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From fusion and quasifission to fission and decay for SHE

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Background

- \Box Heavy-ion fusion reaction
	- \triangleright cold fusion: using ²⁰⁸Pb (²⁰⁹Bi) as target, the heaviest is up to Z=113 so far
	- \triangleright hot fusion: using ⁴⁸Ca as projectile, the heaviest is up to Z=118 so far
- □ Quasifission and fusion-fission are the primary mechanism in preventing SHE formation

S. G. Zhou, Nucl. Phys. Rev 34, 318 (2017)

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	- \triangleright hot fusion: using ⁴⁸Ca as projectile, the heaviest is up to Z=118 so far
- Quasifission and fusion-fission are the primary mechanism in preventing SHE formation
- \Box Fusion probability is several orders of magnitude difference among the different phenomenological models

Synthesis and identification of SHE involves the physical process

S. G. Zhou, Nucl. Phys. Rev 34, 318 (2017)

Content

Part I: cold and hot fusion

Part II: quasifission

Part III: fission

Part IV: alpha decay

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Part I: cold and hot fusion

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Part V: multinucleon transfer reactions

time-dependent Hartree-Fock

$$
S = \int_{t_1}^{t_2} \left\langle \Psi(t) \left| H - i \hbar \partial_t \right| \Psi(t) \right\rangle dt, \qquad H = \sum_{i=1}^A t_i + \sum_{i < i}^A v_{ij} \qquad \qquad \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \mathbb{E}_{t_3} \longrightarrow \mathbb{E}_{t_4} \longrightarrow \mathbb{E}_{t_5} \longrightarrow \mathbb{E}_{t_6} \longrightarrow \mathbb{E}_{t_7} \longrightarrow \mathbb{E}_{t_7} \longrightarrow \mathbb{E}_{t_8} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \mathbb{E}_{t_3} \longrightarrow \mathbb{E}_{t_4} \longrightarrow \mathbb{E}_{t_6} \longrightarrow \mathbb{E}_{t_7} \longrightarrow \mathbb{E}_{t_8} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \mathbb{E}_{t_3} \longrightarrow \mathbb{E}_{t_4} \longrightarrow \mathbb{E}_{t_5} \longrightarrow \mathbb{E}_{t_6} \longrightarrow \mathbb{E}_{t_7} \longrightarrow \mathbb{E}_{t_8} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \mathbb{E}_{t_3} \longrightarrow \mathbb{E}_{t_4} \longrightarrow \mathbb{E}_{t_6} \longrightarrow \mathbb{E}_{t_7} \longrightarrow \mathbb{E}_{t_8} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_9} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \mathbb{E}_{t_3} \longrightarrow \mathbb{E}_{t_1} \longrightarrow \mathbb{E}_{t_2} \longrightarrow \math
$$

$$
H = \sum_{i=1}^A t_i + \sum_{i
$$

 λ λ

 φ_{λ} \qquad

$$
\Psi(\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_A, t) = \frac{1}{\sqrt{A!}} \det \left| \varphi_\lambda(\mathbf{r}_i, t) \right|, \qquad \qquad i\hbar \frac{\partial \varphi_\lambda}{\partial t} = h\varphi_\lambda
$$

Advantages:

- ➢ **Fully microscopic, parameter-free theory in heavy-ion collisions;**
- ➢ **Nuclear structure and reactions in a unified framework (same EDF);**
- ➢ **Dynamical and quantum effects are automatically incorporated**; **Limitations:**
	- ➢ **Quantum tunneling is missing;**

 E_{ϵ}

Part I: cold and hot fusion dynamics

Theoretical framework for fusion evaporation reactions

Fusion evaporation reactions

$$
\sigma_{ER}(E_{c.m.}, x) = \int_0^1 d \cos(\theta_P) \int_0^1 d \cos(\theta_T)
$$

$$
\times \frac{\pi}{k^2} \sum_J (2J+1) T_J(E_{c.m.}, \theta_T, \theta_P)
$$

 $\times P_{\text{CN}}(\theta_P, \theta_T, E_{\text{c.m.}}, J)W_{\text{sur}}(E_{\text{CN}}^*, x, J).$

Capture process: Schrodinger equation for penetration probability

$$
\left[\frac{-\hbar^2}{2\mu}\frac{d^2}{dR^2} + \frac{J(J+1)\hbar^2}{2\mu R^2} + V_{\text{DC-FHF}}(R) - E_{\text{c.m.}}\right]\psi(R) = 0
$$

