# The inclusive $\pi^0$ on the SPD-ECAL endcaps for online polarimetry

Katherin Shtejer Díaz

SPD Physics & MC Meeting 23.10.2024

□ Vertex assumed at (0, 0, 0): Gaussian smeared:  $\sigma_z = 30 \ cm$  and  $\sigma_{x,y} = 0.1 \ cm$ 

#### Generation

- ☐ SpdRoot version 4.1.6
- $pp @ \sqrt{s} = 27 \text{ GeV}$
- ☐ Particle generator: Pythia 8 (number of events:  $1.8 \cdot 10^8$ )
- Minimum Bias

#### Realistic reconstruction

- ☐ Focus on the "ECAL" reconstructed particle
- Identified the cluster to which the particle belongs
- Position and energy taken from cluster<sup>(\*)</sup>
- Selected clusters that belong to the ECAL endcaps
- (\*) Cluster splitting is not available yet.

$$\sqrt{s} = 27 \text{ GeV}$$

$$\mathcal{L} \approx 10^{32} cm^{-2} s^{-1}$$
  $\sigma_{pp} = 40 \text{ mb}$ 

$$\mathcal{R} = \mathcal{L} \cdot \sigma = 4 \cdot 10^6 s^{-1}$$

$$N_{ev} = 1.8 \cdot 10^8$$

$$t_{mc} = N_{ev} \cdot \frac{1}{\mathcal{R}} = 46.8 \, sec$$

#### Analysis

- $\square$  Cuts:  $E_{\rm V} > 400$  MeV,  $p_{\rm T} > 0.5$  GeV/c
- $\square$  Candidates to  $\pi^0$  selected from  $\gamma\gamma$  combinations (invariant mass)
- Photon candidates:
  - ✓ no especial constraint is applied to select photons (i.e. pdg-based filtering)
  - $\checkmark$  candidates to  $\pi^0$  selected from all possible  $\gamma\gamma$  combinations (invariant mass)
- $\Box$  Fitting the invariance mass distribution: gausn + pol2

$$f(x) = [p0] \cdot exp(-0.5 \cdot ((x - [p1])/[p2]) \cdot ((x - [p1])/[p2]))/(sqrt(2\pi) \cdot [p2])$$
$$+ [p3] + [p4] \cdot x + [p5] \cdot pow(x, 2)$$

- ☐ The yield is extracted from the integral of the fit function in certain limits, using the parameters of the Gaussian "signal" peak.
- ☐ The integral errors are calculated using the parameter uncertainties and the covariance matrix obtained from the fit.

$$p^{\uparrow} + p \to \boldsymbol{\pi^0} + X \qquad \phi = 2\pi$$

The cross section of hadron production in polarized  $p^{\uparrow}+p$  collisions, is modified in azimuth.

$$\frac{d\sigma}{d\varphi} = \frac{d\sigma}{d\varphi_0} \left[ 1 + P \cdot A_N \cdot \cos(\varphi + \varphi_0) \right]$$

Azimuthal cosine modulation

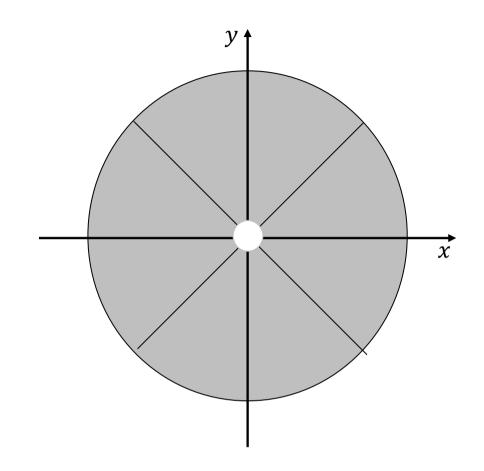
$$N_{\pi^0}(\varphi) = A[1 + P \cdot A_N \cdot \cos(\varphi + \varphi_0)]$$

