Prompt charmonium production at small $p_{\mathcal{T}}$ in the Soft Gluon Resummation approach

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Outline

- Hadronisation model: NRQCD
- Soft Gluon Resummation approach
- InEW, matching of factorisation theorems
- ▶ J/ψ production at small p_T
- ▶ J/ψ polarisation
- Summary

Hadronisation model: NRQCD

• J/ψ wave function as a series with respect to relative constituent quarks velocity v:

$$\begin{split} |J/\psi\rangle &= \mathcal{O}(\upsilon^0) \, |c\bar{c}[{}^3S_1^{(1)}]\rangle + \mathcal{O}(\upsilon^1) \, |c\bar{c}[{}^3P_J^{(8)}]g\rangle + \mathcal{O}(\upsilon^2) \, |c\bar{c}[{}^3S_1^{(1,8)}]gg\rangle + \\ &+ \mathcal{O}(\upsilon^2) \, |c\bar{c}[{}^1S_0^{(8)}]g\rangle + \mathcal{O}(\upsilon^2) \, |c\bar{c}[{}^1D_J^{(1,8)}]gg\rangle + \dots \end{split}$$

- \blacktriangleright Approximate v-scaling due to $v^2\approx 0.2$
- Hard cross section factorisation:

$$d\hat{\sigma}(ab \to \mathcal{C}X) = \sum_{n} d\hat{\sigma}(ab \to c\bar{c}[n]X) \langle \mathcal{O}^{\mathcal{C}}[n] \rangle$$

Nonperturbative (hadronisation) factors:

 $\langle \mathcal{O}^{\mathcal{C}}[n] \rangle$ – long-distance matrix elements (LDME):

color singlet LDMEs — potential models, data for leptonic decay

color octet LDMEs - lattice QCD calculation or experimental data fitting

General remarks on our approximations in calculations of prompt J/ψ production

▶ Direct production: $g + g \rightarrow J/\psi + X$, feed-down contributions from $\psi(2S) \rightarrow J/\psi + X$ and $\chi_{cJ} \rightarrow J/\psi + \gamma$

Prompt = Direct + Feed-down contributions

- We study the direct production & the P-wave feed-down. At $\sqrt{s} = 200$ GeV (PHENIX data), feed-down contribution is estimated at 30%
- We study gluon-gluon fusion & quark-antiquark annihilation, quark-antiquark subprocesses may contribute about 10% to the total cross section at $\sqrt{s} = 200$ GeV (within the Improved Color Evaporation model) [Saleev, Chernyshev, 2022]
- Our preliminary calculations were done in the LO approximation of the pQCD in α_S

TMD factorisation and initial parton transverse momenta

Fransverse Momentum Dependent (TMD) factorization: $q_T, k_T \ll \mu_F \sim M$,

TMD parton distribution functions $F(x, q_T, \mu_F, \zeta) \Rightarrow$ two-scale **Collins-Soper** equations,

$$q_1^{\mu} = x_1 p_1^{\mu} + y_1 p_2^{\mu} + q_{1T}^{\mu}, \qquad q_2^{\mu} = x_2 p_2^{\mu} + y_2 p_1^{\mu} + q_{2T}^{\mu},$$

▶ Preserving $\mathcal{O}(q_T/M)$ terms, neglecting $\mathcal{O}(q_T^2/M^2)$ terms and, therefore, assuming $y_{1,2} \rightarrow 0$:

$$q_1 \approx \left(\frac{x_1\sqrt{s}}{2}, \boldsymbol{q_{1T}}, \frac{x_1\sqrt{s}}{2}\right), \quad q_2 \approx \left(\frac{x_2\sqrt{s}}{2}, \boldsymbol{q_{2T}}, -\frac{x_2\sqrt{s}}{2}\right),$$

▶ Relevant processes only $2 \rightarrow 1$, intermediate $c\bar{c}$ -states can be

- color-octet $^1S_0^{(8)}$, $^3P_{0,2}^{(8)}$ and $^3S_1^{(8)}$ for J/ψ
- color-sinlget $^3P_{0,2}^{(1)}$ and color-octet $^3S_1^{(8)}$ for χ_{cJ}