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$$

Fusion Process: fusion-by-diffusion (FbD) model for fusion probability

$$
P_{\text{CN}}(\theta_P, \theta_T, E_{\text{c.m.}}, J) = \frac{1}{2} \bigg[1 - \text{erf}\bigg(\frac{\Delta V_J(\theta_P, \theta_T)}{T_J(\theta_P, \theta_T)} \bigg) \bigg],
$$

Capture process: Schrodinger equation for penetration probability

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$$

Survival Process: statistical model for survival probability

$$
W_{\text{sur}}(E_{\text{CN}}^*, x, J) = P(E_{\text{CN}}^*, x) \prod_{i}^{x} \left(\frac{\Gamma_n(E_i^*, J)}{\Gamma_n(E_i^*, J) + \Gamma_f(E_i^*, J)} \right),
$$

Fusion dynamics involving halo nuclei

14,15C+²³²Th Fusion dynamics involving halo nuclei

 \triangleright Spherical neutron magic¹⁴C, deformed one-neutron halo ¹⁵C + ²³²Th;

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- ➢ Parameter-free microscopic calculations well reproduce the enhancement of cross sections at sub-barrier; reveal the underlying mechanism of this enhancement, which is driven by the halo structure of ${}^{15}C$;

Fusion dynamics involving halo nuclei

- \triangleright Spherical neutron magic¹⁴C, deformed one-neutron halo ¹⁵C + ²³²Th;
- ➢ Parameter-free microscopic calculations well reproduce the enhancement of cross sections at sub-barrier; reveal the underlying mechanism of this enhancement, which is driven by the halo structure of ${}^{15}C$;
- \triangleright One-neutron transfer probabilities are more sensitive to the orientations of 15C than 232Th;
- \triangleright Notable effect of halo structure on reaction dynamics.

Cold fusion dynamics

Compound-nucleus formation in cold-fusion reactions ⁴⁸Ca, ⁵⁰Ti, ⁵⁴Cr+²⁰⁸Pb

□ Above the capture barrier, our calculations reproduce the measured fusion probability reasonably well.

□ The restrictions from microscopic dynamics theory TDHF improve the predictive power of coupled-channels and diffusion calculation

X. X. Sun and Lu Guo, Phys. Rev. C 105, 054610 (2022)

 \Box The orientation effects of U are self-consistently included in the capture and fusion processes

- \Box The calculated capture cross sections agree well with the experimental data.
- The capture cross sections are strongly dependent on the orientations with several orders of magnitude

X. X. Sun and Lu Guo, Phys. Rev. C 107, 064609 (2023)

- \Box The TDHF evolutions with different orientations and incident energies are used to extract the injection distance, which is the only input of the fusion-by-diffusion model for fusion probabilities.
- \Box The fusion probabilities are strongly dependent on the orientations and the present calculations without any free parameters show that the tip-orientation collision is favorable for both the capture process and the formation of compound nucleus
- \Box our calculations reproduce the experimental evaporation-residue cross sections

X. X. Sun and Lu Guo, Phys. Rev. C 107, 064609 (2023)

Part II: quasifission

Quasifission is the main reaction channel hindering the SHE formation

 40,48 Ca+ 238 U PRL113,182502 (2014) 48 Ti+ 238 U PRL119,222502 (2017)

- ments with $210 \le A \le 214$ (see the text for details).
- \triangleright Experimental indication of shell effects initially came from mass-yield measurements;
- \triangleright To unambiguously confirm shell effects, proton or neutron numbers distributions have to be measured;
- \triangleright Atomic number distribution in the fragments have been measured, confirming the role of Z=82 magic shell;

Quasifission dynamics

First evidence for the impact of tensor force on the dynamical shell effects

$$
\sigma_{\lambda} \propto \int_{b_{\rm min}}^{b_{\rm max}} b \, db \int_0^{\frac{\pi}{2}} d\beta \sin(\beta) P_b^{(\lambda)}(\beta) ,
$$

 \Box the parameter-free microscopic TDHF study;

- \Box the charge distribution shows much better agreement with experiments for the SLy5t
- \Box the prominence of shell effects is manifested not only through shifts in peak positions but also through narrower yield distributions

L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Lett. B 833,137349 (2022) L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Rev. C (in press)

Quasifission dynamics

First evidence for the impact of tensor force on the dynamical shell effects

The inclusion of tensor force causes the spherical shell effects to become more prominent at Z=82, N=126

L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Lett. B 833,137349 (2022) L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Rev. C (in press)

Quasifission dynamics

First evidence that tensor force not only influences shell evolution in nuclear structure, but also in quasifission dynamics

L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Lett. B 833,137349 (2022) L. Li, Lu Guo, K. Godbey, A. S. Umar, Phys. Rev. C (in press)

Part III: fission

gure 6. The proton localization functions C_p for five configurations along the static fission pathway. The corresponding quadrupole moments e displayed on the top of each column. The upper panel shows the results with SLy5 and the bottom panel for SLy5t.