$$A_N = \frac{Amp}{P}$$

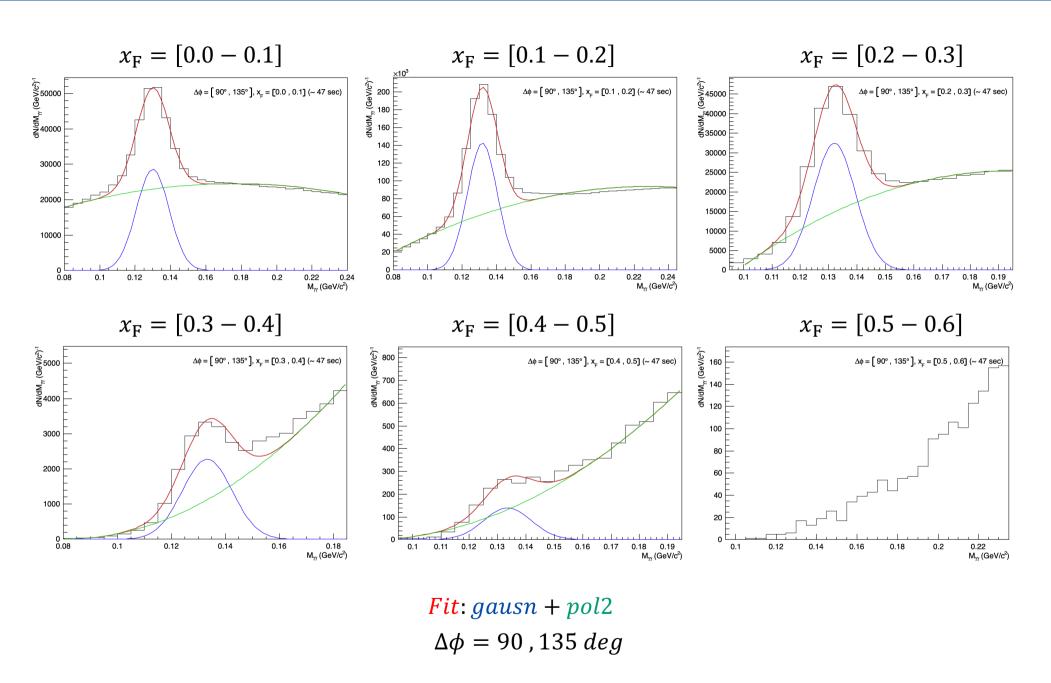
 $N_{\pi^0}(\varphi)$ : Yield of  $\pi^0$ 

P: Beam polarization

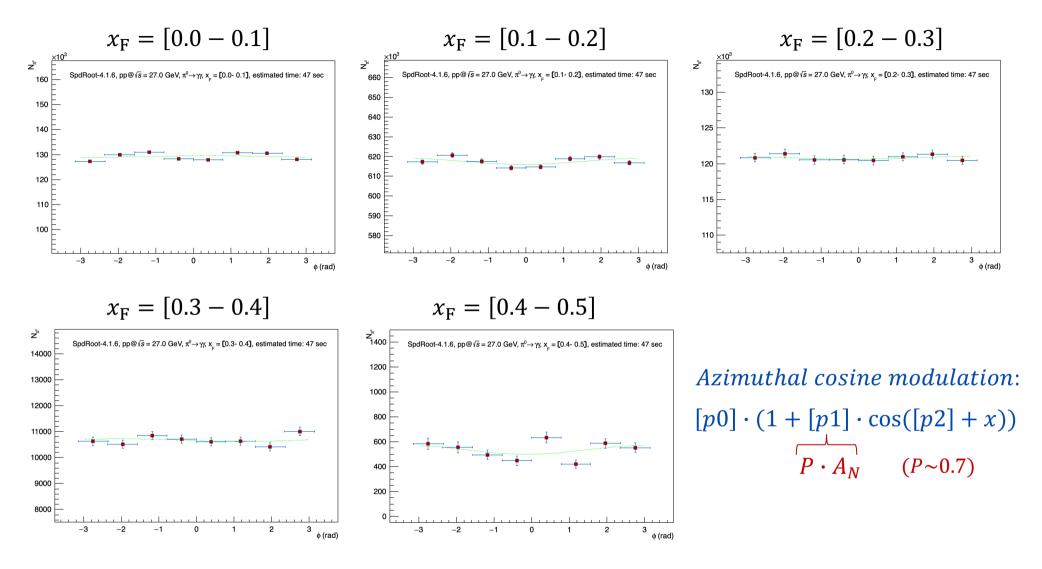
P = 0.7 was assumed



The spin dependent  $\pi^0$  yields for each bin are extracted from the invariant mass spectra in different  $x_{\rm F}$  sub-ranges for each  $\varphi$  bin.



