

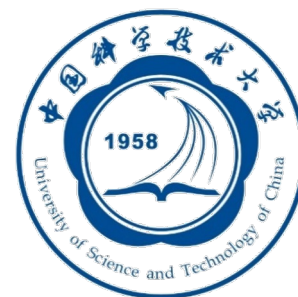


2024

# Heavy Flavor and Quarkonia

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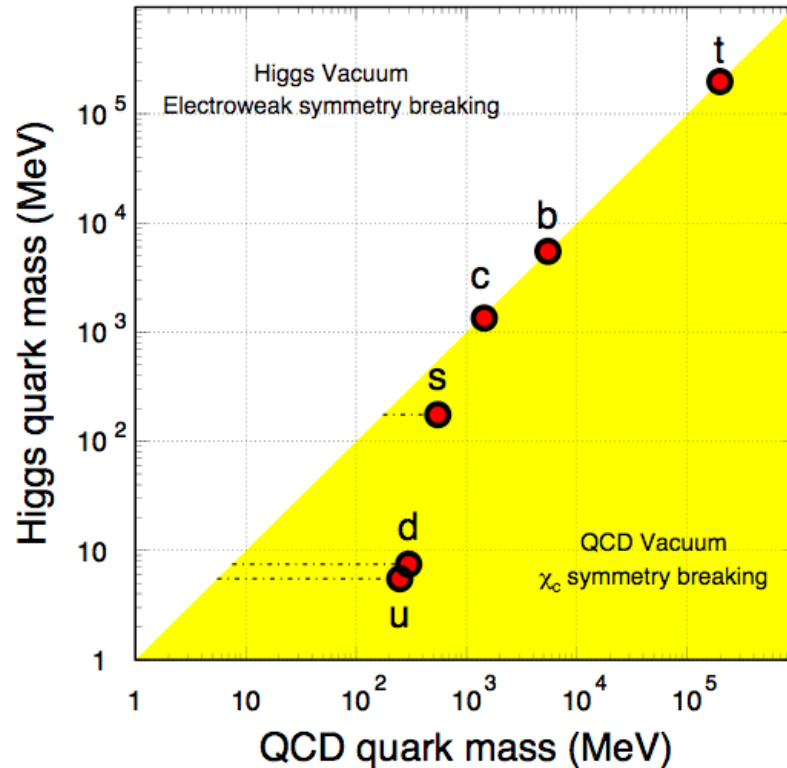


November 25, 2024



# What Does “Heavy” Mean

**Heavy Flavor:** quarks with large masses, usually refer to **charm** and **bottom** quark



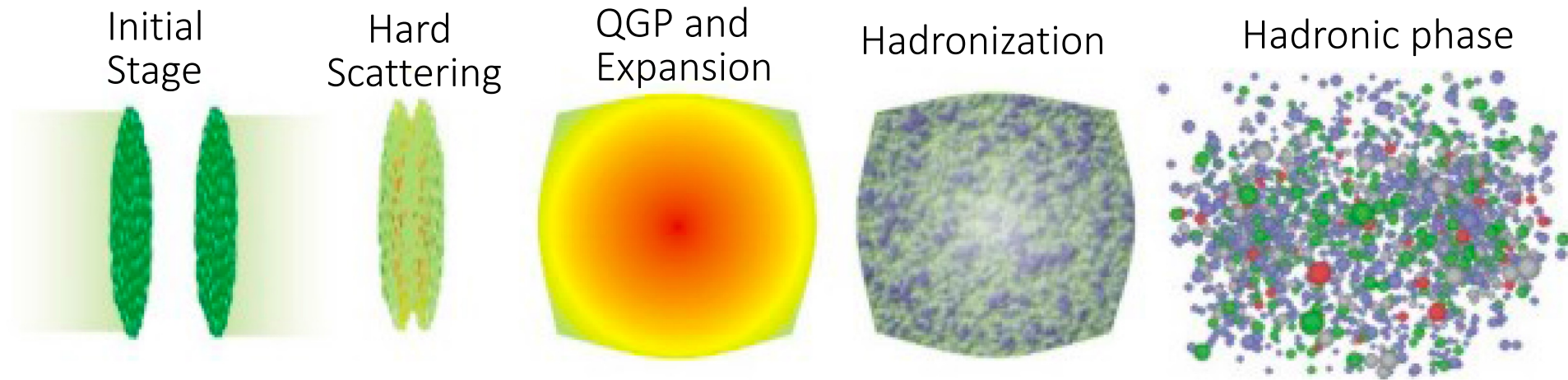
*X. Zhu et al., PLB 647, 366 (2007)*

Strong interactions do not affect heavy quark masses

Top quark has too short lifetime

- $\sim 0.15 \text{ fm}/c \ll \text{QGP formation time}$
- Irrelevant to heavy ion collision physics (?)

# Pros of Being Heavy

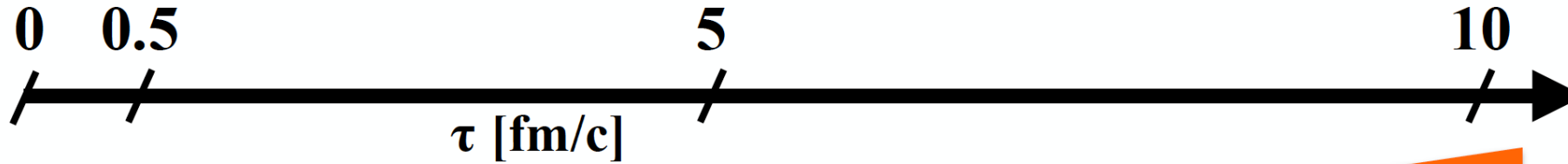


- $t \sim \frac{1}{m_c}, \frac{1}{m_b} < \tau_{0,QGP}$ : Produced early
- $m_c, m_b \gg \Lambda_{QCD}$ : Produced in initial hard scattering calculable in pQCD
- $m_c, m_b \gg T_{QGP}$ : Production in QGP is negligible

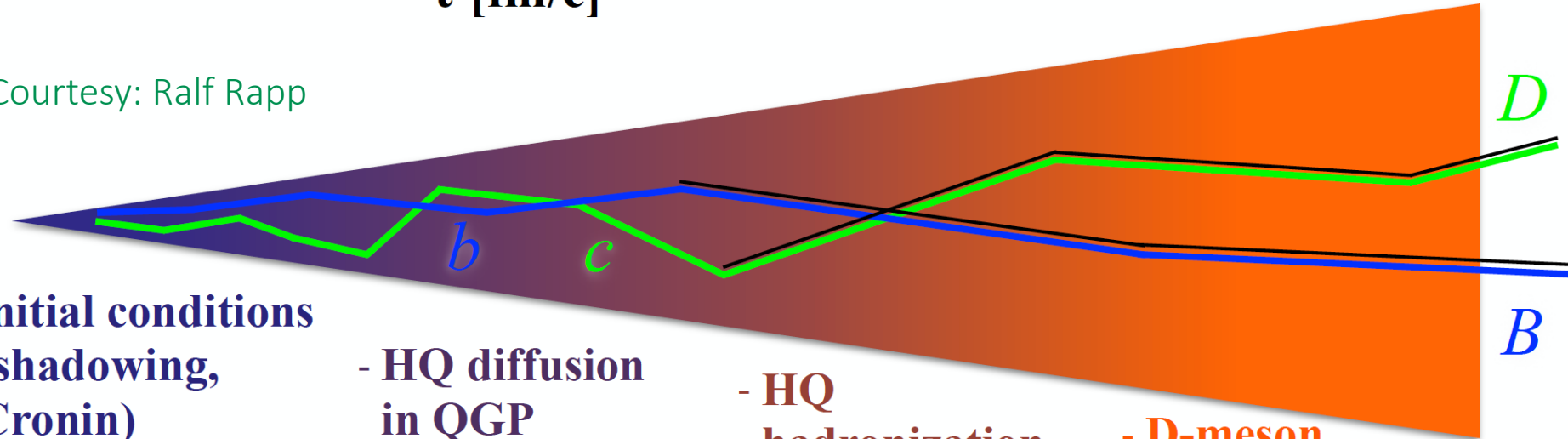
**Penetrating probe:** experience the whole evolution of QGP



# A Journey of HQ in Heavy Ion Collisions



Courtesy: Ralf Rapp



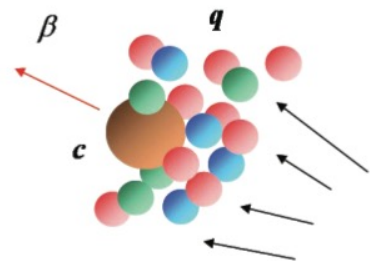
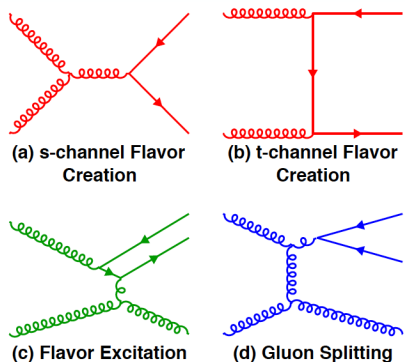
# of heavy quarks is constant throughout the evolution

- initial conditions (shadowing, Cronin)
- Pre-equil. fields

- HQ diffusion in QGP

- HQ hadronization

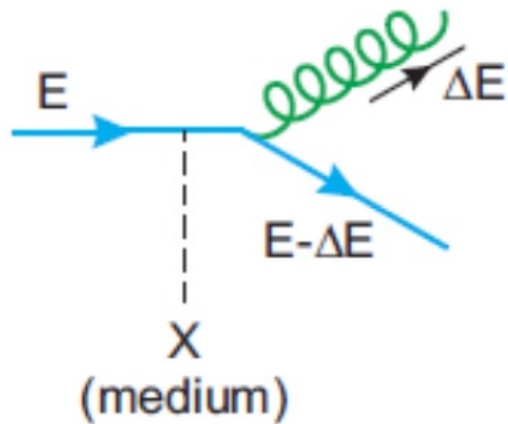
- D-meson diffusion at hadronic phase



Probe the medium with final heavy flavor **hadrons**



## Radiative energy loss Inelastic scattering



*D. d'Enterria and B. Betz,  
Lect. Notes Phys. 785, 285 (2010)*

Energy lost by radiation of gluons induced by interactions with the hot and dense medium

$$\langle \Delta E \rangle \propto \alpha_s C_R \hat{q} L^2$$

$\hat{q}$ : Transport coefficient of the QGP medium

$$\hat{q} = \frac{\langle k_T^2 \rangle}{\lambda}$$

Relate to the energy (gluon) density of the medium

$C_R$ : Casimir coupling factor

4/3 for quark-gluon coupling

3 for gluon-gluon coupling

$$\langle \Delta E \rangle_g > \langle \Delta E \rangle_q$$



# Heavy Quark Energy Loss in QGP

“Dead-Cone” effect: Gluon radiation in vacuum is suppressed for small angle due to kinematical constraints

*Y. Dokshitzer, D. Kharzeev, PLB519, 199 (2001)*

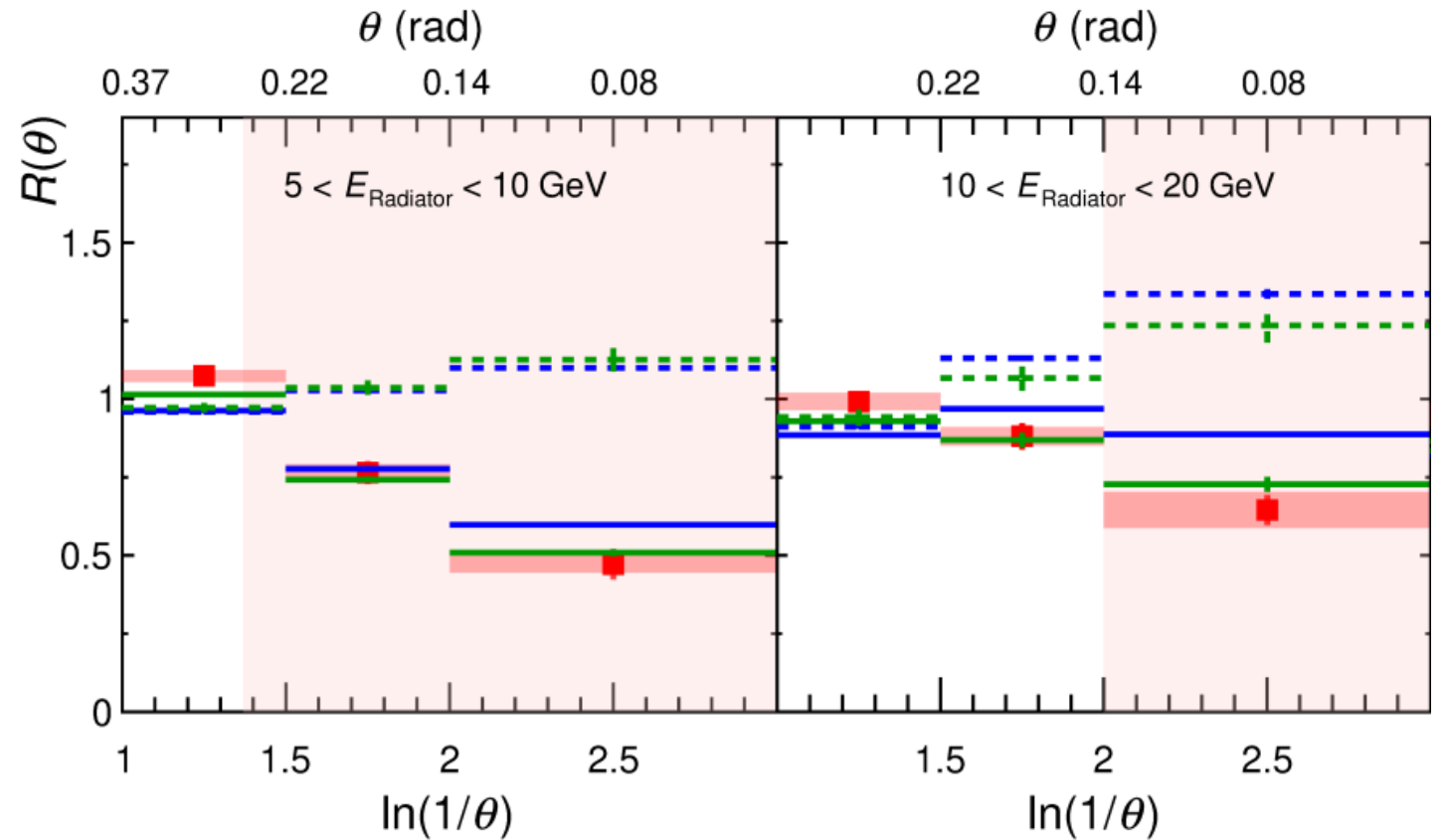
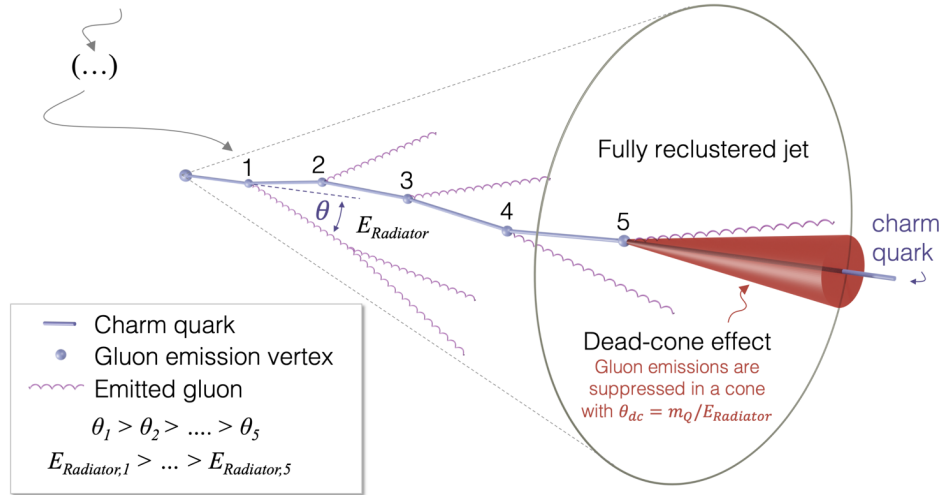


$$\langle \Delta E \rangle = \frac{\langle \Delta E \rangle_0}{(\theta^2 + \theta_0^2)^2} \quad \theta_0 = \frac{1}{\gamma} = \frac{m_q}{E}$$

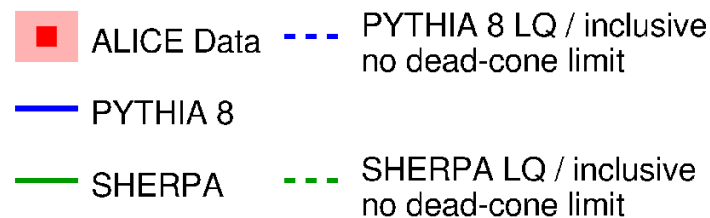
Smaller energy loss for heavy quarks,  
especially when the energy of a parton is close to its mass

$$\langle \Delta E \rangle_g > \langle \Delta E \rangle_{u,d,s} > \langle \Delta E \rangle_c > \langle \Delta E \rangle_b$$

# Direct Observation of “Dead-Cone” Effect



First direct observation of “dead-cone” effect in jets containing a soft  $D^0$ -meson

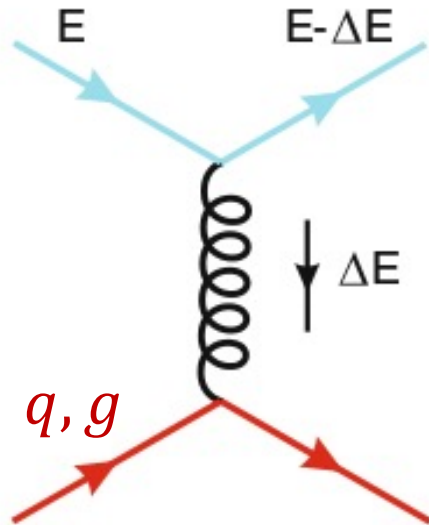


*ALICE, Nature 605, 440 (2022)*


# Elastic Collision of Heavy Quark in QGP

Collisional energy “loss”

Elastic scattering



One-dimensional elastic collision in classic mechanics



The diagram shows two spheres on a horizontal surface. The left sphere is solid black with mass  $M$  and is moving to the right with velocity  $v$ . The right sphere is hollow white with mass  $m$  and is initially at rest. Below the diagram, the energy loss formula is given as  $\frac{\Delta E}{E} = \frac{4k}{(k+1)^2}$ , where  $k = \frac{M}{m} \gg 1$ .

$$\frac{\Delta E}{E} = \frac{4k}{(k+1)^2}, \quad k = \frac{M}{m} \gg 1$$

High-energy object: Lose small energy per collision

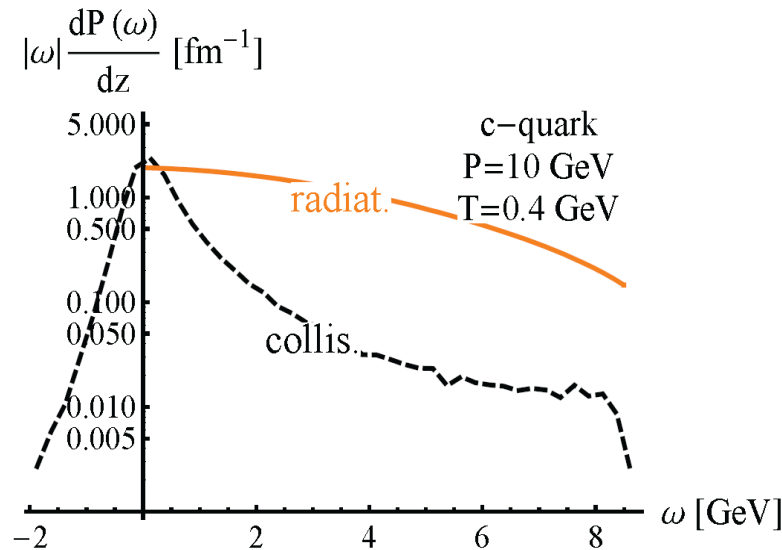
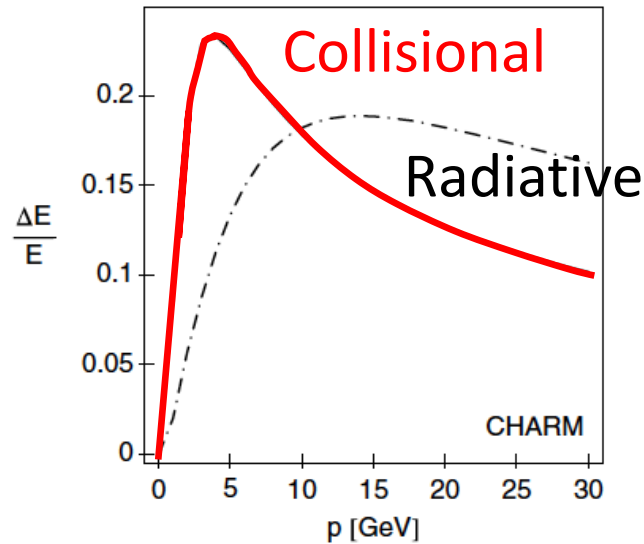
Low-energy object: May gain small energy from the medium  
“Brownian motion”

Large relaxation time for heavy flavor quark, larger or comparable to QGP lifetime

→ Interactions at all stages are important (so-called memory?)



