

Manifestation of clustering in fission of heavy nuclei

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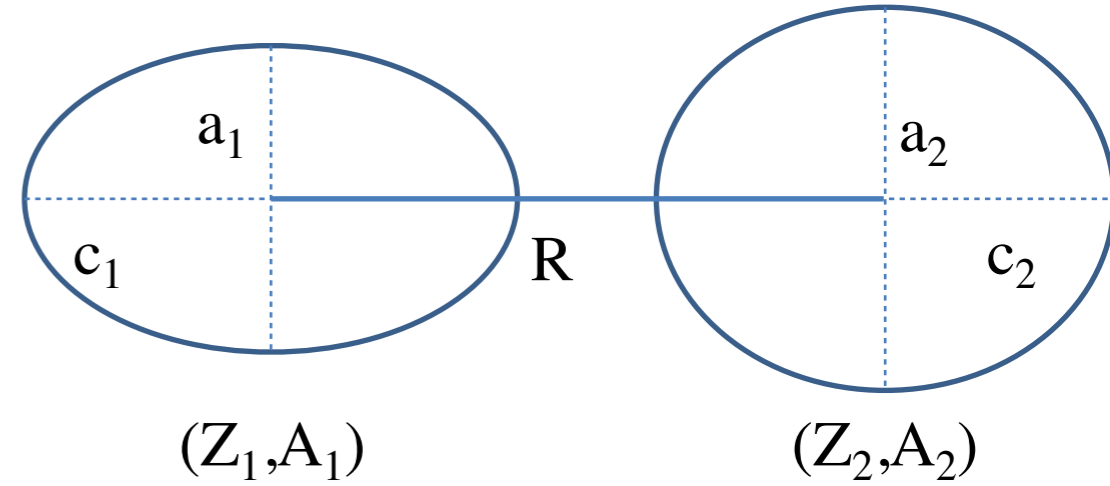
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Contents

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- *Description of the model*
- *Results:*
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Model

Parameters $Z_1, A_1, \beta_1, \beta_2$ and R completely describe the geometry of the system.



The interaction potential between the fragments: $V(R, Z, A, J, \beta_1, \beta_2) = V_N + V_C$

$$V_C(R, Z_1, Z_2, \beta_1, \beta_2) = \frac{e^2 Z_1 Z_2}{R} + \left(\frac{9}{20\pi}\right)^{1/2} \frac{e^2 Z_1 Z_2}{R^3} \sum_{i=1}^2 R_i^2 \beta_i \left[1 + \frac{2}{7} \left(\frac{5}{\pi}\right)^{1/2} \beta_i \right] P_2(\cos\theta_i)$$

$$V_N = \int \rho_1(r_1) \rho_2(R - r_2) F(r_1 - r_2) dr_1 dr_2$$

$$F(r_1 - r_2) = C_0 \left[F_{in} \frac{\rho_1(r_1)}{\rho_{00}} + F_{ex} \left(1 - \frac{\rho_1(r_1)}{\rho_{00}} \right) \right] \delta(r_1 - r_2)$$

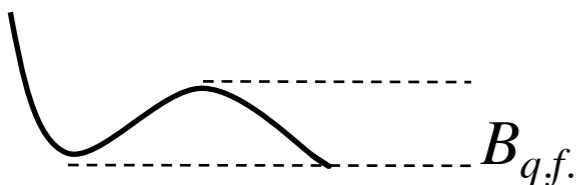
$$\rho_0(r) = \rho_1(r) + \rho_2(R - r)$$

$$F_{in,ex} = f_{in,ex} + f'_{in,ex} \frac{(N - Z)(N_2 - Z_2)}{(N + Z)(N_2 + Z_2)}$$

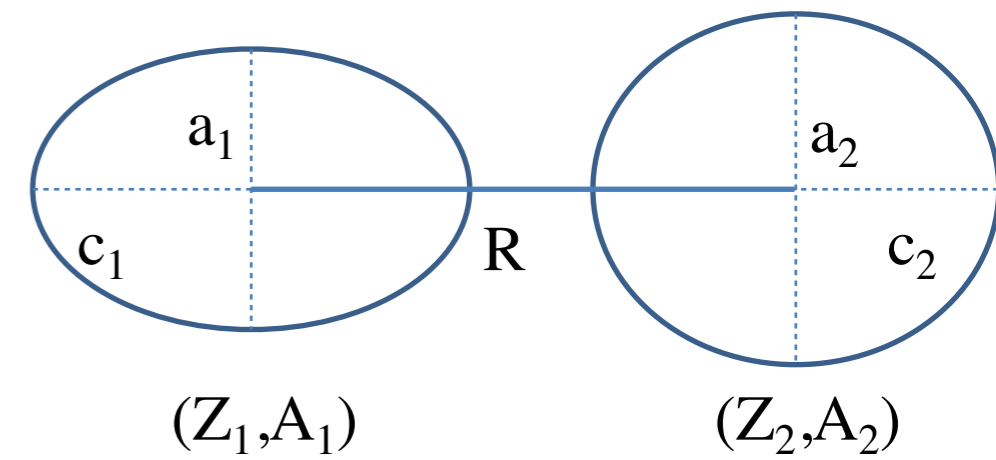
$$C_0 = 300 \text{ MeV fm}^3$$

$$f_{in} = 0.09, f_{ex} = -2.59$$

$$\rho_{00} = 0.17 \text{ fm}^{-3}, a = 0.51 - 0.56 \text{ fm}$$



Model



- The total energy:

$$\begin{aligned}
 U(A_i, Z_i, \beta_i, d) &= \\
 &= U_{macro}(A_i, Z_i, \beta_i, d) + \delta U^{shell}(A_i, Z_i, \beta_i) = \\
 &= \sum_{i=1,2} U_i^{LD}(A_i, Z_i, \beta_i) + \sum_{i=1,2} \delta U_i^{shell}(A_i, Z_i, \beta'_i, E_i^*) + \\
 &+ V_N(A_i, Z_i, \beta_i, d) + V_C(A_i, Z_i, \beta_i, d).
 \end{aligned}$$

$$U_i^{L.D.}(A_i, Z_i, \beta_i) = U_i^{Surface}(A_i, Z_i, \beta_i) + U_i^C(A_i, Z_i, \beta_i) + U_i^{Sym}(A_i, Z_i)$$

- Liquid drop terms:

$$U_i^{sym}(A_i, Z_i) = 27.612 \frac{(N_i - Z_i)^2}{A_i}$$

$$U_i^C(A_i, Z_i, \beta_i) = \frac{3}{5} \frac{(Z_i e)^2}{R_{0,i}} \frac{\beta_i^{1/3}}{\sqrt{\beta_i^2 - 1}} \ln(\beta_i + \sqrt{\beta_i^2 - 1})$$

Model

Surface energy with variable surface tension:

$$U_i^{Surface}(A_i, Z_i, \beta_i) = \sigma_i S_i$$

$$\sigma_i = \sigma_{0,i}(1 + k_i(\beta_i - \beta_i^{g.s.})^2)$$

$$\sigma_{0,i} = 0.9517(1 - 1.7826((N_i - Z_i)^2)/A_i^2)$$

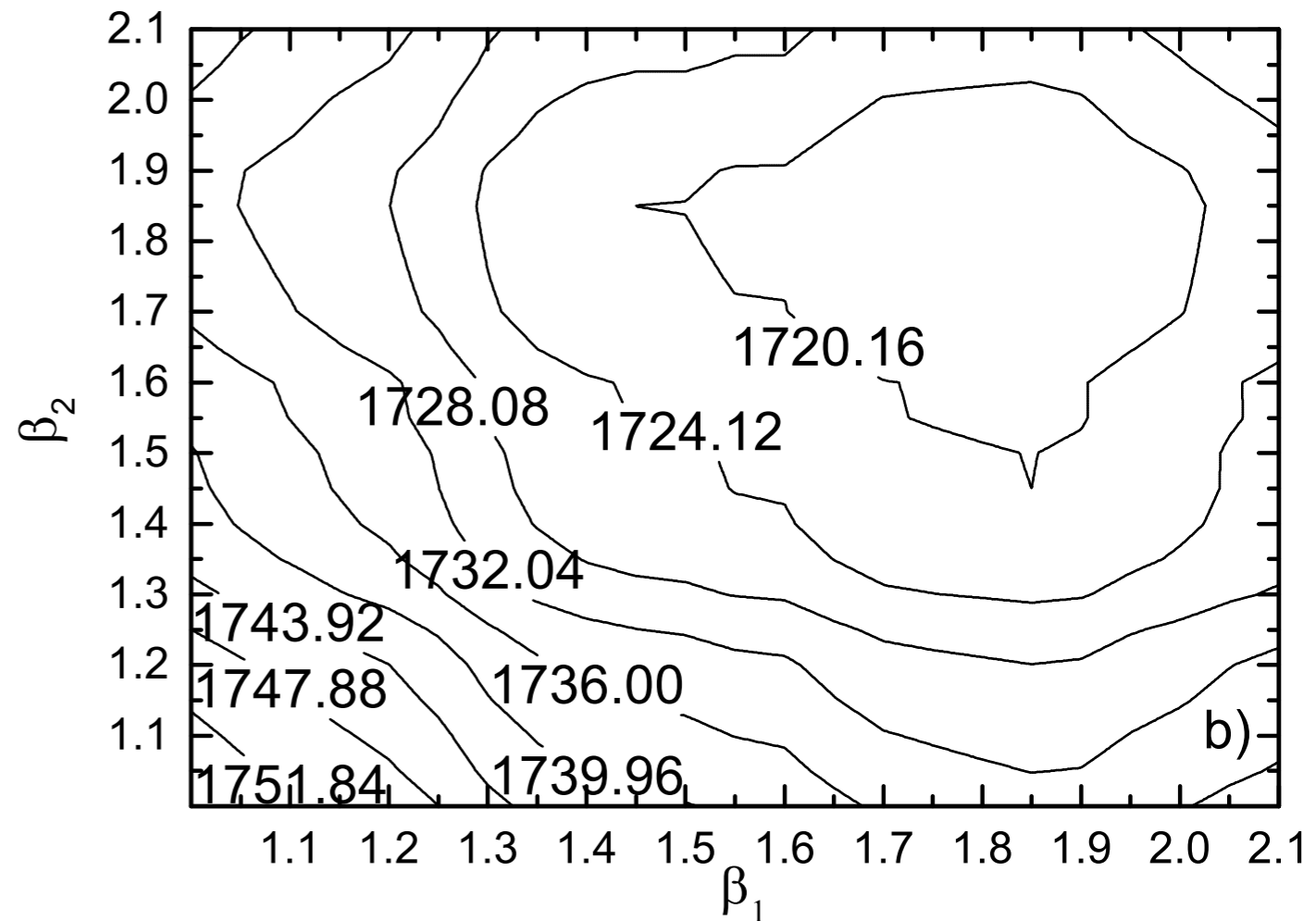
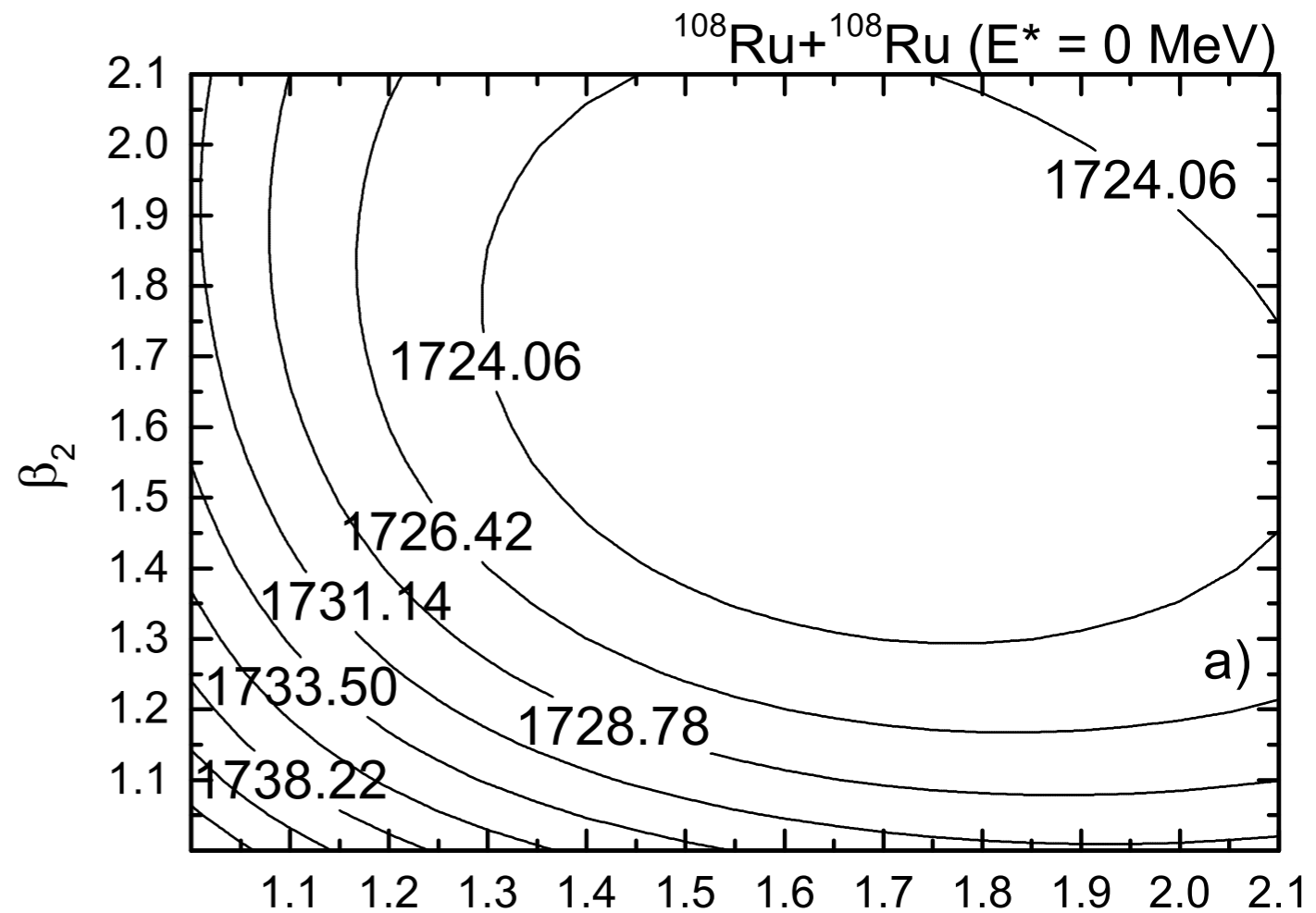
$$k_i = \frac{1}{1 + \exp[-0.063(C_{vib}(Z_i, A_i) - 67)]}$$

$$C_{vib}(Z_i, A_i) = \frac{\hbar\omega_{vib}^i(3Z_i e R_{0,i}^2/(4\pi))^2}{2B(E2)_{vib}^i}$$

$$B(E2)_{vib} = E_{2+}^i B(E2)_{rot}^i / (\hbar\omega_{vib}^i)$$

Shell corrections are calculated as in :

J. Maruhn and W. Greiner, Z. Phys. 251, 431 (1972).



Model

Excitation energy of the scission configuration can be calculated as a sum of the initial excitation energy of the fissioning nucleus and the difference of the potential energies of the fissioning nucleus and scission configuration:

$$E^*(A_i, Z_i, \beta_i, R_m) = E_{CN}^* + [U_{CN}(A, Z, \beta) - U(A_i, Z_i, \beta_i, R_m)].$$

$$T_{DNS}(E^*) = \sqrt{E^*/a}, \quad a = A/12 \text{ MeV}^{-1}$$

Temperature dependence of LD terms:

$$U_i^{sym}(A_i, Z_i, T) = U_i^{sym}(A_i, Z_i, T = 0)(1 + 6 * 10^{-4} E_i^*/A_i),$$

$$U_i^C(A_i, Z_i, \beta_i, T) = U_i^C(A_i, Z_i, \beta_i, T = 0)(1 - 0.12 E_i^*/A_i)$$

$$U_i^{Surf}(A_i, Z_i, \beta_i, T) = U_i^{Surf}(A_i, Z_i, \beta_i, T = 0)(1 + 0.102 E_i^*/A_i).$$

$$k_i(E_i^*) = k_i * \exp[-E_i^*/E_k]$$

Shell damping:

$$\delta U_i^{shell}(A_i, Z_i, \beta_i', E_i^*) = \delta U_i^{shell}(A_i, Z_i, \beta_i', E_i^* = 0) \exp[-E_i^*/E_D]$$

Model

Yields:

Using $P_{Z,A}(E_{CN}^*, \beta_1, \beta_2) \sim \exp\{-U(R_m, Z, A, \beta_1, \beta_2)/T\}$
 $P_{Z,A,\beta_1,\beta_2}^{decay} \sim \exp\{-B_{qf}(Z, A, \beta_1, \beta_2)/T\}$

$$w(A_i, Z_i, \beta_i, E^*) = N_0 \exp\left[-\frac{U(A_i, Z_i, \beta_i, R_m) + B_{qf}(A_i, Z_i, \beta_i)}{T}\right]$$

The different yields can be calculate by integrating over the deformations:

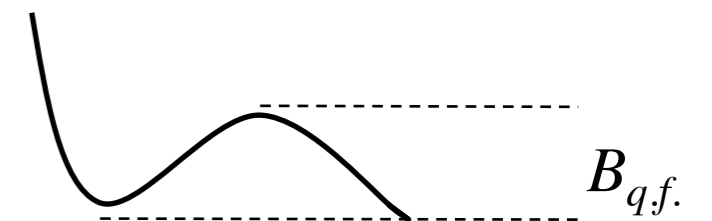
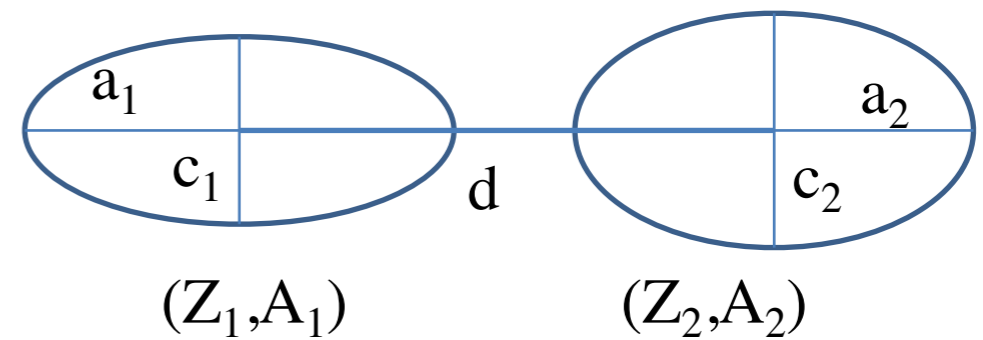
$$Y(A_i, Z_i) = N_0 \int \int w(A_i, Z_i, \beta_1, \beta_2, E^*) d\beta_1 d\beta_2,$$

$$Y(A_i) = N_0 \sum_{Z_i} \int \int w(A_i, Z_i, \beta_1, \beta_2, E^*) d\beta_1 d\beta_2,$$