### Azimuthal cosine modulation of $\pi^0$ yields in $x_F$ intervals



The modulation size is expected to be zero in unpolarized Monte Carlo simulations.

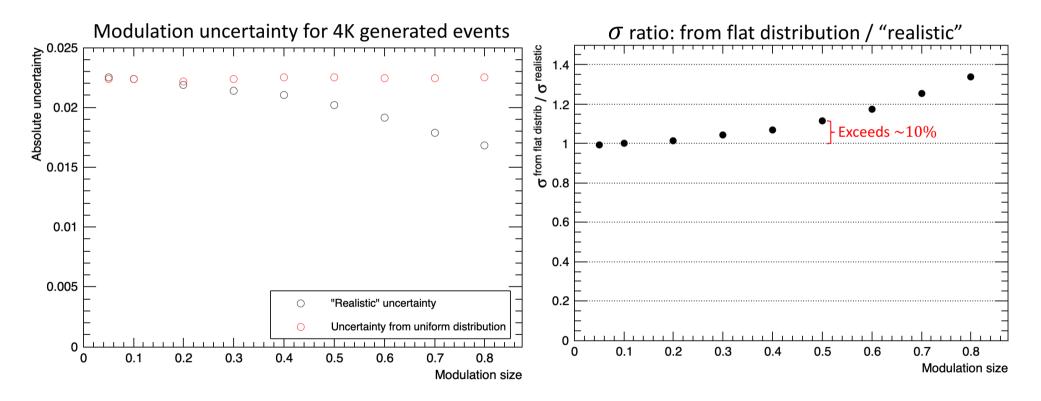
How reliable is to extract the statistical uncertainty of the amplitude modulation of a flat distribution?

#### Statistical uncertainty of uniformly distributed $A_N$ ?

#### Contribution from Igor Denisenko!

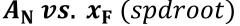
The reliability of extracting the statistical uncertainty of  $A_{\rm N}$  from the fit  $c \cdot (1 + A_{\rm N} \cdot \cos(\phi + b))$  of a flat distribution is evaluated using a toy modelling.

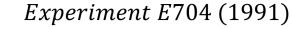
- Two distributions are generated  $\begin{cases} f = 1 + [0] \cdot \cos(x) \cdot [-\pi, \pi] \\ f_0 = 1 \ [-\pi, \pi] \end{cases}$
- Both are fitted with a cosine modulation function
- lacktriangle The  $\sigma_{A_{\mathbf{N}}}$  is extracted in both cases

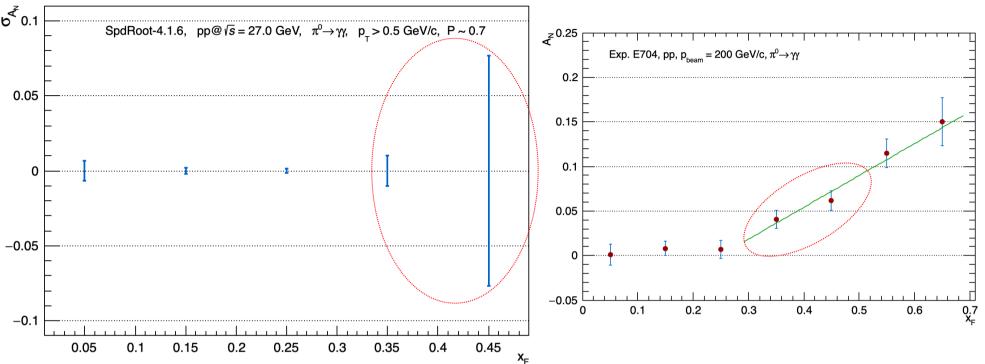


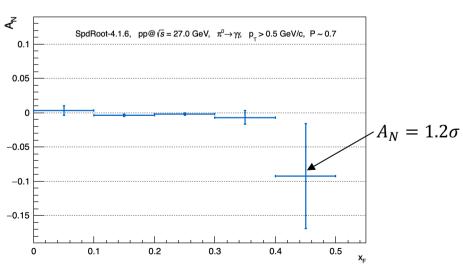
The statistical uncertainty of the amplitude modulation can be reasonably estimated for  $A_N \approx 0$ 