Y. Huang, X. X. Sun, and Lu Guo, Eur. Phys. J. A 60, 100 (2024) Y. Huang, X. X. Sun, and Lu Guo, Phys. Rev. C (accepted)

Figure 7. Similar to Fig. 6 but for the neutron localization functions C_n .

Part IV: alpha decay

 α + ²⁰⁸Pb DC-TDHF

 $T_{1/2} = \frac{\hbar \ln 2}{\Gamma}$

$$
\Gamma_{\alpha} = PF \frac{\hbar^2}{4\mu} \exp \left[-2 \int_{r_2}^{r_3} dr k(r)\right],
$$

$$
k(r)=\sqrt{\frac{2\mu}{\hbar^2}|Q-V(r)|}.
$$

Buck et al, Phys. Rev. Lett. 65, 2975 (1990)

S. C. Li and Lu Guo, Phys. Rev. C (in preparation)

 α + ²⁰⁸Pb DC-TDHF

Preliminary

S. C. Li and Lu Guo, Phys. Rev. C (in preparation)

Part V: multinucleon transfer reaction

producti**on of new exotic nuclei via multi-nucleon transfer reaction**

MNT may be the promising way in..., which are difficult to be produced by other methods.

FIG. 4. Experimentally deduced cross sections for the production of the $N = 126$ isotones as a function of the atomic number. The filled circles are from the present work and the filled stars are from the fragmentation of $^{208}Pb(1 \text{ GeV/nucleon}) + Be [28].$ The solid and dashed lines are to guide the eye.

TDHF studies of multi-nucleon transfer reaction

TDHF gives the average number of transferred nucleons for all reaction channels.

How to obtain the transferred nucleon number for each reaction channel?

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How to obtain the transferred nucleon number for each reaction channel?

particle number projection method (PNP)

particle number projection operator: transfer probability for each channel:

nucleon transfer reaction
ldies of multi-nucleon transfer re
he average number of transferred nucleons f
in the transferred nucleon number for each re
number projection method (PNP)
ber projection operator: transfer proba

$$
\hat{P}_n = \frac{1}{2\pi} \int_0^{2\pi} d\theta e^{i(n-\hat{N}_{V_P})\theta}
$$
is setion of primary products:
DHF+PNP)

transfer cross sction of primary products:

(TDHF+PNP)

ucleon transfer reaction

\nlies of multi-nucleon transfer reaction

\naverage number of transferred nucleons for all reaction channels

\nthe transferred nucleon number for each reaction channels

\nimber projection method (PNP)

\ner projection operator:

\n\n- transfer probability for each channel:
\n- $$
\sum_{n} \sum_{n} \int_{0}^{2\pi} d\theta e^{i(n-\hat{N}_{V_p})\theta} \qquad P_n = \langle \Psi | \hat{P}_n | \Psi \rangle
$$
\n
\nsection of primary products:

\n\n- $$
\sigma_{N,Z}(E) = 2\pi \int_{b_{min}}^{b_{min}} b P_{N,Z}(b, E) db
$$
\n- $$
E_{\text{F+PNP}}
$$
\n
\nExample 2. Show that, 105, 192701 (2010)

 $P_n = \langle \Psi | \hat{P}_n | \Psi \rangle$

C. Simenel, Phys. Rev. Lett. **105**, 192701 (2010)

TDHF studies of multi-nucleon transfer reaction

TDHF gives the average number of transferred nucleons for all reaction channels n
 n n i c n n n n i n c n i c n i c n i c n i c c *n* **c c** *n* **c c** *n* **c c** *n* **c c** *n* **c** *n* **c c** *c n* **c** *n* **c** *n* **c** *c* **Cleon transfer reaction**

so of multi-nucleon transfer reaction

verage number of transferred nucleons for all reaction channels

beer projection method (PNP)

brojection operator: transfer probability for each channel:

How to obtain the transferred nucleon number for each reaction channel?