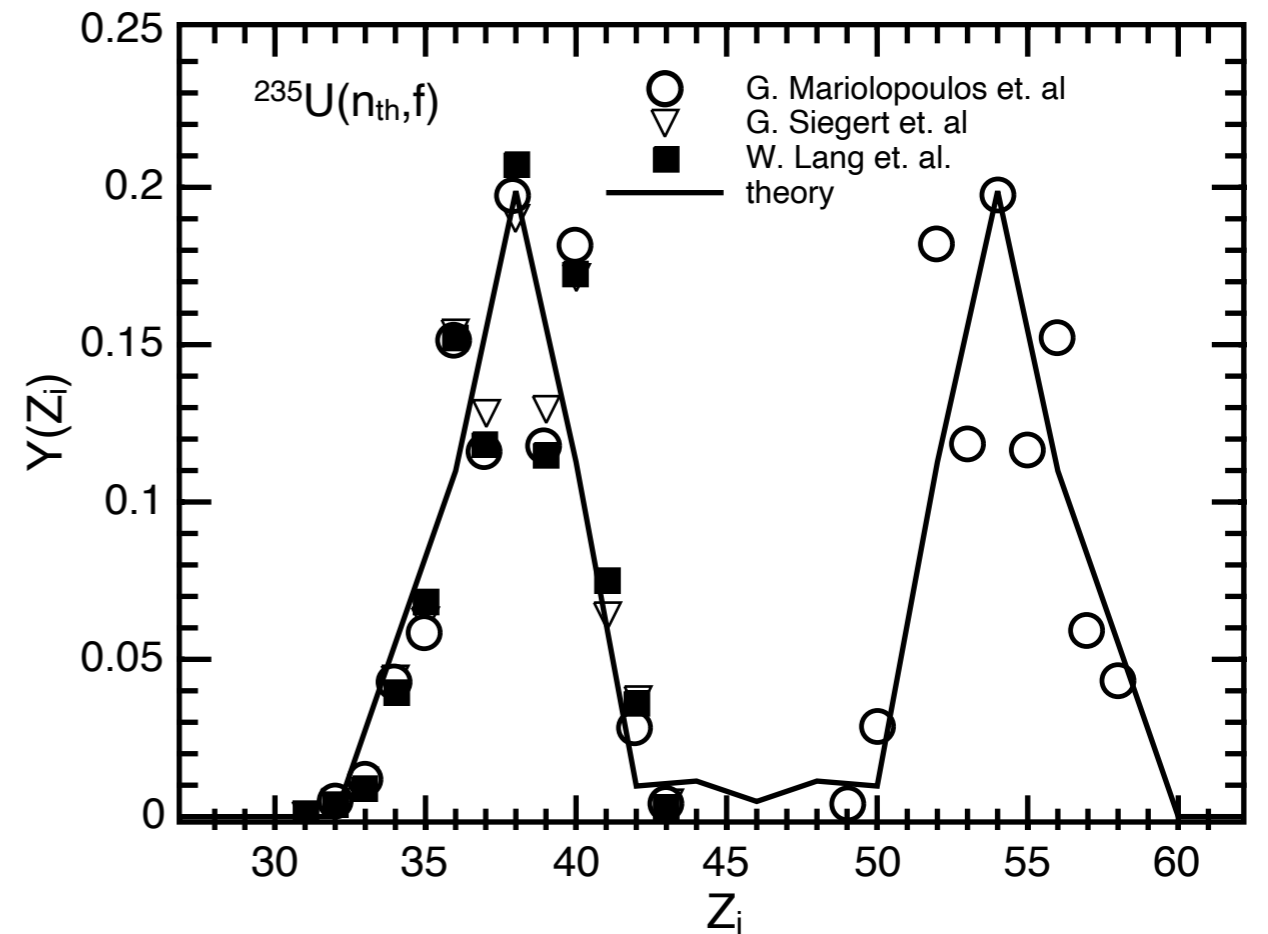
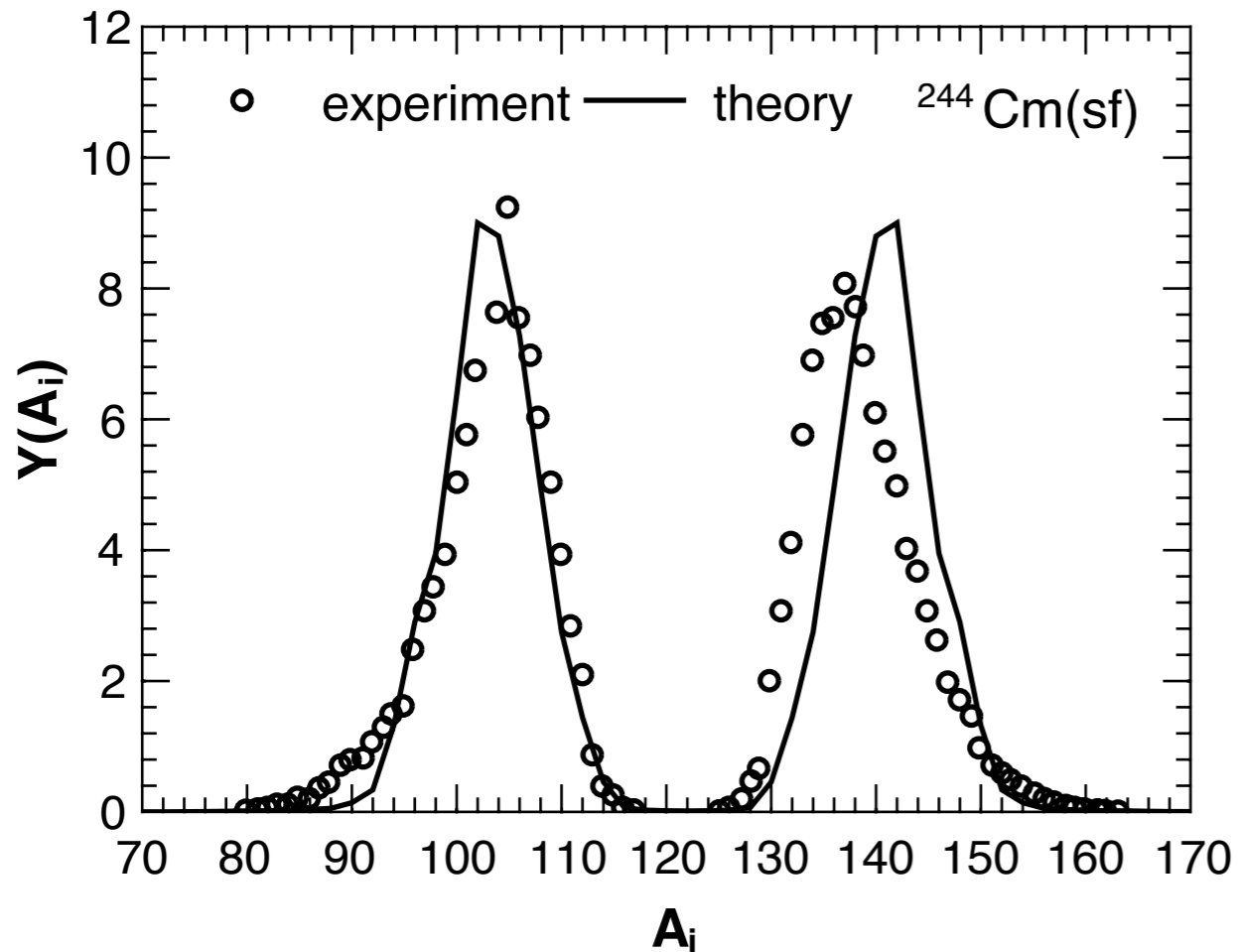
$$Y(Z_i) = N_0 \sum_{A_i} \int \int w(A_i, Z_i, \beta_1, \beta_2, E^*) d\beta_1 d\beta_2,$$

$$TKE(A_i, Z_i) = V_c(A_i, Z_i, \beta_1, \beta_2) + V_n(A_i, Z_i, \beta_1, \beta_2)$$

$$\langle TKE \rangle (A_i) = \frac{\sum_{Z_i} \int TKE(A_i, Z_i, \beta_1, \beta_2) w(A_i, Z_i, \beta_1, \beta_2, E^*) d\beta_1 d\beta_2}{\sum_{Z_i} \int w(A_i, Z_i, \beta_1, \beta_2, E^*) d\beta_1 d\beta_2}$$

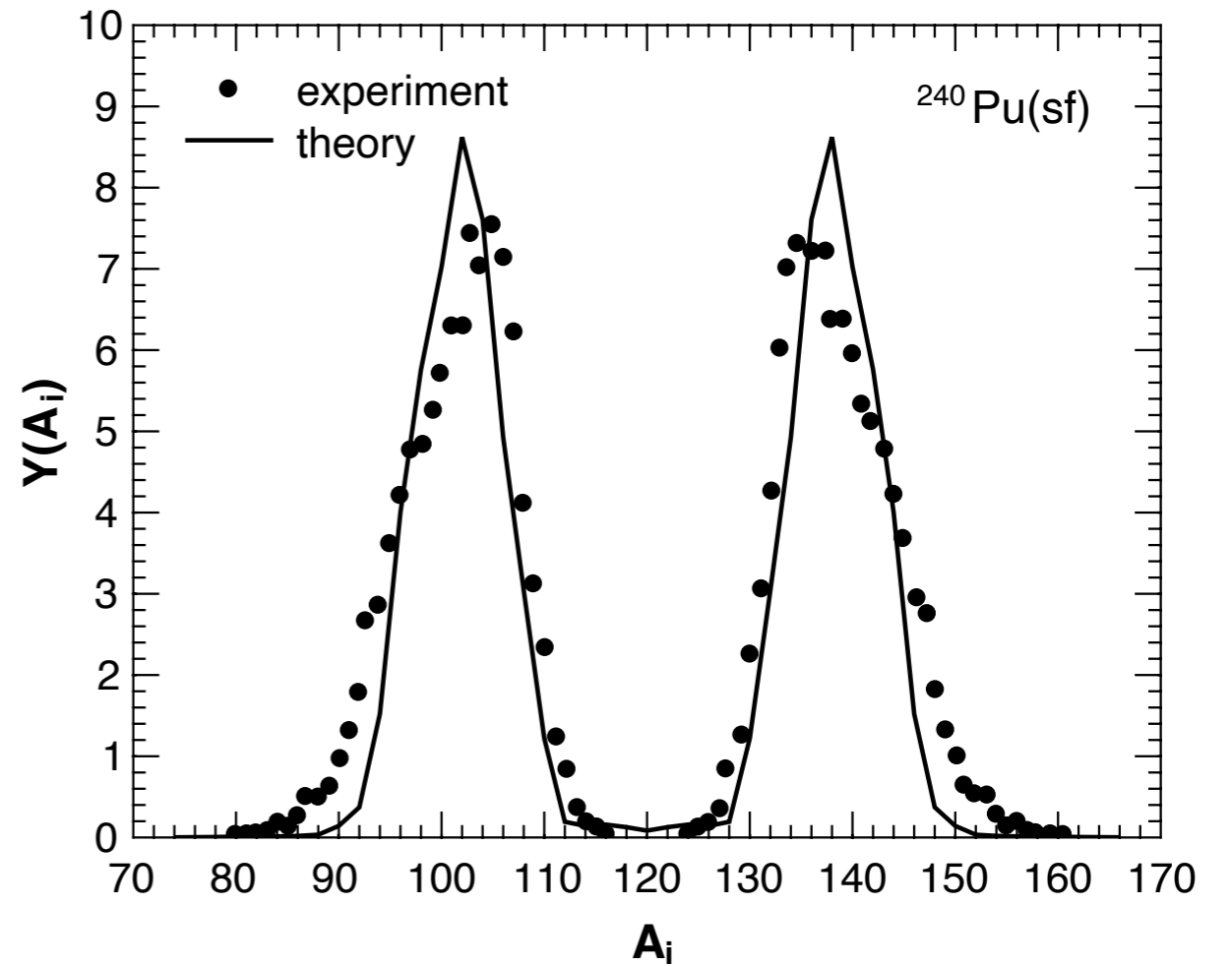
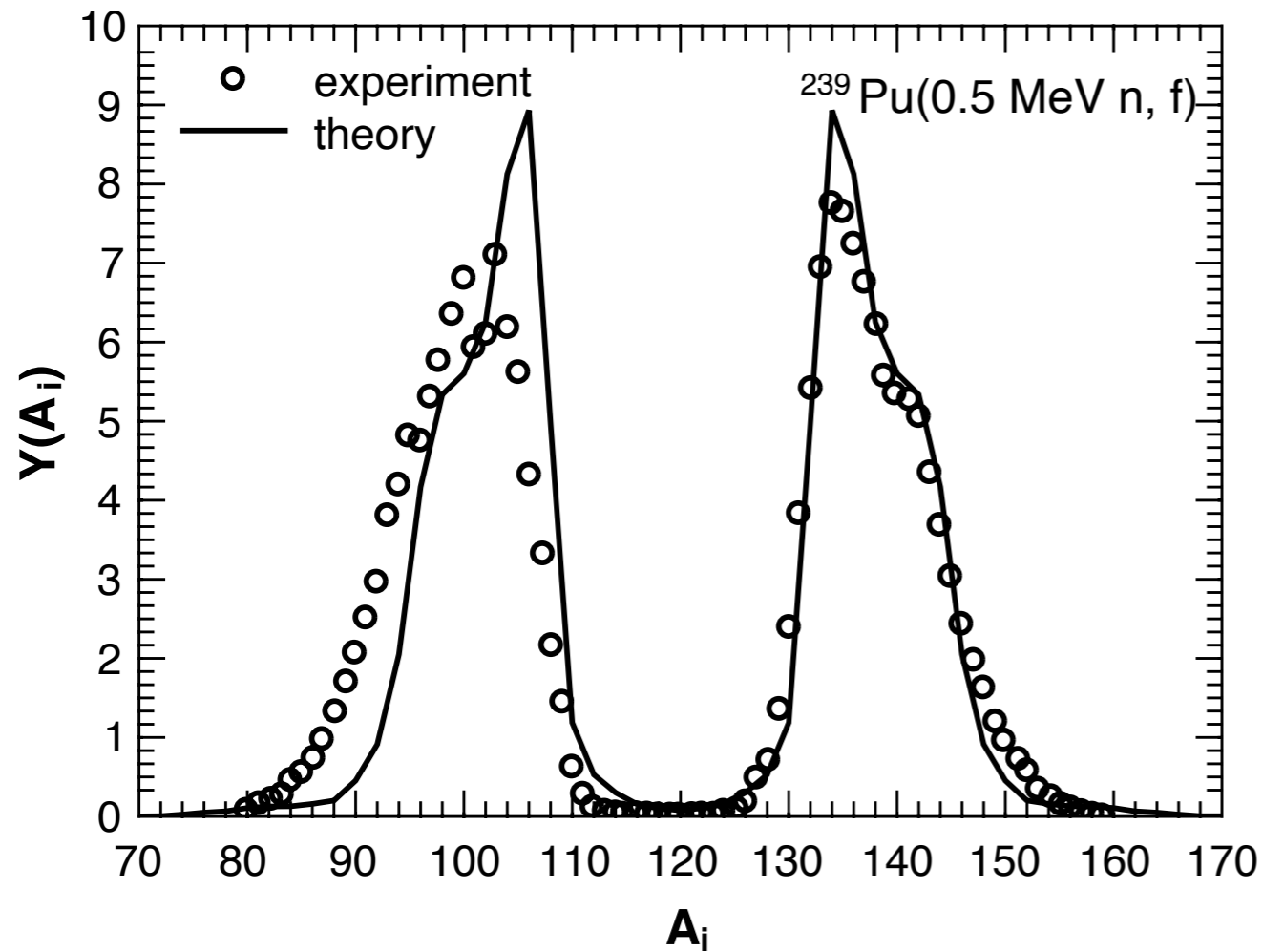


Results



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- G. Siegert, H. Wollnik, J. Grief, R. Decker, G. Fiedler, and B. Pfeiffer, Phys.Rev.C 14, 1864 (1976);
- G. Mariolopoulos, Ch. Hamelin, J. Blachot, J.P. Boucqet, B. Brissot, J. Crancon, H. Nifenecker, and Ch. Ristori, Nucl. Phys. A 361, 213 (1981).
- T.R. England and B.F. Rider, LANL report no. LA-UR-94-3106 (1994).
- P. Schillebeeckx, C. Wagemans, A.J. Deruytter, and R. Barthelemy, Nucl. Phys. A 545, 623 (1992).
- R. Schmidt and H. Henschel, Nucl. Phys. A 395, 15 (1983).

Results

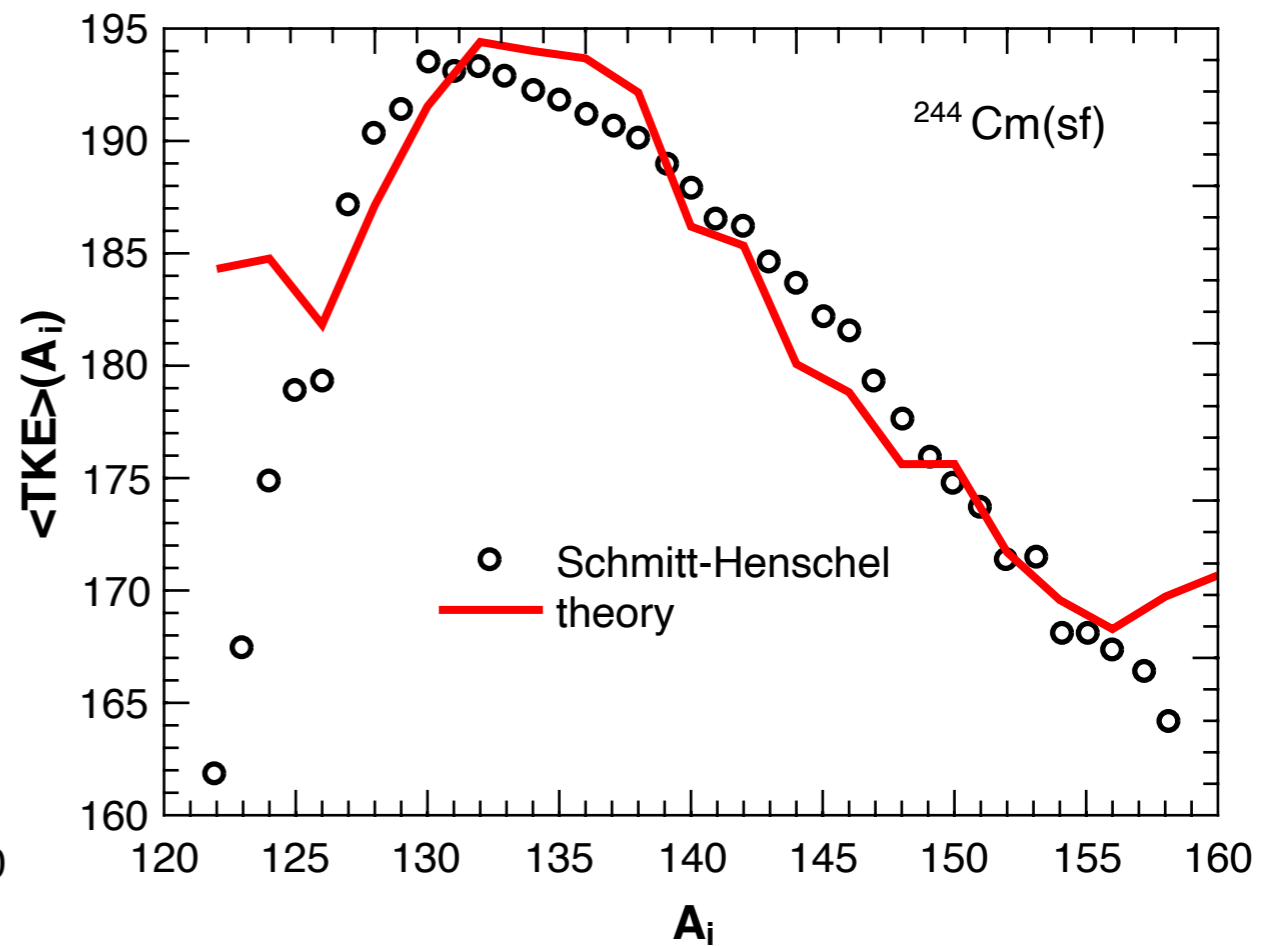
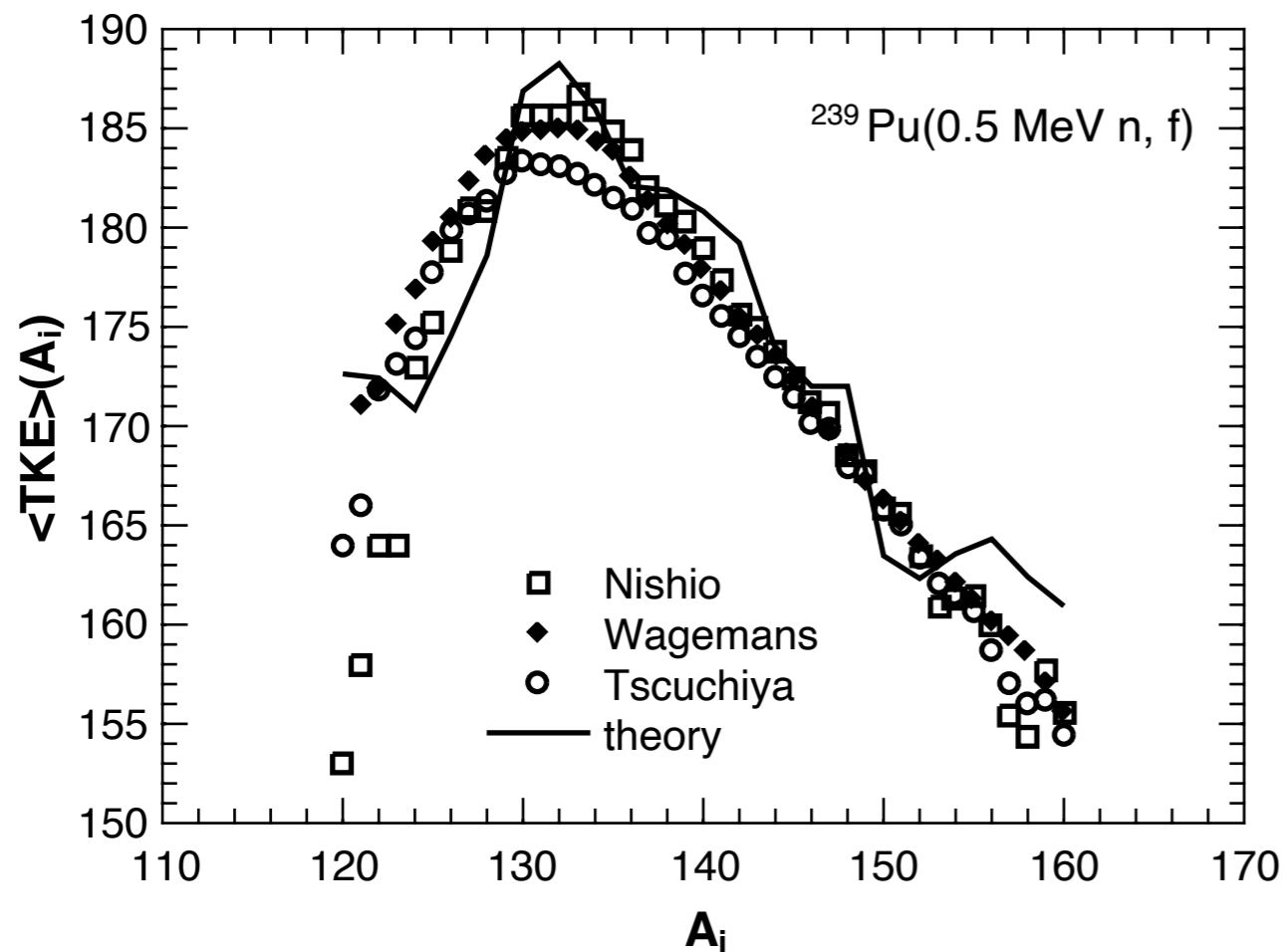
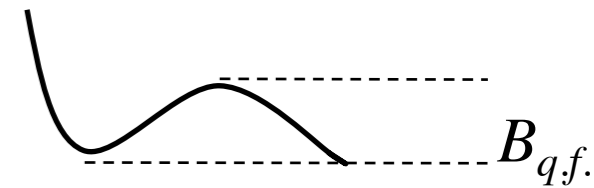


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- T.R. England and B.F. Rider, LANL report no. LA-UR-94-3106 (1994).
- P. Schillebeeckx, C. Wagemans, A.J. Deruytter, and R. Barthelemy, Nucl. Phys. A 545, 623 (1992).