TMD factorisation and TMD PDFs

General formula of TMD factorization [TMD Handbook, arXiv:2304.03302]:

$$\frac{d\sigma}{d\boldsymbol{p_T}} = \sigma_0 \int d\boldsymbol{q_{1T}} \, d\boldsymbol{q_{2T}} \, F(x_1, \boldsymbol{q_{1T}}, \boldsymbol{\mu_F}, \zeta_1) \, F(x_2, \boldsymbol{q_{1T}}, \boldsymbol{\mu_F}, \zeta_2) \, \delta(\boldsymbol{q_{1T}} + \boldsymbol{q_{2T}} - \boldsymbol{p_T})$$

 \blacktriangleright To implement CS evolution, the transfer to impact parameter b_T space by 2D Fourier transform is done:

$$\frac{d\sigma}{d\boldsymbol{p_T}} = \sigma_0 \int \frac{d\boldsymbol{b_T}}{(2\pi)^2} e^{i\boldsymbol{p_T}\boldsymbol{b_T}} \, \hat{F}(x_1, \boldsymbol{b_T}, \boldsymbol{\mu_F}, \zeta_1) \, \hat{F}(x_2, \boldsymbol{b_T}, \boldsymbol{\mu_F}, \zeta_2)$$

• σ_0 is calculated as series in small α_S

Soft Gluon Resummation approach

Soft and collinear gluon resummation approach by [J. Collins, D. Soper, 1981]:

$$\frac{d\sigma(J/\psi)}{d\boldsymbol{p_T}} = \sigma_0 \int_0^\infty db_T \, b_T \, J_0(\boldsymbol{p_T} \boldsymbol{b_T}) \, e^{-S_P(\boldsymbol{b_T}, \boldsymbol{\mu_F}, Q)} \, e^{-S_{NP}(\boldsymbol{b_T})} \, \hat{F}(\boldsymbol{x_1}, \boldsymbol{\mu_{b^*}'}, \boldsymbol{b_T'}) \, \hat{F}(\boldsymbol{x_2}, \boldsymbol{\mu_{b^*}'}, \boldsymbol{b_T'}) \, \hat{$$

Sudakov factor in LL-LO perturbative calculations [J. Collins, D. Soper (1982)]:

$$S_{P}(b_{T},\mu_{F},Q) = \frac{C_{A}}{\pi} \int_{\mu_{b}^{2}}^{Q^{2}} \frac{d\mu'^{2}}{\mu'^{2}} \alpha_{s}(\mu') \left[\ln \frac{Q^{2}}{\mu'^{2}} - \left(\frac{11 - 2N_{f}/C_{A}}{6}\right) \right] + \mathcal{O}(\alpha_{s})$$

Sudakov factor expression is valid only on region $b_0/Q \leq b_T \leq b_{T, \max}$ which is being controlled with [D. Boer, W. J. den Dunnen (2014); J. Collins, D. Soper, G. Sterman (1985)]

$$\mu_b \to \mu_b' = \frac{Q b_0}{Q b_T + b_0} \quad \text{and} \quad b_T^*(b_T) = \frac{b_T}{\sqrt{1 + (b_T/b_{T,\,\text{max}})^2}}$$

Soft Gluon Resummation approach

Master formula for soft gluon resummation:

$$\frac{d\sigma(J/\psi)}{d\boldsymbol{p_T}} = \sigma_0 \int_0^\infty db_T \, b_T \, J_0(\boldsymbol{p_T} b_T) \, e^{-S_P(b_T, \mu_F, Q)} \, e^{-S_{NP}(b_T)} \, \hat{F}(x_1, \mu_{b^*}', b_T^*) \, \hat{F}(x_2, \mu_{b^*}', b_T^*)$$

Nonperturbative quark factor obtained in SIDIS data fitting: [S. Aybat, T. Rogers (2011)]:

$$S_{NP}(b_T, Q) = \left[g_1 \ln \frac{Q}{2Q_{NP}} + g_2 \left(1 + 2g_3 \ln \frac{10xx_0}{x_0 + x}\right)\right] b_T^2$$

- it should be Casimir-scaled by ${\cal C}_A/{\cal C}_F$ for gluons

 \blacktriangleright In the leading order of α_S , the perturbative tail of TMD PDF is expressed with collinear PDF:

$$\hat{F}(x,\mu_{b^*}',b_T^*) = f(x,\mu_{b^*}') + \mathcal{O}(\alpha_S) + \mathcal{O}(b_T\Lambda_{\mathsf{QCD}})$$

Matching of small- p_T and high- p_T regions within Inverse-Error Weighting Scheme