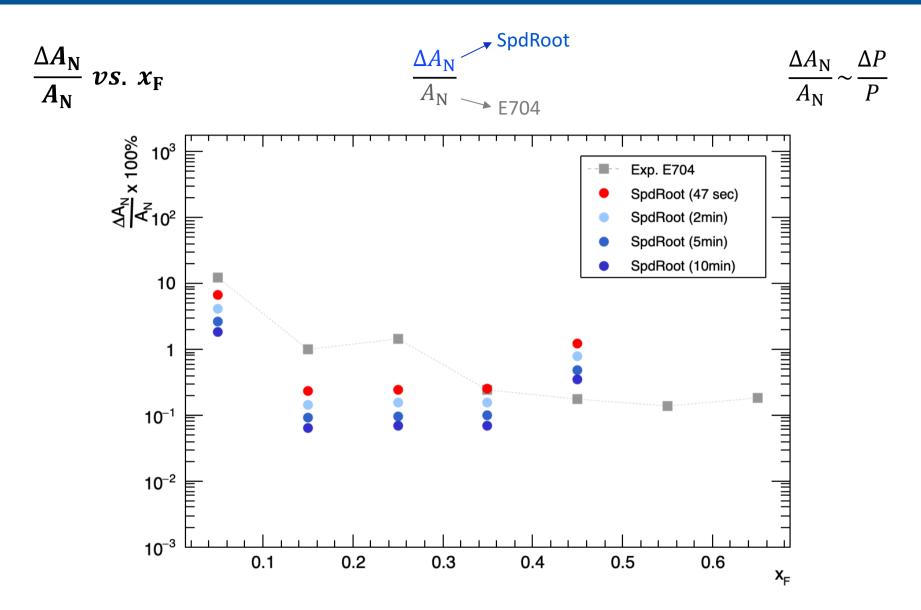








$$\frac{\Delta A_{\rm N}}{A_{N}} \sim \frac{\Delta P}{P}$$

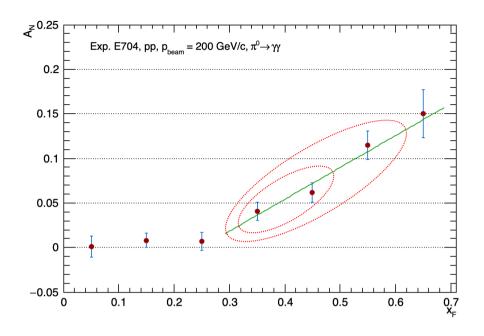


Better precision of the polarization measurement expected at:  $0.1 < x_F < 0.2 \ (\sqrt{s} = 27 \text{ GeV})$ 

#### Estimated relative error of the polarization

$$\frac{\Delta A_{\rm N}}{A_N} \sim \frac{\Delta P}{P}$$

$$\frac{\Delta P}{P} = \frac{1}{\sqrt{\sum_{i} \left(\frac{A_{N_i}}{\Delta A_{N_i}}\right)^2}}$$



Taking **3** experimental points (0.3  $\leq x_F <$  0.6):

$$\frac{\Delta P}{P} = 0.0998 \rightarrow 9.9 \% \text{ (Experiment E704)}$$

Taking **2** experimental points (0.3  $\leq x_F <$  0.5):

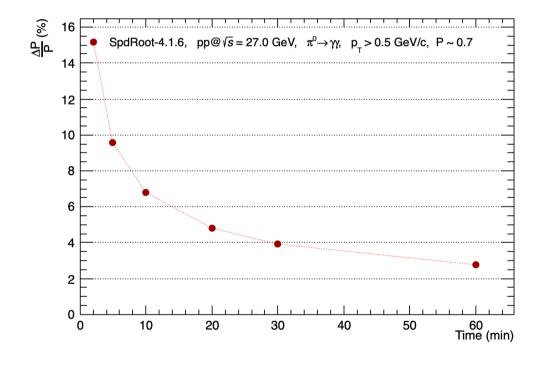
$$\frac{\Delta P}{P} = 0.1434 \rightarrow 14.3 \% \text{ (Experiment E704)}$$

The error of the beam polarization in the experiment **E704** is estimated in **10%** (FERMILAB-Pub-91/15-E[E581,E704])

### Estimated relative error of the polarization

Estimation of the statistical accuracy of the beam polarization measurement, with  $pp \to \pi^0 X$  at  $\sqrt{s}=27$  GeV, in SPD ECAL endcaps.