# Elastic vs. Inelastic Collisions in QGP



## Elastic collision:

- Dominant at low- $p_T$
- Not always energy “loss”
- Responsible for heavy quark collectivity

## Inelastic collision:

- More important at high- $p_T$
- Main contributor of energy loss (jet quenching)

*M. Djordjevic, PRC74, 064907 (2006)*

Both have much less effect on bottom than charm



# Heavy Quarks Probe QGP

HF hadrons production affected by interaction between heavy quark and medium

They are sensitive to:

- Properties of the hot and dense medium
- Mechanism of heavy quark and medium interaction

Low- $p_T$ :

- Thermalization of heavy quarks
- Diffusion coefficient

High- $p_T$ :

- Energy loss of heavy quarks
- Transport coefficient

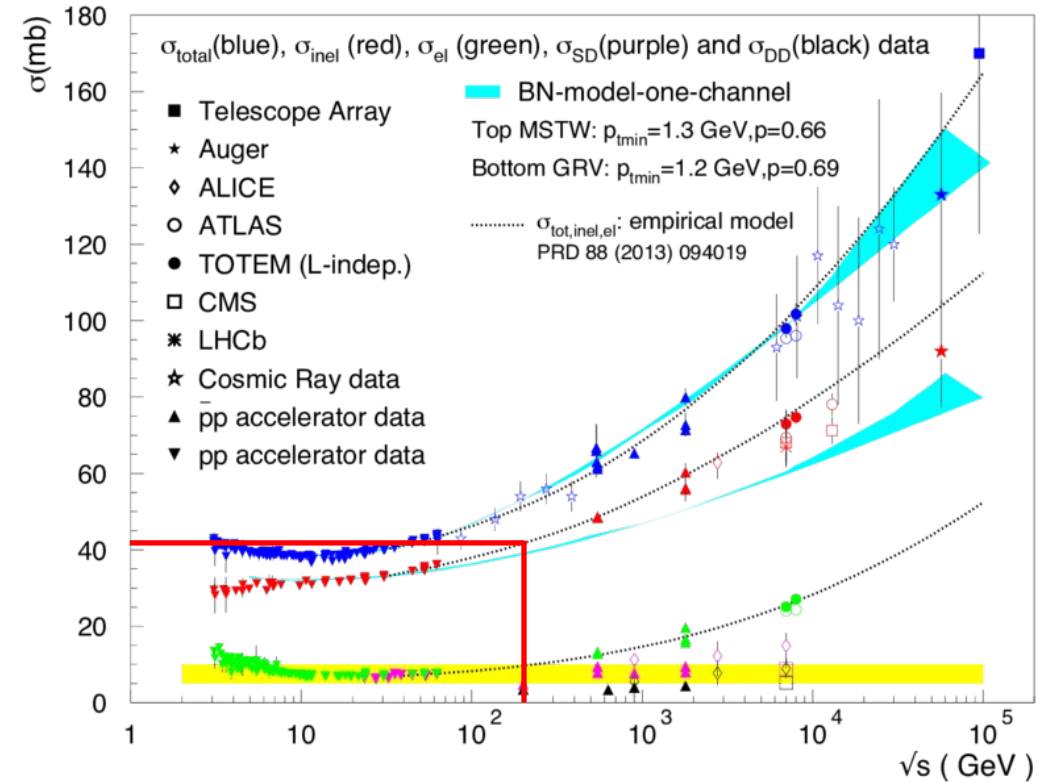
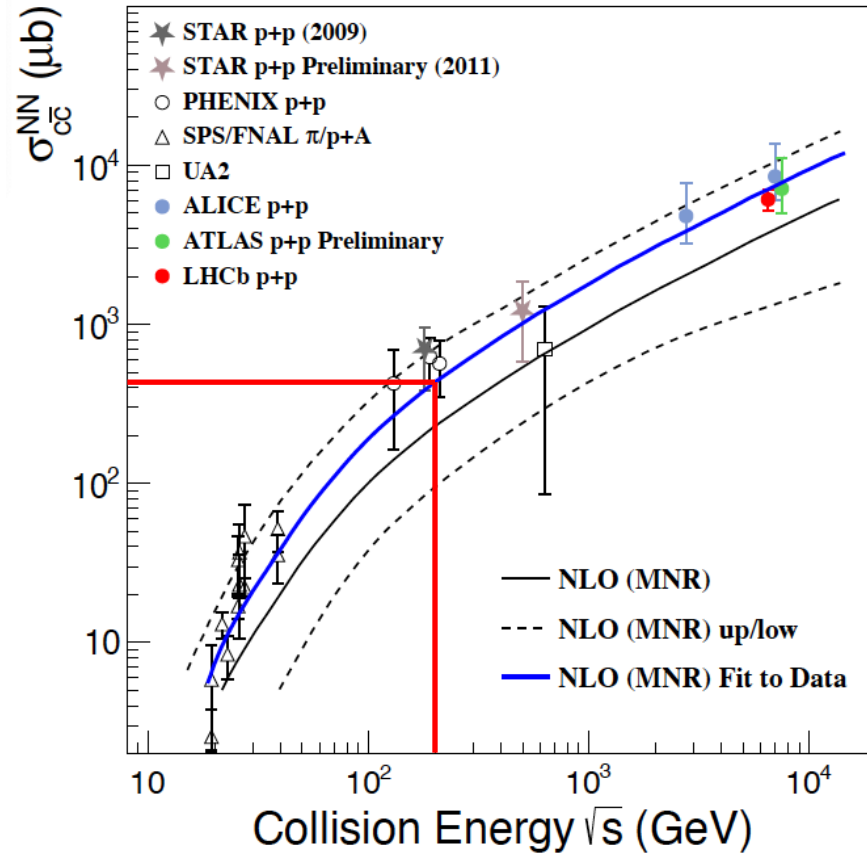
Hadronization mechanism is essential for experiment and theory comparisons

# How to Measure Heavy Flavor Hadrons





# Charm Quark Pair Production Cross Section

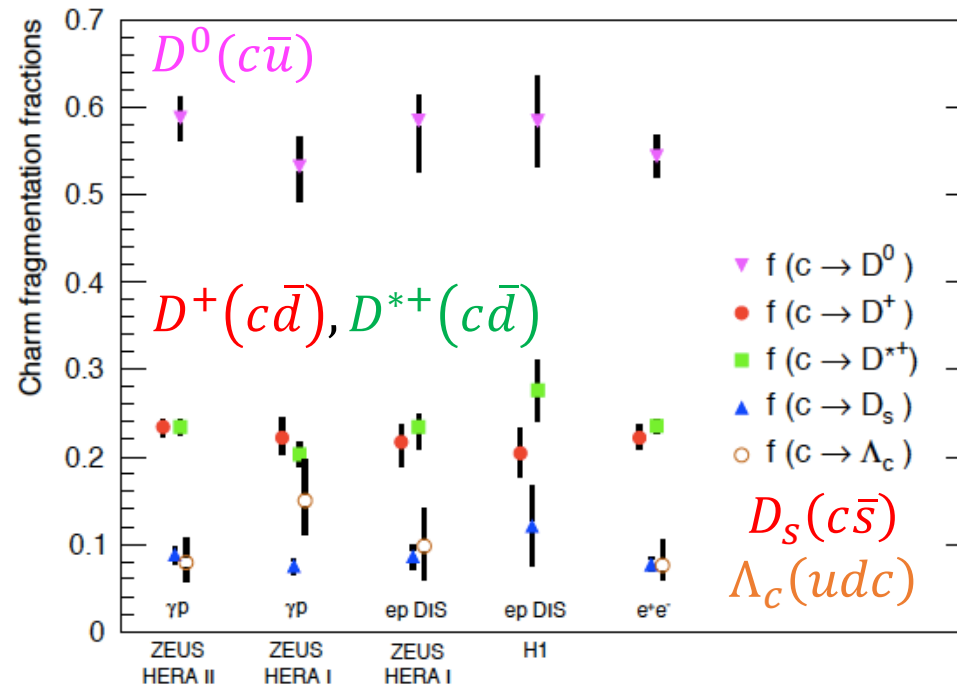


400  $\mu\text{b}/40 \text{ mb} = 1 \text{ } c\bar{c}$  per 100 Min. Bias p+p events @RHIC  
~1/5 produced in one unit of rapidity at mid-rapidity

~0.6  $c\bar{c}$  at mid-rapidity **per MB Au+Au collisions** at RHIC, 10x more at LHC



# Charmed Hadrons



**QUESTION:**  
Why  $D^0$  is significantly higher than  $D^+$ ?

ZEUS, JHEP 1309, 058 (2013)

Hadron	Abundance (fragmentation)
$D^0$	56%
$D^+$	24%
$D_s$	10%
$\Lambda_c$	10%

Short lifetime → Decay before tracker

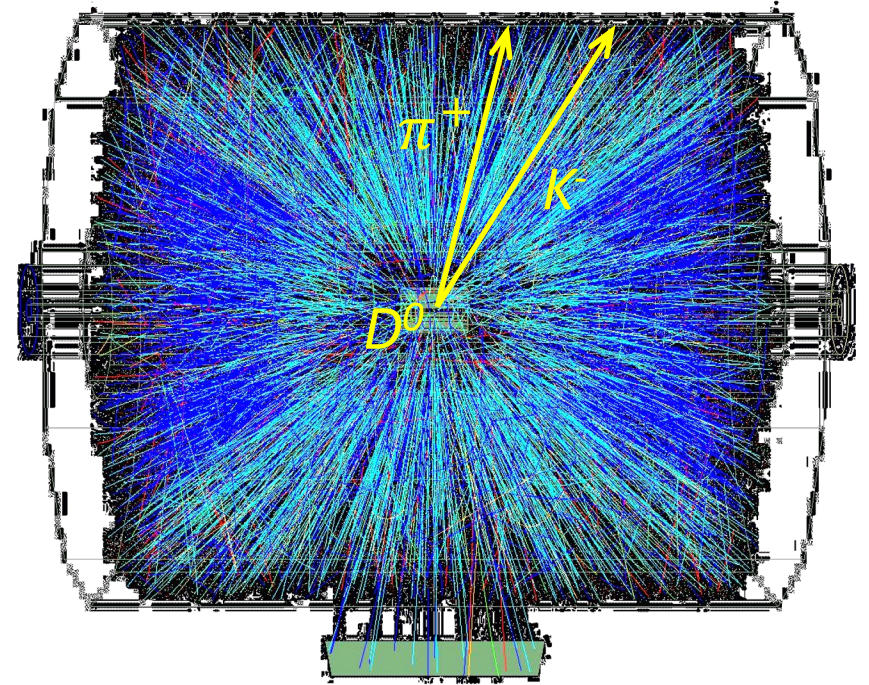
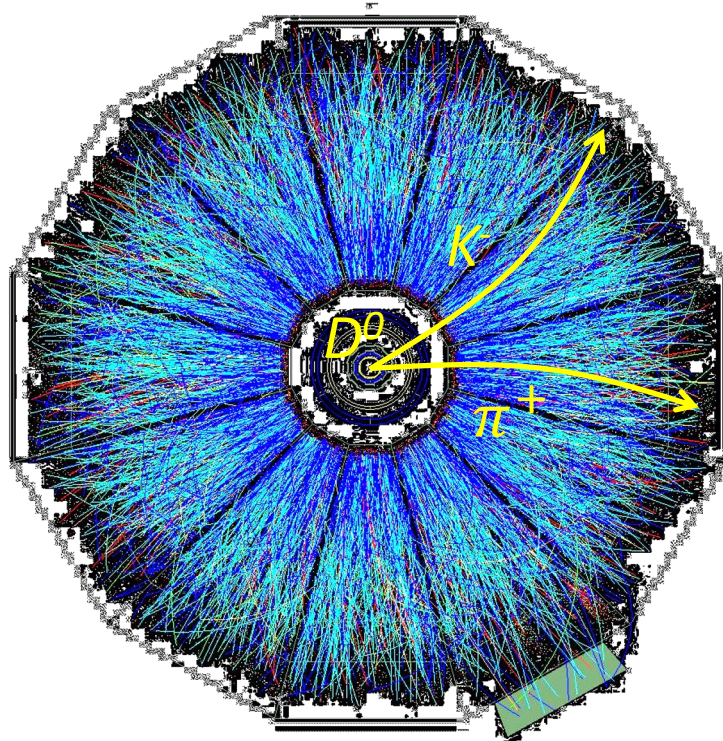
- Can not be detected directly
- Has to be reconstructed from decay products

# $D^0$ -meson Reconstruction

Golden channel:

$$D^0 \rightarrow K^- \pi^+$$

Br.  $\sim 3.9\%$



- Need to pair all  $K^-$  with all  $\pi^+$  in the same event

$$dN/dy(\pi^+) \sim 100 \quad dN/dy(K^-) \sim 20$$

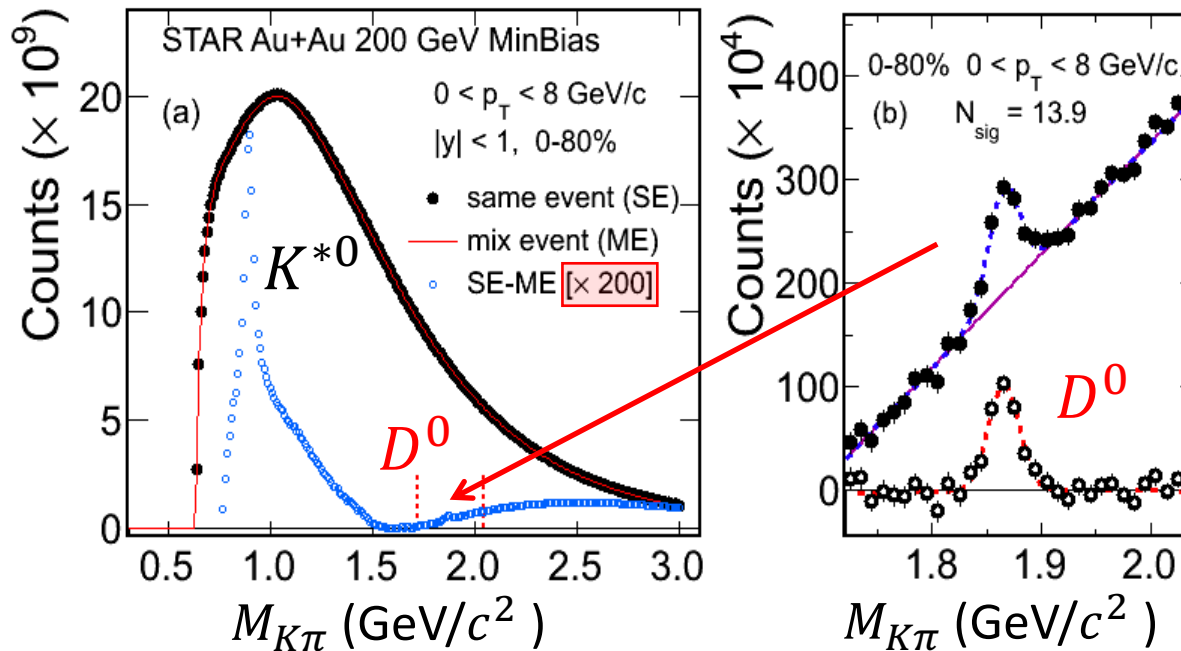
STAR, PRC79, 034909 (2009)

Signal-to-background ratio: 0.01 vs. 2000

Better if concentrate on narrow mass window

- Huge random combinatorial background may be *estimated* with the same data

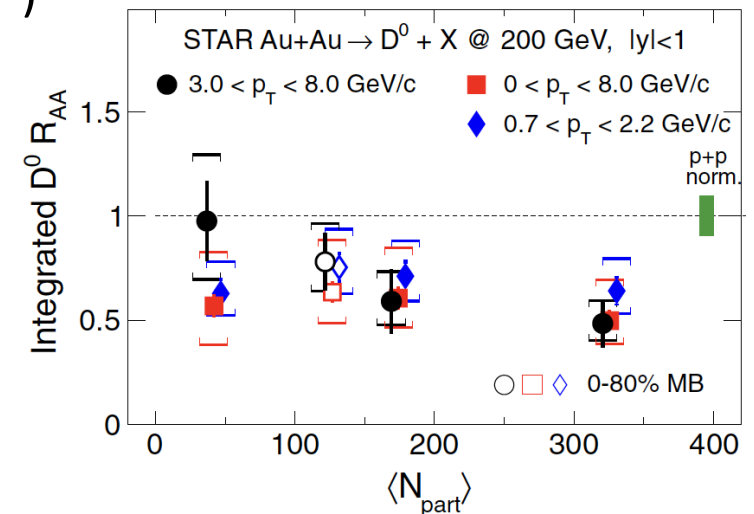
# D<sup>0</sup>-meson Reconstruction



STAR, PRL113, 142301 (2014)

It is doable, but very challenging

- Rely on good particle (mainly kaon) identification
- The statistical and systematic uncertainties are large





# Topological Reconstruction

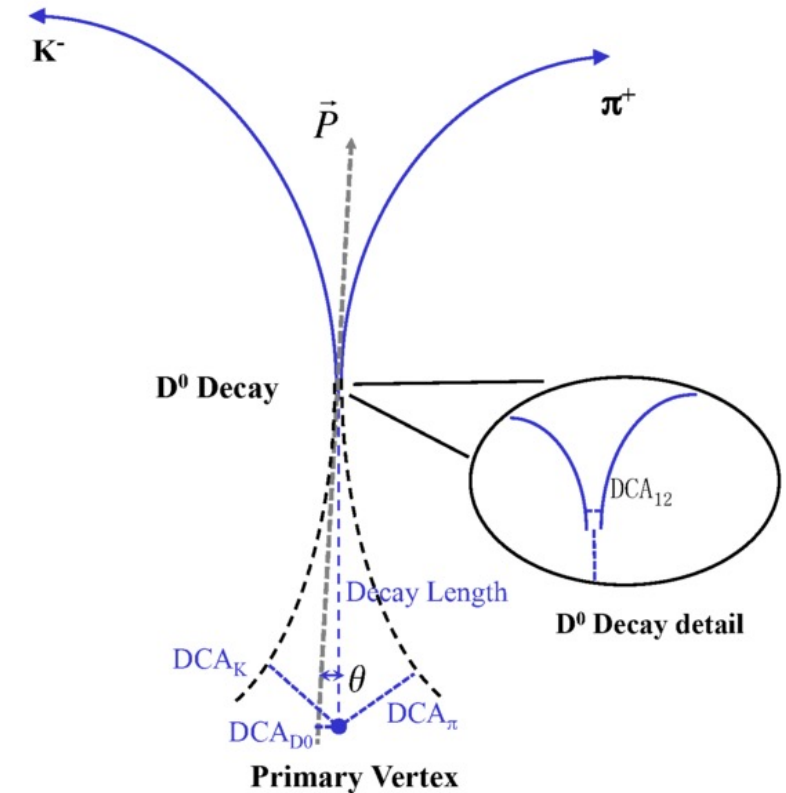
- Kaons and pions from D decay are originated from secondary vertex, *a little bit* away from primary vertex

Hadron	Abundance (fragmentation)	$c\tau$ ( $\mu\text{m}$ )
$D^0$	56%	123
$D^+$	24%	312
$D_s$	10%	150
$\Lambda_c$	10%	60

Less than 1/2  
thickness of  
a fingernail!!!

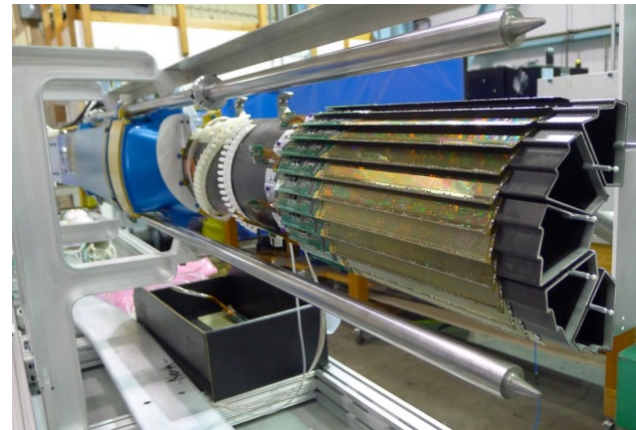
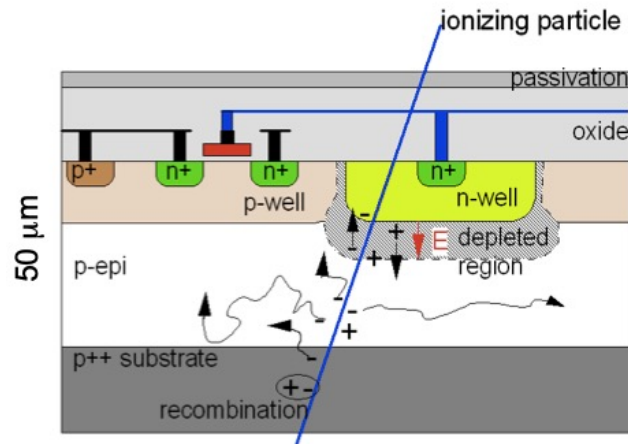
- Majority of kaons and pions are promptly produced and originate from primary vertex
- If the detectors can tell the tiny difference, significant backgrounds can be rejected
- Pointing resolution of tracker is crucial for this purpose

**QUESTION:**  
*Does this method work  
for low- $p_T$  D-mesons?*

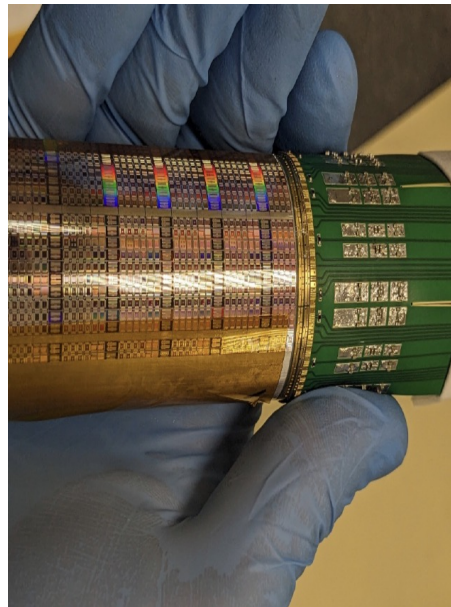


# Silicon Pixel Detectors

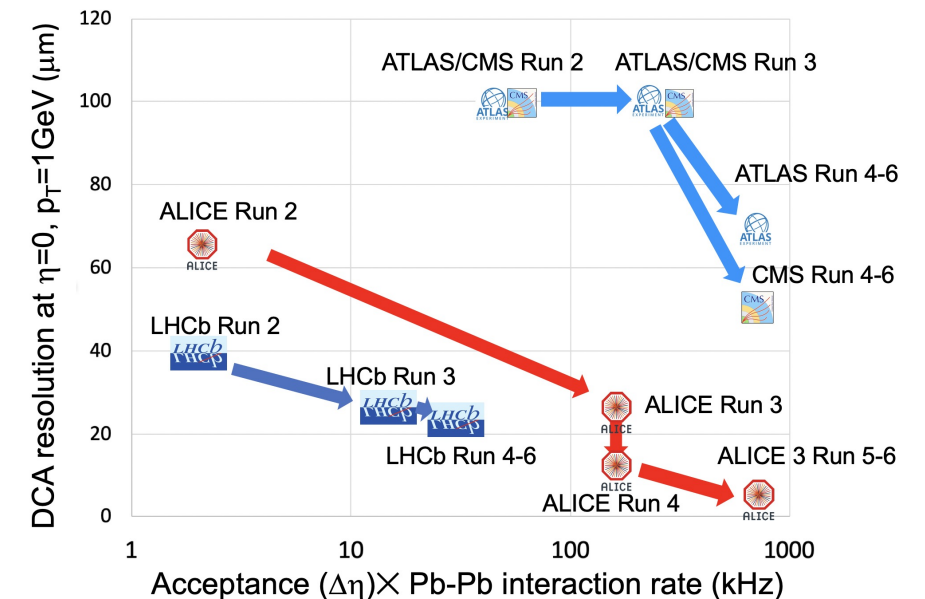
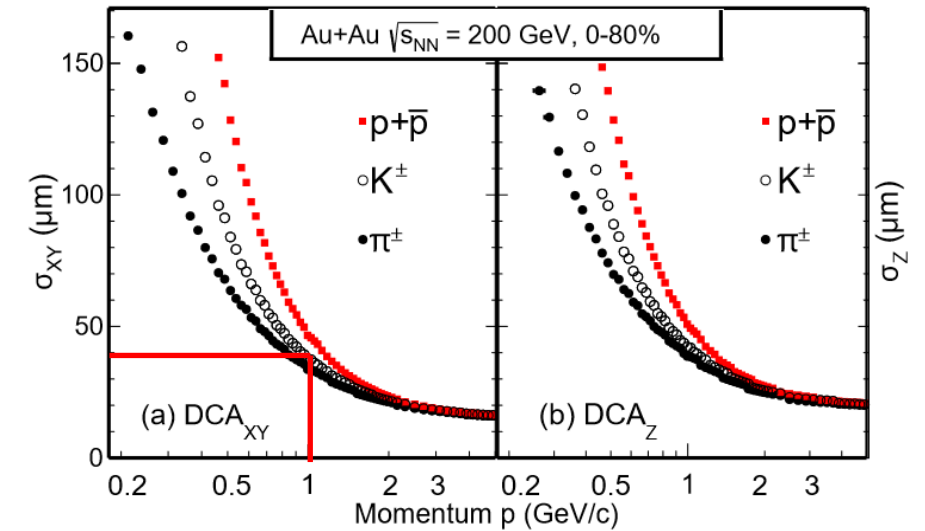
## Monolithic Active Pixel Sensor (MAPS)



STAR Heavy Flavor Tracker



ALICE Inner Tracking System



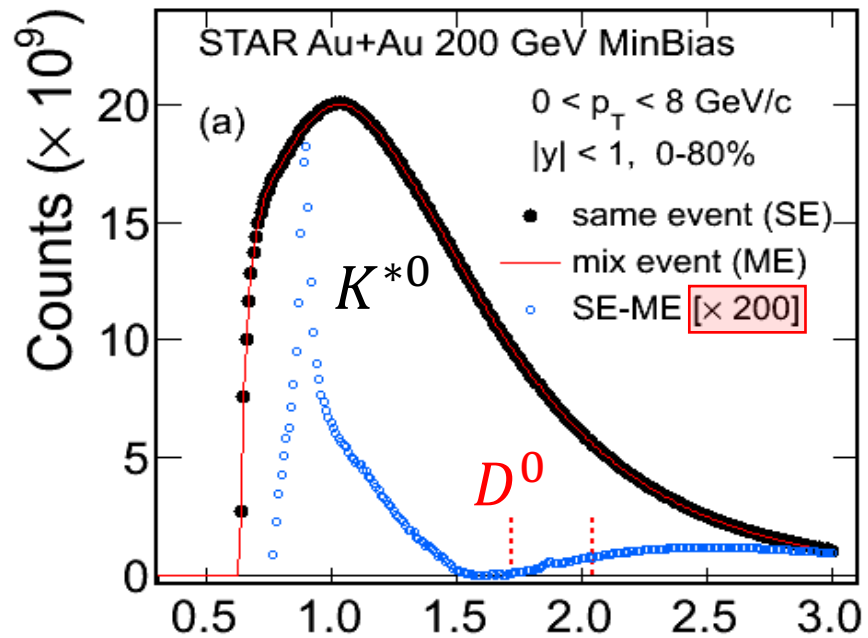


# $D^0$ Signal with Topological Cuts

2010+2011

820M events

w/o HFT



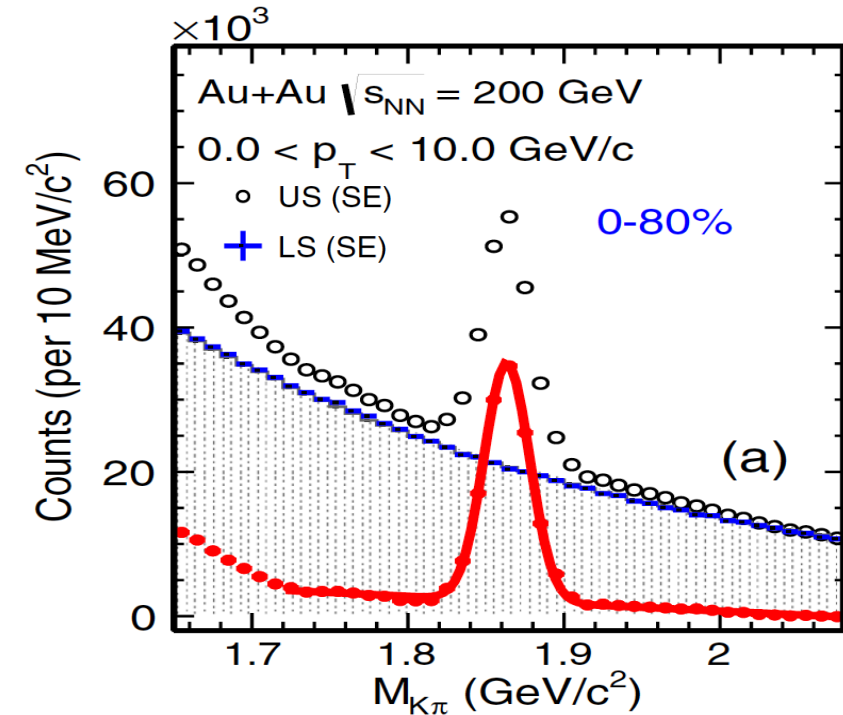
STAR, PRL113, 142301 (2014)