Matched cross-section as a weighed sum of CPM and TMD terms [M. Echevarria, T. Kasemets, J.-P. Lansberg, C. Pisano, A. Signori (2018)]:

$$d\sigma = \mathcal{W} \, d\sigma^{\mathsf{TMD}} + \mathcal{Z} \, d\sigma^{\mathsf{CPM}}$$

Normalised weights for each of the two terms:

$$\begin{split} \mathcal{W} &= \frac{\Delta \mathcal{W}^{-2}}{\Delta \mathcal{W}^{-2} + \Delta \mathcal{Z}^{-2}}, \qquad \mathcal{Z} = \frac{\Delta \mathcal{Z}^{-2}}{\Delta \mathcal{W}^{-2} + \Delta \mathcal{Z}^{-2}}\\ \Delta \mathcal{W} &= \left(\frac{p_T}{Q}\right)^2 + \left(\frac{m}{Q}\right)^2, \qquad \Delta \mathcal{Z} = \left(\frac{M}{p_T}\right)^2 \left(1 + \ln^2 \frac{Q_T}{p_T}\right) \end{split}$$

Uncertainty due to the matching procedure:

$$\Delta d\sigma = \frac{d\sigma}{\sqrt{\Delta \mathcal{W}^{-2} + \Delta \mathcal{Z}^{-2}}} = \frac{\Delta \mathcal{W} \cdot \Delta \mathcal{Z}}{\sqrt{\Delta \mathcal{W}^{2} + \Delta \mathcal{Z}^{2}}} d\sigma$$

Extraction of CO LDME at $p_T < 1~{\rm GeV}~({\rm SGR})$ and at $p_T > 5~{\rm GeV}~({\rm CPM})$



Extraction of CO LDME at $p_T < 1$ GeV (SGR) and at $p_T > 5$ GeV (CPM)



• χ_{cJ} decays contribution $\approx 7\%$

Extraction of CO LDME at $p_T < 1~{\rm GeV}$ (SGR) and at $p_T > 5~{\rm GeV}$ (CPM)

CO LDME	Our fit	LO CPM [Cho, Leibovich (1996)]	NLO CPM, global fit [Butenschön, Kniehl (2011)]
$\langle \mathcal{O}^{J/\psi} [{}^1S_0^{(8)}] angle$, GeV 3	$(9.7 \pm 0.5) \cdot 10^{-2}$	_	$(3.0 \pm 0.4) \cdot 10^{-2}$
$\langle \mathcal{O}^{J/\psi}[^{3}P_{0}^{(8)}] angle$, GeV 5	$(1.3 \pm 0.2) \cdot 10^{-2}$	_	$(-9.1 \pm 1.6) \cdot 10^{-3}$
$M_3^{J/\psi}$, GeV 3	$(1.1 \pm 0.1) \cdot 10^{-1}$	$(6.6 \pm 1.5) \cdot 10^{-2}$	$(1.8 \pm 0.6) \cdot 10^{-2}$
$\langle \mathcal{O}^{J/\psi}[^3S_1^{(8)}] angle$, GeV 3	$(2.0 \pm 1.6) \cdot 10^{-3}$	$(6.6 \pm 2.1) \cdot 10^{-3}$	$(1.7 \pm 0.5) \cdot 10^{-3}$
$\langle \mathcal{O}^{\chi_{c0}}[^3S_1^{(8)}]\rangle$, GeV^3	$(8.6 \pm 2.9) \cdot 10^{-3}$	$(3.3 \pm 0.5) \cdot 10^{-3}$	_
$\chi^2/n.d.f.$	0.76	0.9	3.74

Polarisation test for CO LDME at PHENIX



Predictions for SPD NICA using the Soft Gluon Resummation approach



• quark-antiquark annihilation contribution < 4%

• χ_{cJ} decays contribution $\approx 9\%$

Summary

- \blacktriangleright We have used the Soft Gluon Resummation approach to calculate small- $p_T~J/\psi$ production in the TMD factorisation
- > CO LDMEs of NRQCD are necessary to describe J/ψ production using the TMD and CPM factorisations, where they are major contributions
- ▶ Soft Gluon Resummation approach for gluon and quark TMD PDF satisfyingly describe experimental data for unpolarised and polarised J/ψ production at $\sqrt{s} = 200$ GeV in the TMD domain of $p_T < 1$ GeV
- We have tested the factorisation matching scheme InEW for description of intermediate p_T domain, calculation agrees with experimental data within the uncertainty band
- We estimate the perspective region for the extraction of gluon TMD PDF in the J/ψ production as $p_T \leqslant 1$ GeV