Estimated time	$\frac{\Delta P}{P}$
2 min	15.1 %
5 min	9.6 %
10 min	6.8 %
20 min	4.8 %
30 min	3.9 %
1 h	2.8 %



#### Summary of the first part

- ✓ The energy and position of  $\pi^0$  decayed photons in the endcaps of the SPD ECAL are quantities which are accessible online, with no necessity of particle identification or vertex reconstruction.
- ✓ The accuracy of the beam polarization has been estimated for pp collisions at  $\sqrt{s} = 27$  GeV by Monte Carlo simulations based on SpdRoot-4.1.6
- ✓ Based on the azimuthal asymmetry of  $\pi^0$  detected in the ECAL-endcaps, the accuracy of the beam polarization has been estimated at **9.6% for 5 min.** of data taking, assuming an **average polarization of 0.7**.

### Another approach: projected statistical uncertainty of $A_N$

Raw asymmetry:

$$A_{N}(\phi) = \frac{1}{P\langle|\cos(\phi)|\rangle} \frac{N^{\uparrow}(\phi) - \mathcal{R} \cdot N^{\downarrow}(\phi)}{N^{\uparrow}(\phi) + \mathcal{R} \cdot N^{\downarrow}(\phi)}$$

 $N(\phi)$ : counts in  $\phi$  bins

**P**: beam polarization

 $\frac{1}{\langle |cos(\phi)| \rangle}$ : azimuthal acceptance correction factor

 $\langle |cos(\phi)| \rangle = \frac{\int_{\phi_1}^{\phi_2} \cos(\phi) d\phi}{\phi_2 - \phi_1}$ : average of the cosine of azimuth in the  $\phi$  bin

 $\mathcal{R} = \mathcal{L}^{\uparrow}/\mathcal{L}^{\downarrow}$ : relative luminosity



$$\sigma_{A_{N}}(\phi) = \frac{1}{P\langle|\cos(\phi)|\rangle} \frac{1}{\sqrt{2N}}$$

$$\mathcal{R} = \mathcal{L}^{\uparrow}/\mathcal{L}^{\downarrow} \text{: relative luminosity}$$

$$\mathbf{Statistical \ uncertainty \ of \ } A_{\mathbf{N}} \text{:} \boxed{ \begin{aligned} \sigma_{A_{\mathbf{N}}}(\phi) &= \frac{1}{P\langle|\cos(\phi)|\rangle} \frac{1}{\sqrt{2N}} \end{aligned}} \quad \begin{cases} \mathcal{R} \sim 1 \\ N^{\uparrow} \sim N^{\downarrow} &= N \\ \sigma_{N} &= \sqrt{N} \text{ : Poisson distribution of } N \end{aligned}}$$

0°3 counts/(0.001 GeV/c²) 00 01 05 05

0<sup>-3</sup> counts/(0.01 GeV/c<sup>2</sup>

 $_{100}$  (b)  $\eta \rightarrow \gamma \gamma$ 

0.2 0.3 0.4 0.5 0.6

M<sub>ss</sub> [GeV/c<sup>2</sup>]

 $M_{yy}$  [GeV/c<sup>2</sup>]