Results (and more modelling)

$$\langle TKE \rangle(A_i) = \frac{\sum_{Z_i} \int d\beta_L d\beta_H TKE(A_i, Z_i, \beta_i) w(A_i, Z_i, \beta_i, E^*)}{\sum_{Z_i} \int d\beta_L d\beta_H w(A_i, Z_i, \beta_i, E^*)}$$

$$TKE(\{A_i, Z_i, \beta_i\}) = V^N(\{A_i, Z_i, \beta_i\}) + V^C(\{A_i, Z_i, \beta_i\})$$



- R. Schmidt and H. Henschel, Nucl. Phys. A 395, 15 (1983).
- C. Wagemans, E. Allaert, A. Deruytter, R. Barthelemy, and P. Schillebeeckx, Phys. Rev. C 30, 218 (1984).
- T.R. England and B.F. Rider, LANL report no. LA-UR-94-3106 (1994).

Results (and more modelling)

The probability of emitting exactly x neutrons (Jackson model):

$$P(x) = 1 - e^{-\Delta_x} \left(1 + \sum_{k=1}^{2x-3} \frac{(\Delta_x)^k}{k!} \right)$$

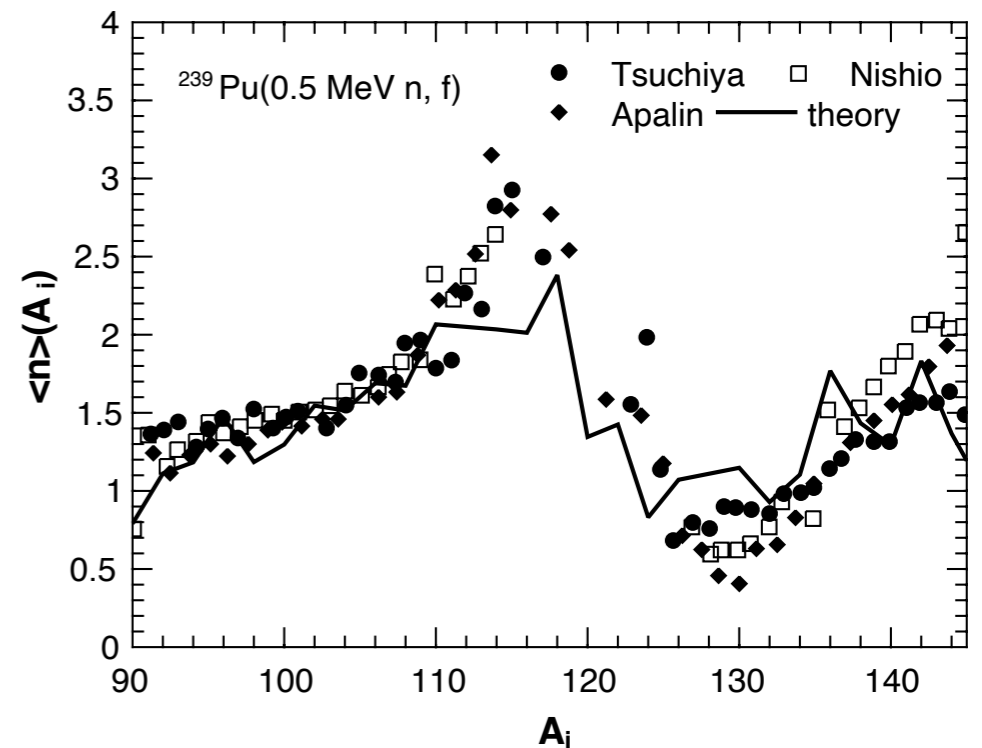
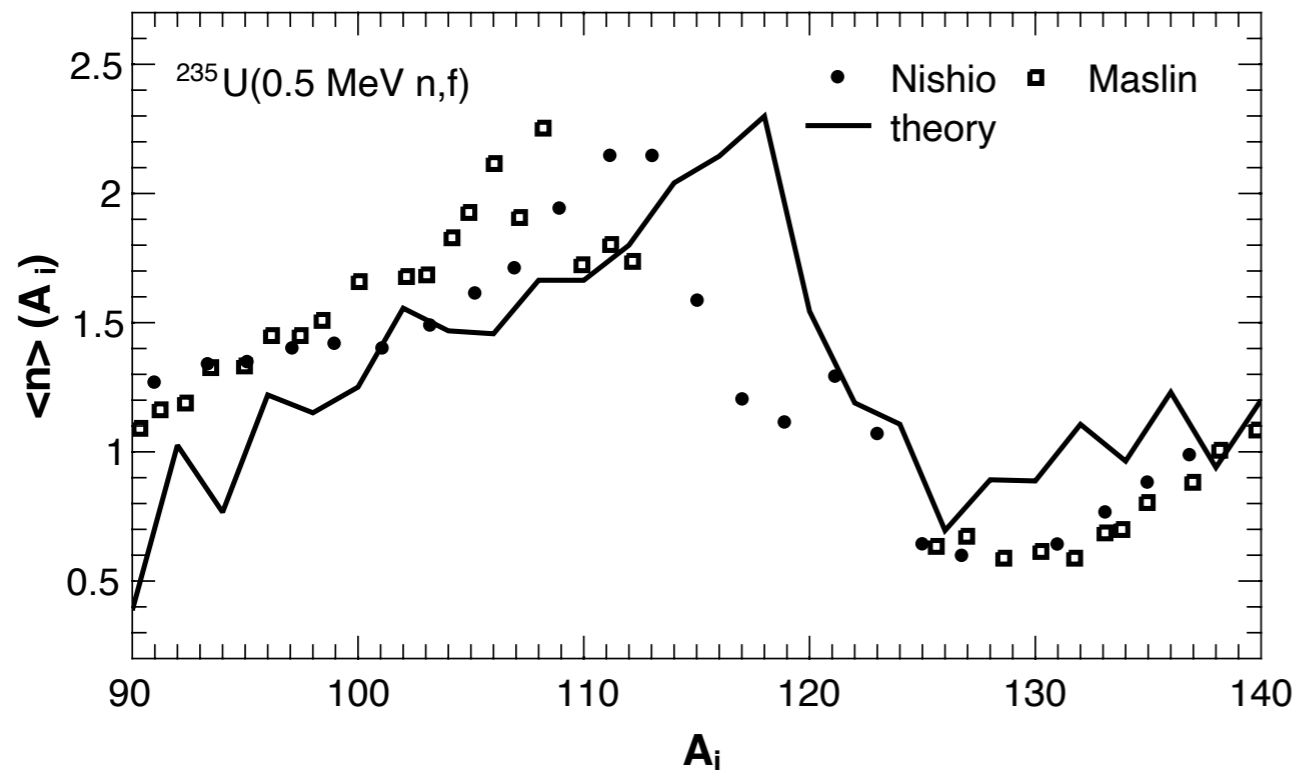
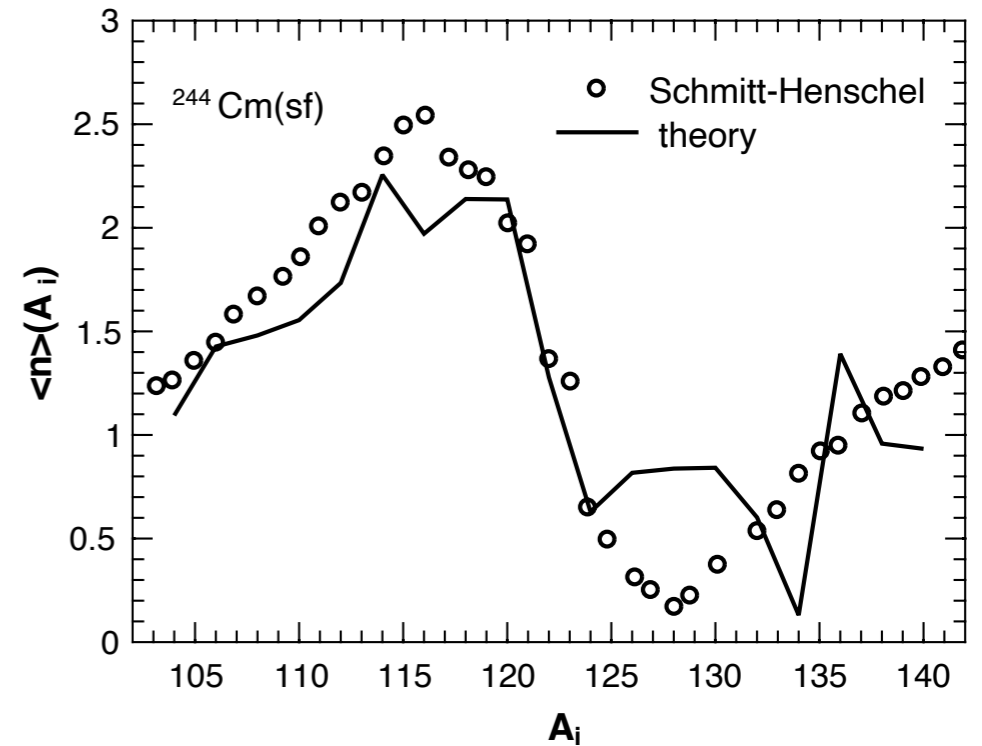
$$P_{xn}(E_i^*) = P(x) - P(x+1)$$

$$P(x+1) = 1 - e^{-\Delta_{x+1}} \left(1 + \sum_{k=1}^{2x-1} \frac{(\Delta_{x+1})^k}{k!} \right)$$

$$\Delta_x = \left(E_i^* - \sum_{k=1}^x B_k \right) / T_i$$

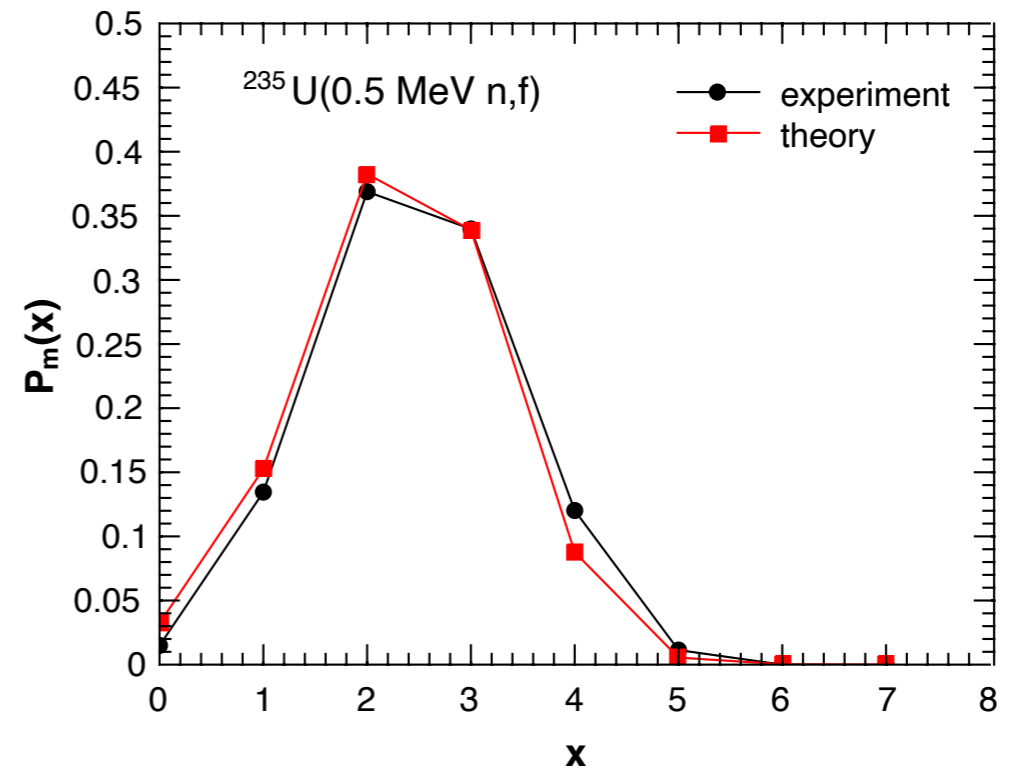
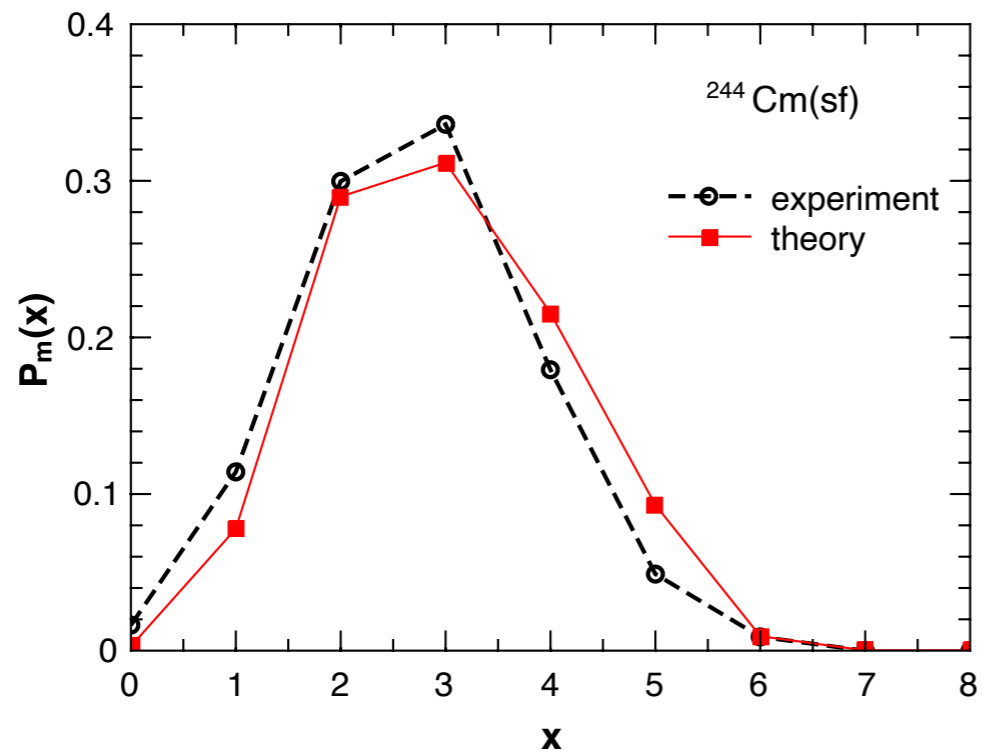
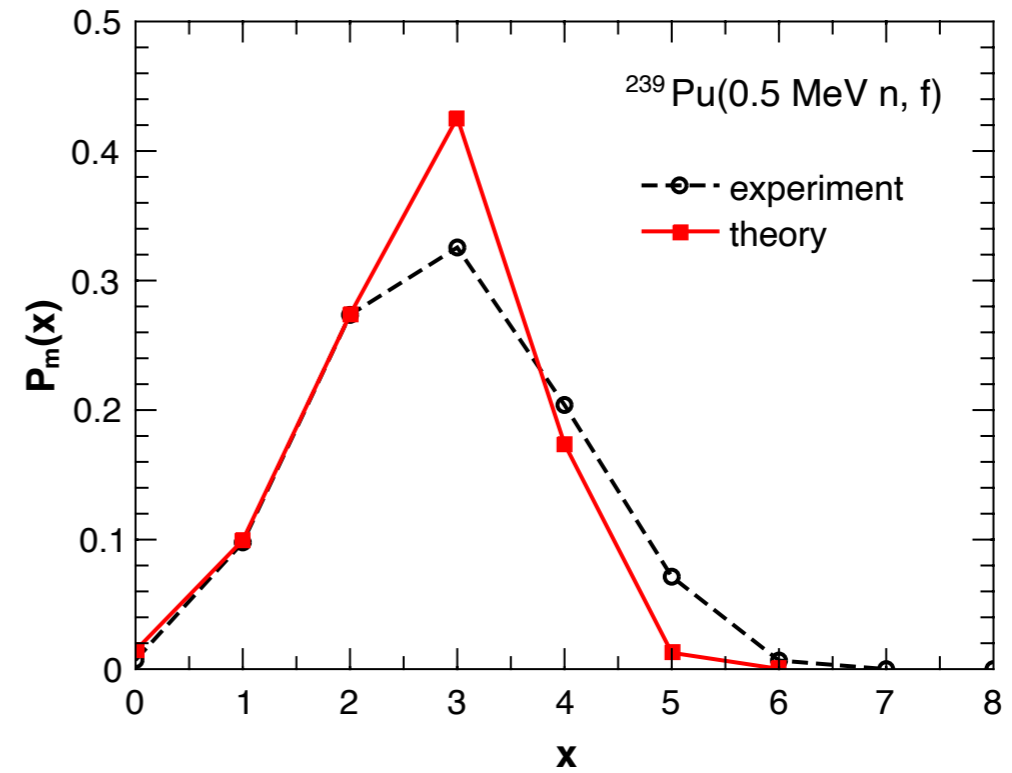
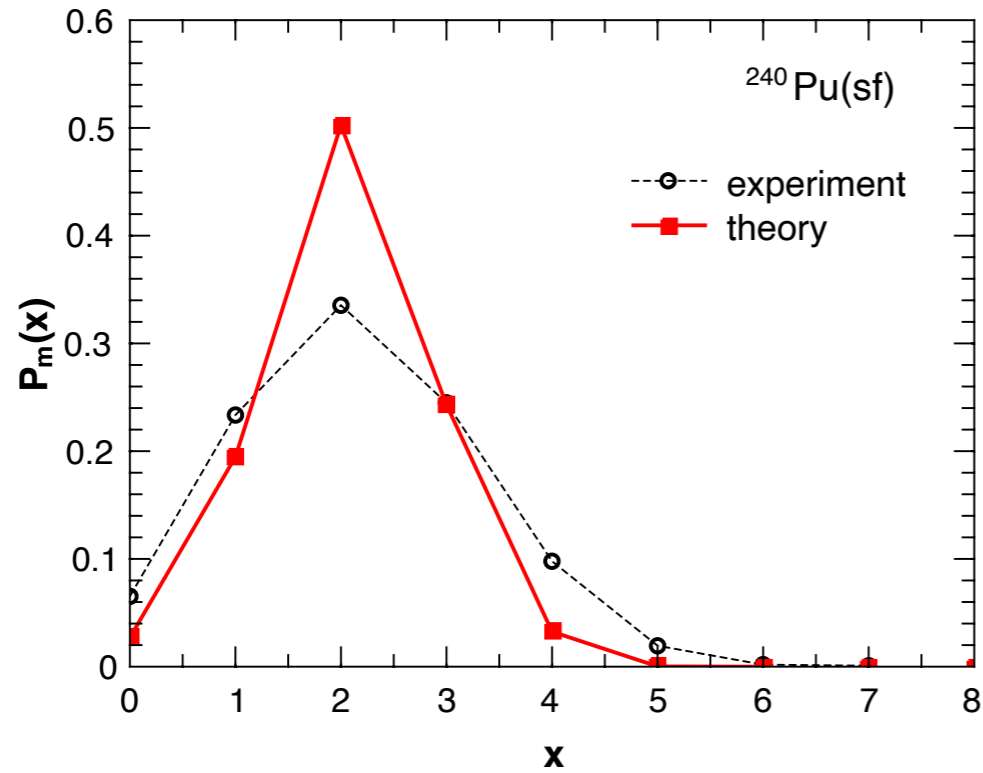
$$\rho_L(E_L^*) \sim \rho_L(E_L^*) \rho_H(E^* - E_L^*), \quad \rho_i(E_i^*) \sim \exp[2(a_i E_i^*)^{1/2}]$$

$$P_L(x) = \frac{\int_0^{E^*} dE_L^* P_{xn}^L(E_L^*) \rho_L(E_L^*) \rho_H(E^* - E_L^*)}{\int_0^{E^*} dE_L^* \rho_L(E_L^*) \rho_H(E^* - E_L^*)}$$