2014

900M events

w/ HFT



STAR, PRC 99, 034908 (2019)

S/B improved by  $O(10^4)$





# Feeddown Contribution

The measured charm hadrons have feeddown contributions

$$\text{Inclusive } D = \text{prompt } D + \text{non-prompt } D$$

**Prompt D:** D from vertex not distinguishable from primary vertex

- Direct D
- Strong/radiative decay products of heavier particle  
i.e.  $D^{*+} \rightarrow D^0 + \pi^+$ ,  $D^{*+} \rightarrow D^+ + \pi^0$ ,  $D^{*0} \rightarrow D^0 + \pi^0/\gamma$

$D^0, D^\pm, D^{*0}, D^{*\pm}$  reflect similar physics, unnecessary to separate prompt

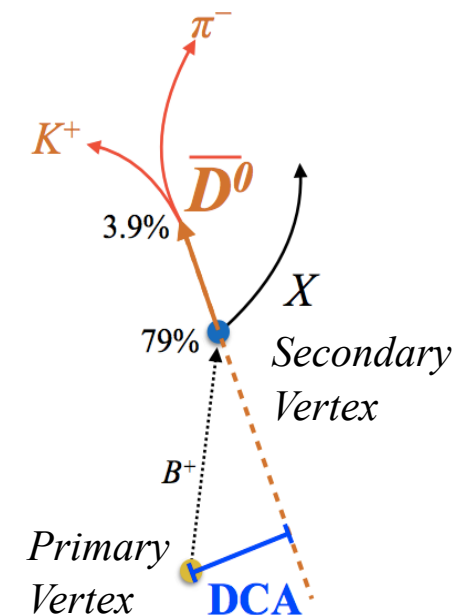
**Non-prompt D:** D from secondary vertex

- Weak decay products of heavier particle, mainly from B-hadrons

Flavor changed far later than QGP lifetime ( $\tau \sim 500 \mu\text{m}/c$ )

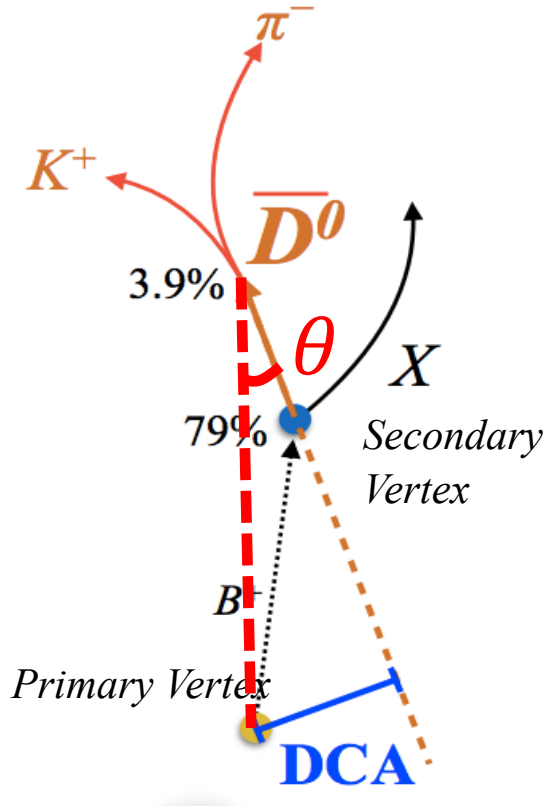
➔ Reflect interaction of bottom instead of charm quark with the medium

**A typical (good) proxy of B-hadrons**



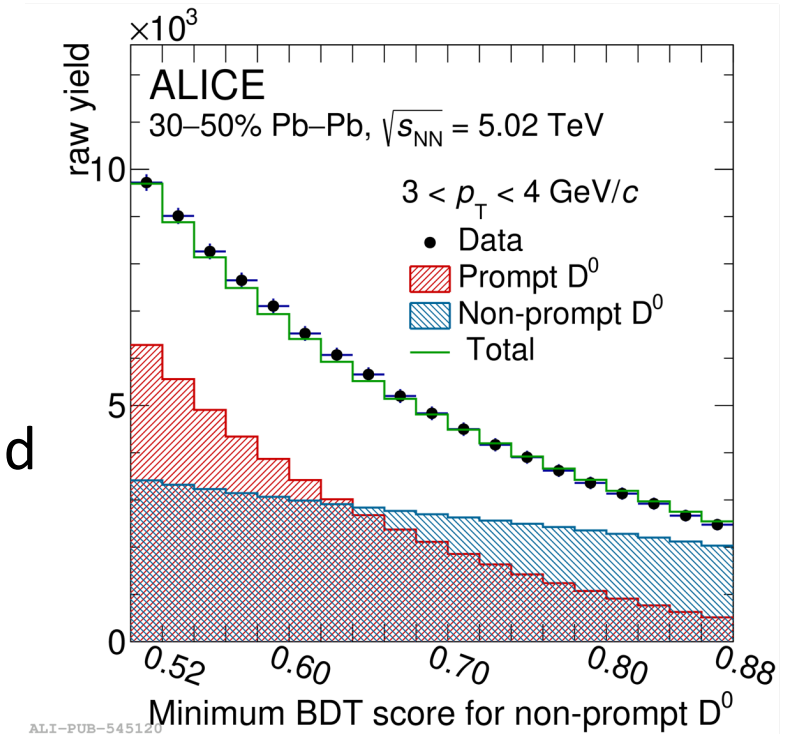
# Prompt and Non-prompt Separation

Machine-learning approach with 3-class classification



Training variables mainly based on:

- DCA of  $D^0$  daughters and  $D^0$
- Distance between  $D^0$  decay vertex and primary vertex
- Pointing angle between line of flight and momentum reconstructed



Non-prompt D yield obtained by combining measurements of *non-prompt fraction + inclusive measurements*

ALICE, EPJ C83, 1123 (2023)

# Experimental Results

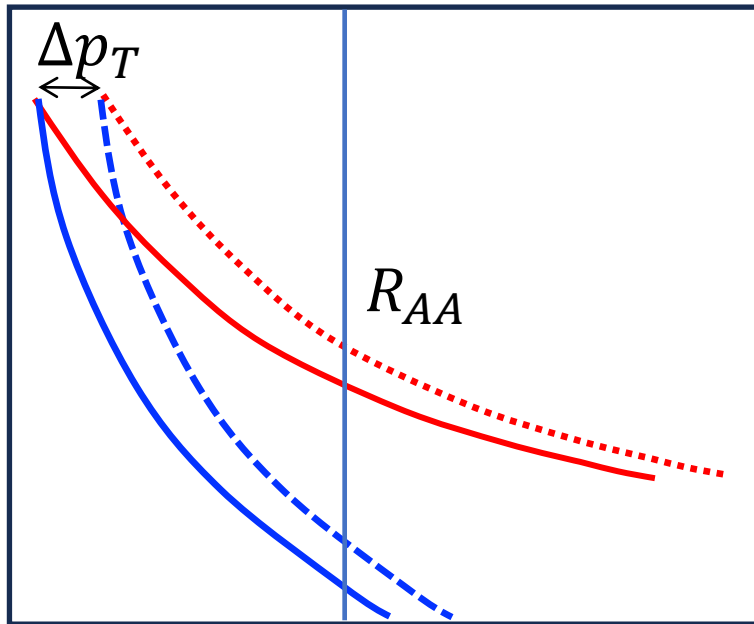


# Observable: Nuclear Modification Factor

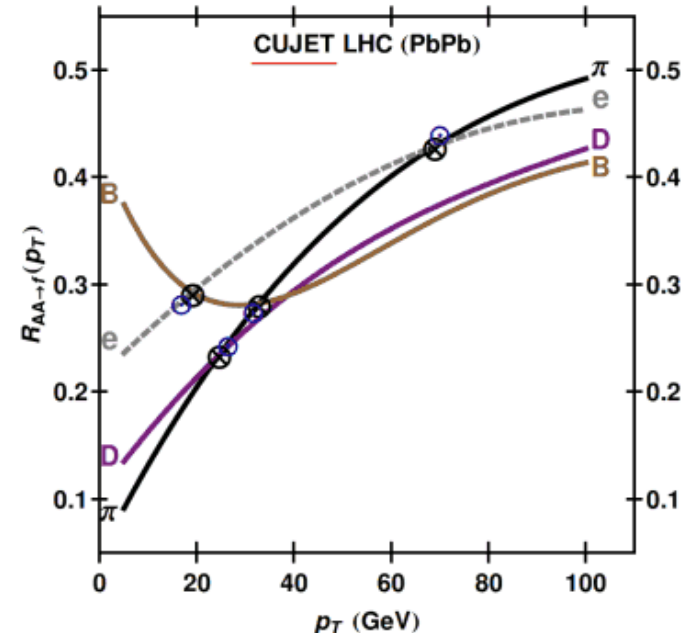
Nuclear Modification Factor: 
$$R_{AA} = \frac{dN_{AA}/dp_T dy}{\langle N_{bin} \rangle dN_{pp}/dp_T dy}$$

Conversion from  $R_{AA}$  measurements to energy loss is not so straight forward

- $R = \frac{f(p_T - \Delta p_T)}{f(p_T)}$  depends not only on  $\Delta p_T$ , but also the shape of  $f(p_T)$



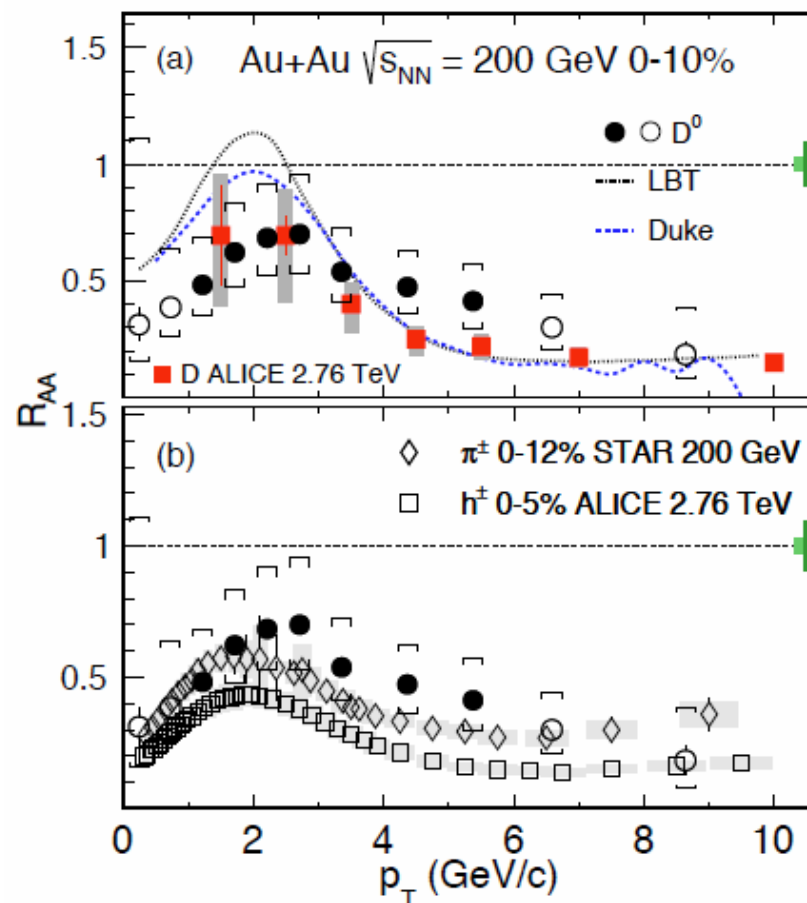
- And diluted by hadronization process ( $q \rightarrow h$ )



$R_{AA}(D) > R_AA(B)$   
@low/intermediate  $p_T$

*A. Buzzatti et al.,  
NPA904, 779c (2013)*

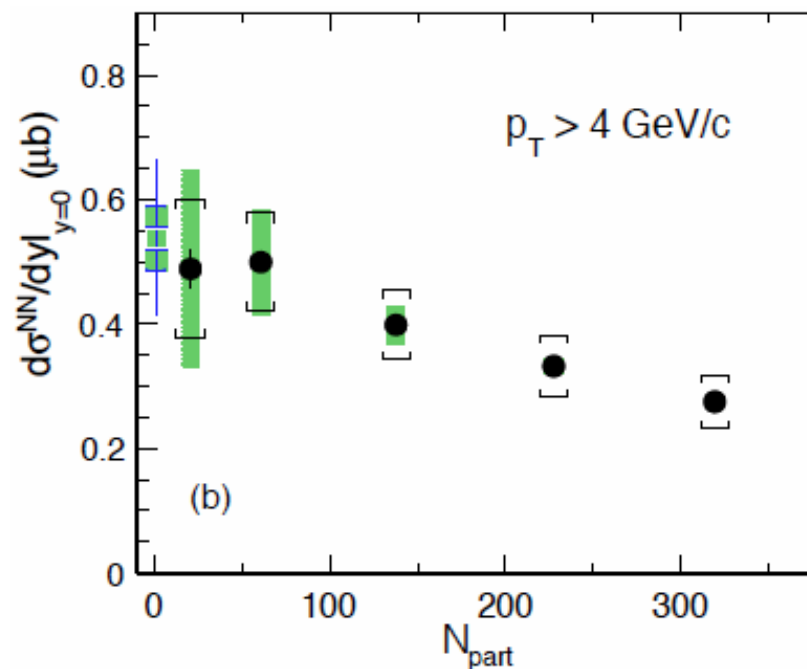
# D<sup>0</sup> Suppression at RHIC



STAR: PRC 99, 034908 (2019)

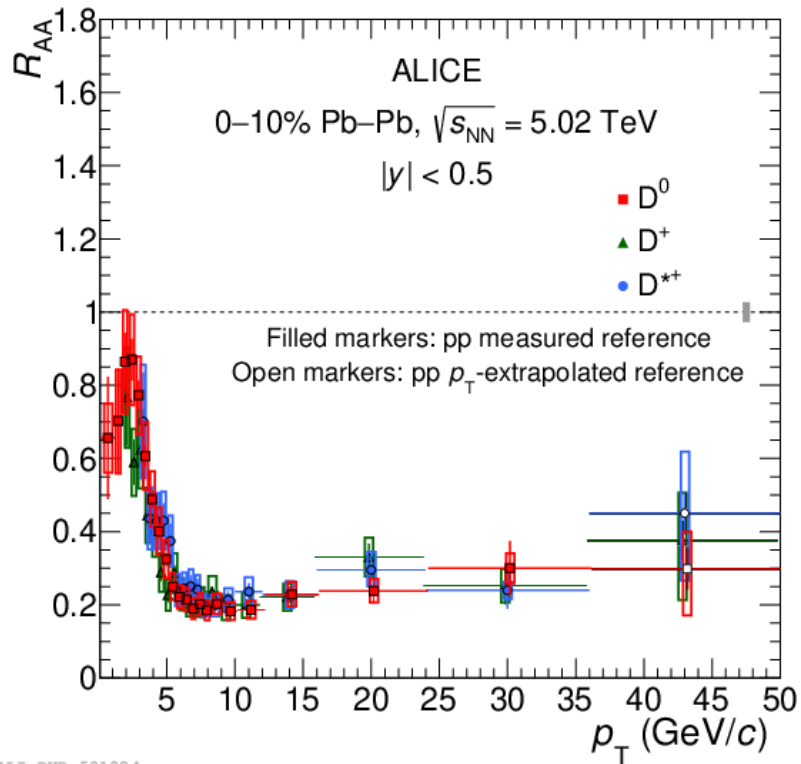
LBT: PRC 97 (2018) 014907

Duke: PRC 94 (2016) 014909

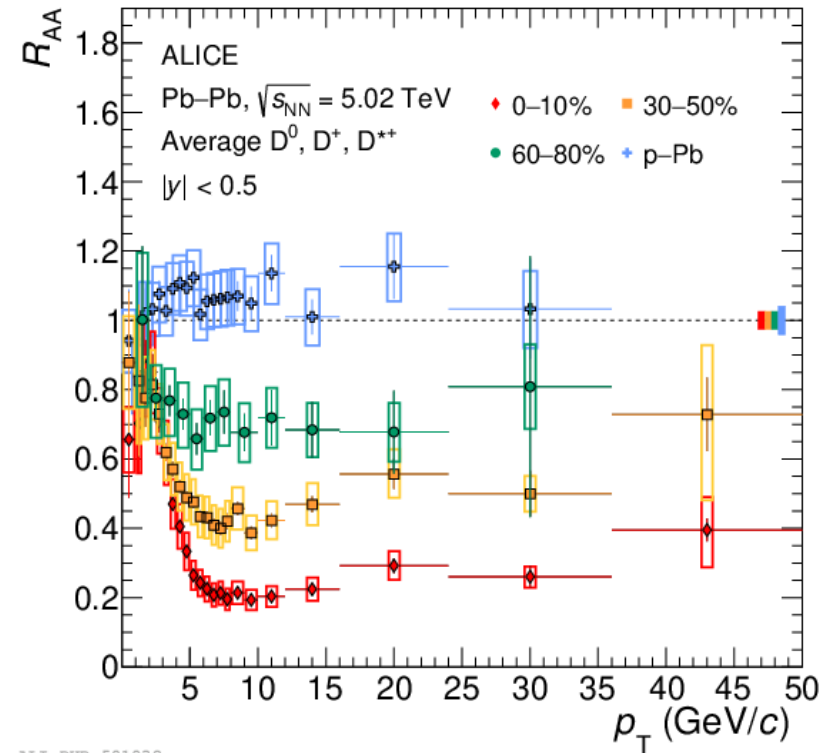


- Significant suppression of D<sup>0</sup> yield at high- $p_T$  observed in 200 GeV Au+Au collisions
- Stronger suppression towards central collisions
- Described by theoretical calculations

# D-mesons Suppression at LHC



ALI-PUB-501924

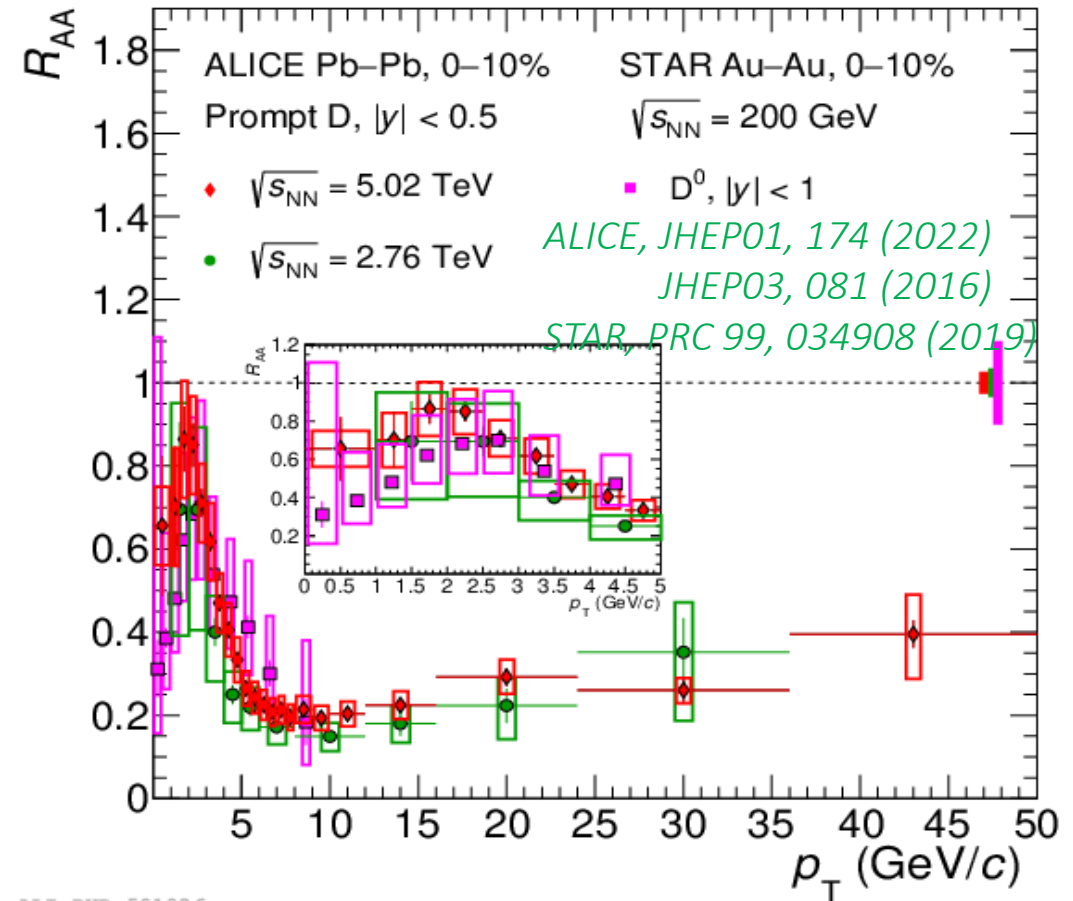
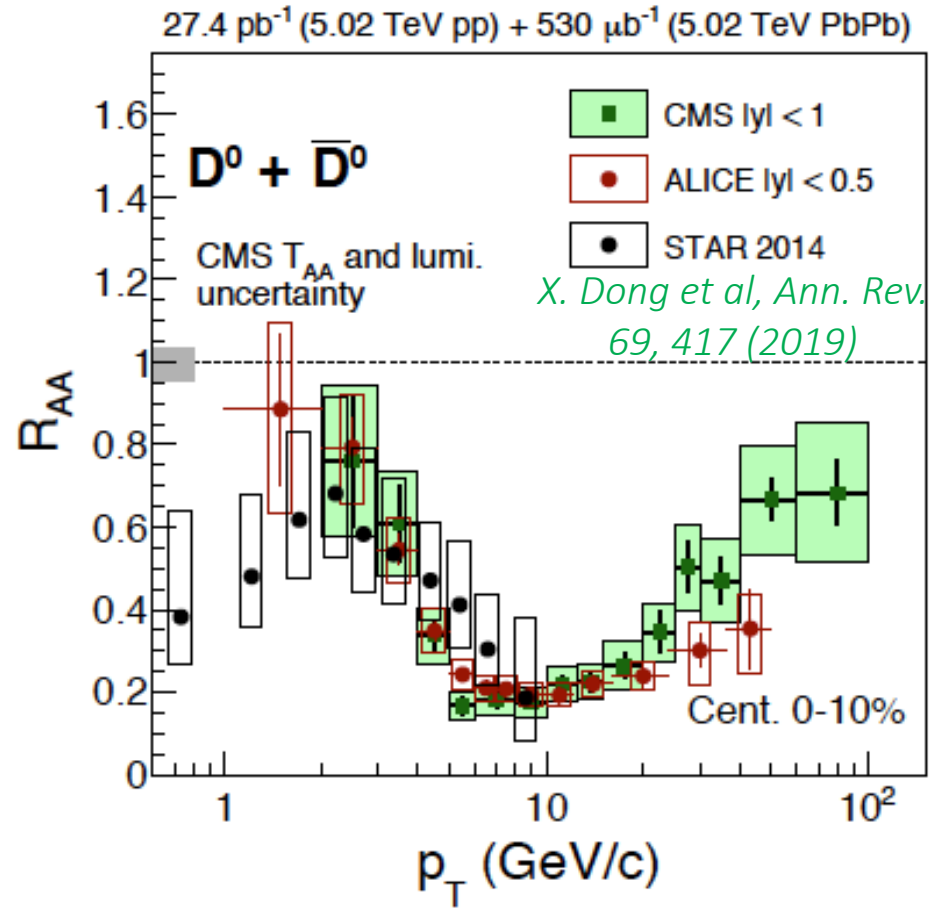


ALI-PUB-501932

ALICE, JHEP01, 174 (2022)

- Prompt  $D^0, D^+, D^{*+}$  same suppression
- Dramatic decrease from 3-6 GeV/c
- Slight increase at high  $p_T$
- Consistent with unity in pPb
- Clear increasing suppression towards central collisions at intermediate/high  $p_T$   
 $\leftarrow$  Increasing energy density, size, lifetime