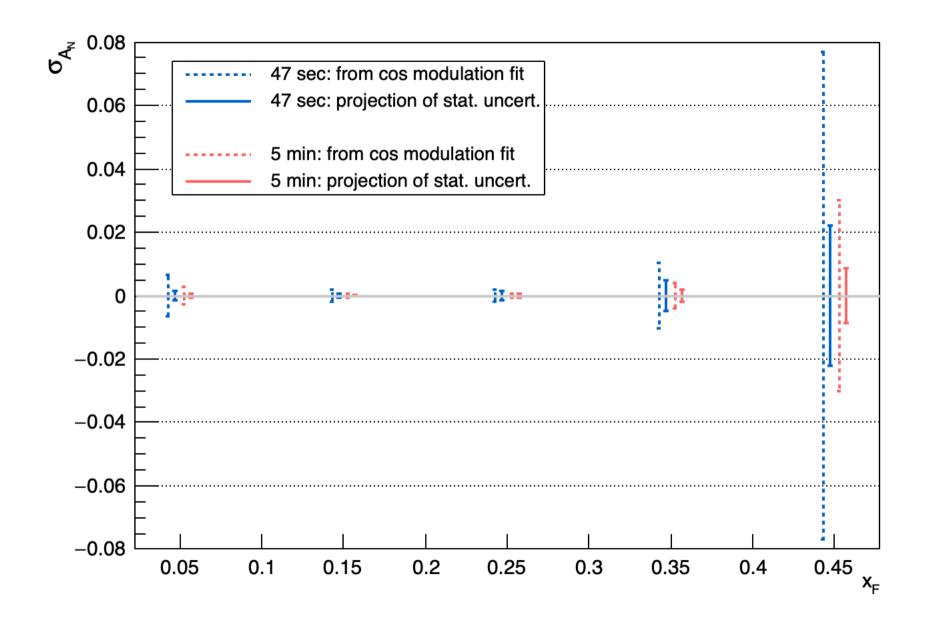
**PHENIX** 

0.25

The statistical uncertainties estimated independently for each  $\phi$  bin,  $\sigma_{A_N}(\phi)$ , can be averaged as:

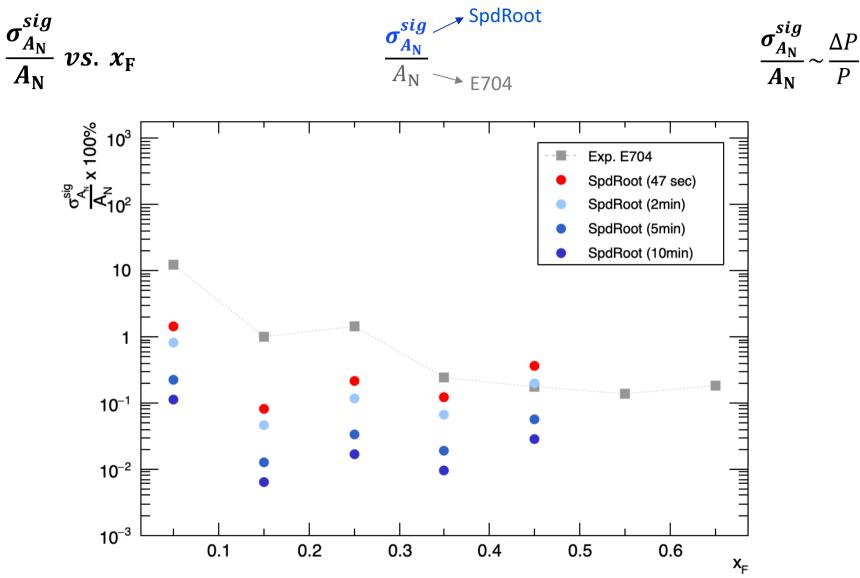
$$\sigma_{A_N^{sig}}(x_F) = \frac{1}{\sqrt{\sum_{i=1}^8 \frac{1}{\sigma_{A_N^{sig}}^2(\phi_i)}}}$$

Amaresh's Analysis Note "Prospects of Open-Charm Asymmetry Measurements at the SPD" (indico.jinr.ru/event/4594/attachments/18860/32246/D Meson Report.pdf)



#### Another approach: projected statistical uncertainty of $A_N$

#### Relative error for $A_N$ , $pp @ \sqrt{s} = 27 \text{ GeV}$



Better precision of the polarization measurement expected at:  $0.1 < x_F < 0.2 \ \ (\sqrt{s} = 27 \ {\rm GeV})$ 

## Method: Calculation of the projected statistical uncertainty of $A_{\rm N}$ .

Estimated time	$\frac{\Delta P}{P}$
2 min	6.4 %
5 min	1.8 %
10 min	0.9 %
20 min	0.5 %
30 min	0.3 %
1 h	0.2 %

Correction for the background is needed!

#### Method:

Extracting the  $A_{\rm N}$  from the modulation amplitude of the cosine function.

Estimated time	$\frac{\Delta P}{P}$
2 min	15.1 %
5 min	9.6 %
10 min	6.8 %
20 min	4.8 %
30 min	3.9 %
1 h	2.8 %