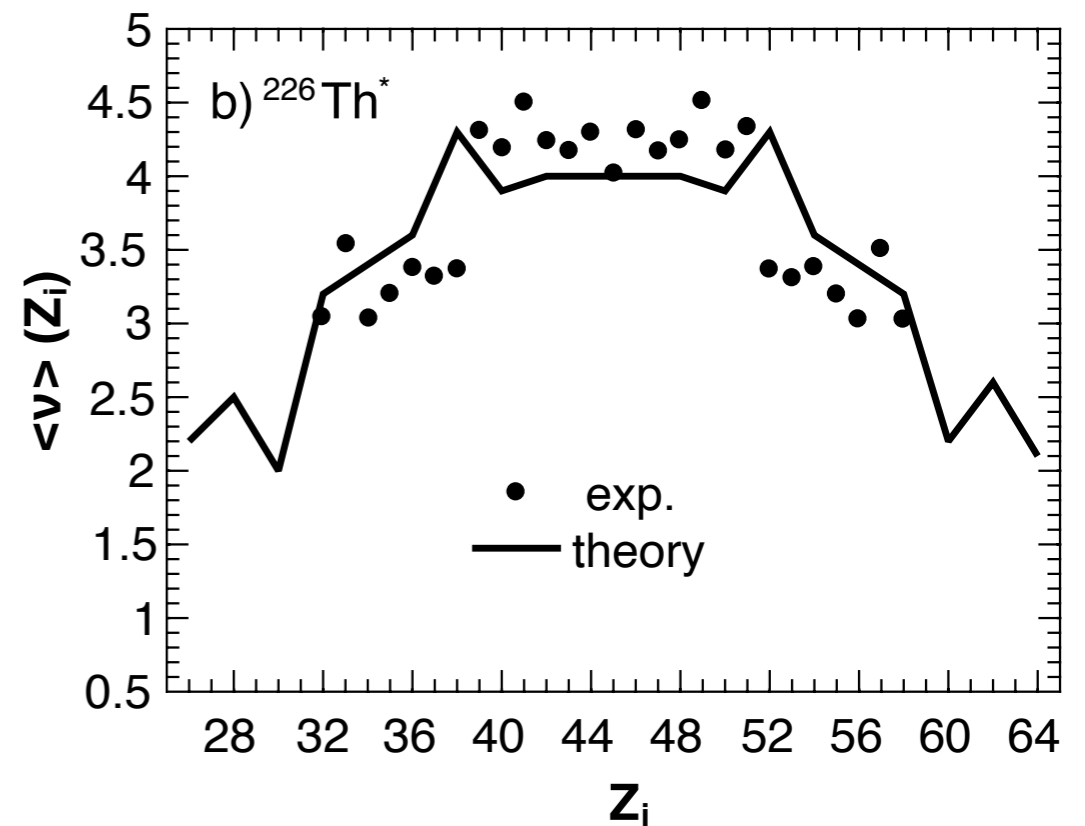
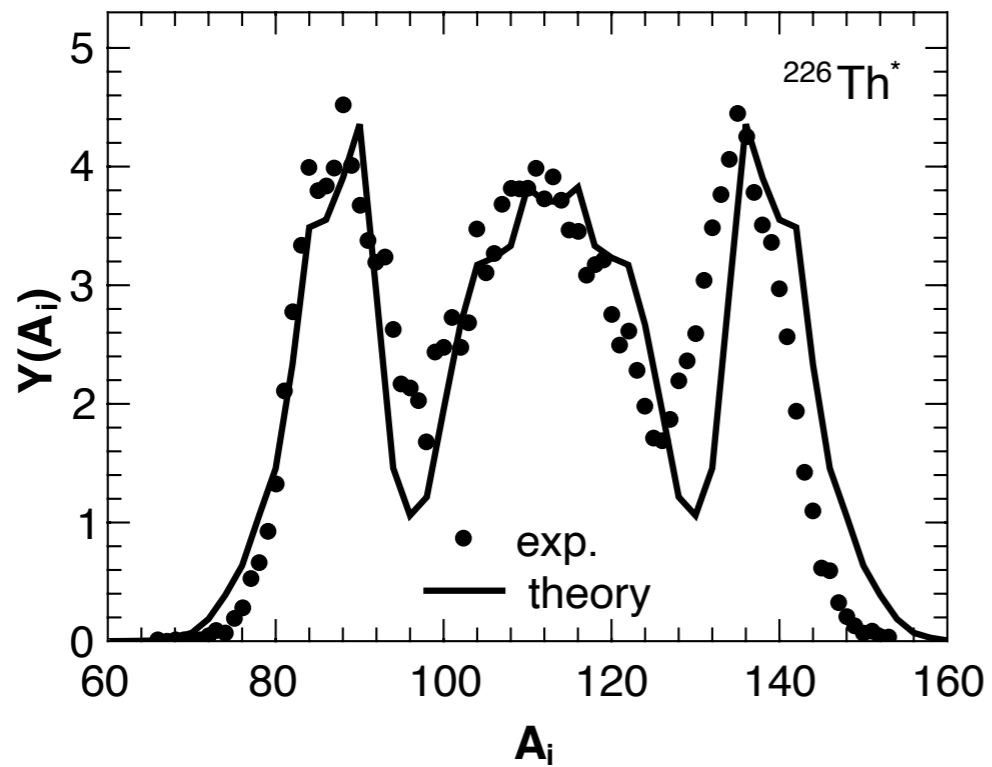
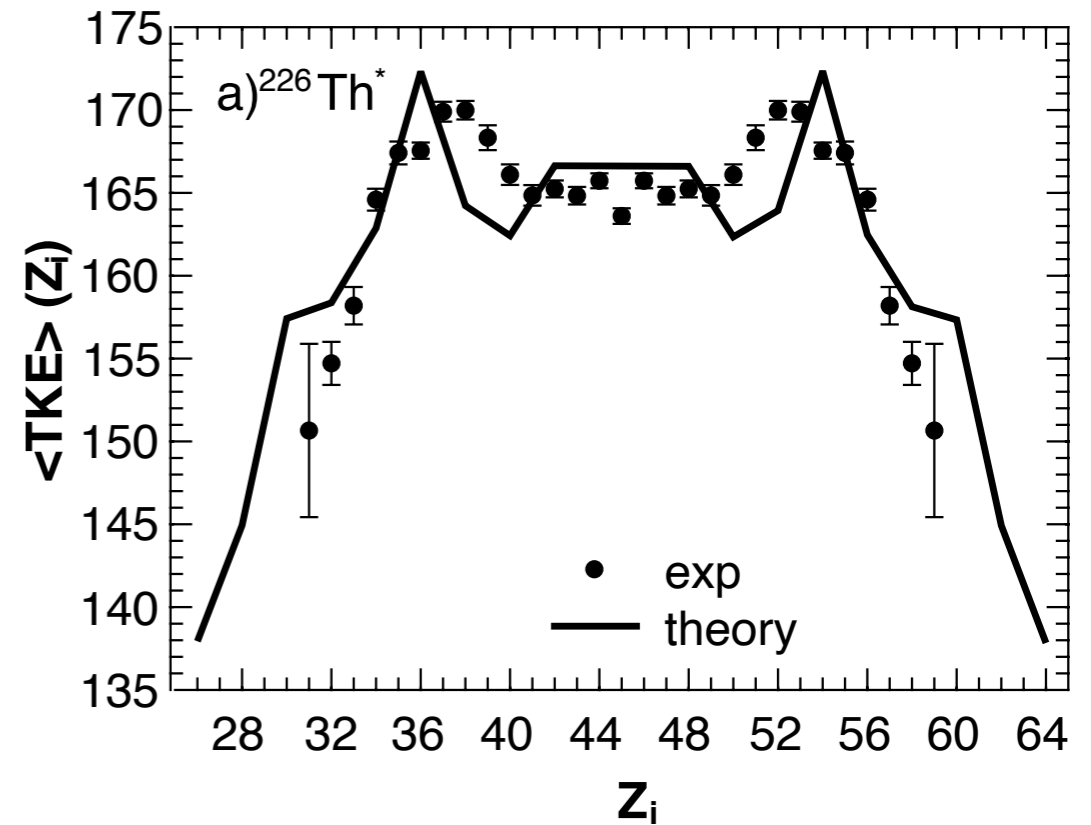
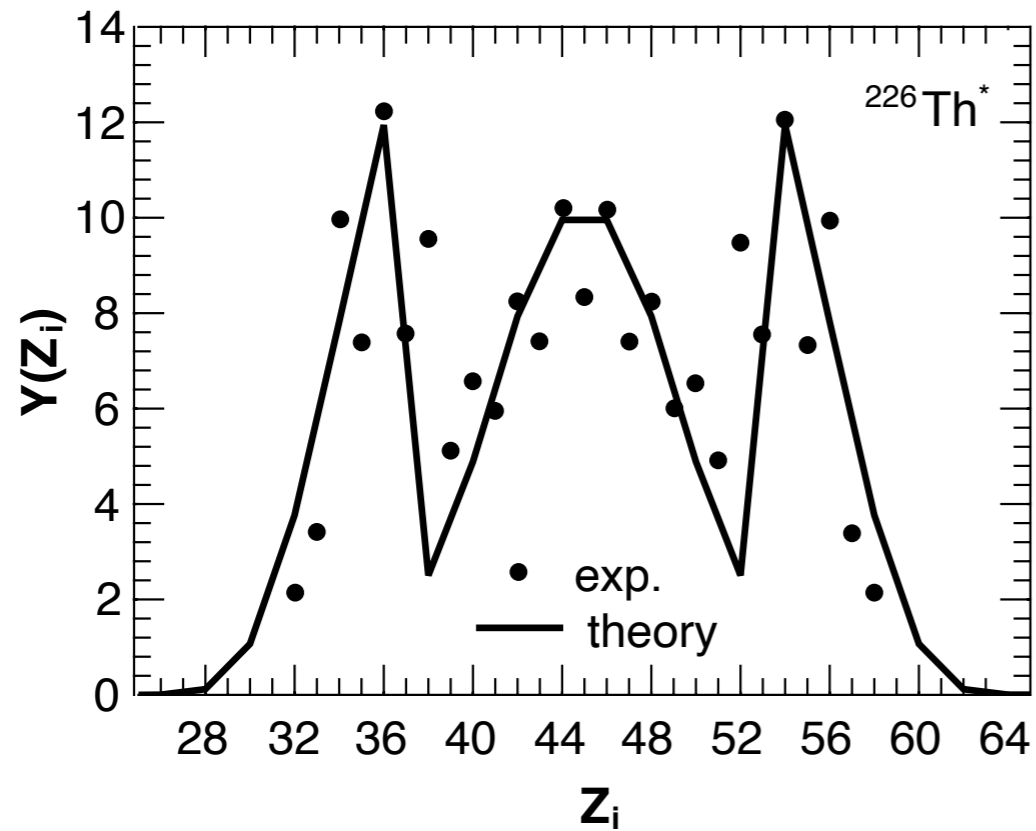
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- R. Schmidt and H. Henschel, Nucl. Phys. A 395, 29 (1983).

Results (and more modelling)



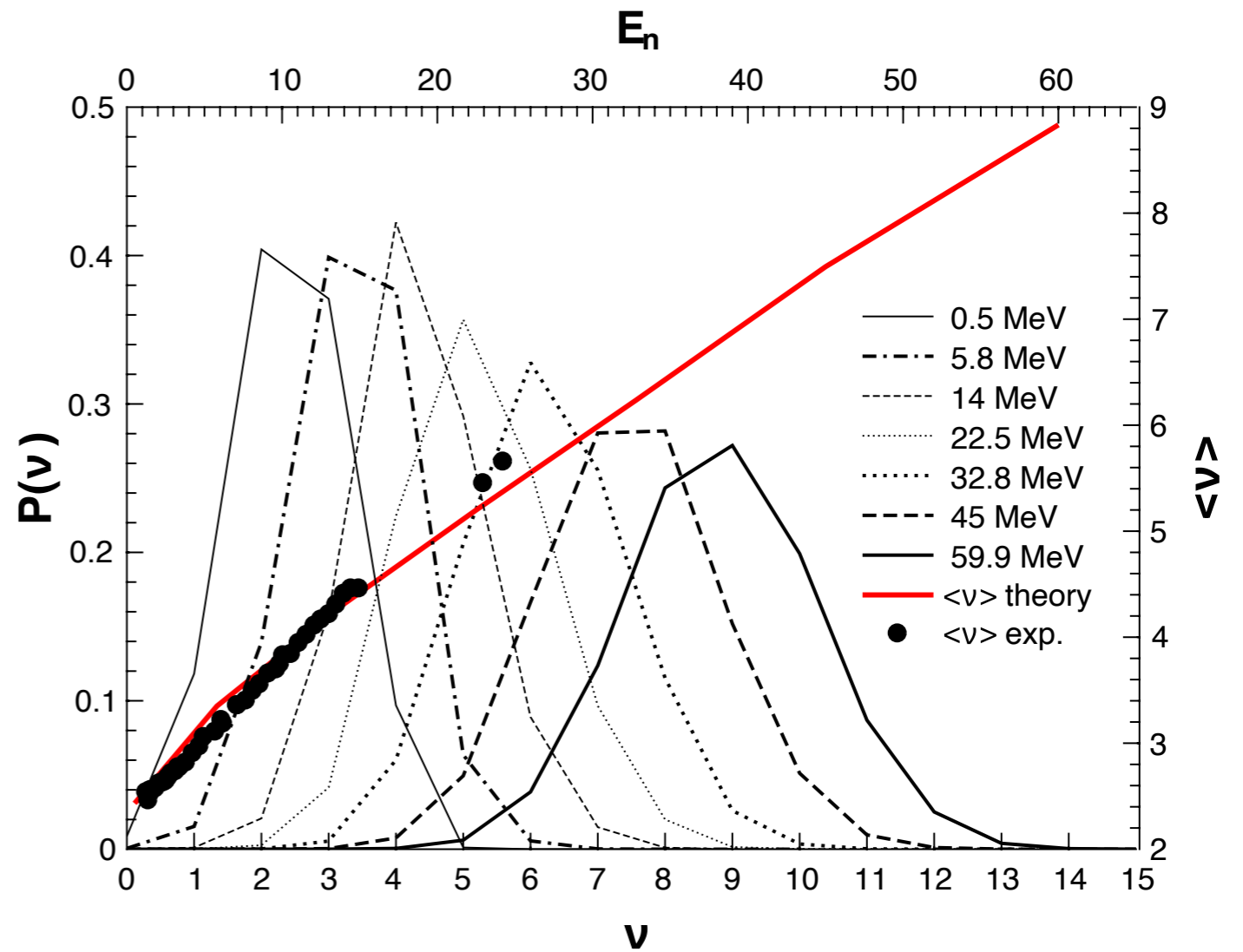
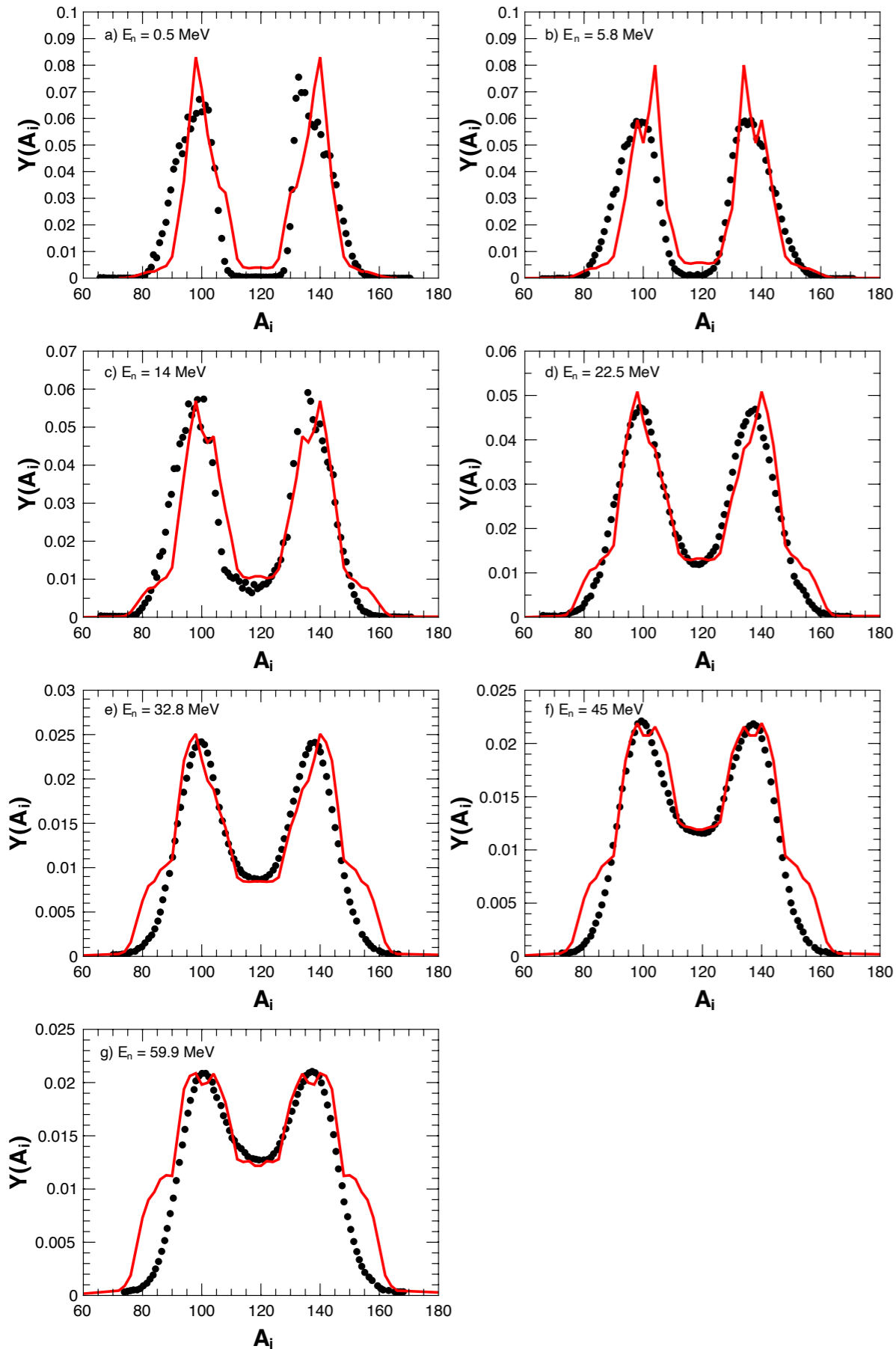
- B.C. Diven, H.C. Martin, R.F. Taschek, and J. Terrell, Phys. Rev. 101, 1012 (1956).
- N. E. Holden and M. S. Zucker, Brookhaven National Laboratory Report No. BNL-NCS-35513 (1985).
- Z. Huanqiao, L. Zuhua, D. Shengyao, and L. Shaoming, Nucl. Sci. Eng. 86, 315 (1984).