# Energy Dependence



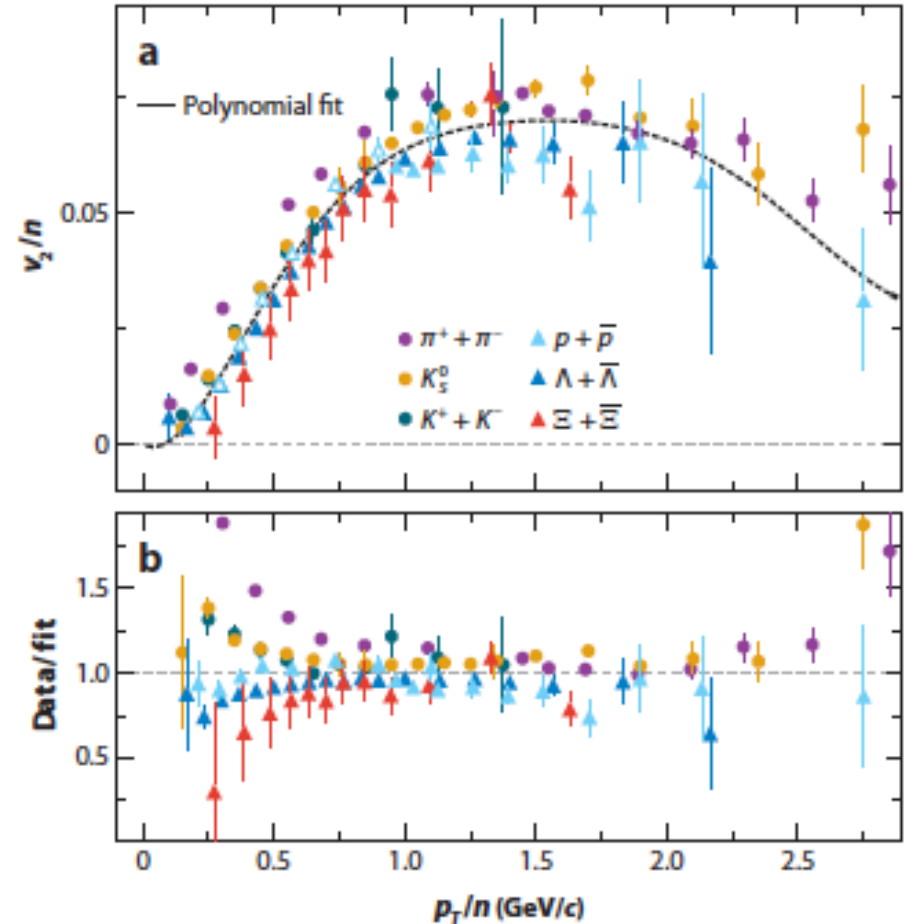
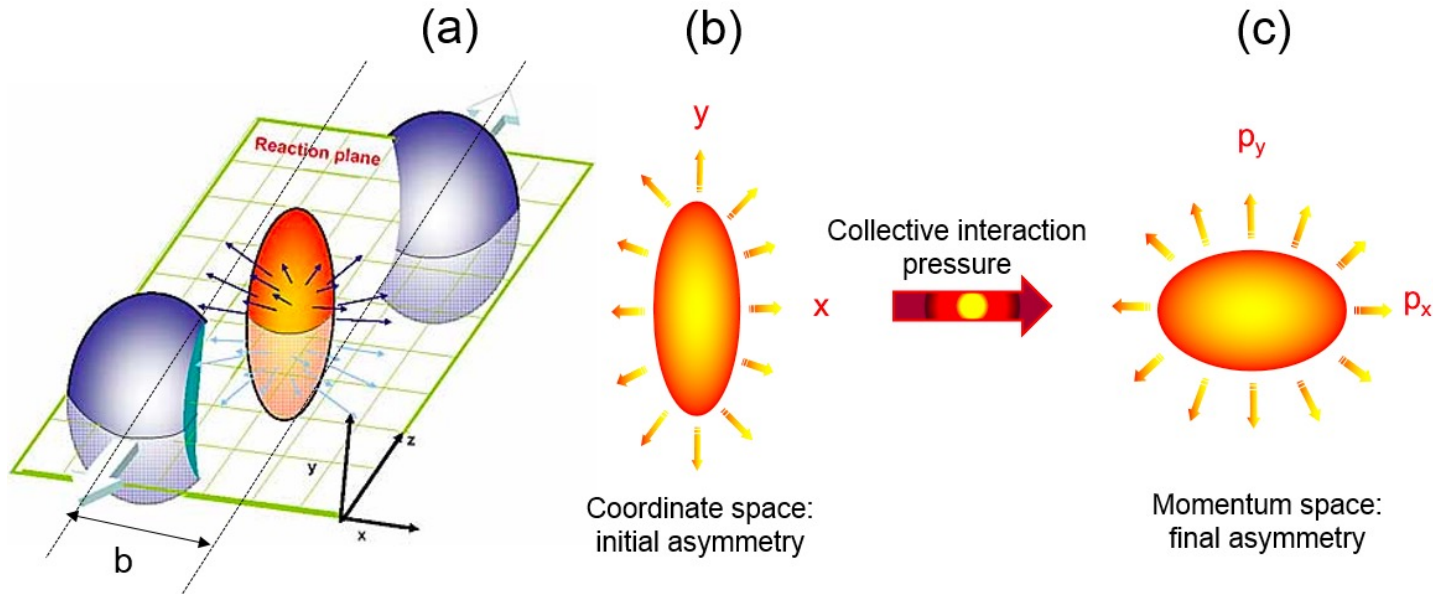
ALI-PUB-501936

200 GeV ~ 2.76 TeV = 5.02 TeV

Counterbalance of temperature and medium density vs. p<sub>T</sub> spectrum steepness



# Observable: Collectivity

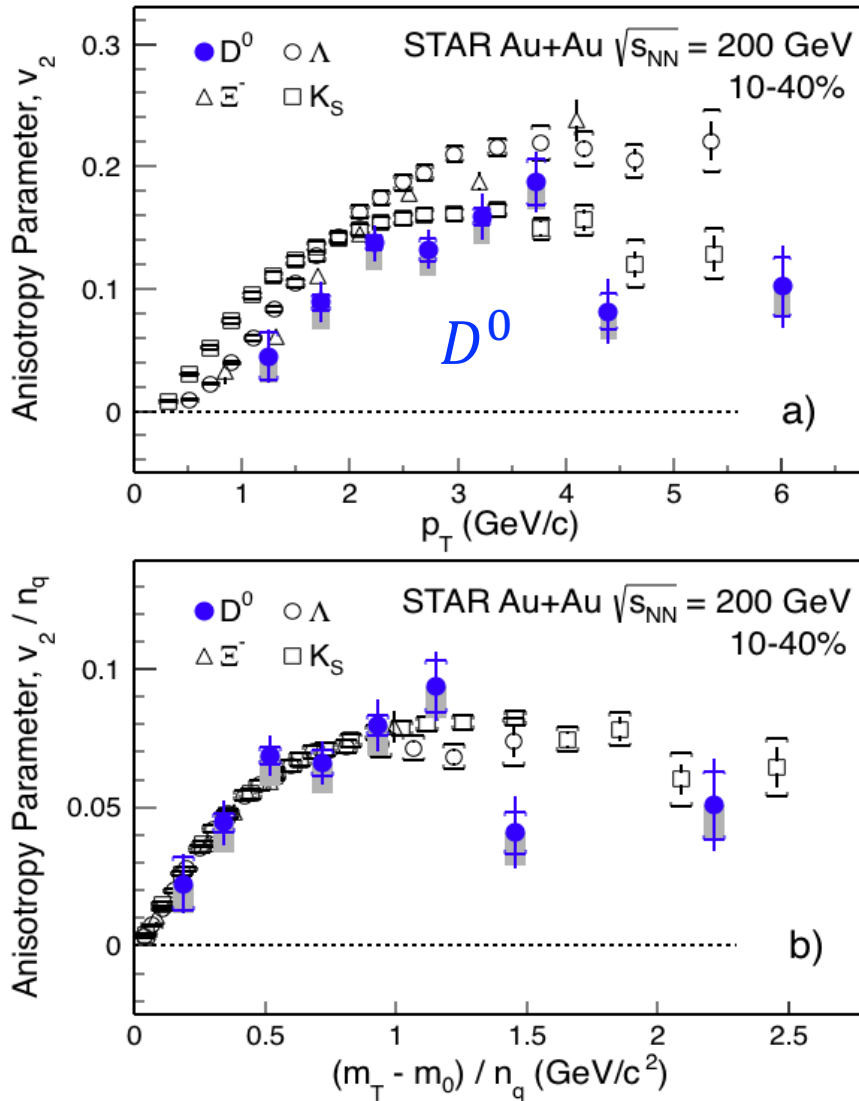


Number-of-Constituent Quark scaling  
 → Partonic flow + Coalescence



# Elliptic/Triangular Flow of D-meson

STAR, PRL  
118, 212301  
(2017)

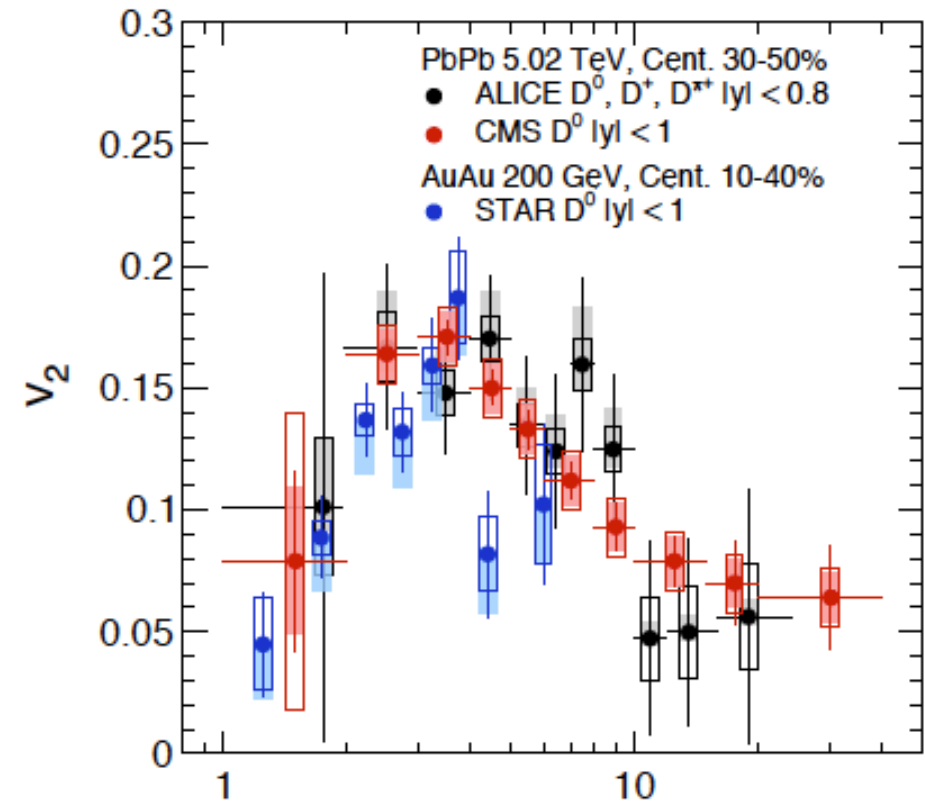
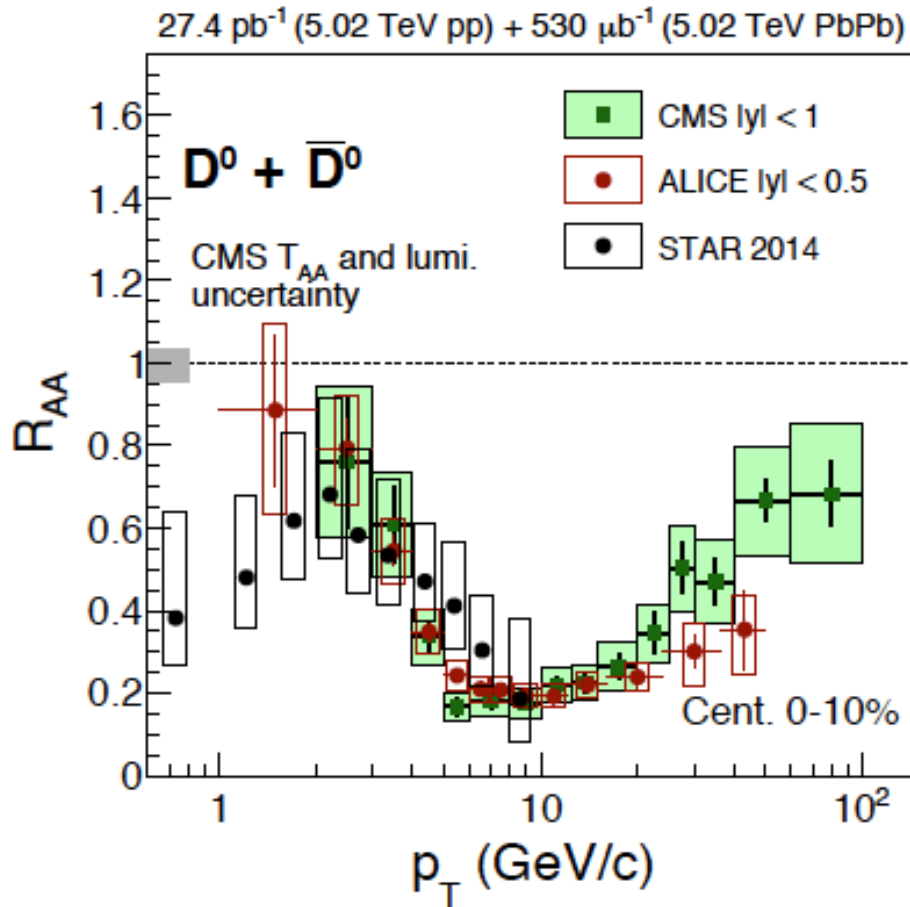


Significant  $v_2$  observed for  $D^0$ -meson

- Mass ordering at low  $p_T$   
– hydrodynamics behavior
- Follow mesons' flow at intermediated  $p_T$   
– quark coalescence



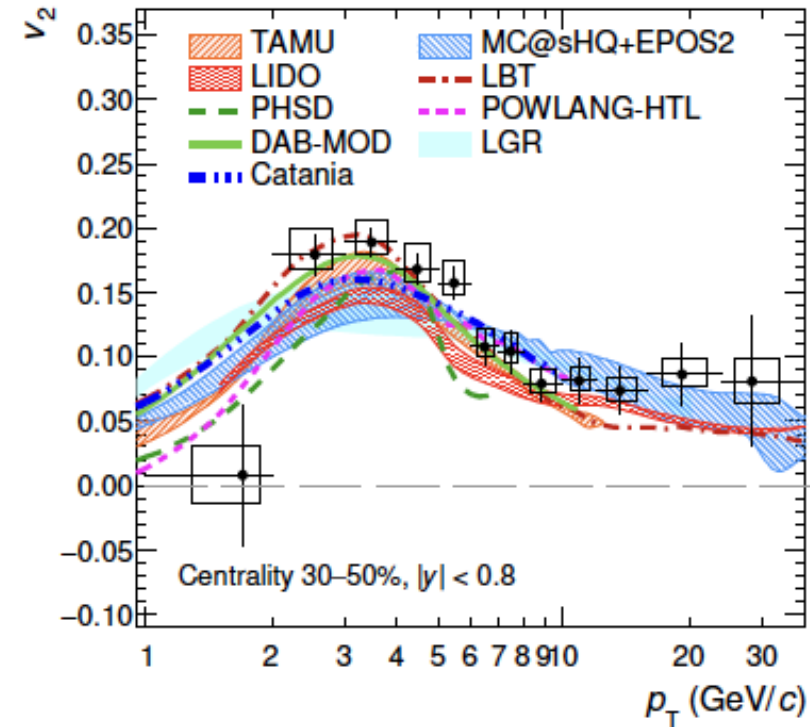
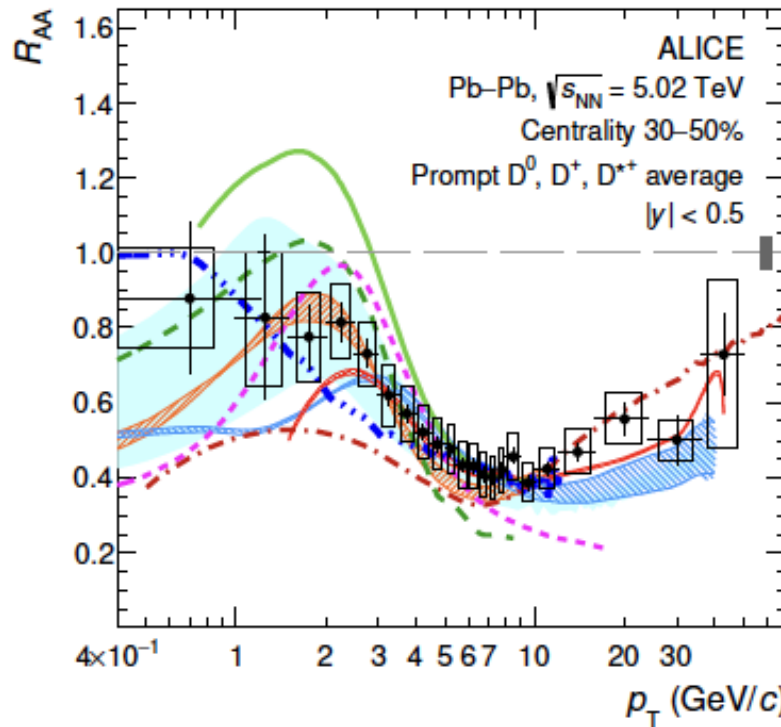
# Energy Dependence



*X. Dong et al, Ann. Rev. 69, 417 (2019)*

No obvious energy dependence for both  $R_{AA}$  and  $v_2$

# Comparison with Theories



ALICE, JHEP01, 174 (2022)

- TAMU, MC@sHQ+EPOS2, LIDO, LGR and Catania able to describe  $R_{AA}$  and  $v_2$  simultaneously  
All includes charm quark diffusion in the medium and quark coalescence hadronization
- Charm quark spatial diffusion coefficient is constrained by the comparisons



# Spatial Diffusion Coefficient

The heavy quark distribution function can be written as (Fokker-Planck equation):

*X. Dong et al, Ann. Rev. 69, 417 (2019)*

$$\frac{\partial}{\partial t} f_Q = \frac{\partial}{\partial p} A(p) p f_Q + \frac{\partial^2}{\partial^2 \vec{p}} B(p) f_Q$$

Spatial diffusion coefficient

$$D_s = \frac{T}{m_Q A(p=0)}$$

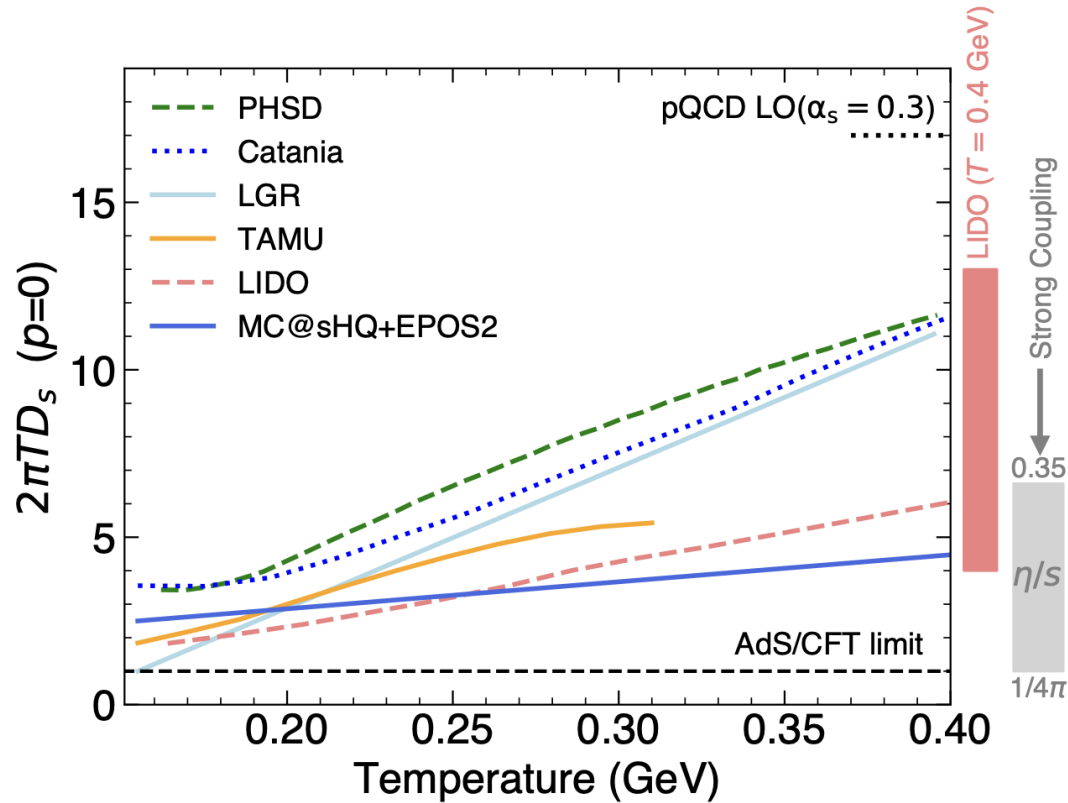
Relaxation time

$$\tau_{relax} = \frac{1}{A} = \frac{m_Q}{T} D_s = 2\pi T D_s \times \frac{m_Q}{2\pi T^2}$$

$$\frac{m_c}{2\pi T_c^2} \sim 2 \text{ fm}/c$$



# Charm Quark Diffusion Coefficient



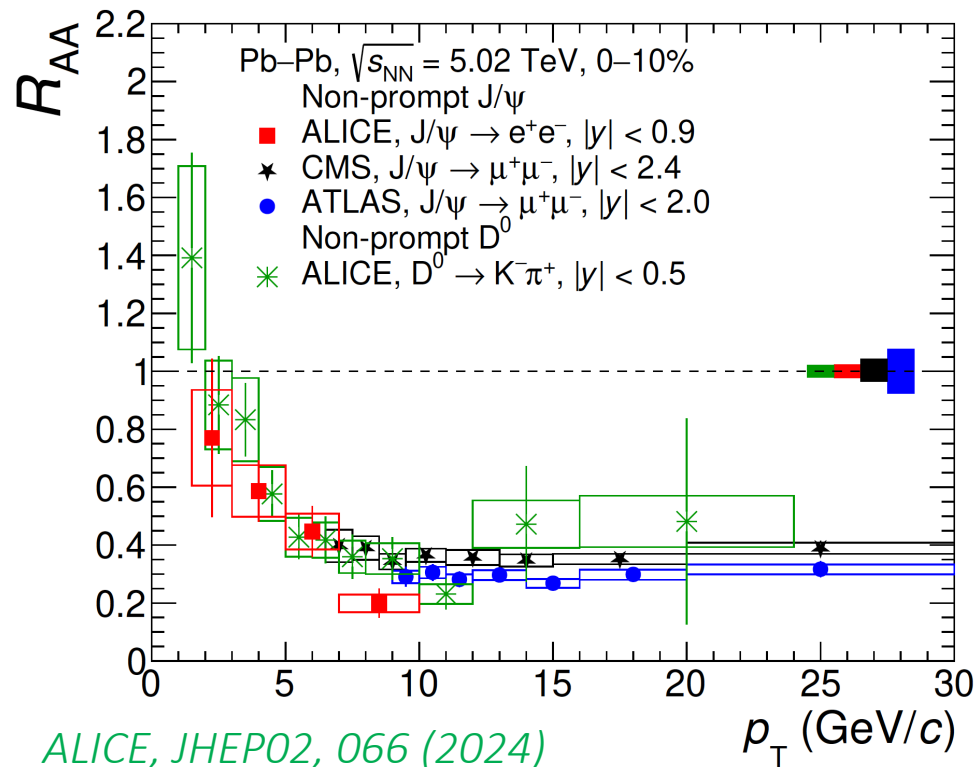
$$1.5 < 2\pi T D_s < 4.5 @ T_c$$

$$\tau_{relax} = (3 - 9) fm/c \lesssim \tau_{QGP}$$

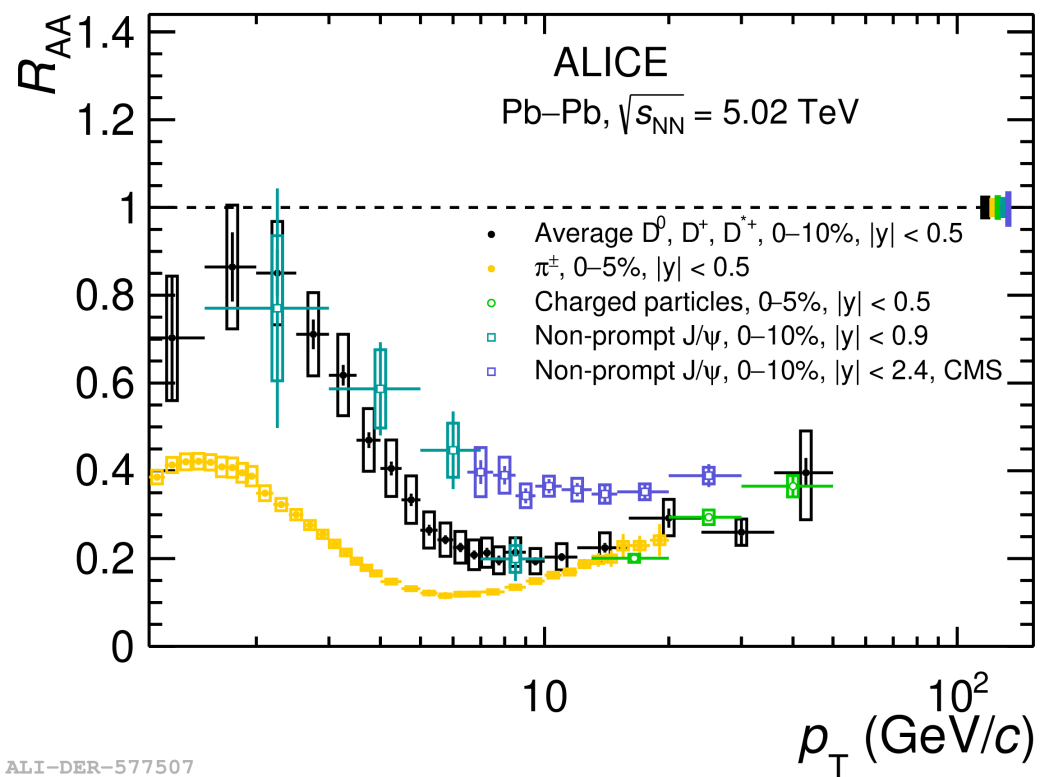
ALICE, arXiv:2211.04384

Charm is fully thermalized in QGP

# Suppression of Bottom



- Measured via non-prompt  $D^0$  and J/ $\psi$ 
  - $D^0$  has better statistics
  - J/ $\psi$  has better kinematics
- Strong suppression observed