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indles of multi-nucleon transfer reaction

\nudies of multi-nucleon transfer re

\nthe average number of transferred nucleons from the transfered nucleons of number projection method (PNP)

\nwhere projection operator: transfer proba

\n
$$
\hat{P}_n = \frac{1}{2\pi} \int_0^{2\pi} d\theta e^{i(n-\hat{N}_{V_P})\theta}
$$
\nas section of primary products:

\nDHF+PNP

\nall decay method (GEMINI++)

\narison with experimental measurement

\n=
$$
\sum_{N \geq N'} \sum_{Z \geq Z'} P_{N,Z} P_{\text{decay}}(E_{N,Z}^*, J_{N,Z}, N, Z)
$$
\npass section of secondary products:

\nDHF+PNP+GENINI)

transfer cross sction of primary products: (TDHF+PNP)

particle in proposition in the proposition operator: transfer probability for each channel: $\mathbf{D} = \Delta \mathbf{U} \mathbf{I} + \mathbf{\hat{D}} \mathbf{I} \mathbf{I} \mathbf{I}$

$$
\Gamma_n = \langle \mathbf{T} | \Gamma_n | \mathbf{T} \rangle
$$

$$
\sigma_{N,Z}(E) = 2\pi \int_{b_{\min}}^{b_{\text{cut}}} b P_{N,Z}(b,E) db
$$

Statistical decay method (GEMINI++)

direct comparison with experimental measurement

$$
P_{{}_{N',Z'}} = \sum_{N \geq N'}\sum_{Z \geq Z'} P_{_{N,Z}} P_{\text{decay}}(E_{_{N,Z}}^*, {J}_{_{N,Z}}, {N}, {Z}; {N}', {Z}')
$$

transfer cross section of secondary products:

(TDHF+PNP+GEMINI)

$$
\sigma_{_{N',Z'}}(E)\,{=}\,2\pi{\textstyle \int}_{b_{\rm min}}^{b_{\rm cut}}\!bP_{_{N',Z'}}(b,E)db
$$

- ➢Parameter-free and fully microscopic calculations;
- ➢ TDHF+PNP+GEMINI method well accounts for the experiment.;

Z. J. Wu and Lu Guo, Phy. Rev. C 100, 014612 (2019)

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- ➢ TDHF+PNP+GEMINI method well accounts for the experiment.;

Z. J. Wu and Lu Guo, Phy. Rev. C 100, 014612 (2019)

Production of proton-rich nuclei in ¹⁰⁰Sn region

- ¹⁰⁰Sn has an unexpected large Gamow-Teller strength measured ;
- \Box Experiments identify a new island of cluster radioactivity;
- End point of astrophysical rp-process (Sn-Sb-Te cycle);

Conventional experimental methods:

- \Box Fragmentation process
- \square Fusion-evaporation reaction
- Fusion: ⁵⁴Fe(⁵⁸Ni,4n)¹⁰⁸Xe, ¹⁰⁸Xe→¹⁰⁴Te→¹⁰⁰Sn α-decay chain;
- The production cross section of $108Xe$ is less than 1nb; PRL121, 182501 (2018);

Nature 486, 341 (2012); PRL116, 162501 (2016); PRL 97, 082501 (2006); PRL121, 182501 (2018);

Explore new possibility to reach and beyond proton drip line

 σ of Sn isotopes in MNT are several orders of magnitudes higher, indicating an enormous advantage to discover new protonrich isotopes;

Combined effect of several factors, i.e., shell structure, neck dynamics, charge equilibrium;

Z. J. Wu, Lu Guo, Zhong Liu, GX Peng, Phys. Lett. B 825, 136886 (2022)

Production of proton-rich nuclei in ¹⁰⁰Sn region

MNT reaction ⁵⁸Ni+¹¹²Sn: energy dependence

 \Box About 19 new proton-rich isotopes with cross sections of larger than 1 nb are predicted to be produced

Z. J. Wu, Lu Guo, Zhong Liu, GX Peng, Phys. Lett. B 825, 136886 (2022)

Possible production of heavy and superheavy elements $238U+248Cm$

With an appropriate reaction system, the projectile evolves to the doubly magic nuclide and transfer neutrons to the target, and consequently very neutron-rich heavy nuclides and particularly very neutron-rich superheavy nuclides are expected to be produced.

Possible production of heavy and superheavy elements $238U+248Cm$ Ternary quasifission

N=89.66 N=112.01 N=97.42

Summary

- ➢ Develop the microrscopic theoretical methods to study SHE production
- \triangleright The fusion mechanism involving the exotic nuclei
- \triangleright Cold fusion reactions and hot fusion reactions for the production of superheavy
- ➢ Quasifission mechanism and shell effects
- ➢ Multinucleon transfer reactions

Perspectives

- ➢ basis space, especially with pairing (TDHFB, TDHF+BCS)
- \triangleright beyond the independent-particle approximation (TDRPA)
- \triangleright the role of two-body collisions and fluctuation effects

Thanks for your attention!

University of Chinese Academy of Sciences