Results - thorium isotopes

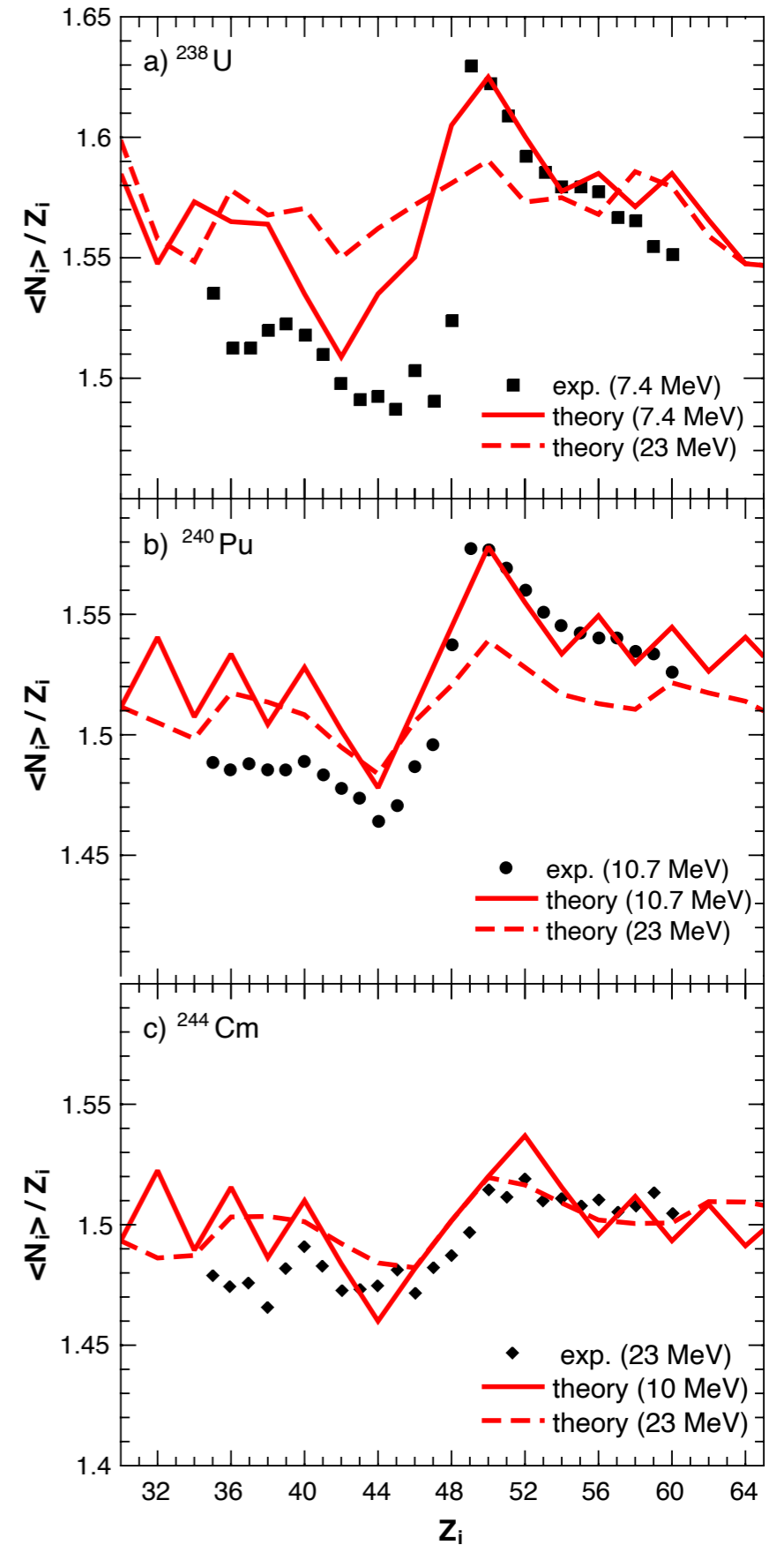
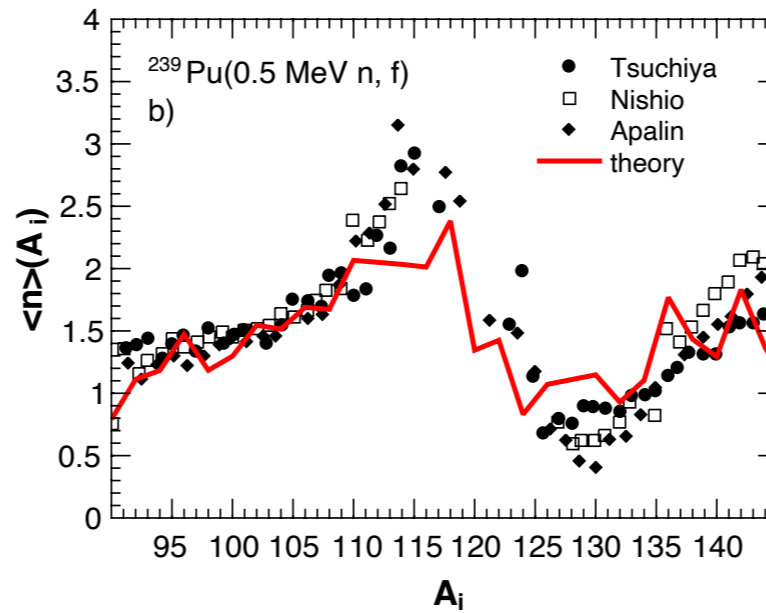
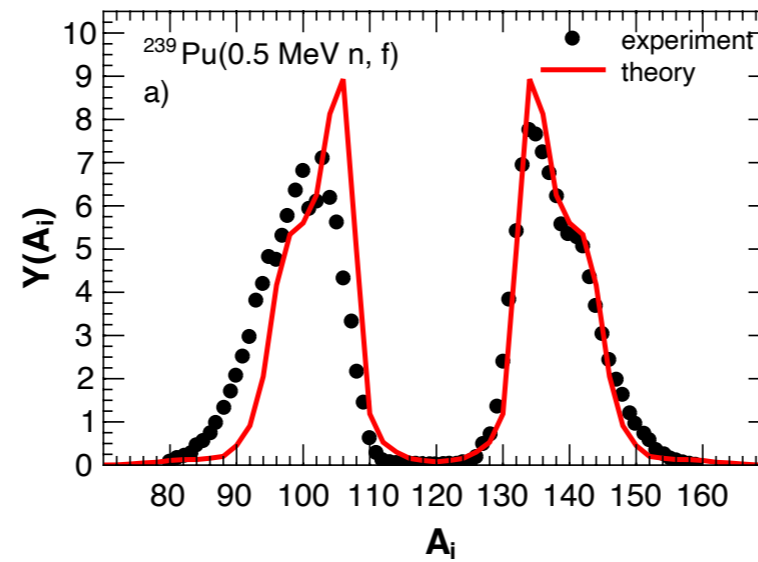
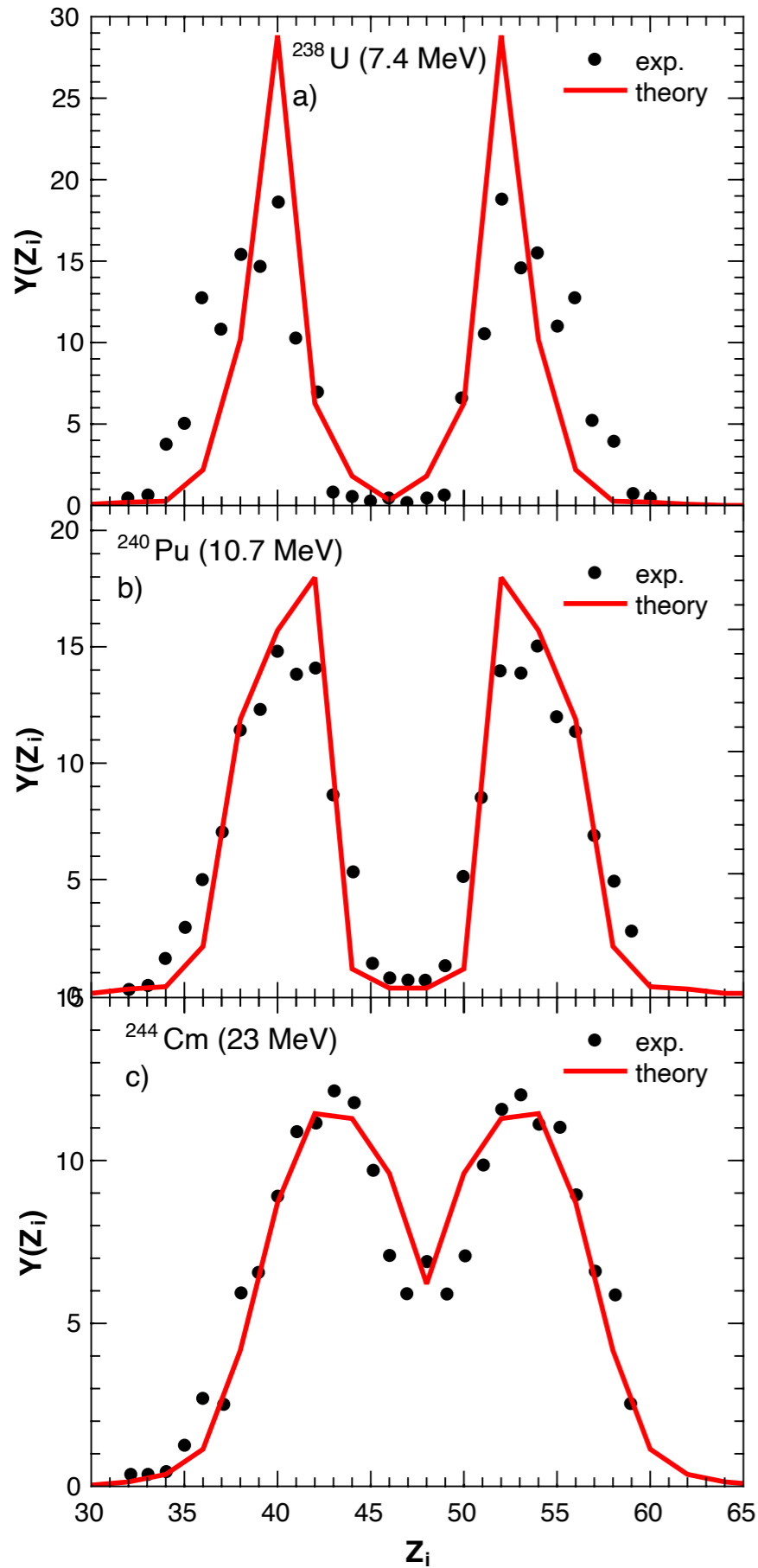


- K.-H. Schmidt et al., Nucl. Phys. A 665, 221 (2000); 693, 169 (2001).
- A. Chatillon et al., Phys. Rev. Lett. 124, 202502 (2020)

What about excitation energy of CN? $^{238}\text{U}(n,f)$



Results - $\langle N/Z \rangle$ ratios in primary fragments



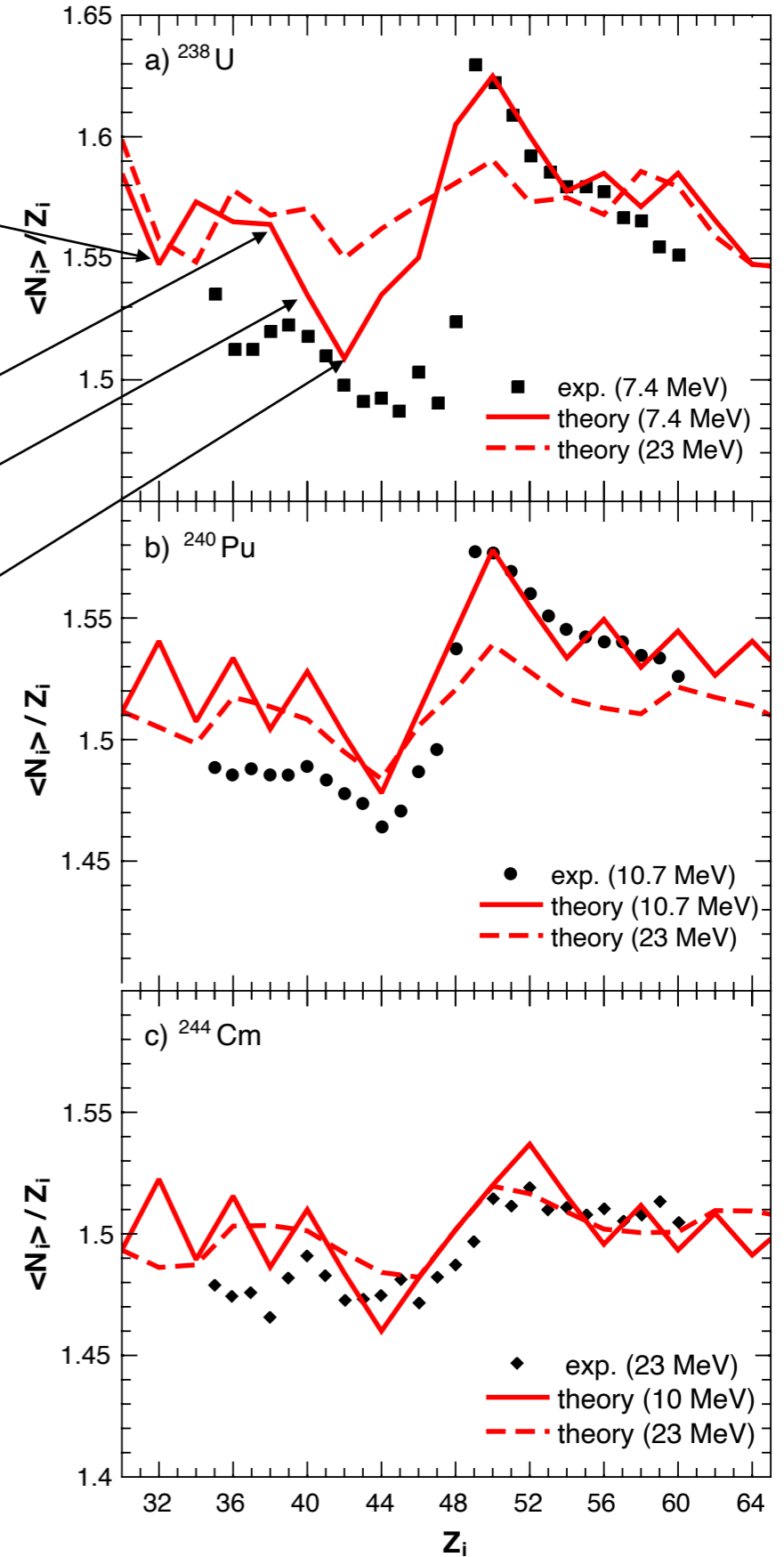
Results - $\langle N/Z \rangle$ ratios in primary fragments

$Z_L=32 \rightarrow {}^{82}\text{Ge} (N_L=50) + {}^{156}\text{Nd}$

$Z_L=38 \rightarrow {}^{102}\text{Sr} + {}^{136}\text{Xe} (N_H=82), {}^{100}\text{Sr} + {}^{138}\text{Xe} (N_H=84), {}^{98}\text{Sr} + {}^{140}\text{Xe} (N_H=86)$

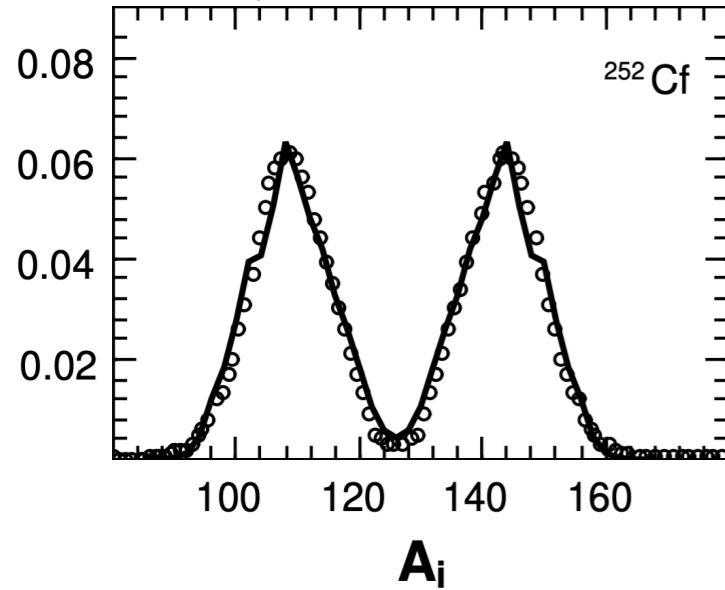
$Z_L = 40 \rightarrow {}^{102}\text{Zr} + {}^{136}\text{Te}, {}^{104}\text{Zr} + {}^{134}\text{Te}, {}^{106}\text{Zr} + {}^{132}\text{Te}$

$Z_L = 42 \rightarrow {}^{106}\text{Mo} + {}^{132}\text{Sn}$ (double magic)



Focus: ^{252}Cf

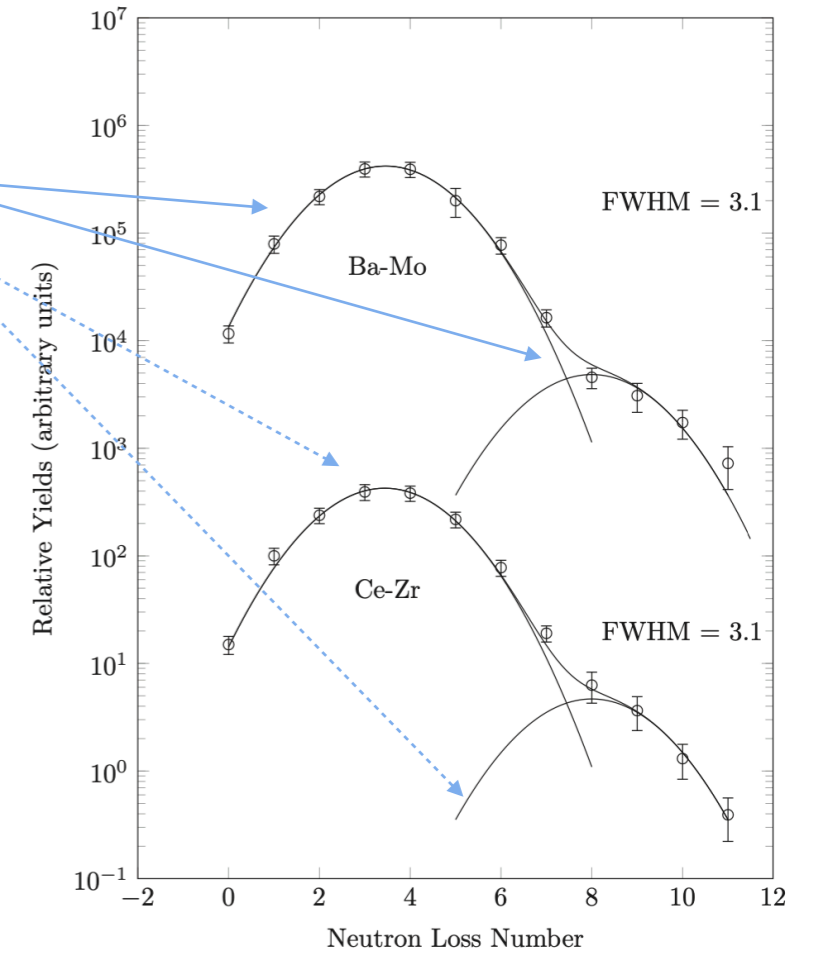
Paşca et. al, PRC **99**, 064611 (2019)



Two fission modes?

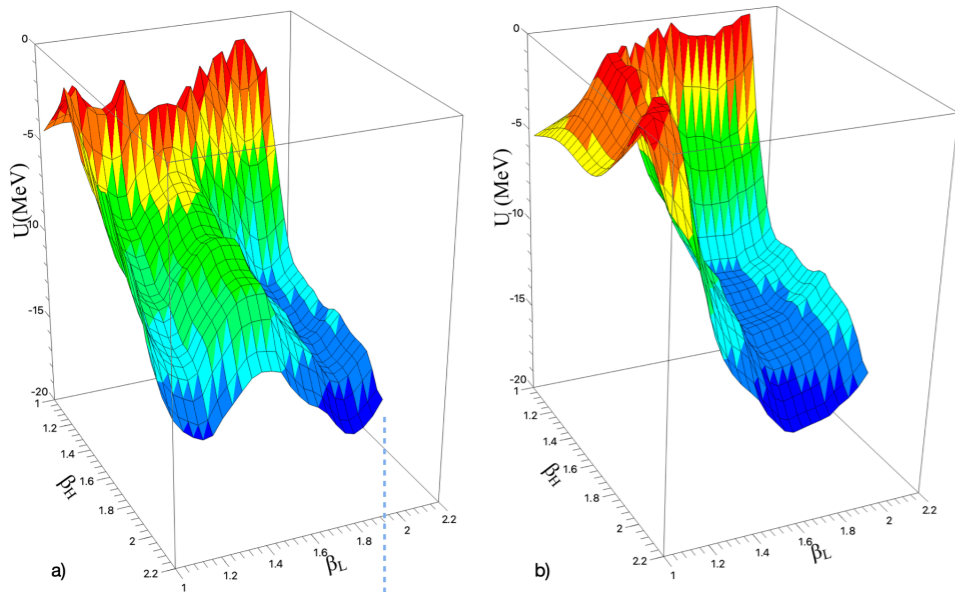
Ba-Mo: 1st mode: $\langle TKE \rangle = 189 \pm 1$ MeV
 $\langle \nu \rangle = 3 - 4$

2nd mode: $\langle TKE \rangle = 153 \pm 3$ MeV
 $\langle \nu \rangle \simeq 8$

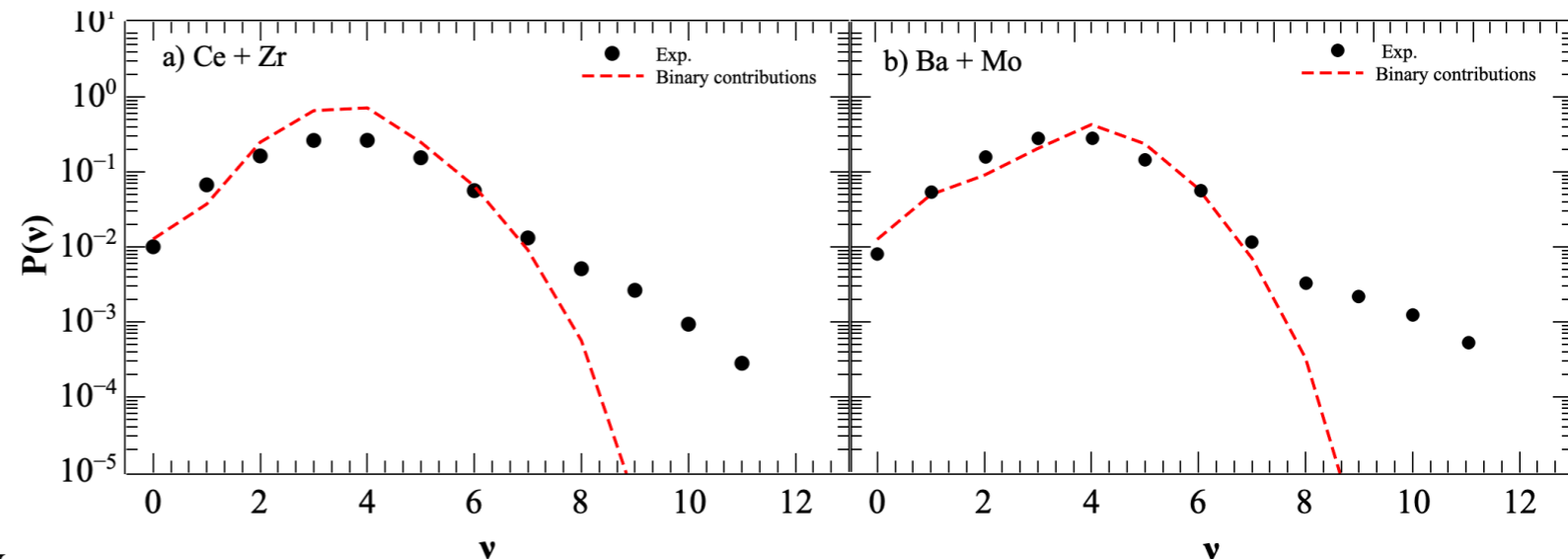


G.M. Ter-Akopian et.al., PRC**55** 1146 (1997): “(...) data clearly indicate that most of the neutrons in the high neutron emission events come out of the Ba fragments – not the Mo fragments.”

B. M. Musangu et.al., PRC **101**, 034610 (2020)
 see also: G.M. Ter-Akopian et.al., PRL**77**, 32 (1996)
 and G.M. Ter-Akopian et.al., PRC**55** 1146 (1997)



Theory: $\langle TKE \rangle = 192.4$ MeV (Ba-Mo)
 $\langle TKE \rangle = 189.1$ MeV (Ce-Zr)
 $\langle \nu \rangle \simeq 3.5$



Theory: $TKE(\beta_1 = \beta_2 = 2.2) = 173.4$ MeV (Ba-Mo) and
 171.0 MeV (Ce-Zr), respectively.