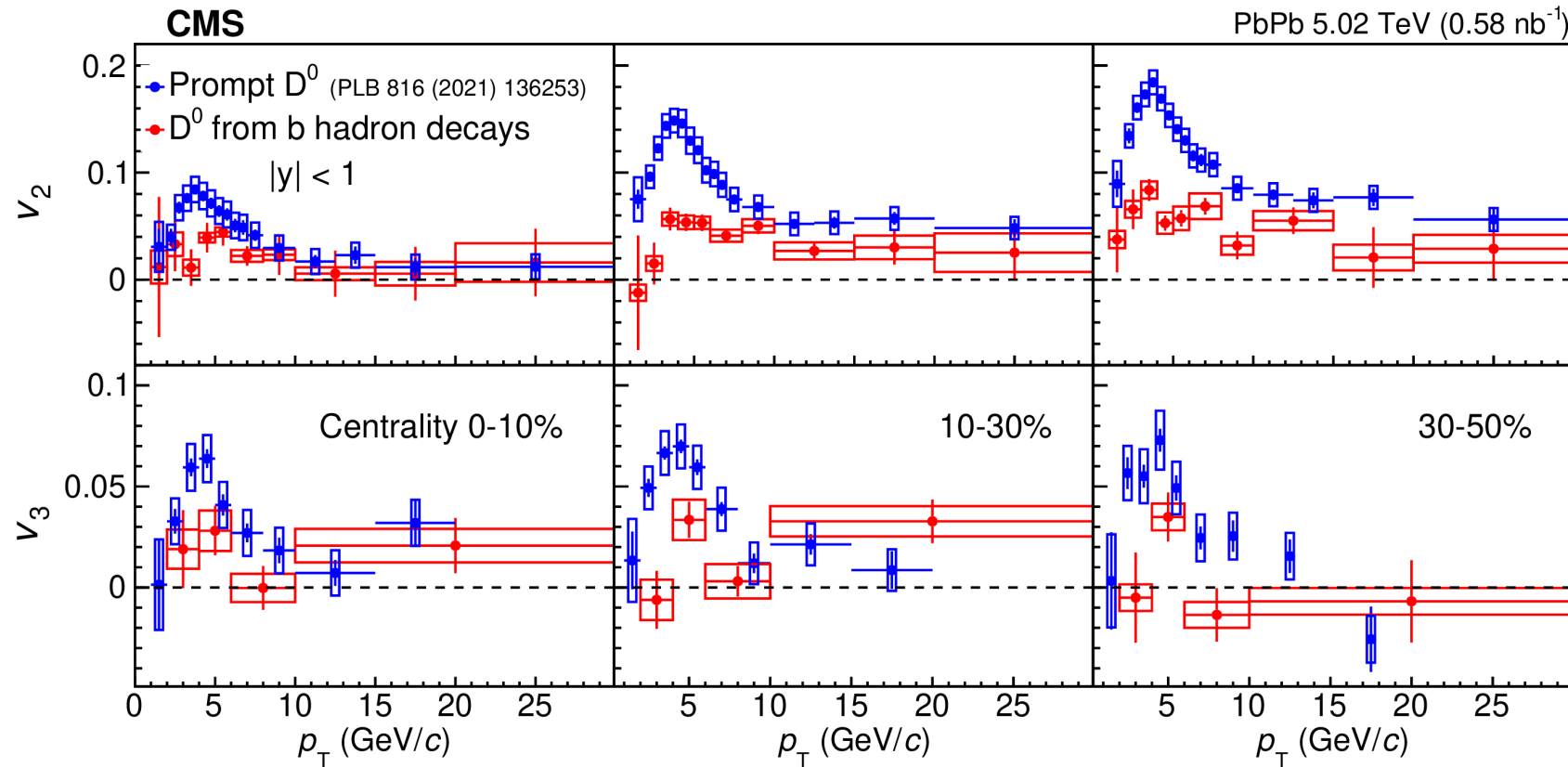


- Clear mass hierarchy at intermediate  $p_T$ 

$$R_{AA}(B) > R_{AA}(D) > R_{AA}(\text{light hadrons})$$
- Converge at high  $p_T$



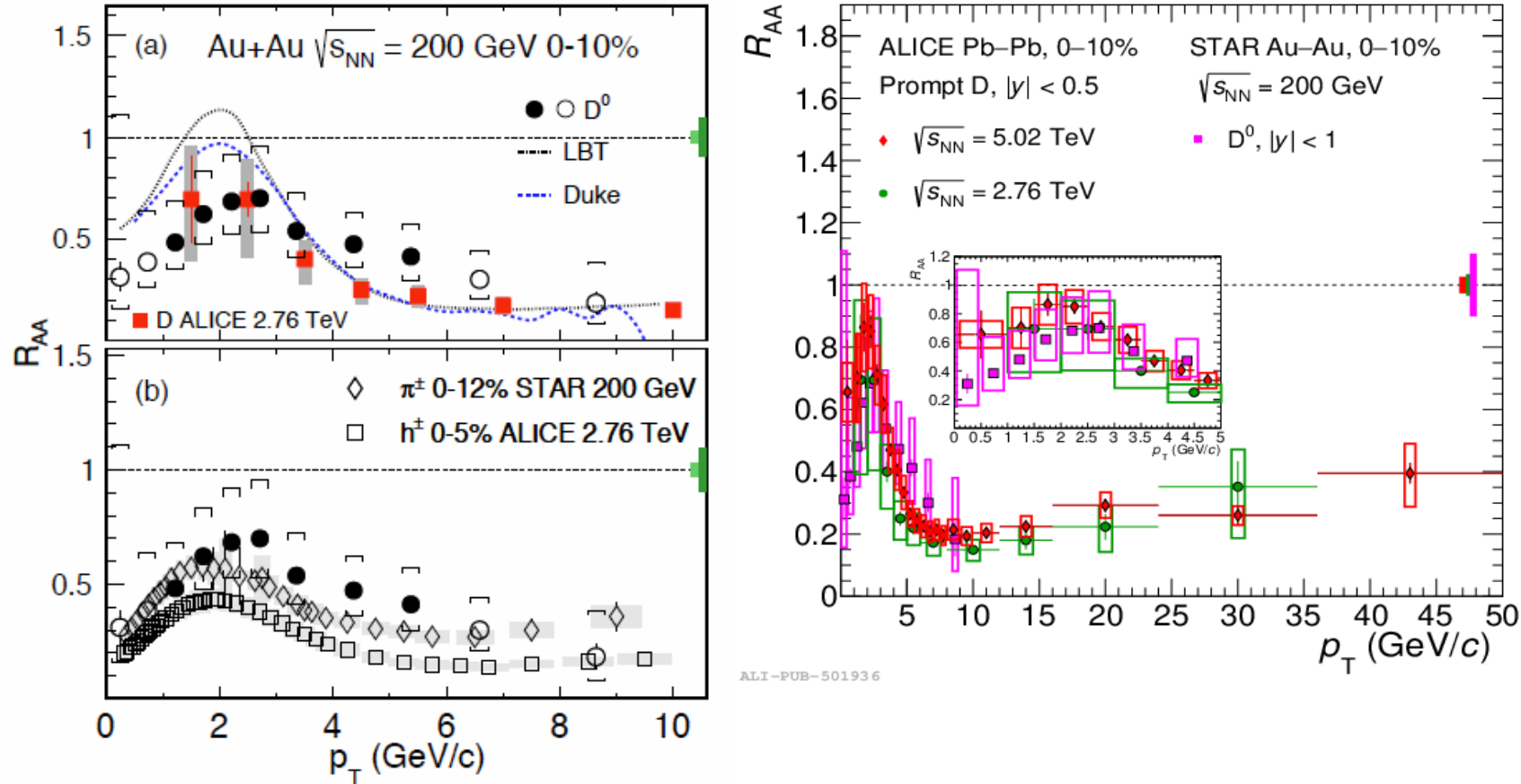
# Collectivity of Bottom



*CMS, PLB850, 138389 (2024)*

- Significant flow of D from B decay
- Much smaller than prompt D
- Different degree of thermalization of charm and bottom

# Heavy Flavor Hadronization in QGP

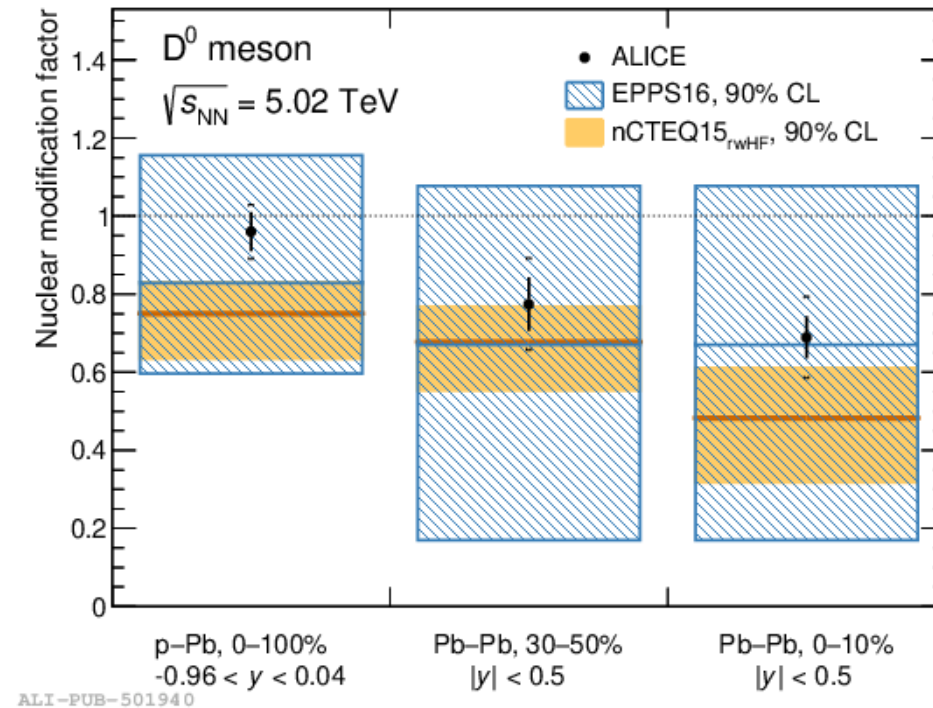


**QUESTION:**  
 Light hadrons have similar behavior, why we don't worry?

D is suppressed in all  $p_T$  at both RHIC and LHC  $\rightarrow N^{AA} < N_{bin} \times N^{pp}$

Contradict to  $N_{bin}$  scaling of  $\#ccbar$  ?!!

# Cold Nuclear Matter Effect?



ALICE, PLB839, 137796 (2023)

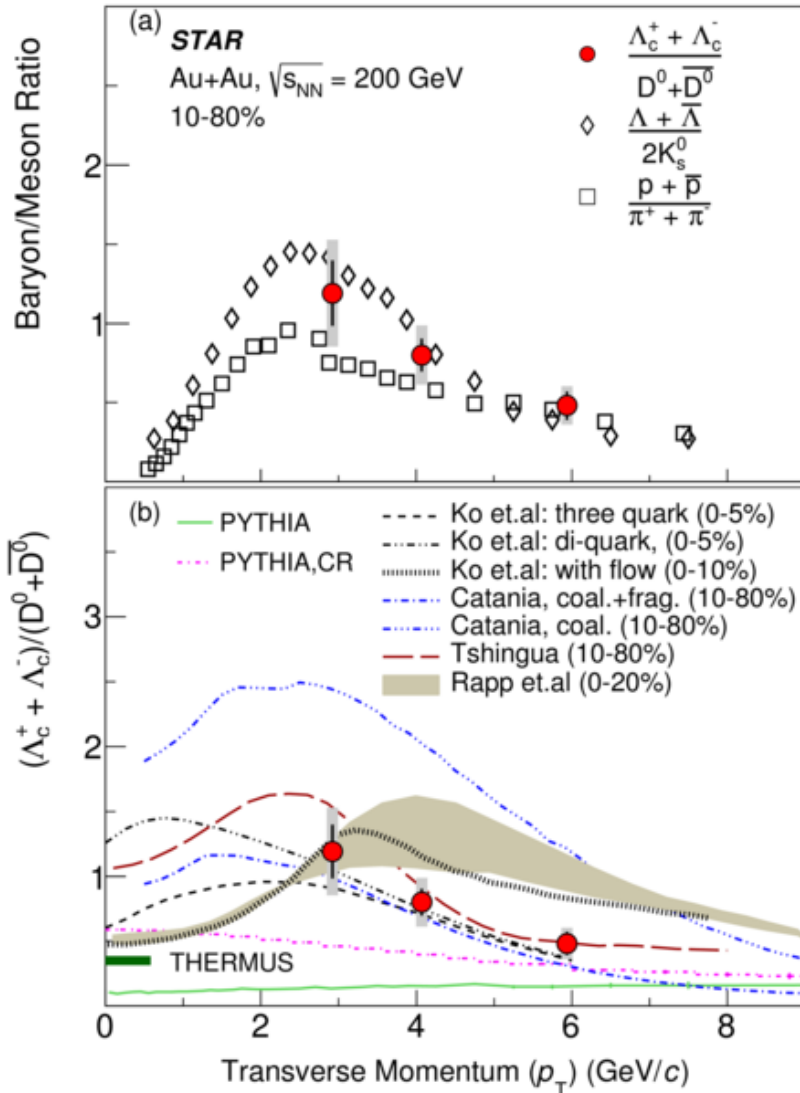
~30% suppression in  
central PbPb collisions

Decreasing trend from pPb to central PbPb

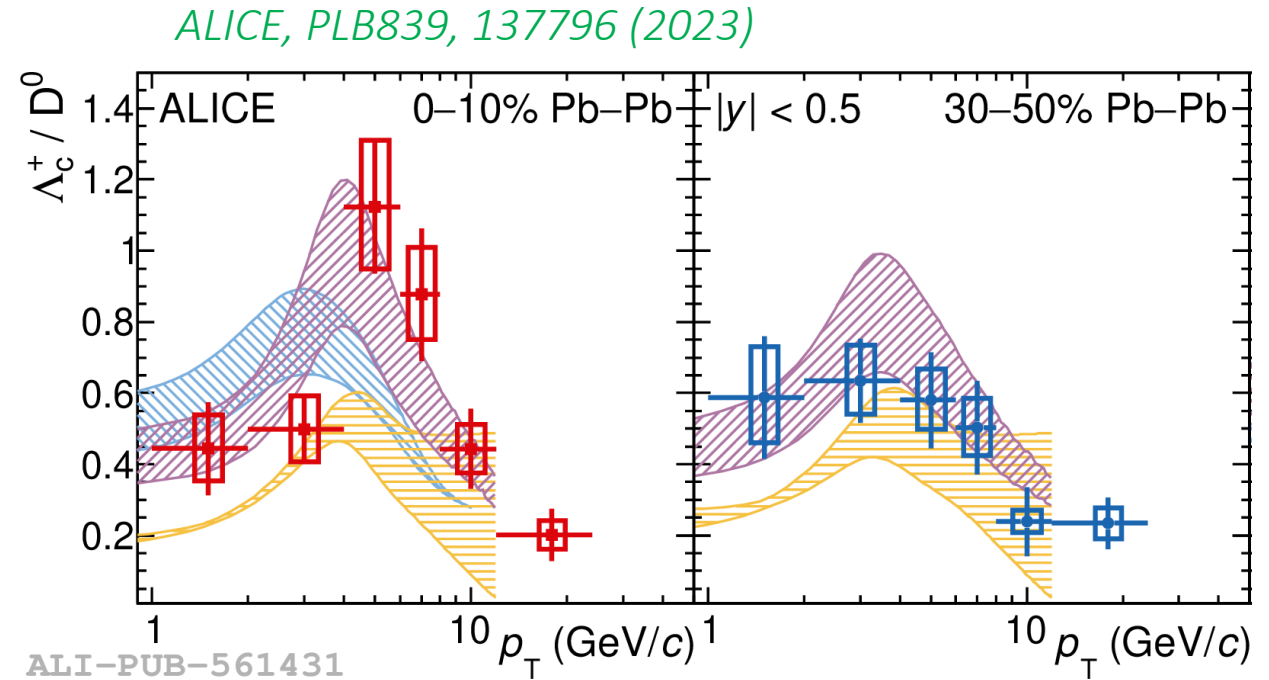
Consistent with models implemented shadowing effect within large uncertainties

Purely from shadowing? Is there any other effects on top of it? --> Particle ratios

# Charmed Baryon/Meson Ratio



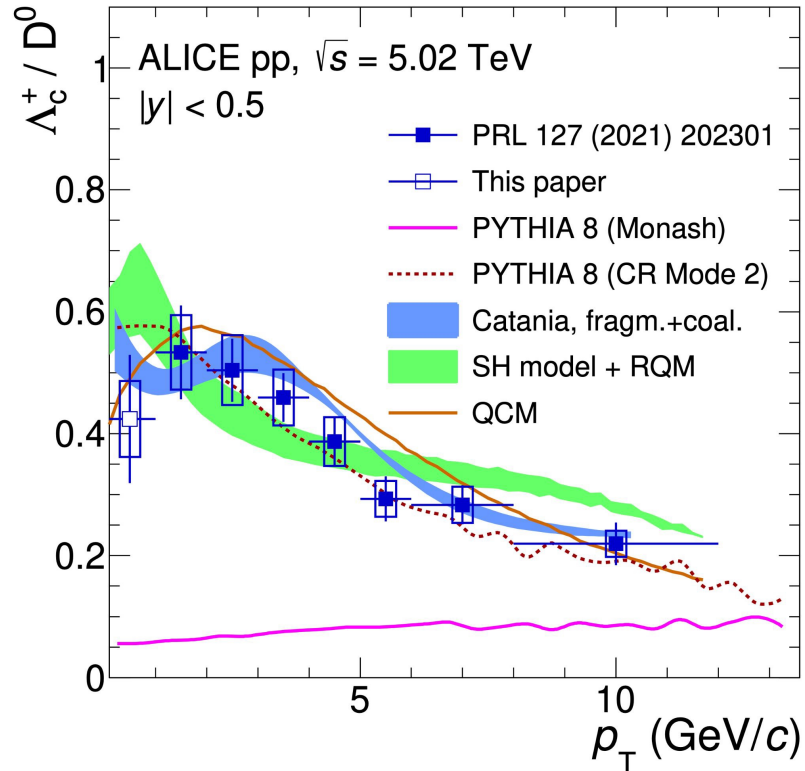
STAR, PRL124, 172301 (2020)



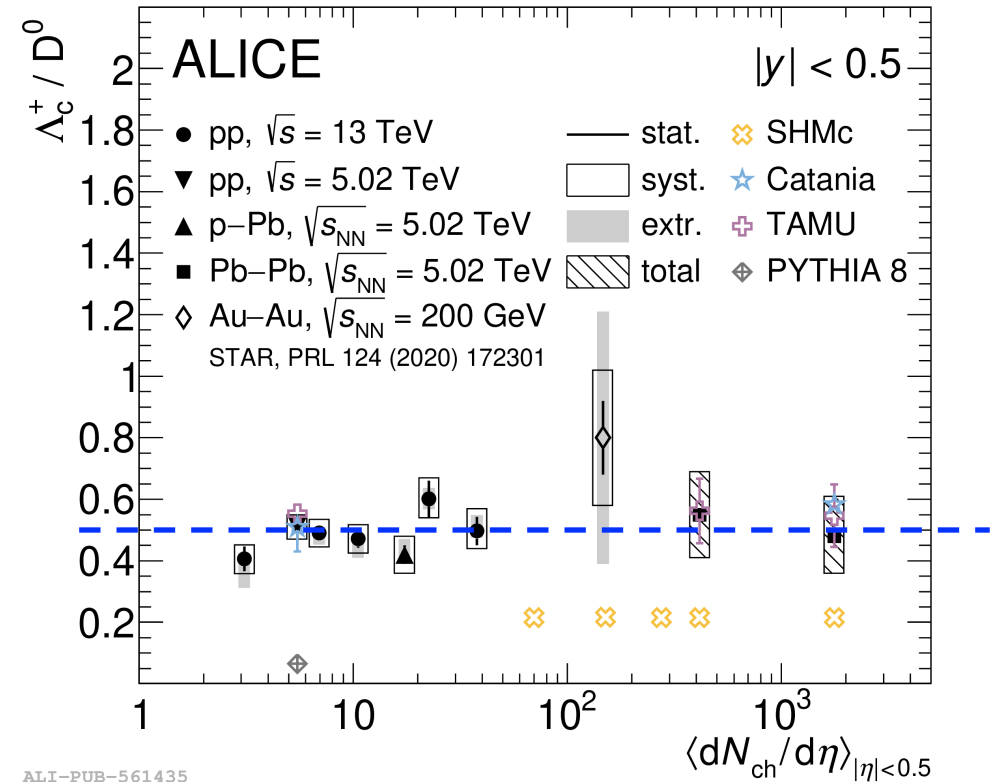
- Charmed baryon/meson ratio is similar to light hadrons
- Significantly higher than PYTHIA (constrained by ee/ep)
- Model including coalescence describe the enhanced ratios

# Ratios in Small System

ALICE, PRC107, 064901 (2023)



ALICE, PLB839, 137796 (2023)



- Significantly larger than default PYTHIA
- Qualitatively described by models

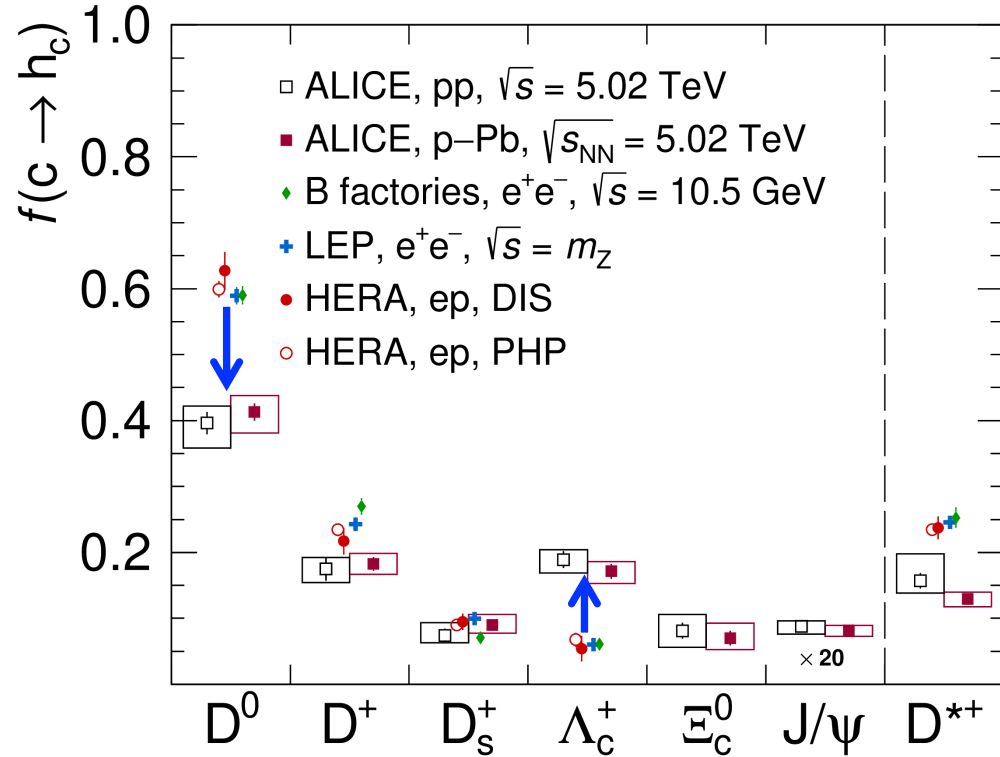
Smooth trend from pp to central AA

Same hadronization in pp and AA?

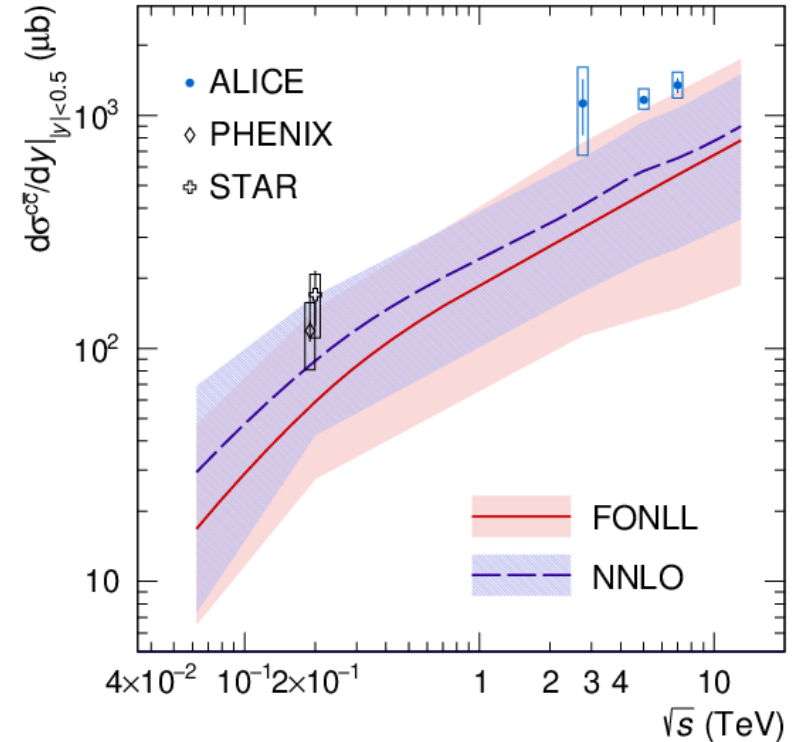
Coalescence  $\leftrightarrow$  QGP



# Relative Abundance in Small System



ALI-PUB-570972



ALI-PUB-500755

*ALICE, arXiv:2405.14571*  
*ALICE, PRD105, L011103 (2022)*

Suppression of mesons and enhancement of baryons

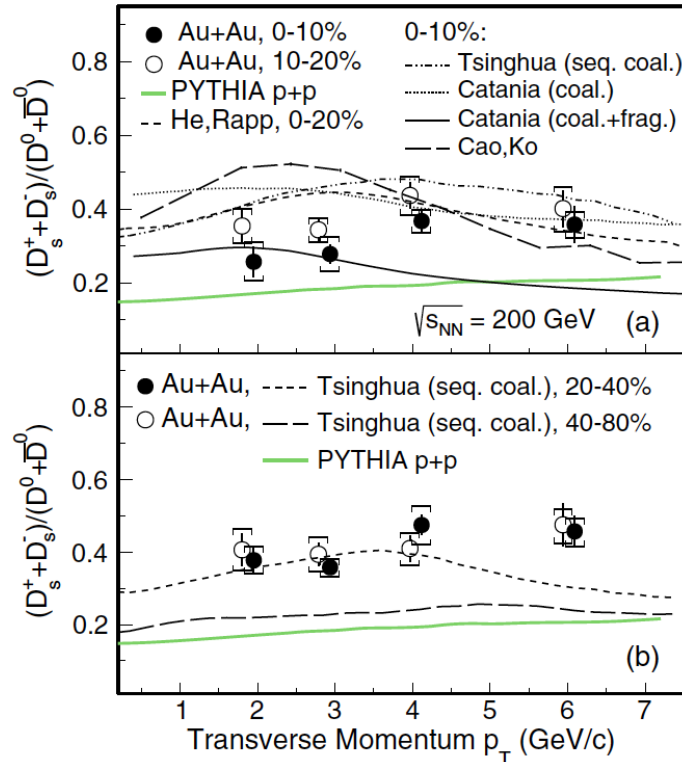
Total cross-section consistent with FONLL

**Redistribution** of charm quarks among hadrons

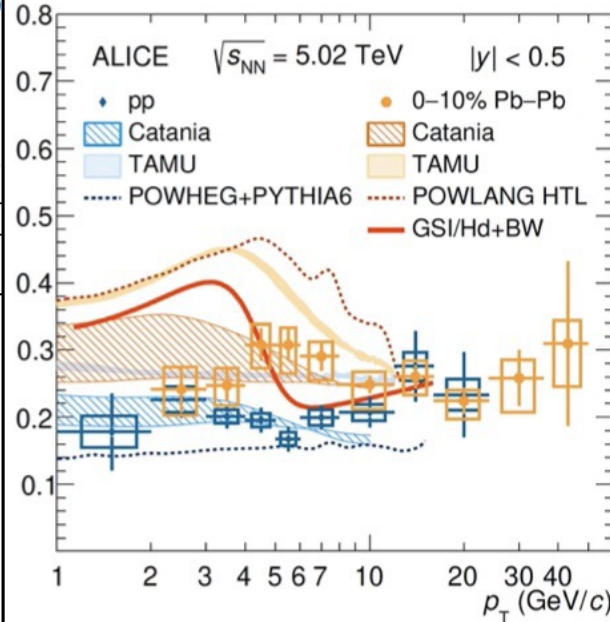




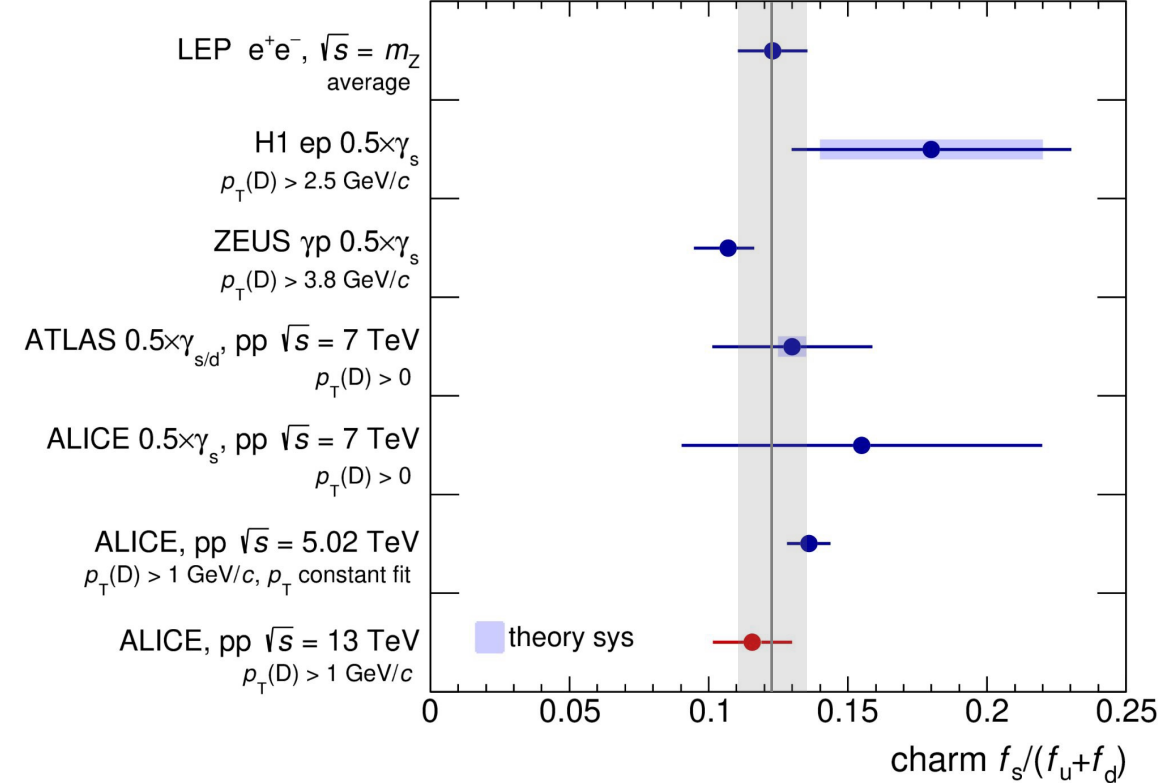
# Ds/D0 Ratio



STAR, PRL127, 092301 (2021)



ALICE, PLB827, 136986 (2022)



ALI-PUB-567901

$D_s/D^0$  in AA  $>$   $D_s/D^0$  in pp

$D_s/D^0$  in pp  $=$   $D_s/D^0$  in ee

Coalescence + **Strangeness enhancement** in QGP



# Summary of Open Heavy Flavor

Heavy quarks are unique probes of QGP due to their large masses

Extensive experimental studies have been conducted thanks to start-of-the-art silicon tracker

Experimental and theoretical studies show that

- Charm quark exhibit significant energy loss and collective motion
- Mass dependence of yield suppression and collective flow is observed
- Quark coalescence plays an important role in heavy quark hadronization at low/mid  $p_T$
- Dimensionless spatial diffusion coefficient of charm in the medium is constrained by comparing experimental results and theoretical calculations

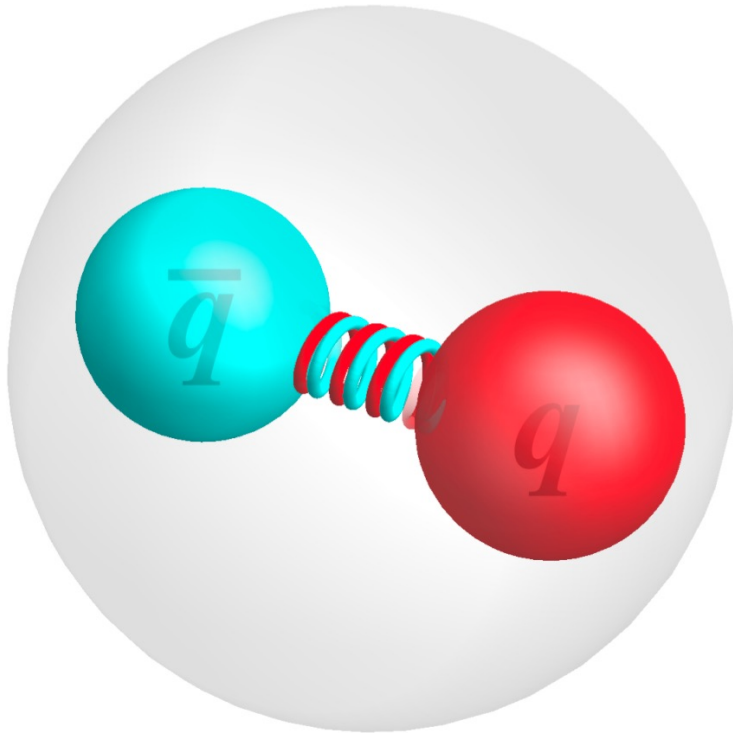


# Quarkonium

Bound state of quark and its own antiquark, usually refer to heavy quark

One of the simplest systems in QCD

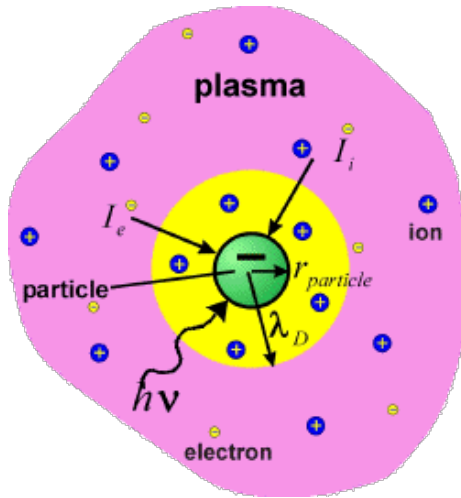
Analogue to hydrogen in atomic physics (QED)



State	$J/\psi$	$\chi_c$	$\psi'$	$\Upsilon$	$\chi_b$	$\Upsilon'$	$\chi'_b$	$\Upsilon''$
Mass [GeV]	3.10	3.53	3.68	9.46	9.99	10.02	10.26	10.36
$\Delta E$ [GeV]	0.64	0.20	0.05	1.10	0.67	0.54	0.31	0.20
$\Delta M$ [GeV]	0.02	-0.03	0.03	0.06	-0.06	-0.06	-0.08	-0.07
$r_0$ [fm]	0.50	0.72	0.90	0.28	0.44	0.56	0.68	0.78

Please see also Enrico Scomparin's lecture at Quark Matter 2023: [slides](#)

# Debye Screening in Plasma

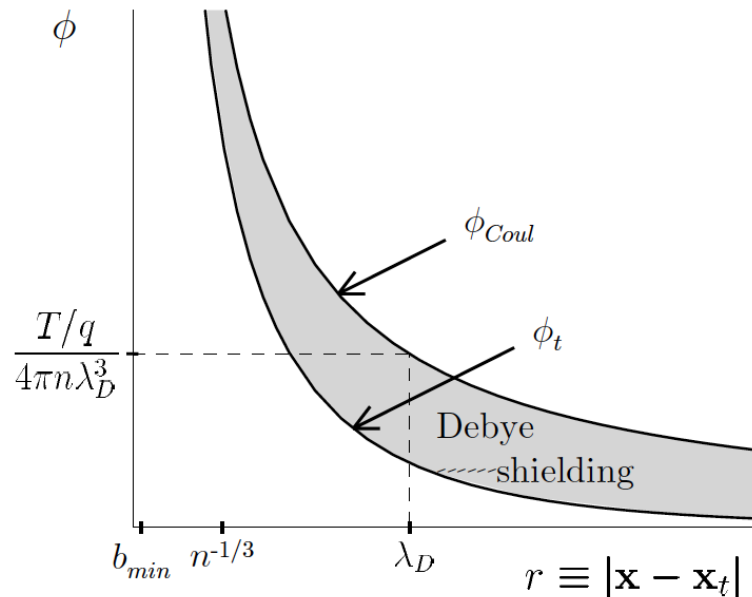


Potential of point charge in vacuum:

$$V = \frac{Q}{4\pi\epsilon_0 r}$$

Potential of test charge in a plasma:

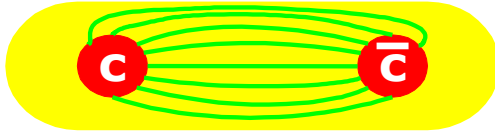
$$V = \frac{Q}{4\pi\epsilon_0 r} e^{-r/\lambda_D}$$



Electromagnetic interaction limited in the Debye radius

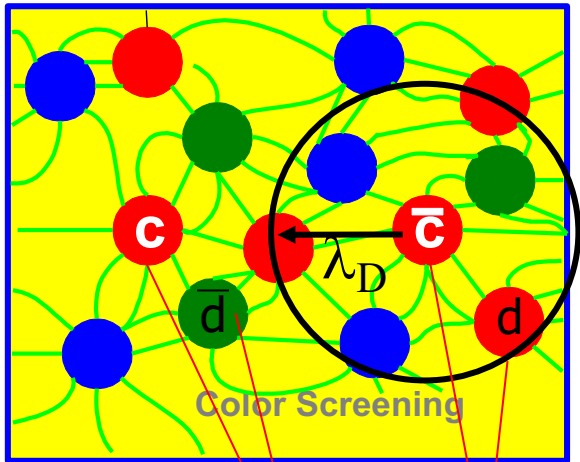
# Debye Screening of Strong Interaction in QGP

Vacuum



$$V(r) = \sigma r - \frac{\alpha}{r}$$

QGP

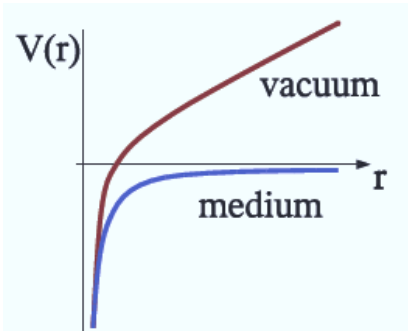


$$\frac{\sigma}{\mu} \{1 - e^{-\mu r}\} - \frac{\alpha}{r} e^{-\mu r}$$

$$\mu = 1/\lambda_D$$

$D^+$

$D^-$



Strong interaction between heavy quark and its antiquark is reduced in the deconfined medium due to the surrounding free quarks and gluons

Bound state will be dissociated into open heavy flavor hadrons when the Debye radius is smaller than the size of the bound state

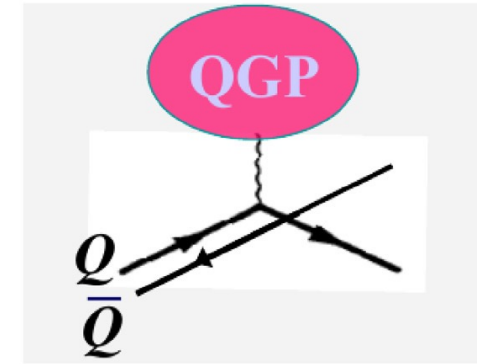
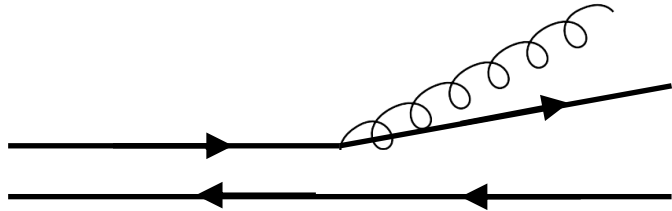
Suppression of quarkonium in relativistic heavy ion collisions should provide a “smoking-gun” signature of QGP formation

**QUESTION:**

*Will light hadrons such as  $\phi$  have the same effect?*

*T. Matsui, H. Satz, PLB174, 416 (1986)*

# Dynamic Dissociation in QGP



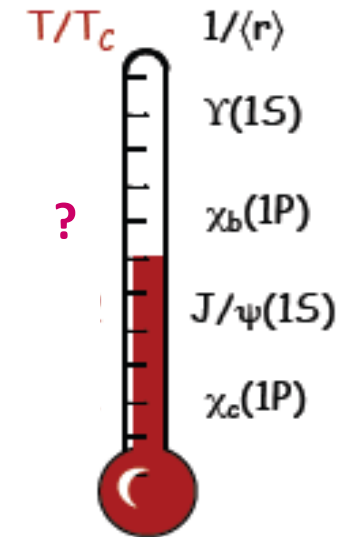
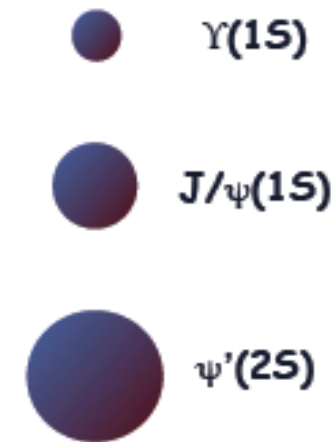
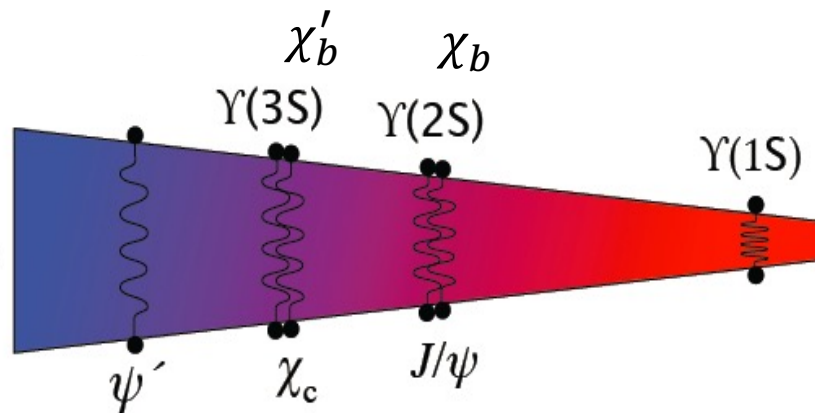
*M. He, H. van Hees and R. Rapp, PNP130, 104020 (2023)*

Quarkonium may absorb a gluon or interact with partons in QGP and dissociated

Dissociation rate depends also on QGP temperature, binding energy of quarkonium etc

# Quarkonium Suppression: QGP Thermometer

## Plasma thermometer



Debye radius is inversely proportional to the temperature of QGP

Different quarkonium states dissociate at different temperatures

→ Sequential melting

By measuring sequential melting, one get some information of QGP temperature



# Quarkonium Reconstruction

$$J/\psi \rightarrow e^+e^- \text{ or } \mu^+\mu^-$$

Branching ratio  $\sim 6\%$

$$\psi(2S) \rightarrow e^+e^- \text{ or } \mu^+\mu^-$$

Branching ratio  $\sim 0.8\%$

$$\psi(2S) \rightarrow J/\psi \pi^+ \pi^-$$

Branching ratio  $\sim 35\%$

$$35\% \times 6\% = 2.1\%$$

$$\Upsilon(1S) \rightarrow e^+e^- \text{ or } \mu^+\mu^-$$

Branching ratio  $\sim 2.4\%$

$$\Upsilon(2S) \rightarrow e^+e^- \text{ or } \mu^+\mu^-$$

Branching ratio  $\sim 1.9\%$

$$\Upsilon(3S) \rightarrow e^+e^- \text{ or } \mu^+\mu^-$$

Branching ratio  $\sim 2.2\%$

**QUESTION:**

*Why the dilepton decay branching ratios are different?*

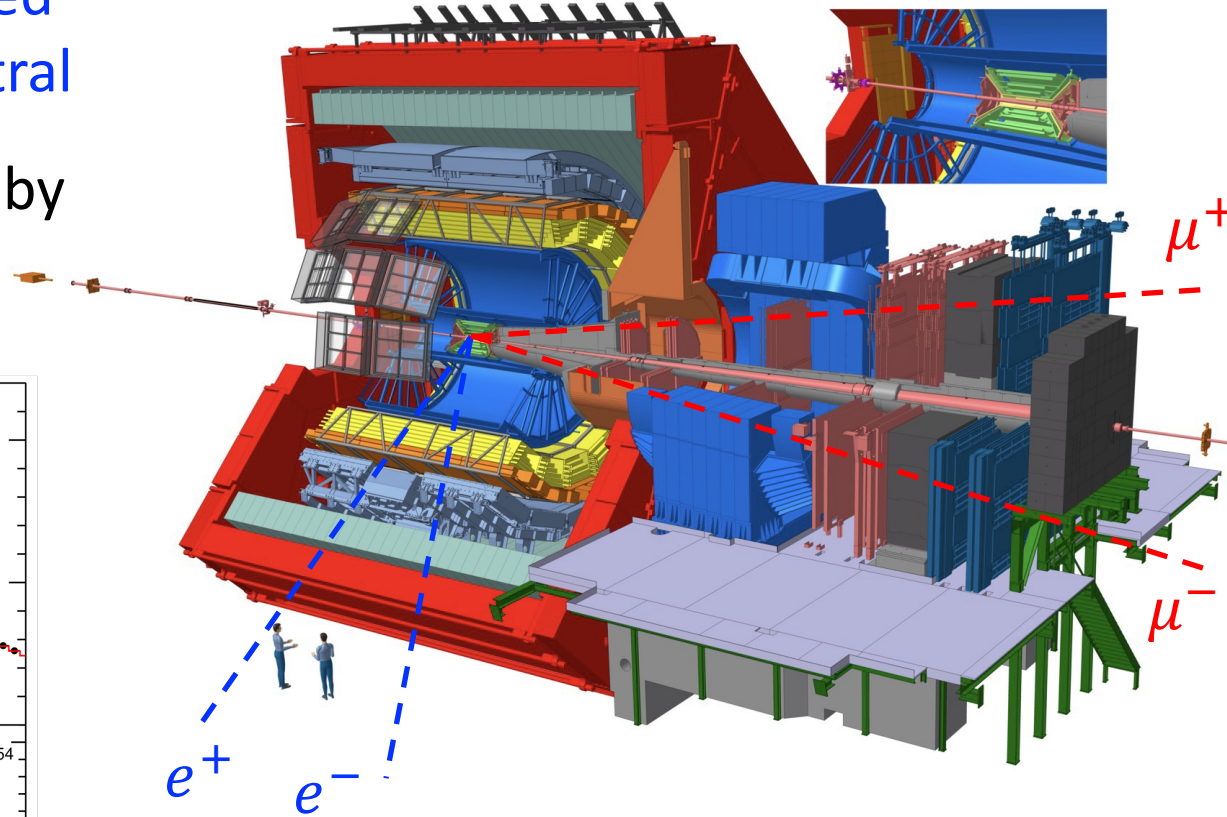
Typically need 1 million p+p events to reconstruct one  $J/\psi$   
Much less for other quarkonium states