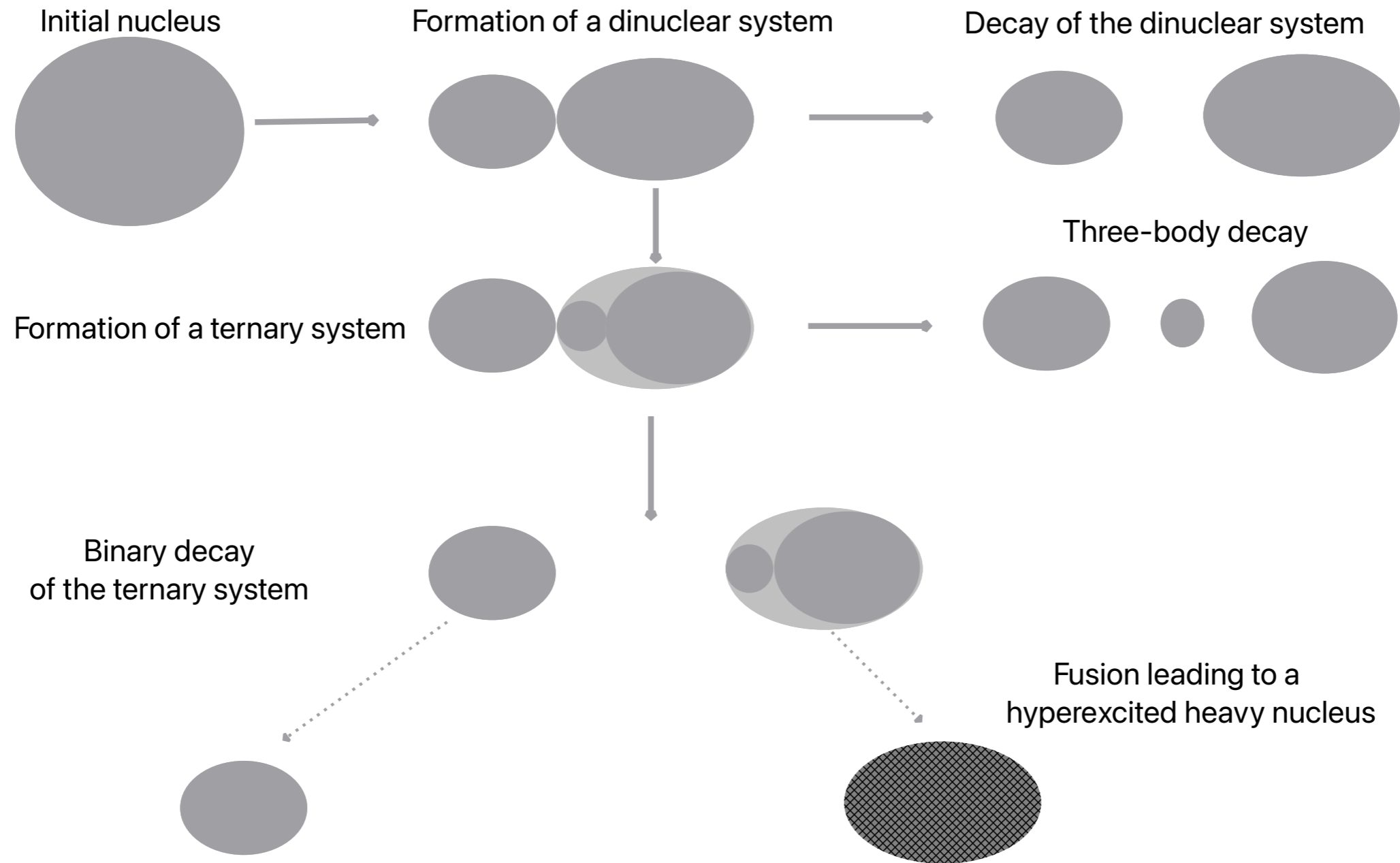
Maximum excitation energy: $\epsilon_{Ba}^* = E_{Ba}^* + U_{Ba}^{def} = 28.1$ MeV
 and $\epsilon_{Ce}^* = E_{Ce}^* + U_{Ce}^{def} = 24.1$ MeV, respectively.

To obtain $TKE \sim 153$ MeV one needs $\beta_{Ba,Mo} \geq 2.5$ (>3 PRC55 1146 (1997)) - unrealistic.

Focus: ^{252}Cf

Competing mechanisms?

Idea supported by: I. Tsekhanovich *et. al.*, PRC67 034610 (2003), Yu.V. Pyatkov *et.al.*, Physics of Atomic Nuclei 86/4 (2023), ...

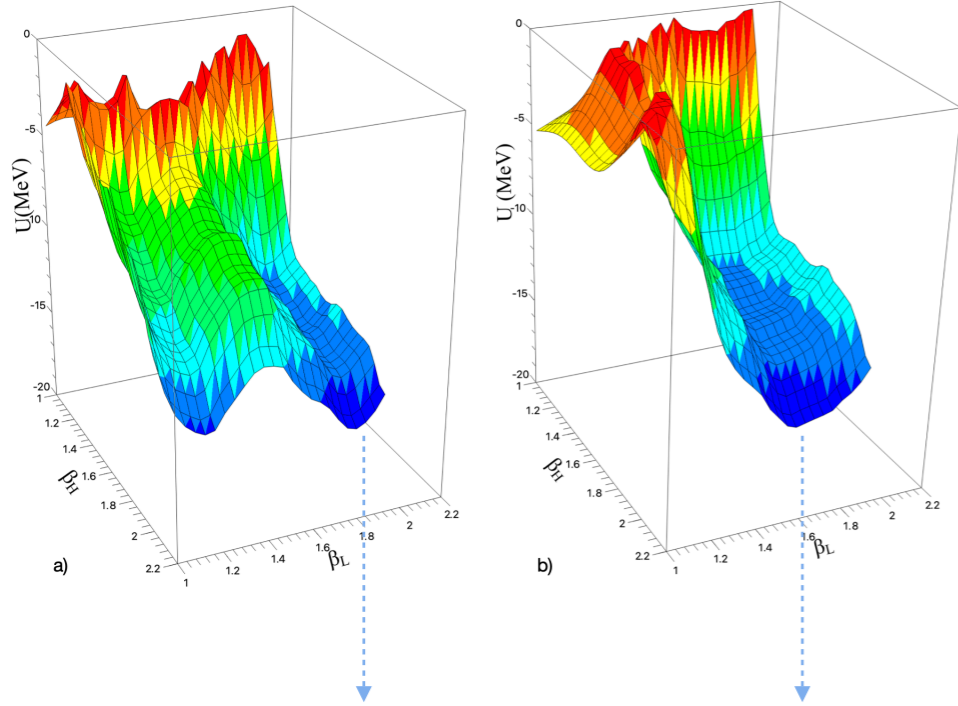


Schematic illustration of normal binary decay and ternary fission processes, as well as the ternary cluster mechanism leading to a hyperexcited heavy fragment.

Focus: ^{252}Cf

Competing mechanisms?

Idea supported by: I. Tsekhanovich *et. al.*, PRC67 034610 (2003), Table I (right)



Ba (left) and Ce (right) nuclei have sufficient ϵ^* to be energetically feasible to be represented as clusters.

Notable third fragments: α , ^{10}Be , ^{14}C , ^{20}O , $^{24,26}\text{Ne}$ (neutron rich nuclei).

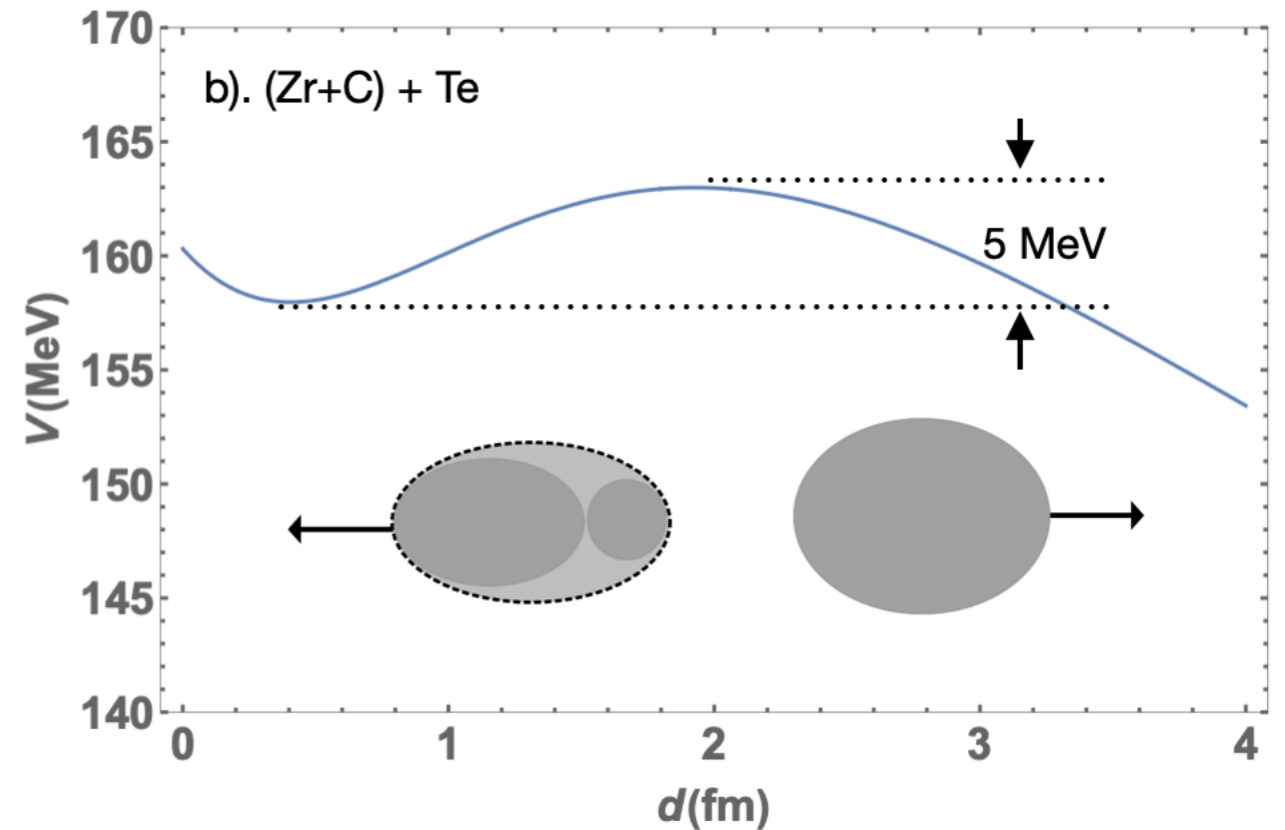
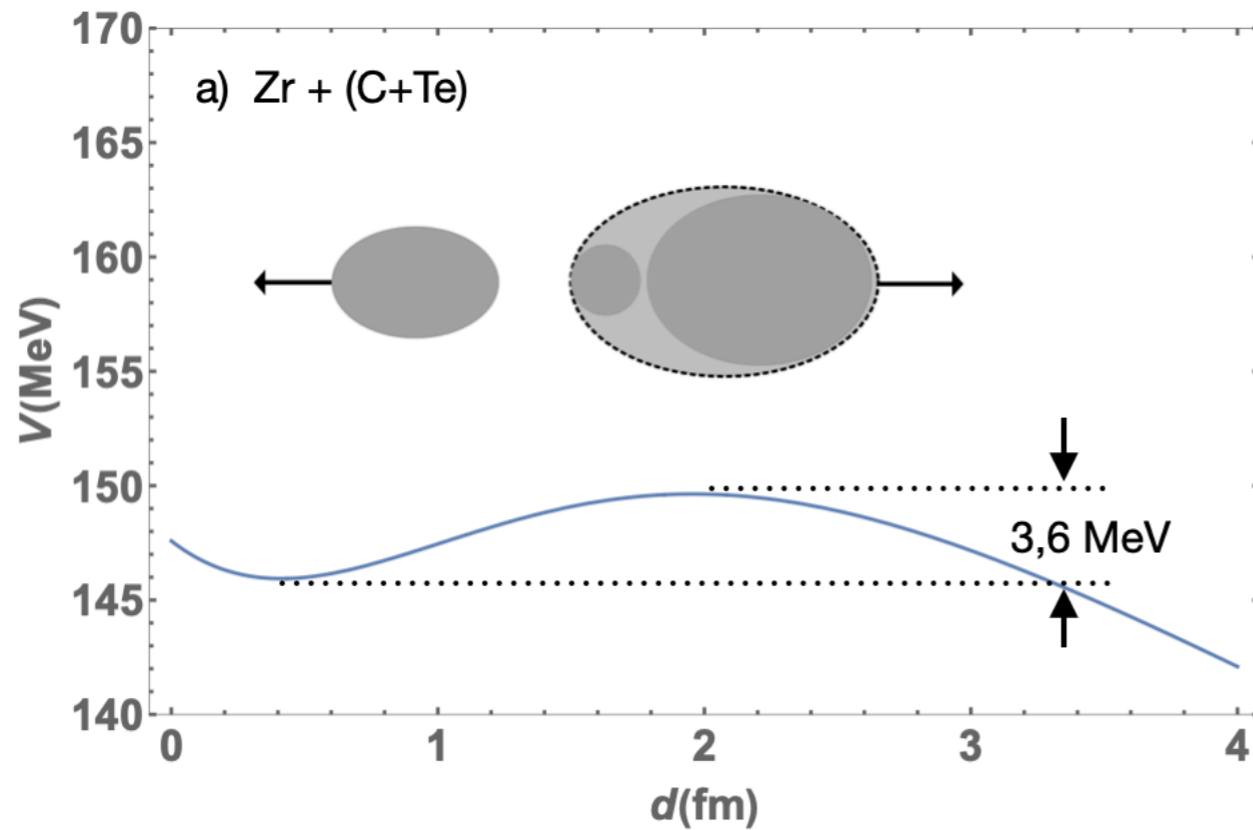
TABLE I. Mean kinetic energy \bar{E} , width of energy distribution σ_E and absolute yield of ternary particles. Values given in *italics* are those from an “enforced” fit, as explained in the text; they should be considered as preliminary.

	\bar{E} (MeV)	σ_E (MeV)	Yield
^8Li	15.1 ± 1.4	7.1 ± 1.3	$(2.6 \pm 0.7) \times 10^{-6}$
^9Li	12.5 ± 0.9	5.5 ± 1.0	$(3.8 \pm 1.0) \times 10^{-6}$
^{10}Be	17.5 ± 0.4	7.7 ± 0.6	$(3.8 \pm 0.7) \times 10^{-5}$
^{11}Be	16.5 ± 1.3	7.4 ± 0.9	$(4.7 \pm 1.2) \times 10^{-6}$
^{12}Be	15.1 ± 1.1	7.1 ± 1.1	$(2.7 \pm 0.7) \times 10^{-6}$
^{12}B	21.8 ± 0.8	8.2 ± 1.8	$(1.5 \pm 0.4) \times 10^{-6}$
^{13}B	20.1 ± 1.1	8.1 ± 0.9	$(2.4 \pm 0.6) \times 10^{-6}$
^{14}B	17.0 ± 1.2	7.3 ± 0.7	$(1.4 \pm 0.4) \times 10^{-7}$
^{15}B	16.8 ± 1.9	7.0 ± 1.0	$(9.1 \pm 4.1) \times 10^{-8}$
^{14}C	27.0 ± 0.3	9.9 ± 0.5	$(1.3 \pm 0.2) \times 10^{-5}$
^{15}C	25.1 ± 0.5	8.9 ± 0.7	$(5.3 \pm 1.1) \times 10^{-6}$
^{16}C	24.4 ± 1.1	9.6 ± 1.2	$(4.8 \pm 1.1) \times 10^{-6}$
^{17}C	21.3 ± 1.7	8.3 ± 0.9	$(7.5 \pm 2.8) \times 10^{-7}$
^{18}C	20.4 ± 2.8	8.5 ± 1.4	$(2.4 \pm 0.7) \times 10^{-7}$
^{16}N	25.9 ± 2.2	9.8 ± 1.7	$(1.5 \pm 0.4) \times 10^{-7}$
^{17}N	25.0 ± 1.6	9.4 ± 1.2	$(8.1 \pm 2.0) \times 10^{-7}$
^{18}N	23.8 ± 1.5	9.9 ± 1.2	$(4.5 \pm 1.1) \times 10^{-7}$
^{20}N	<i>fixed</i>	7.0 ± 0.9	1.3×10^{-8}
^{21}N	<i>fixed</i>	<i>fixed</i>	3.4×10^{-9}
^{20}O	31.4 ± 1.7	10.6 ± 1.9	$(2.5 \pm 0.7) \times 10^{-6}$
^{21}O	24.2 ± 1.2	10.7 ± 0.7	$(6.4 \pm 1.3) \times 10^{-7}$
^{22}O	33.0 ± 7.4	14.3 ± 4.2	$(4.2 \pm 1.6) \times 10^{-7}$
^{24}O	<i>fixed</i>	9.5 ± 3.2	5.8×10^{-8}
^{20}F	25.4 ± 3.3	<i>fixed</i>	9.7×10^{-9}
^{21}F	26.5 ± 2.1	9.8 ± 1.3	$(1.6 \pm 0.4) \times 10^{-7}$
^{22}F	33.8 ± 10.5	12.2 ± 4.6	$(1.4 \pm 0.8) \times 10^{-7}$
^{24}F	26.3 ± 2.8	12.1 ± 2.0	$(8.3 \pm 4.0) \times 10^{-8}$
^{24}Ne	33.9 ± 2.9	14.2 ± 1.9	$(2.4 \pm 0.6) \times 10^{-7}$
^{27}Ne	35.9 ± 5.9	<i>fixed</i>	2.0×10^{-8}
^{28}Ne	<i>fixed</i>	<i>fixed</i>	1.8×10^{-8}
^{27}Na	38.4 ± 8.2	16.3 ± 4.5	$(8.2 \pm 3.2) \times 10^{-8}$
^{28}Na	<i>fixed</i>	<i>fixed</i>	1.0×10^{-7}
^{30}Na	31.7 ± 8.6	11.9 ± 6.1	$(2.2 \pm 2.2) \times 10^{-8}$
^{30}Mg	34.9 ± 3.7	13.0 ± 1.8	$(1.3 \pm 0.4) \times 10^{-7}$
^{32}Mg	<i>fixed</i>	10.8 ± 2.7	3.7×10^{-8}
^{34}Mg	<i>fixed</i>	<i>fixed</i>	1.0×10^{-9}
^{30}Al	<i>fixed</i>	<i>fixed</i>	9.0×10^{-9}
^{32}Al	<i>fixed</i>	<i>fixed</i>	1.1×10^{-8}
^{33}Al	<i>fixed</i>	<i>fixed</i>	1.8×10^{-8}
^{32}Si	<i>fixed</i>	12.0 ± 1.7	8.9×10^{-9}
^{33}Si	<i>fixed</i>	11.3 ± 1.4	1.5×10^{-8}
^{34}Si	<i>fixed</i>	11.3 ± 1.3	2.2×10^{-8}
^{37}Si	<i>fixed</i>	<i>fixed</i>	2.0×10^{-9}
^{39}P	<i>fixed</i>	<i>fixed</i>	$< 5.6 \times 10^{-9}$
^{37}S	<i>fixed</i>	<i>fixed</i>	4.7×10^{-9}
^{40}S	<i>fixed</i>	<i>fixed</i>	$< 3.3 \times 10^{-9}$

Focus: ^{252}Cf

Competing mechanisms?

Idea supported by: I. Tsekhanovich *et. al.*, PRC67 034610 (2003), Yu.V. Pyatkov *et.al.*, Physics of Atomic Nuclei 86/4 (2023), ...



Interaction potentials favor the formation of the cluster composing of the heaviest nucleus with the third particle.

G.M. Ter-Akopian *et.al.*, PRC55 1146 (1997): “(...) data clearly indicate that most of the neutrons in the high neutron emission events come out of the Ba fragments – not the Mo fragments.”

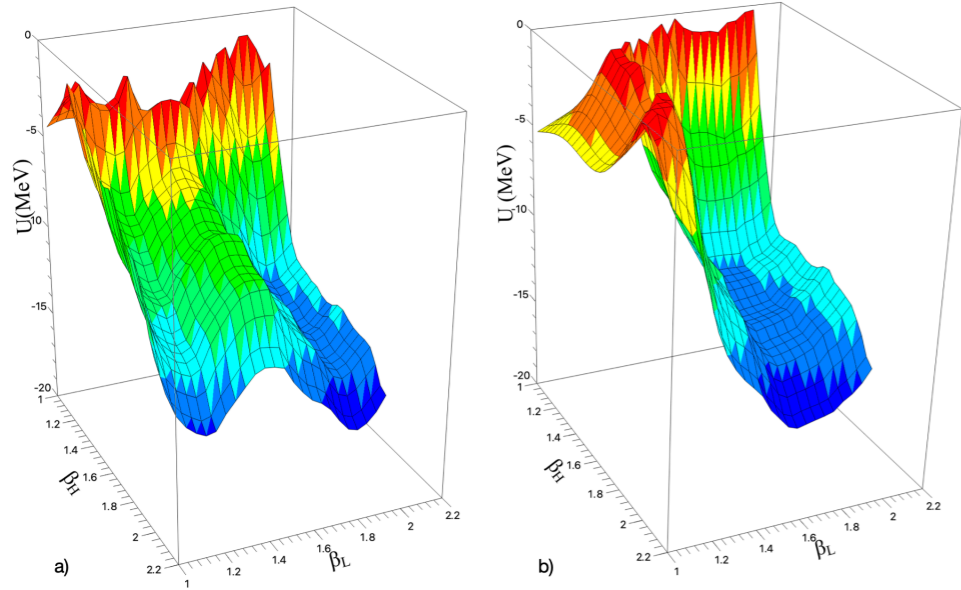
Focus: ^{252}Cf

Final excitation energy of the heavy fragment is: $\varepsilon^{**} = \varepsilon^* + \Delta V - Q$,

where $\varepsilon^* = E_i^* + U^{def}$;

ΔV - the reduction of the interaction potential due to the third fragment;

Q is the fusion Q -value.



$^{148}\text{Ce} + ^{104}\text{Zr}$

$$\varepsilon_{\text{Ce}}^* = E_{\text{Ce}}^* + U_{\text{Ce}}^{def} = 24.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
$(^{144}\text{Ba} + \alpha) + ^{104}\text{Zr}$	8.3	1.1	31.3	162.0
$(^{138}\text{Xe} + ^{10}\text{Be}) + ^{104}\text{Zr}$	13.9	2.9	35.1	140.5
$(^{134}\text{Te} + ^{14}\text{C}) + ^{104}\text{Zr}$	30.8	-9.2	64.0	150.3
$(^{128}\text{Sn} + ^{20}\text{O}) + ^{104}\text{Zr}$	39.6	-9.2	72.7	153.6
$(^{122}\text{Cd} + ^{26}\text{Ne}) + ^{104}\text{Zr}$	47.4	-9.9	81.4	144.5
$(^{124}\text{Cd} + ^{24}\text{Ne}) + ^{104}\text{Zr}$	45.8	-11.0	80.9	144.2

$^{146}\text{Ba} + ^{106}\text{Mo}$

$$\varepsilon_{\text{Ba}}^* = E_{\text{Ba}}^* + U_{\text{Ba}}^{def} = 28.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
$(^{142}\text{Xe} + \alpha) + ^{106}\text{Mo}$	12.7	2.0	38.8	166.6
$(^{136}\text{Te} + ^{10}\text{Be}) + ^{106}\text{Mo}$	13.7	3.2	38.6	144.4
$(^{132}\text{Sn} + ^{14}\text{C}) + ^{106}\text{Mo}$	30.5	-8.5	67.1	155.9
$(^{126}\text{Cd} + ^{20}\text{O}) + ^{106}\text{Mo}$	39.0	-3.5	70.6	160.6
$(^{120}\text{Pd} + ^{26}\text{Ne}) + ^{106}\text{Mo}$	45.0	-4.7	77.8	157.1
$(^{122}\text{Pd} + ^{24}\text{Ne}) + ^{106}\text{Mo}$	44.8	-5.6	78.5	158.0

Large variations in ΔV , especially for heavy third nuclei.

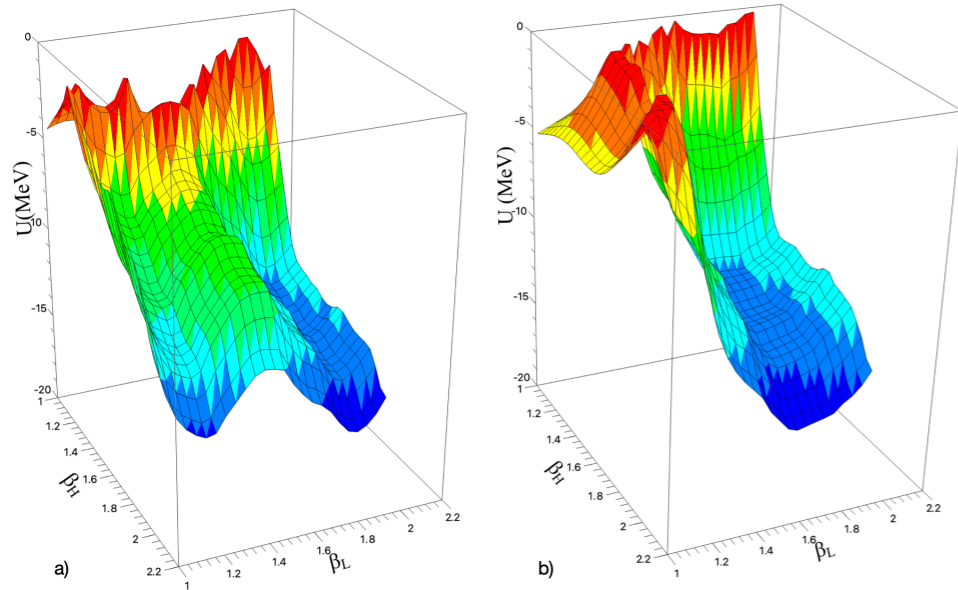
Focus: ^{252}Cf

Final excitation energy of the heavy fragment is: $\varepsilon^{**} = \varepsilon^* + \Delta V - Q$,

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ΔV - the reduction of the interaction potential due to the third fragment;

Q is the fusion Q -value.



$^{148}\text{Ce} + ^{104}\text{Zr}$

$\varepsilon_{\text{Ce}}^* = E_{\text{Ce}}^* + U_{\text{Ce}}^{def} = 24.1 \text{ MeV}$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
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$^{146}\text{Ba} + ^{106}\text{Mo}$

$\varepsilon_{\text{Ba}}^* = E_{\text{Ba}}^* + U_{\text{Ba}}^{def} = 28.1 \text{ MeV}$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
$(^{142}\text{Xe} + \alpha) + ^{106}\text{Mo}$	12.7	2.0	38.8	166.6
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$(^{122}\text{Pd} + ^{24}\text{Ne}) + ^{106}\text{Mo}$	44.8	-5.6	78.5	158.0

Large variations in ΔV , especially for heavy third nuclei.

Large negative Q -values enhance ε^{**} , especially for heavy third nuclei.

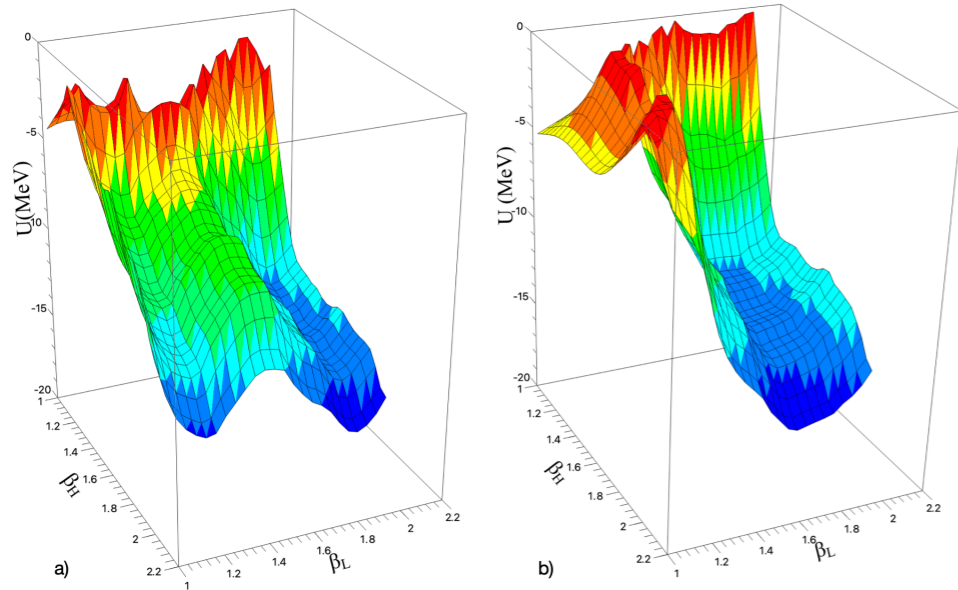
Focus: ^{252}Cf

Final excitation energy of the heavy fragment is: $\varepsilon^{**} = \varepsilon^* + \Delta V - Q$,

where $\varepsilon^* = E_i^* + U^{def}$;

ΔV - the reduction of the interaction potential due to the third fragment;

Q is the fusion Q -value.



$$^{148}\text{Ce} + ^{104}\text{Zr} \quad \varepsilon_{\text{Ce}}^* = E_{\text{Ce}}^* + U_{\text{Ce}}^{def} = 24.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
$(^{144}\text{Ba} + \alpha) + ^{104}\text{Zr}$	8.3	1.1	31.3	162.0
$(^{138}\text{Xe} + ^{10}\text{Be}) + ^{104}\text{Zr}$	13.9	2.9	35.1	140.5
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$(^{124}\text{Cd} + ^{24}\text{Ne}) + ^{104}\text{Zr}$	45.8	-11.0	80.9	144.2

$$^{146}\text{Ba} + ^{106}\text{Mo} \quad \varepsilon_{\text{Ba}}^* = E_{\text{Ba}}^* + U_{\text{Ba}}^{def} = 28.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ε_H^{**} (MeV)	TKE (MeV)
$(^{142}\text{Xe} + \alpha) + ^{106}\text{Mo}$	12.7	2.0	38.8	166.6
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Large variations in ΔV , especially for heavy third nuclei.

Large negative Q -values enhance ε^{**} , especially for heavy third nuclei.

Extremely large excitation energy for the heavy fragment after heavy cluster's fusion.

Focus: ^{252}Cf

Competing mechanisms?

The probability of emitting exactly ν neutrons, taking into account binary and tripartition contribution, can be written as:

$$P(\nu) = \lambda_d P_\nu(A_i, Z_i, \beta_i, \varepsilon^*) + \lambda_t P_\nu(A_i, Z_i, \beta_i, \varepsilon^{**})$$

$$P_\nu(A_i, Z_i, \beta_i, \varepsilon) = \sum_{\nu_L=0}^{\nu} \int_0^{\varepsilon - U_H^{\text{def}} - U_L^{\text{def}}} dE P_C(E) P_{\nu_L}(U_L^{\text{def}} + E) P_{\nu-\nu_L}(\varepsilon - U_L^{\text{def}} - E)$$

The probability of emitting the third cluster is the product of the probabilities for formation of ternary system and three body decay.

↳ λ_t can be estimated using the experimental values.

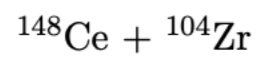
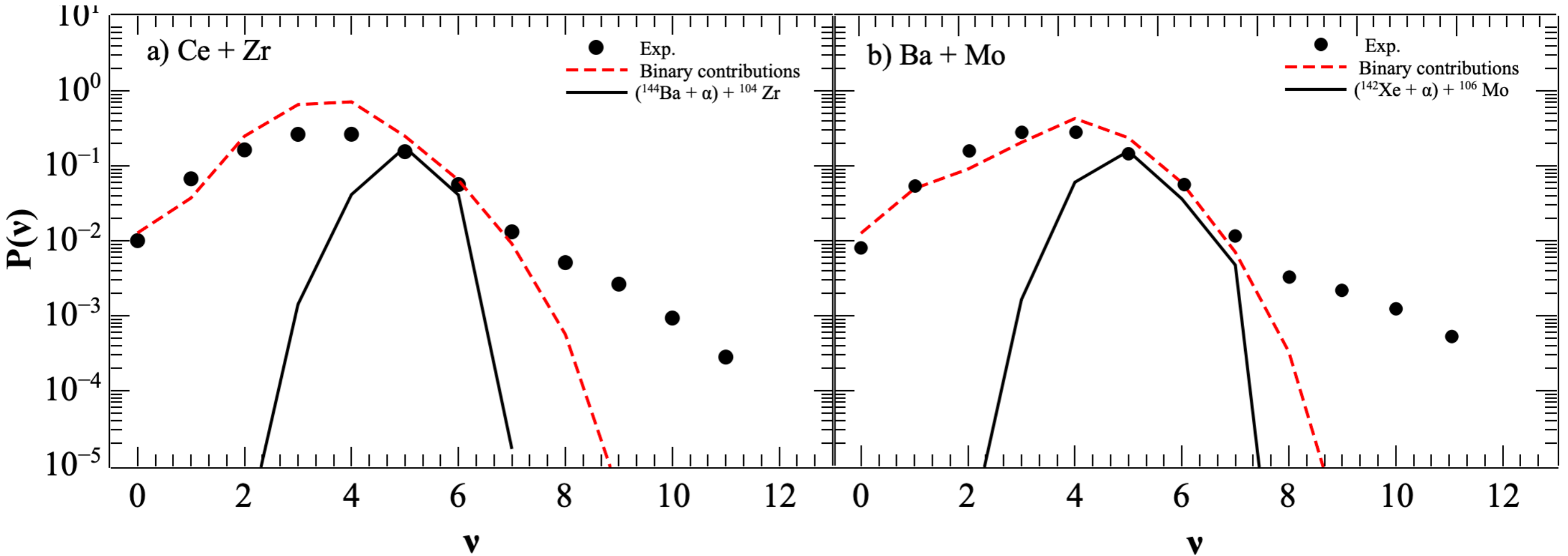
↳ our estimates: $< 10^{-2}$ compared to binary decay.

TABLE I. Mean kinetic energy \bar{E} , width of energy distribution σ_E and absolute yield of ternary particles. Values given in *italics* are those from an “enforced” fit, as explained in the text; they should be considered as preliminary.

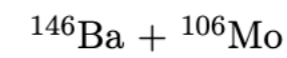
	\bar{E} (MeV)	σ_E (MeV)	Yield
^8Li	15.1 ± 1.4	7.1 ± 1.3	$(2.6 \pm 0.7) \times 10^{-6}$
^9Li	12.5 ± 0.9	5.5 ± 1.0	$(3.8 \pm 1.0) \times 10^{-6}$
^{10}Be	17.5 ± 0.4	7.7 ± 0.6	$(3.8 \pm 0.7) \times 10^{-5}$
^{11}Be	16.5 ± 1.3	7.4 ± 0.9	$(4.7 \pm 1.2) \times 10^{-6}$
^{12}Be	15.1 ± 1.1	7.1 ± 1.1	$(2.7 \pm 0.7) \times 10^{-6}$
^{12}B	21.8 ± 0.8	8.2 ± 1.8	$(1.5 \pm 0.4) \times 10^{-6}$
^{13}B	20.1 ± 1.1	8.1 ± 0.9	$(2.4 \pm 0.6) \times 10^{-6}$
^{14}B	17.0 ± 1.2	7.3 ± 0.7	$(1.4 \pm 0.4) \times 10^{-7}$
^{15}B	16.8 ± 1.9	7.0 ± 1.0	$(9.1 \pm 4.1) \times 10^{-8}$
^{14}C	27.0 ± 0.3	9.9 ± 0.5	$(1.3 \pm 0.2) \times 10^{-5}$
^{15}C	25.1 ± 0.5	8.9 ± 0.7	$(5.3 \pm 1.1) \times 10^{-6}$
^{16}C	24.4 ± 1.1	9.6 ± 1.2	$(4.8 \pm 1.1) \times 10^{-6}$
^{17}C	21.3 ± 1.7	8.3 ± 0.9	$(7.5 \pm 2.8) \times 10^{-7}$
^{18}C	20.4 ± 2.8	8.5 ± 1.4	$(2.4 \pm 0.7) \times 10^{-7}$
^{16}N	25.9 ± 2.2	9.8 ± 1.7	$(1.5 \pm 0.4) \times 10^{-7}$
^{17}N	25.0 ± 1.6	9.4 ± 1.2	$(8.1 \pm 2.0) \times 10^{-7}$
^{18}N	23.8 ± 1.5	9.9 ± 1.2	$(4.5 \pm 1.1) \times 10^{-7}$
^{20}N	<i>fixed</i>	7.0 ± 0.9	1.3×10^{-8}
^{21}N	<i>fixed</i>	<i>fixed</i>	3.4×10^{-9}
^{20}O	31.4 ± 1.7	10.6 ± 1.9	$(2.5 \pm 0.7) \times 10^{-6}$
^{21}O	24.2 ± 1.2	10.7 ± 0.7	$(6.4 \pm 1.3) \times 10^{-7}$
^{22}O	33.0 ± 7.4	14.3 ± 4.2	$(4.2 \pm 1.6) \times 10^{-7}$
^{24}O	<i>fixed</i>	9.5 ± 3.2	5.8×10^{-8}
^{20}F	25.4 ± 3.3	<i>fixed</i>	9.7×10^{-9}
^{21}F	26.5 ± 2.1	9.8 ± 1.3	$(1.6 \pm 0.4) \times 10^{-7}$
^{22}F	33.8 ± 10.5	12.2 ± 4.6	$(1.4 \pm 0.8) \times 10^{-7}$
^{24}F	26.3 ± 2.8	12.1 ± 2.0	$(8.3 \pm 4.0) \times 10^{-8}$
^{24}Ne	33.9 ± 2.9	14.2 ± 1.9	$(2.4 \pm 0.6) \times 10^{-7}$
^{27}Ne	35.9 ± 5.9	<i>fixed</i>	2.0×10^{-8}
^{28}Ne	<i>fixed</i>	<i>fixed</i>	1.8×10^{-8}
^{27}Na	38.4 ± 8.2	16.3 ± 4.5	$(8.2 \pm 3.2) \times 10^{-8}$
^{28}Na	<i>fixed</i>	<i>fixed</i>	1.0×10^{-7}
^{30}Na	31.7 ± 8.6	11.9 ± 6.1	$(2.2 \pm 2.2) \times 10^{-8}$
^{30}Mg	34.9 ± 3.7	13.0 ± 1.8	$(1.3 \pm 0.4) \times 10^{-7}$
^{32}Mg	<i>fixed</i>	10.8 ± 2.7	3.7×10^{-8}
^{34}Mg	<i>fixed</i>	<i>fixed</i>	1.0×10^{-9}
^{30}Al	<i>fixed</i>	<i>fixed</i>	9.0×10^{-9}
^{32}Al	<i>fixed</i>	<i>fixed</i>	1.1×10^{-8}
^{33}Al	<i>fixed</i>	<i>fixed</i>	1.8×10^{-8}
^{32}Si	<i>fixed</i>	12.0 ± 1.7	8.9×10^{-9}
^{33}Si	<i>fixed</i>	11.3 ± 1.4	1.5×10^{-8}
^{34}Si	<i>fixed</i>	11.3 ± 1.3	2.2×10^{-8}
^{37}Si	<i>fixed</i>	<i>fixed</i>	2.0×10^{-9}
^{39}P	<i>fixed</i>	<i>fixed</i>	$< 5.6 \times 10^{-9}$
^{37}S	<i>fixed</i>	<i>fixed</i>	4.7×10^{-9}
^{40}S	<i>fixed</i>	<i>fixed</i>	$< 3.3 \times 10^{-9}$

Focus: ^{252}Cf

Competing mechanisms? - Neutron multiplicity



$$\epsilon_{\text{Ce}}^* = E_{\text{Ce}}^* + U_{\text{Ce}}^{\text{def}} = 24.1 \text{ MeV}$$



$$\epsilon_{\text{Ba}}^* = E_{\text{Ba}}^* + U_{\text{Ba}}^{\text{def}} = 28.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ϵ_H^{**} (MeV)	TKE (MeV)
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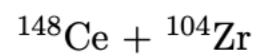
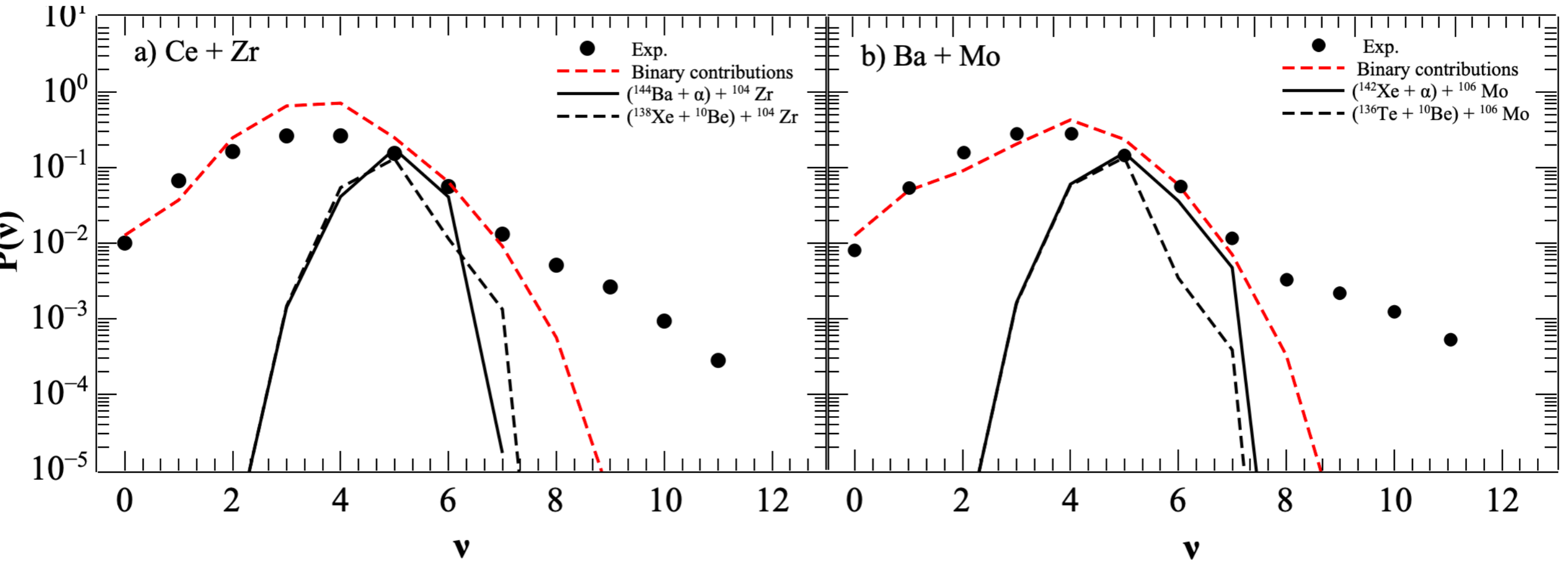
$(^{144}\text{Ba} + \alpha) + ^{104}\text{Zr}$	8.3	1.1	31.3	162.0
$(^{138}\text{Xe} + ^{10}\text{Be}) + ^{104}\text{Zr}$	13.9	2.9	35.1	140.5
$(^{134}\text{Te} + ^{14}\text{C}) + ^{104}\text{Zr}$	30.8	-9.2	64.0	150.3
$(^{128}\text{Sn} + ^{20}\text{O}) + ^{104}\text{Zr}$	39.6	-9.2	72.7	153.6
$(^{122}\text{Cd} + ^{26}\text{Ne}) + ^{104}\text{Zr}$	47.4	-9.9	81.4	144.5
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	ΔV (MeV)	Q (MeV)	ϵ_H^{**} (MeV)	TKE (MeV)
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$(^{142}\text{Xe} + \alpha) + ^{106}\text{Mo}$	12.7	2.0	38.8	166.6
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$(^{132}\text{Sn} + ^{14}\text{C}) + ^{106}\text{Mo}$	30.5	-8.5	67.1	155.9
$(^{126}\text{Cd} + ^{20}\text{O}) + ^{106}\text{Mo}$	39.0	-3.5	70.6	160.6
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Focus: ^{252}Cf