Good **trigger** and **PID** detectors are crucial

# Quarkonium in ALICE @ Run2

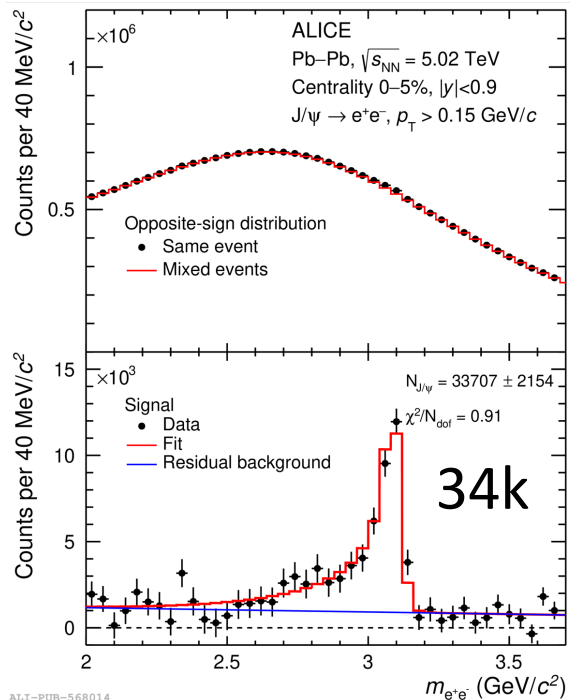
22  $\mu\text{b}^{-1}$  MB triggered  
105  $\mu\text{b}^{-1}$  0-10% central

Electrons identified by  
dE/dx in TPC



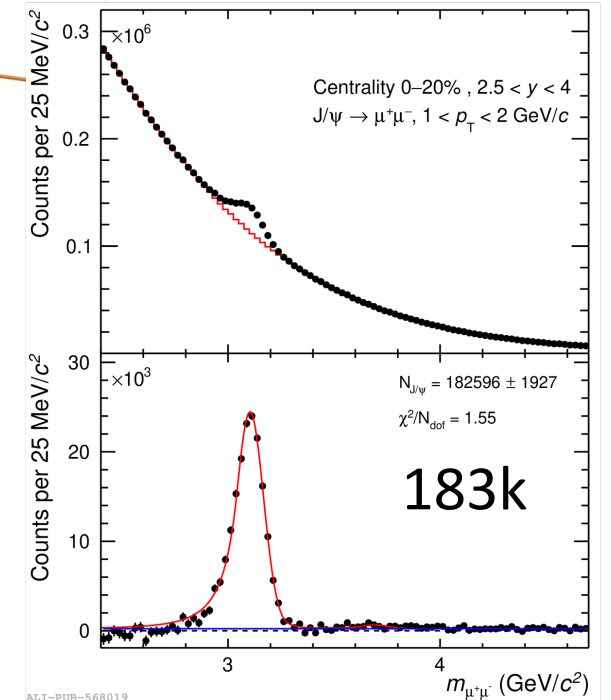
756  $\mu\text{b}^{-1}$   $\mu^+\mu^-$  trigger

Hadrons suppressed by  
absorbers ( $10+7.2 \lambda_I$ )



$e^+e^-$  at mid-rapidity  
 $\mu^+\mu^-$  at forward rapidity

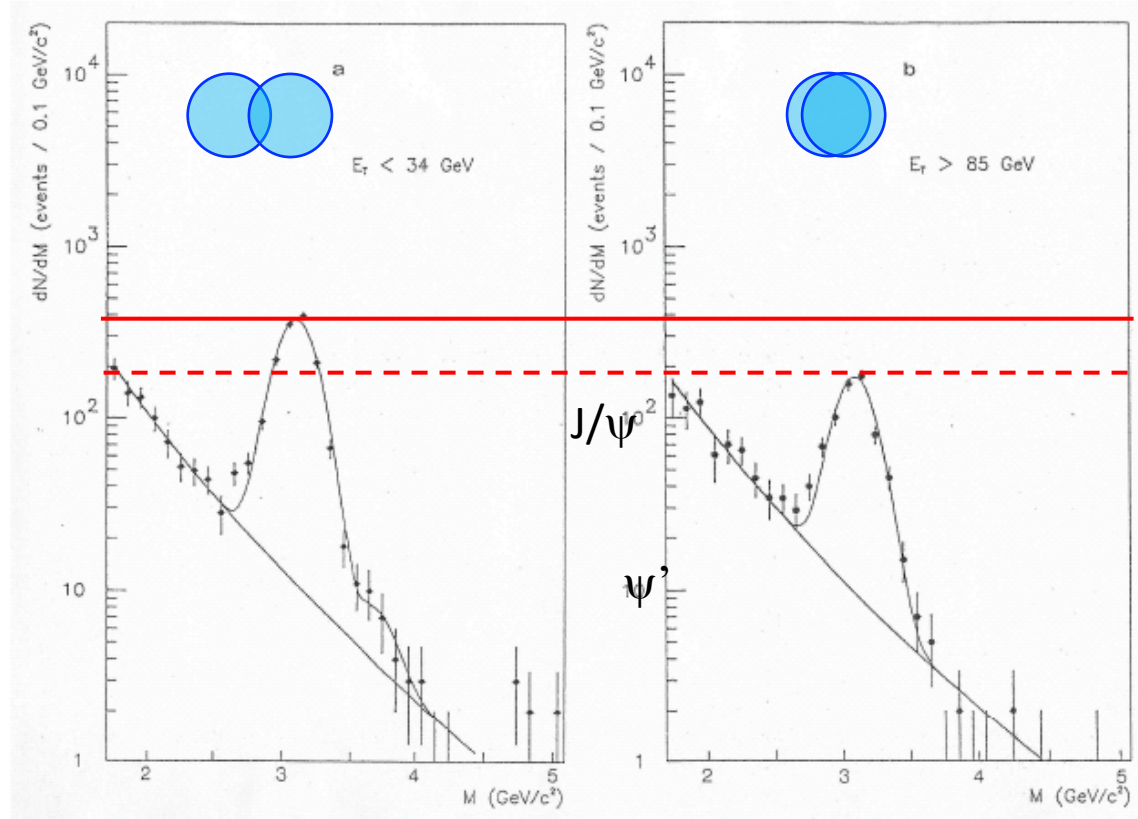
ALICE, JHEP02, 066 (2024)





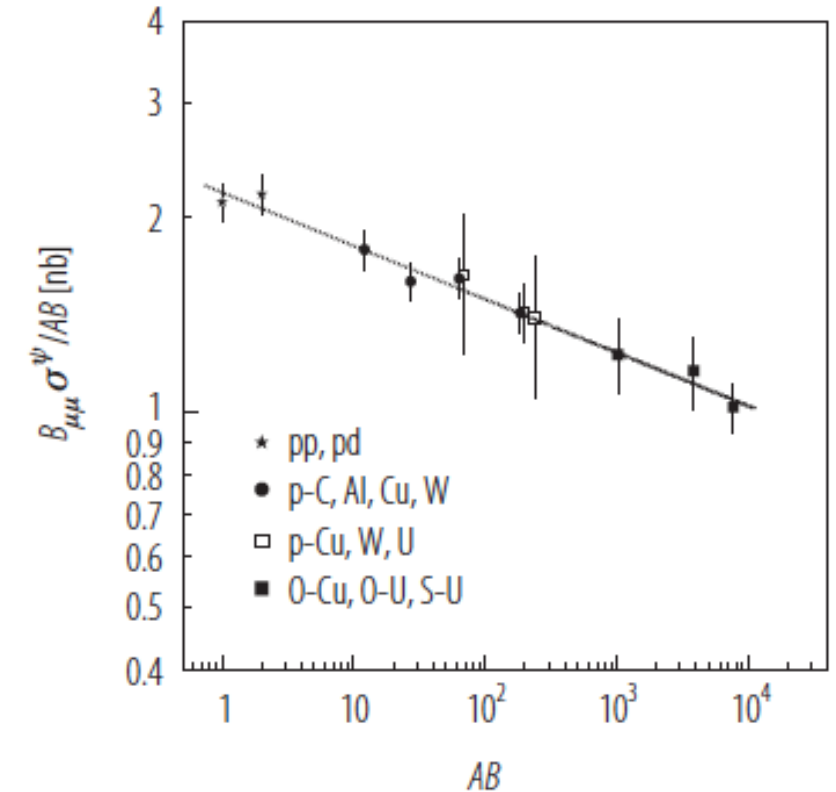
# First Observation of $J/\psi$ Suppression

200 AGeV O+U collisions



*NA38, PLB220, 471 (1989)*

Figure 1: First observation of the  $J/\psi$  suppression effect in O(200 AGeV)-U collisions experiment at CERN-SPS. When comparing the invariant-mass spectrum of muon pair peripheral collisions (characterized by a small transverse energy  $E_T < 34$  GeV; left panel) central collisions (at high transverse energy,  $E_T > 85$  GeV; right panel), a reduction of  $t$  over the Drell-Yan continuum is apparent (from [8]).

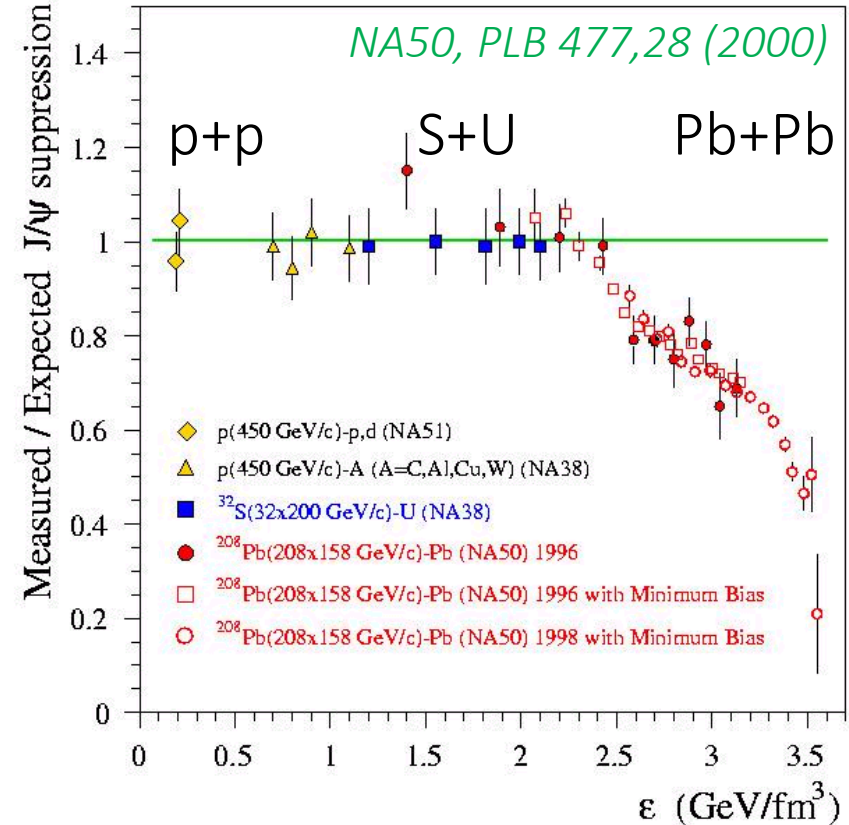
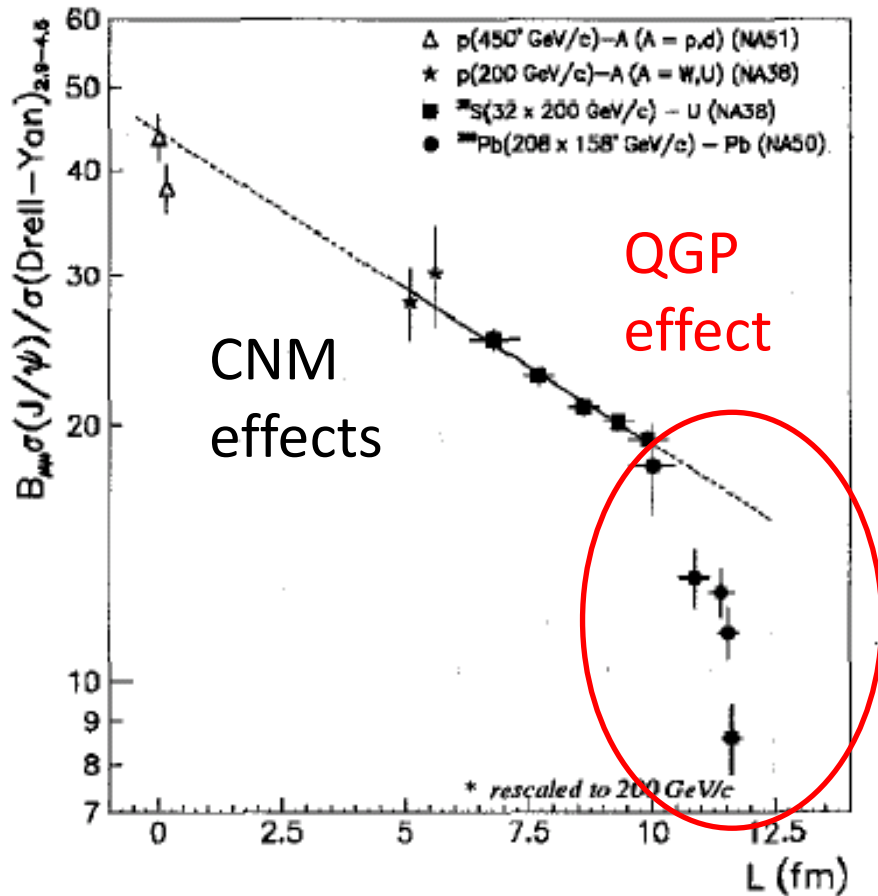


But observed also in p+A collisions

Explained by Cold Nuclear Matter (CNM effects)



# Anomalous $J/\psi$ Suppression at SPS

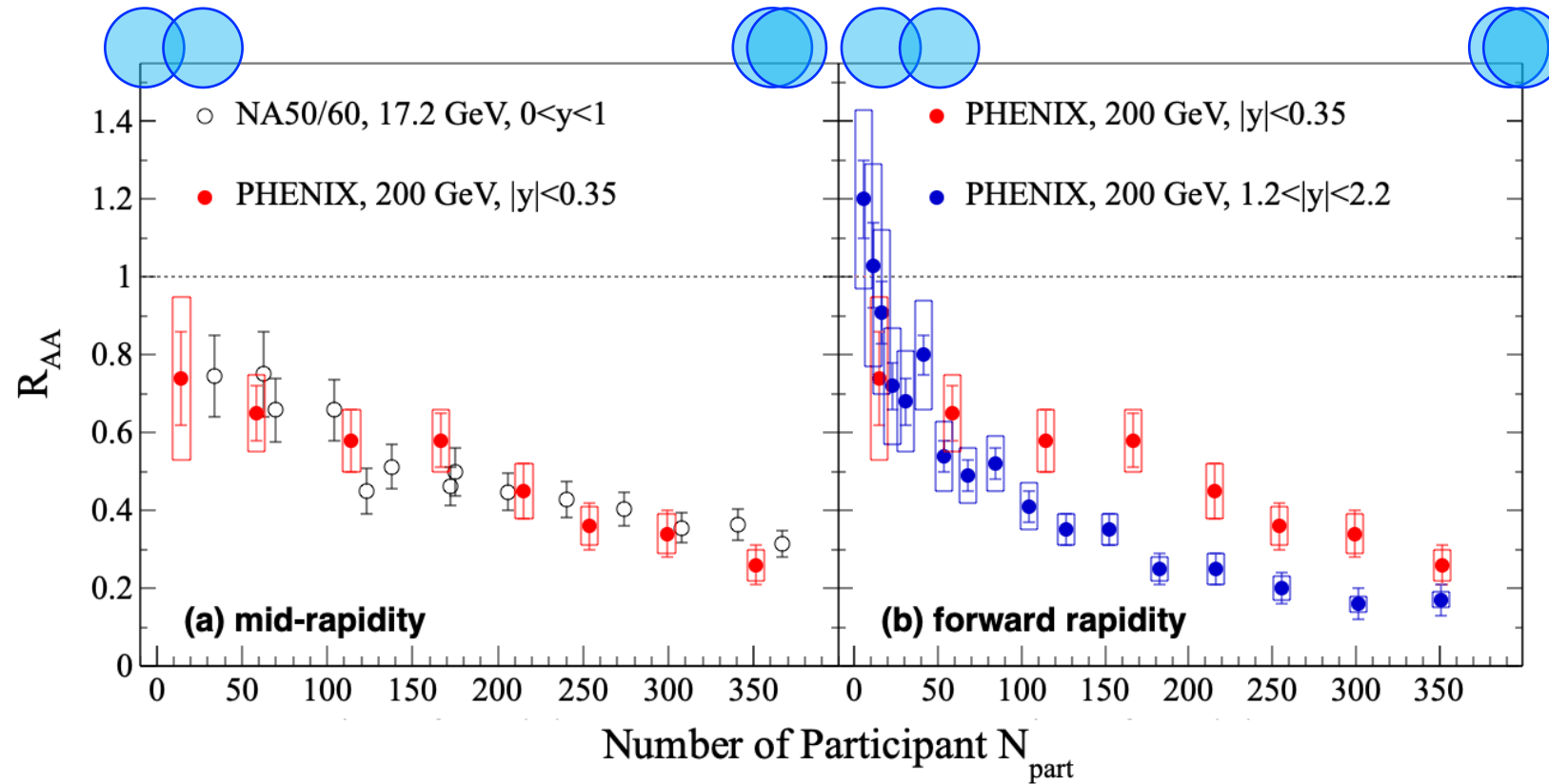


Beyond normal suppression  
observed in Pb+Pb

**“Anomalous” suppression**

gluons. Therefore, we must conclude that the  $J/\psi$  suppression pattern observed in our data provides significant evidence for deconfinement of quarks and gluons in the Pb-Pb collisions probed by NA50.

# Suppression at RHIC

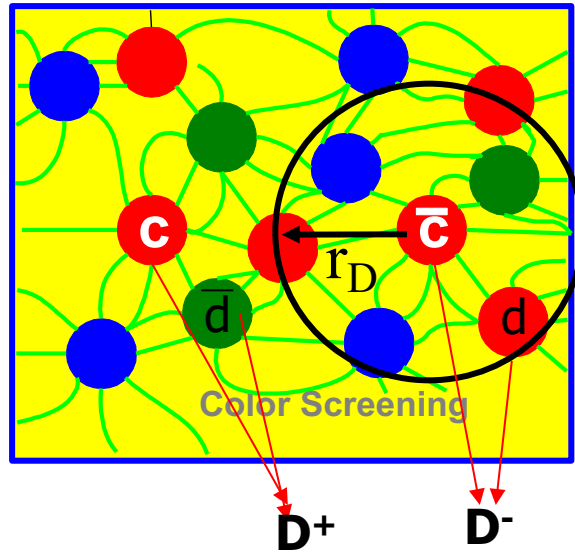


J/ $\psi$  suppression in 200GeV Au+Au at RHIC **similar** as  
J/ $\psi$  suppression in 17.2GeV Pb+Pb at SPS  
Despite the increase of energy by a factor of 10+

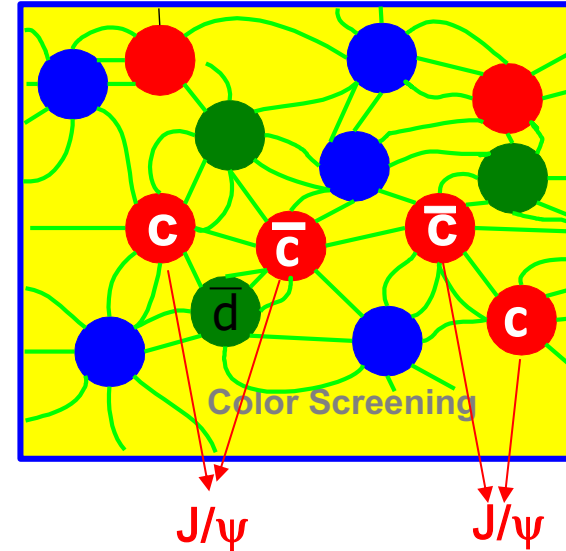
Stronger suppression at  
forward than at mid-rapidity

**Puzzle!!**

# Color Screening vs. Regeneration



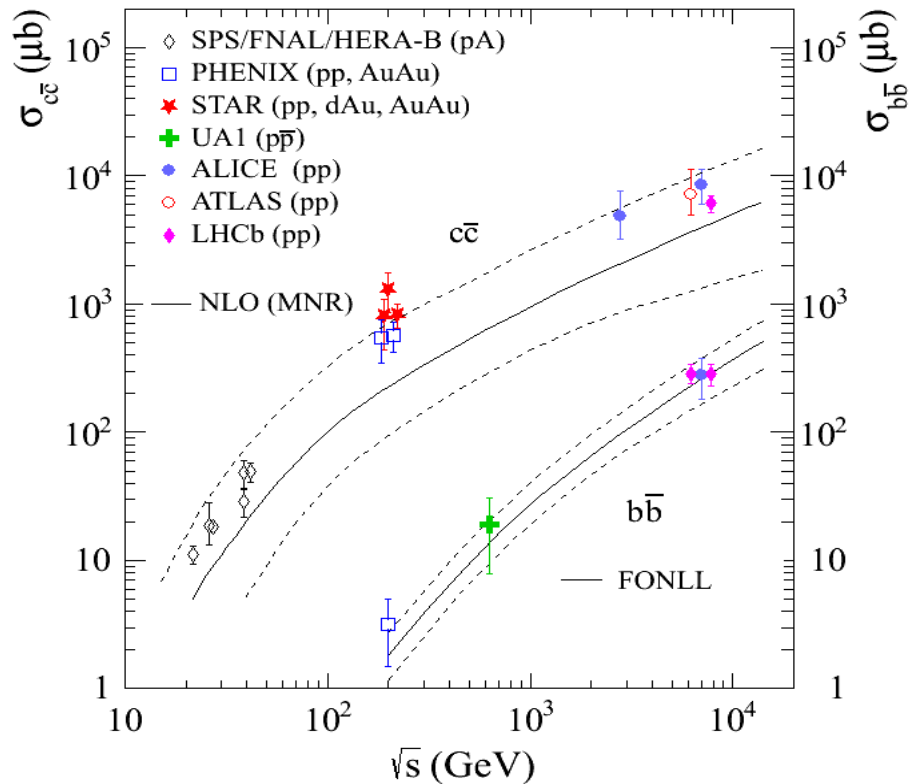
Quarkonium melting in QGP



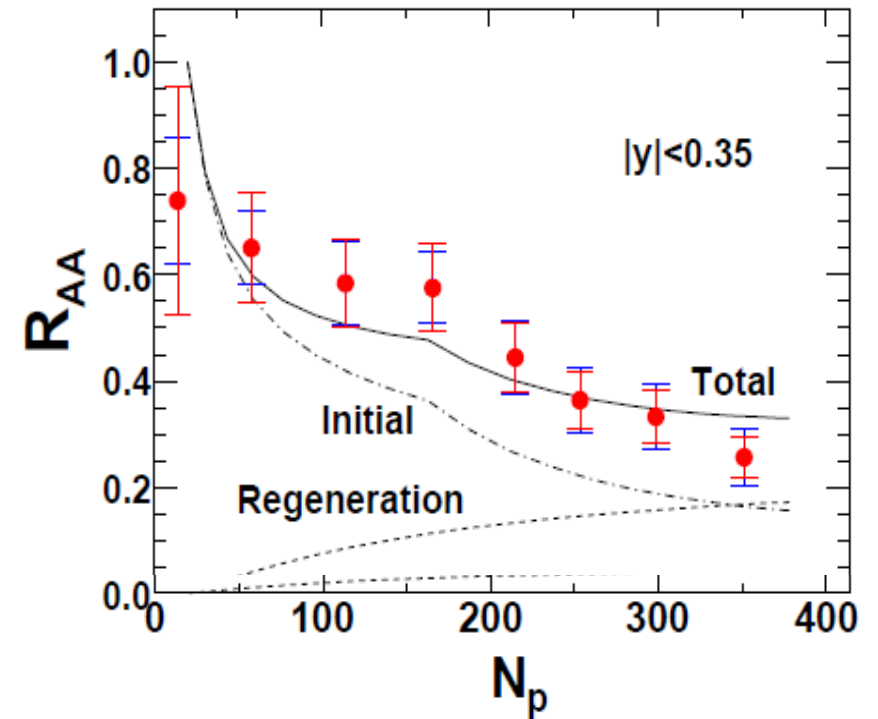
Quarkonium regeneration in QGP

QGP formation is the prerequisite of both effects

# Puzzle Solved



*Z. Qu, Y. Liu, N. Xu, P. Zhuang, NPA830, 335c (2009)*



Charm cross-section increases with energy

More regeneration at higher energy and in central collisions

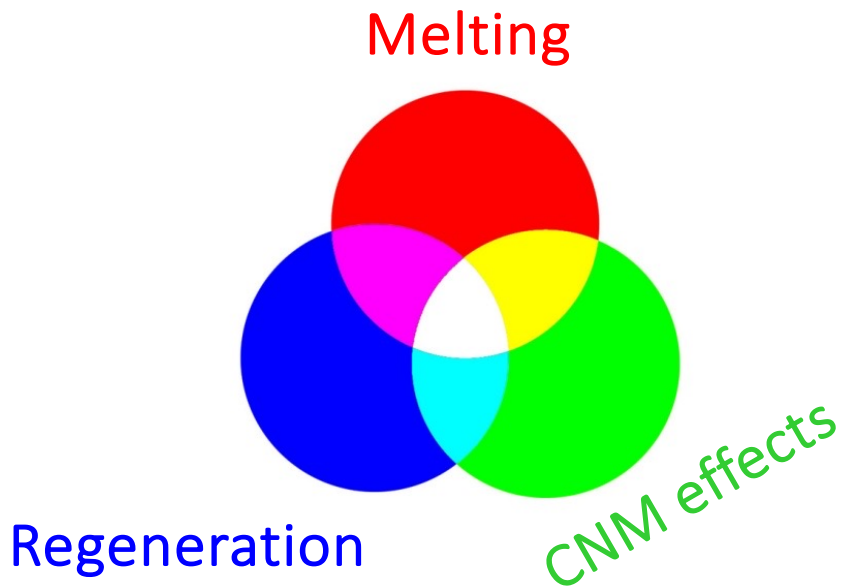
Theoretical calculations with regeneration can describe both RHIC and SPS data

In that the end?