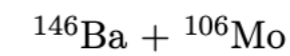
Competing mechanisms? - Neutron multiplicity



$$\epsilon_{Ce}^* = E_{Ce}^* + U_{Ce}^{def} = 24.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ϵ_H^{**} (MeV)	TKE (MeV)
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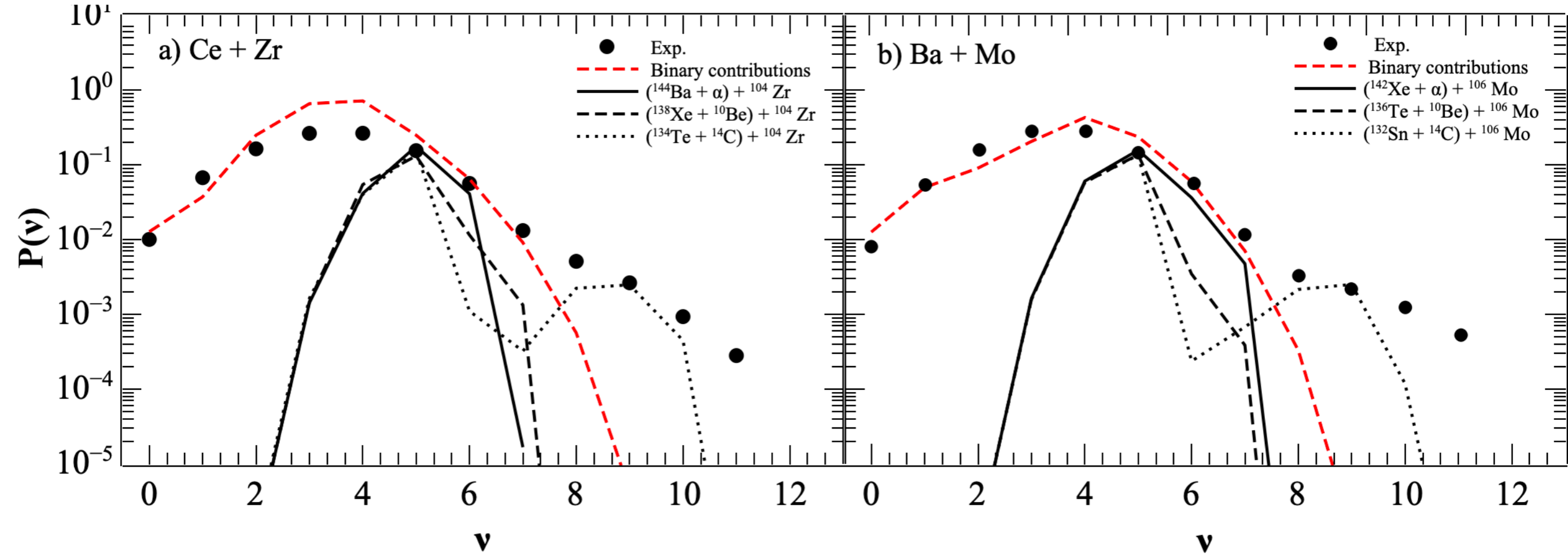
$$\epsilon_{Ba}^* = E_{Ba}^* + U_{Ba}^{def} = 28.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ϵ_H^{**} (MeV)	TKE (MeV)
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Focus: ^{252}Cf

Competing mechanisms? - Neutron multiplicity



$$^{148}\text{Ce} + ^{104}\text{Zr} \quad \epsilon_{\text{Ce}}^* = E_{\text{Ce}}^* + U_{\text{Ce}}^{\text{def}} = 24.1 \text{ MeV}$$

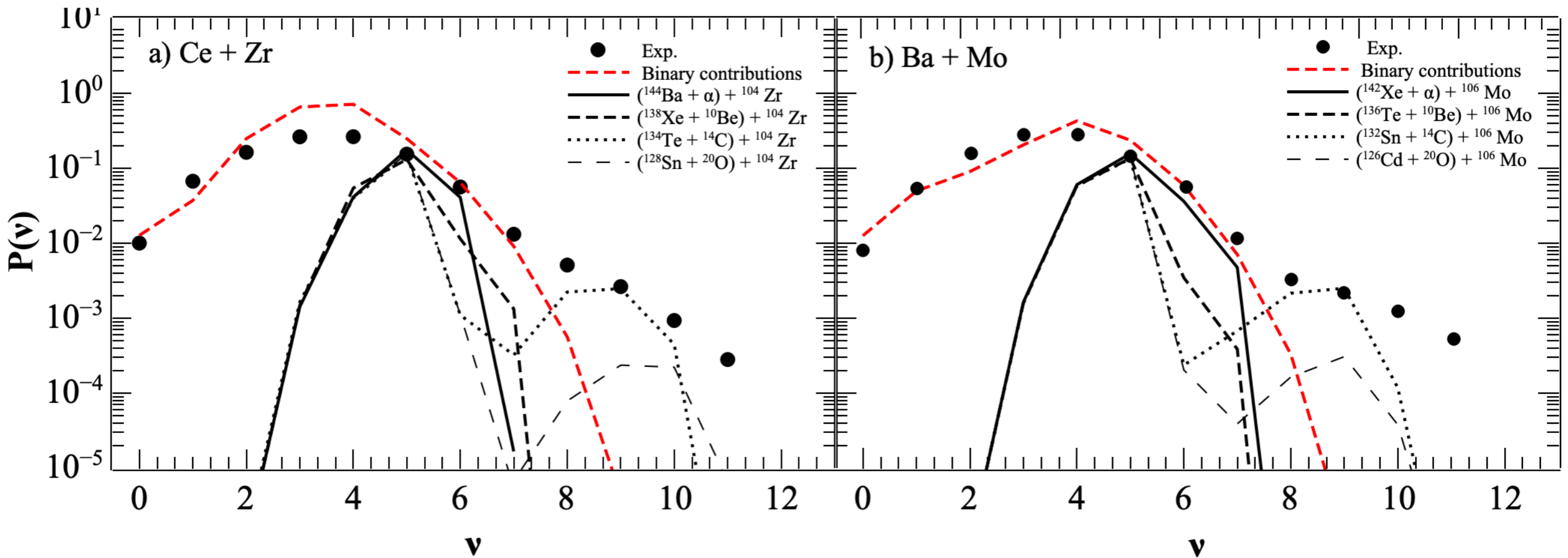
$$^{146}\text{Ba} + ^{106}\text{Mo} \quad \epsilon_{\text{Ba}}^* = E_{\text{Ba}}^* + U_{\text{Ba}}^{\text{def}} = 28.1 \text{ MeV}$$

	ΔV (MeV)	Q (MeV)	ϵ_H^{**} (MeV)	TKE (MeV)
$(^{144}\text{Ba} + \alpha) + ^{104}\text{Zr}$	8.3	1.1	31.3	162.0
$(^{138}\text{Xe} + ^{10}\text{Be}) + ^{104}\text{Zr}$	13.9	2.9	35.1	140.5
$(^{134}\text{Te} + ^{14}\text{C}) + ^{104}\text{Zr}$	30.8	-9.2	64.0	150.3
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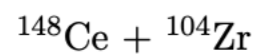
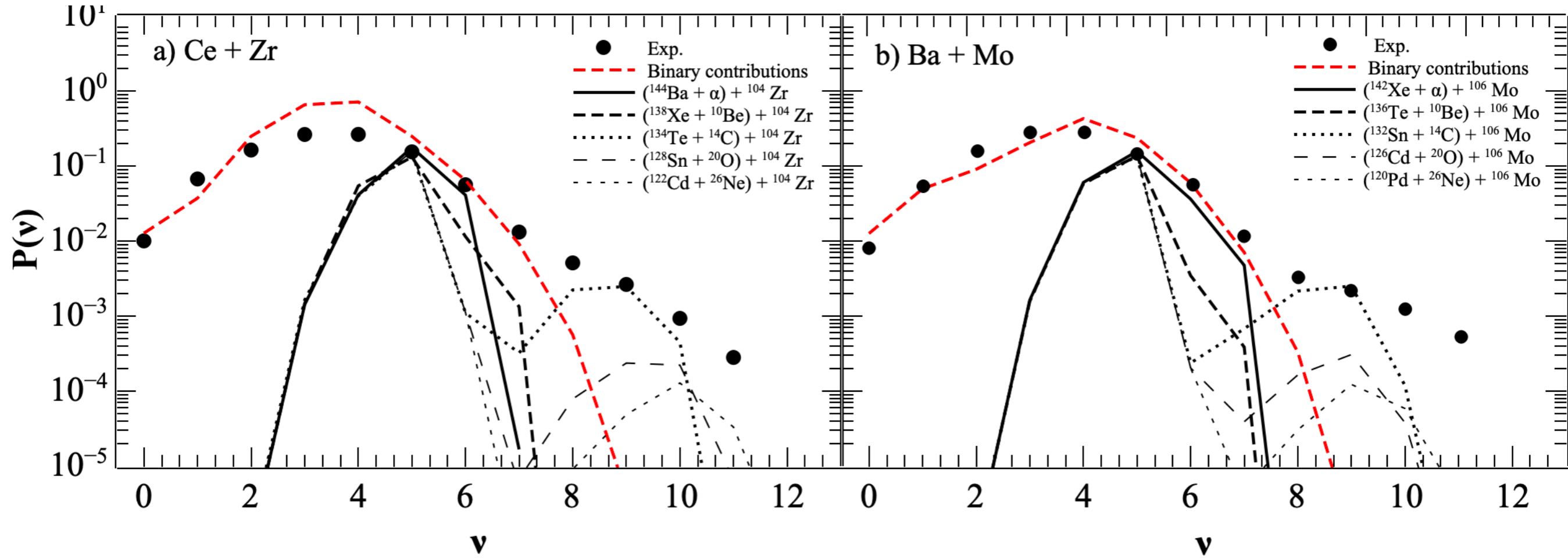
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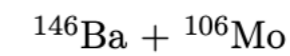
Competing mechanisms? - Neutron multiplicity



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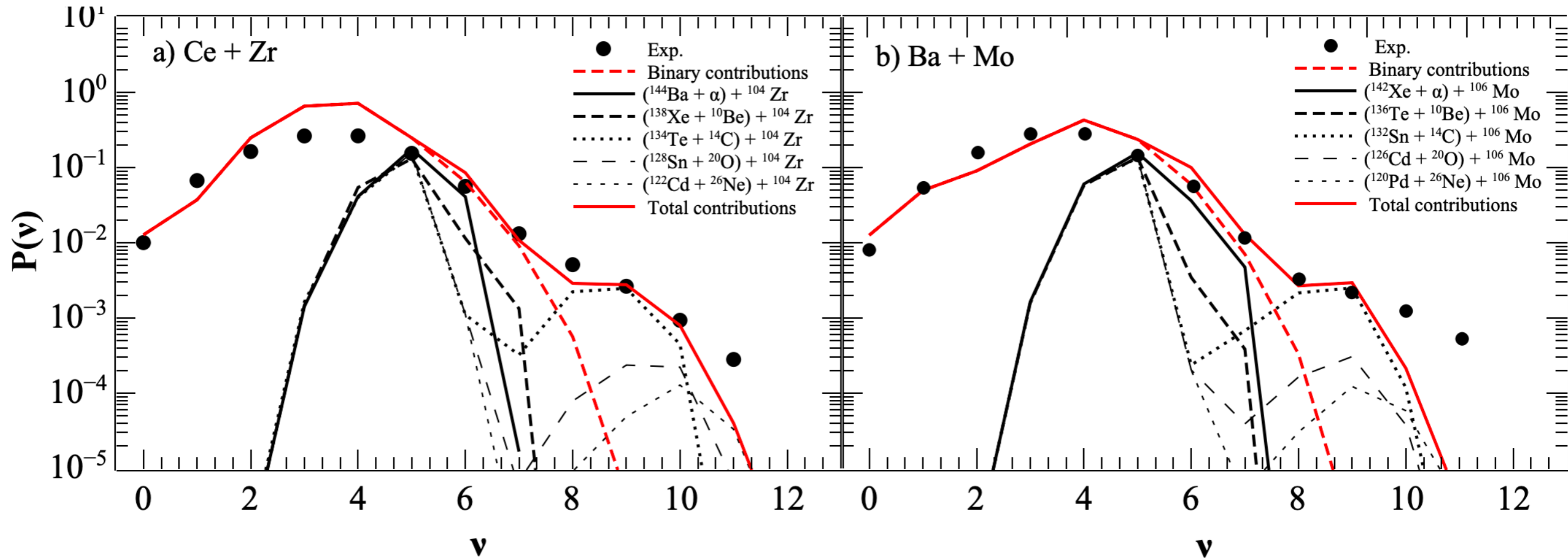
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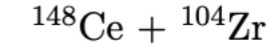
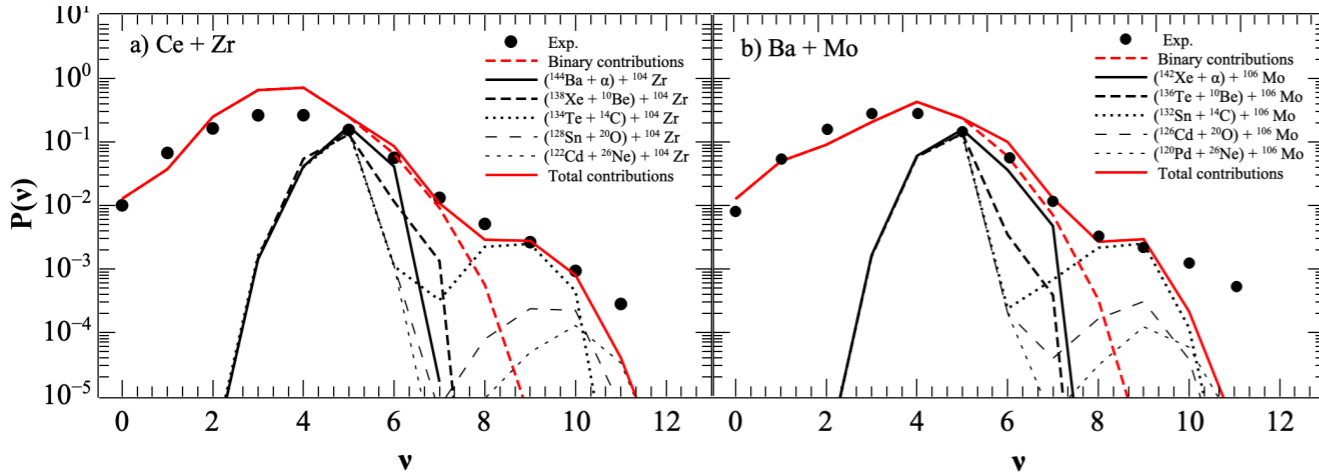
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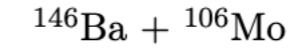
Focus: ^{252}Cf

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$\epsilon_{Ce}^* = E_{Ce}^* + U_{Ce}^{def} = 24.1 \text{ MeV}$

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Other observables in the 2nd mode for Ba-Mo:

Total kinetic energy: $\langle TKE \rangle_{exp} = 153 \pm 3 \text{ MeV}$
 $\langle TKE \rangle_{th} = 155,9 \text{ MeV}$

Neutron multiplicity: $\langle \nu \rangle_{exp} = 8$
 $\langle \nu \rangle_{th} \simeq 9$

Excitation energy: $\left(\frac{E_{Ba,mode1}^*}{E_{Ba,mode2}^*} \right)_{exp} = 2.6$
G.M. Ter-Akopian et.al., PRL77, 32 (1996)

$\left(\frac{E_{Ba,mode1}^*}{E_{Ba,mode2}^*} \right)_{th} = 2.4$

Focus: ^{252}Cf

Other fragmentations - is the second mode present?

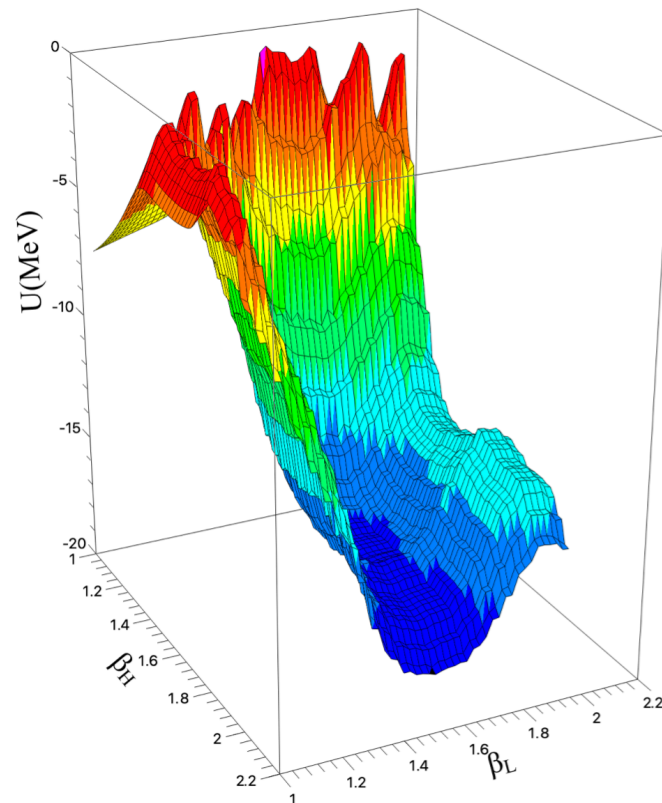
Nd+Sr fragmentation:

- Binary decay contribution much narrower than experimental data;
 - ↳ PES minimum is much narrower than previous cases.
 - ↳ example $P(\nu = 7)$ is at least two orders of magnitude less than previous cases.

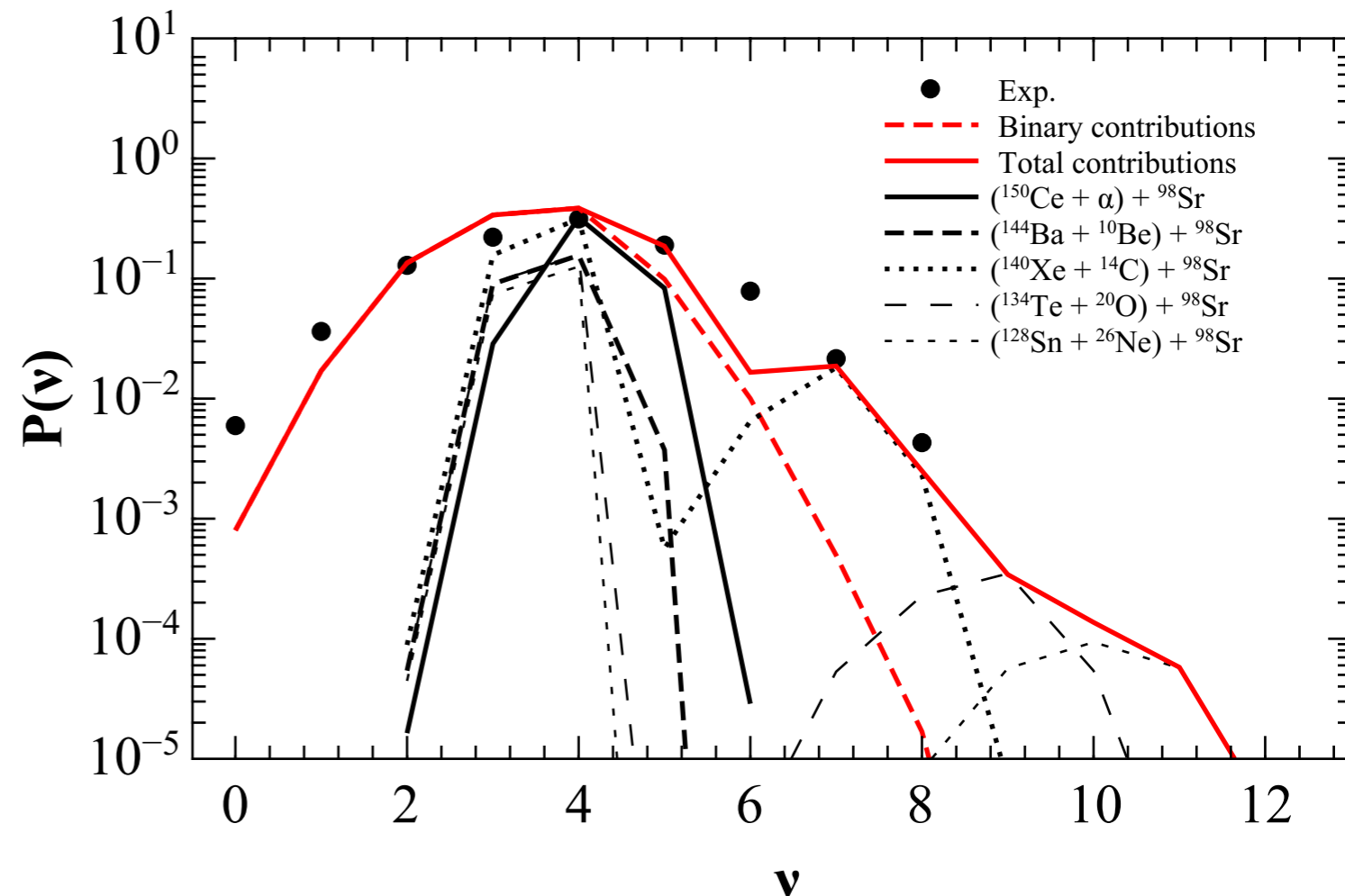
→ Most probable cluster is $\text{Nd} \rightarrow \text{Xe} + \text{C}$.

→ 2nd mode overlaps significantly with 1st mode.

→ O , Ne clusters offer just a “tail” due to low probabilities.



Configuration	ΔV	Q	ϵ^{**}
$(^{150}\text{Ce} + \alpha) + ^{98}\text{Sr}$	9,7	4,12	29,4
$(^{144}\text{Ba} + ^{10}\text{Be}) + ^{98}\text{Sr}$	14,7	6,53	31,9
$(^{140}\text{Xe} + ^{14}\text{C}) + ^{98}\text{Sr}$	28,75	-4,3	50,09
$(^{134}\text{Te} + ^{20}\text{O}) + ^{98}\text{Sr}$	41,8	-13,1	78,6
$(^{128}\text{Sn} + ^{26}\text{Ne}) + ^{98}\text{Sr}$	56,04	-17,2	97



Conclusions

- The mass, charge, TKE and neutron distributions can be calculated within the same model.
- Clusterisation of the heavy fragment can lead to a highly excited fission fragment.
- Competing mechanisms are possible, and they provide key signatures in $\langle n \rangle$ and TKE values.
- Evidence for cluster effects for can be “hidden” within existing experimental data.
- Question : Is this mechanism more common than previously thought ?

Thank you!

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