# Quarkonium in Heavy-Ion Collisions

Quarkonium production in heavy-ion collisions are the interplay of **color-screening/melting**, **regeneration** in QGP and **CNM effects**



How to disentangle them?

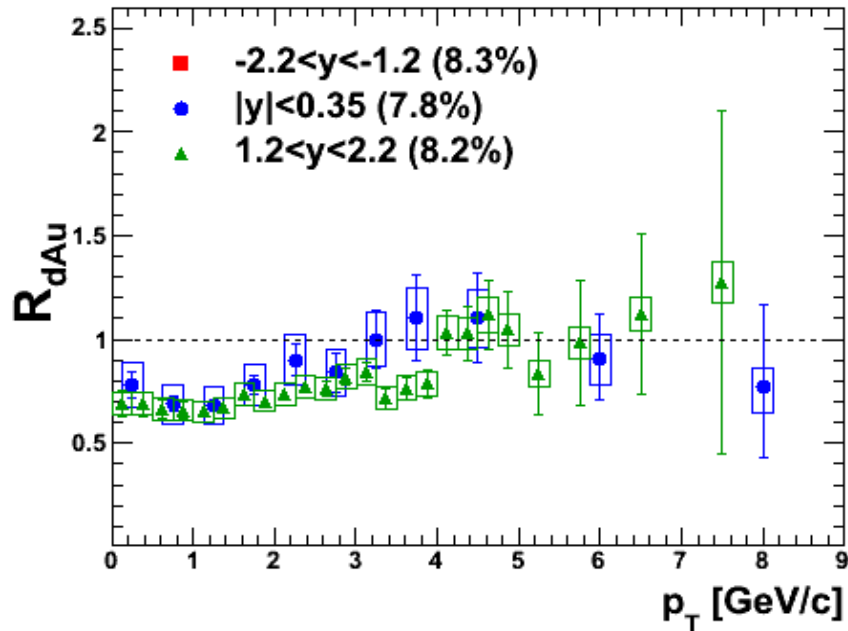
How to prob QGP with quarkonium?

Each of the effects have different dependence on

- $p_T$
- energy
- quarkonium size



# Move to High $p_T$ to Study Color Screening



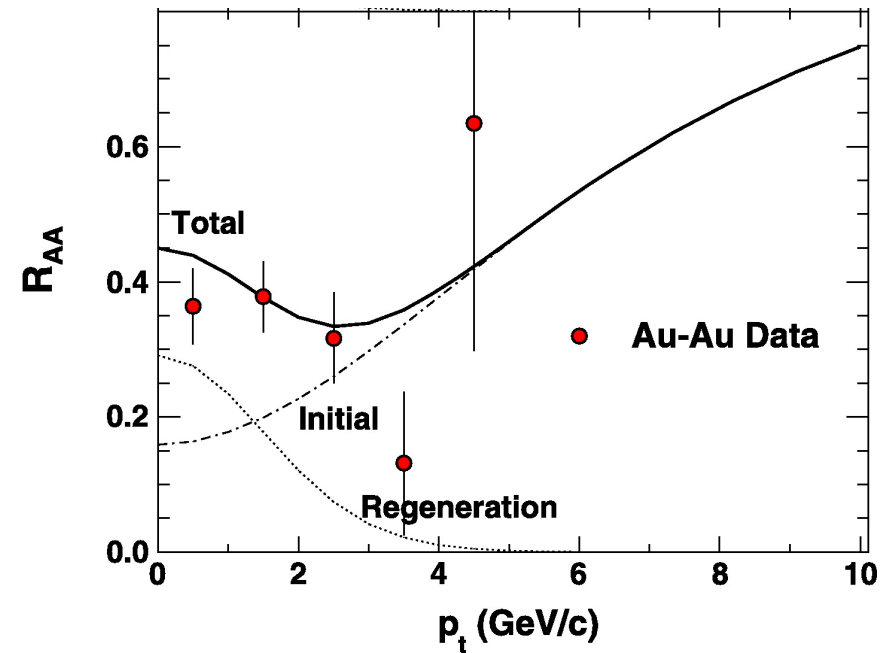
At  $p_T > 4$  GeV/c,

CNM is negligible

High- $p_T$   $J/\psi$ : clean probe of color screening

Very challenging measurement:

- Only  $< 1\%$  of  $J/\psi$  are at high- $p_T$

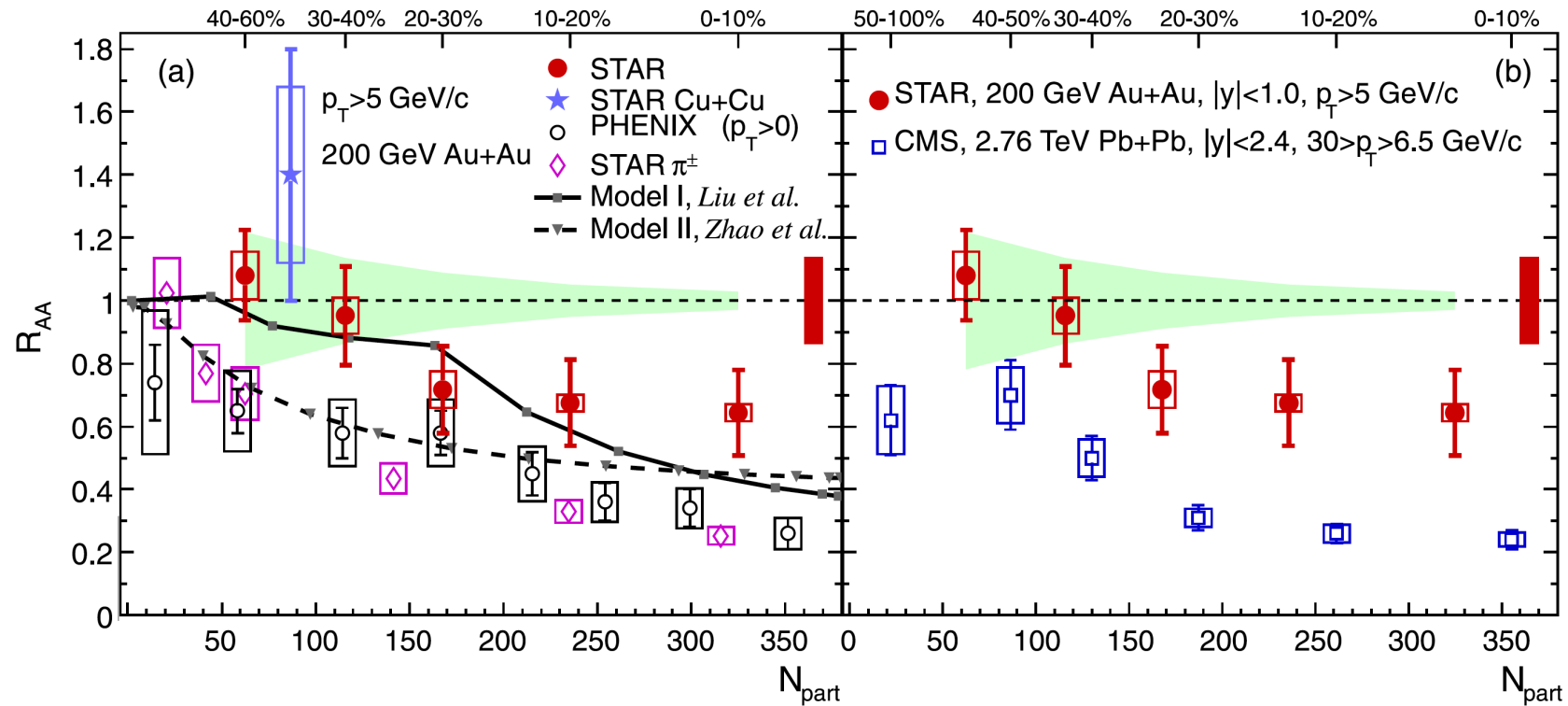


At  $p_T > 5$  GeV/c,

regeneration is negligible

# Measurements in Au+Au Collisions

STAR, PLB722, 55 (2013)



Significant suppression of high- $p_T$   $J/\psi$  observed in central Au+Au collisions

“Points to the color screening feature”



# Improved Measurements with Muons

Au+Au @ 200 GeV, Inclusive  $J/\psi$

★ STAR:  $J/\psi \rightarrow \mu^+ \mu^-$ ,  $|y| < 0.5$

□ Systematic uncertainty

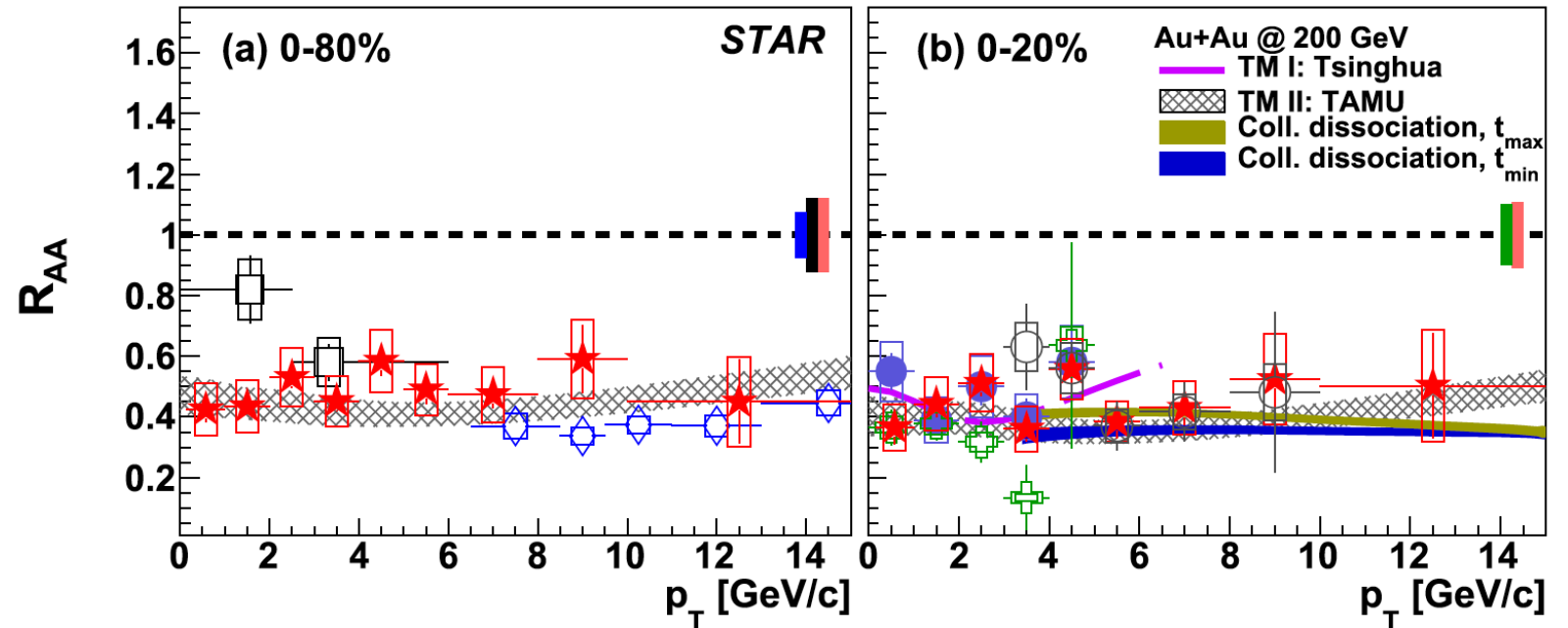
+ PHENIX:  $J/\psi \rightarrow e^+ e^-$ ,  $|y| < 0.35$

○ ● STAR:  $J/\psi \rightarrow e^+ e^-$ ,  $|y| < 1$

Pb+Pb @ 2.76 TeV

□ ALICE: Inclusive  $J/\psi$ , 0-40%,  $|y| < 0.8$

◇ CMS: Prompt  $J/\psi$ , 0-100%,  $|y| < 2.4$



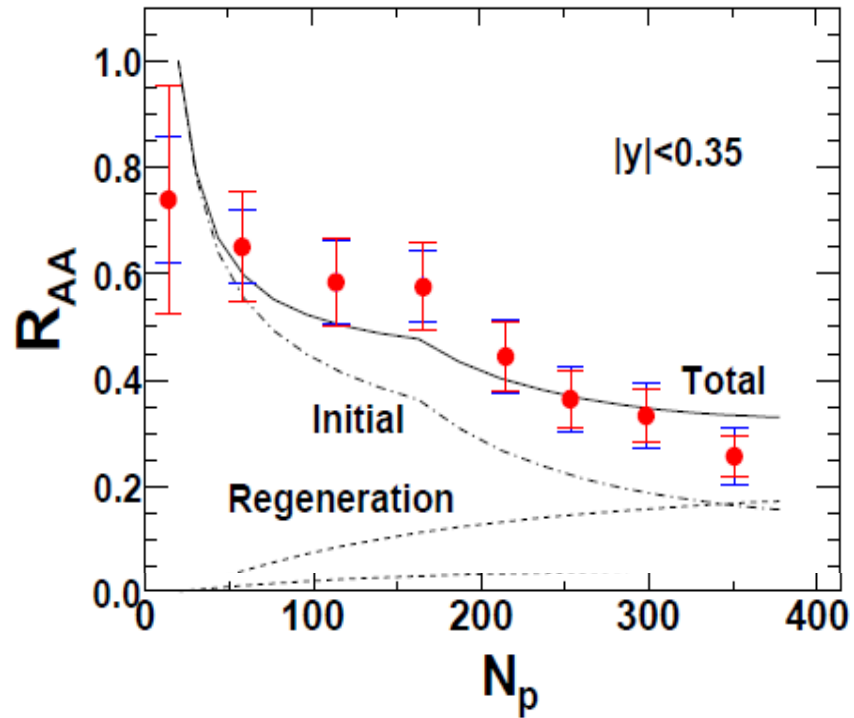
STAR, PLB797, 134917 (2019)

Significantly improves precision, and extends to low and higher  $p_T$

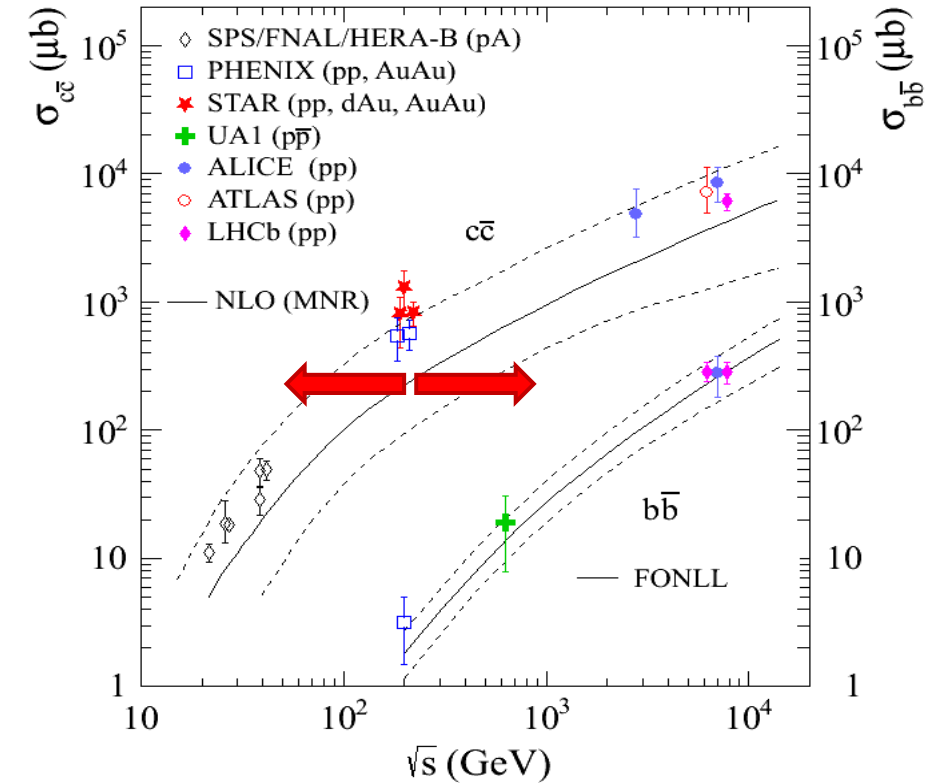
Significant suppression observed at high  $p_T$

“Providing strong evidence for the color-screening in the deconfined medium”

# Energy Dependence of Charm Yield



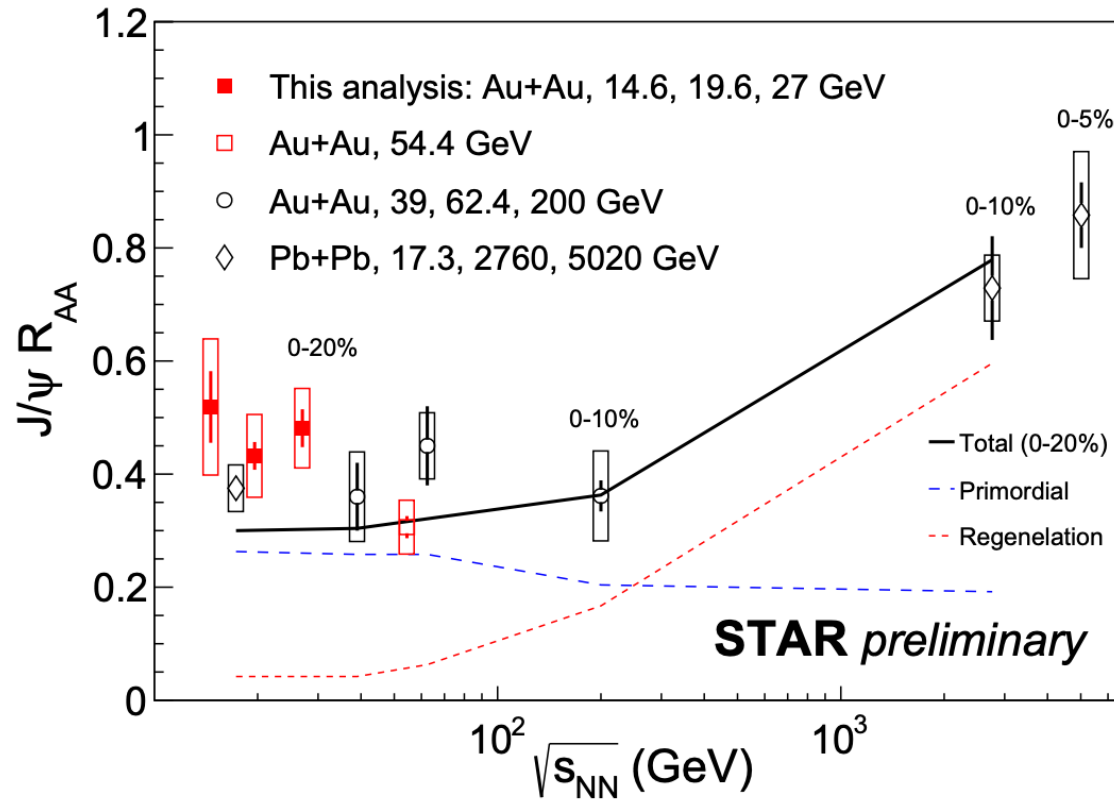
Theory predicts comparable contribution of color screening and regeneration in central collisions at RHIC



It's contribution should be less important at lower energy but dominant at LHC energy



# Energy Dependence of $J/\psi$ Suppression



NA50, PLB 477, 28 (2000)  
Wei Zhang, QM 2023  
STAR, PLB 771, 13 (2017)  
Kaifeng Shen, SQM 2021  
ALICE, PLB 734, 314 (2014)  
ALICE, PLB 849, 138451 (2024)

SPS



RHIC



LHC

CNM + screening

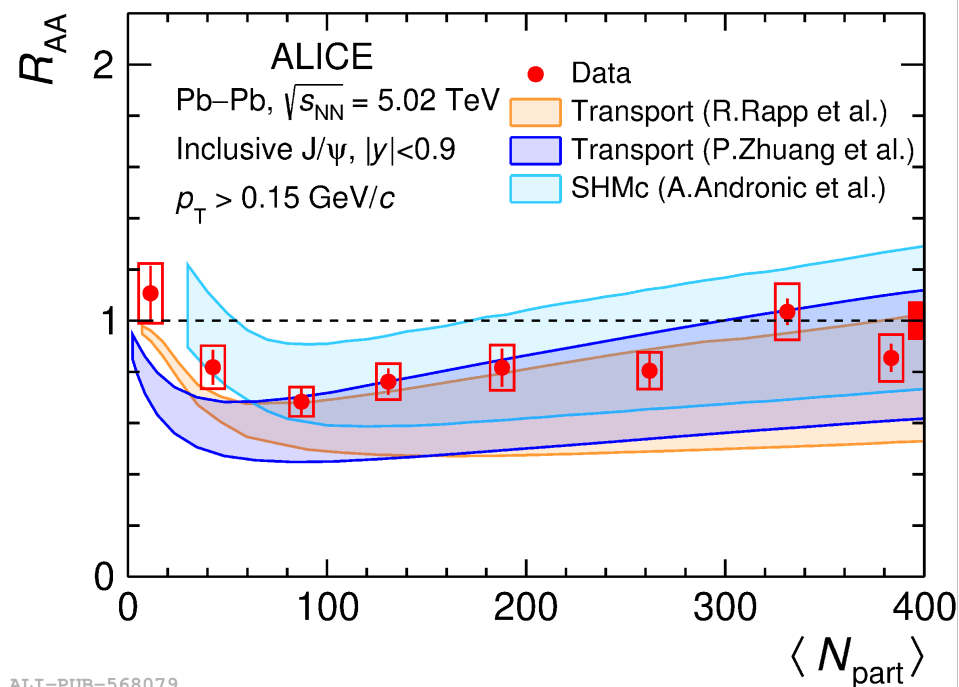
CNM + screening + regeneration

Regeneration domain



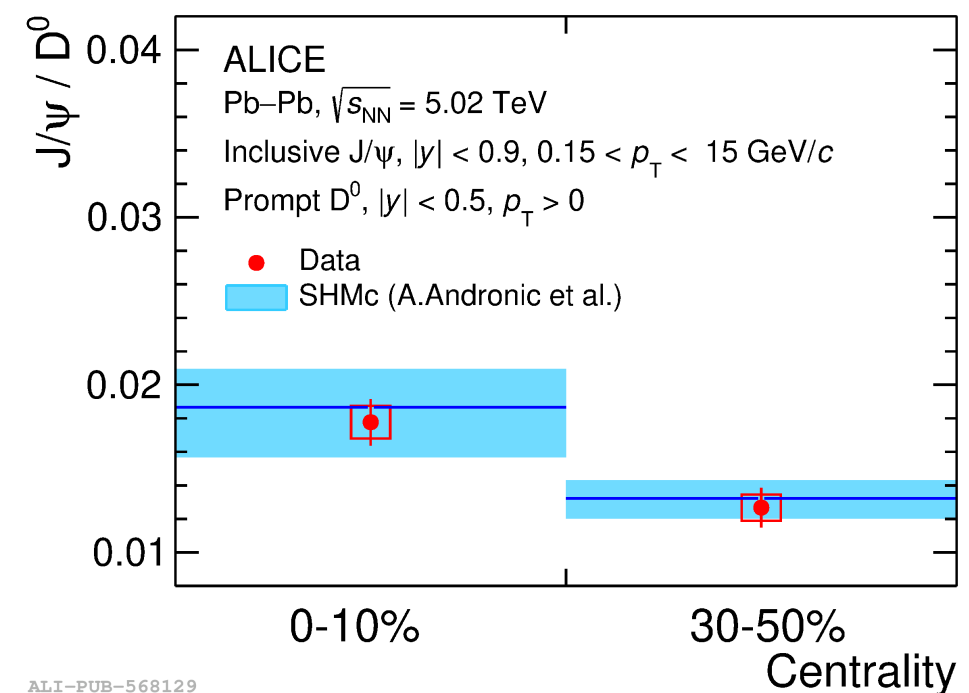
# J/ $\psi$ Yield at LHC

ALICE, PLB 849, 138451 (2024)



## Centrality dependence:

- Increase towards central collisions

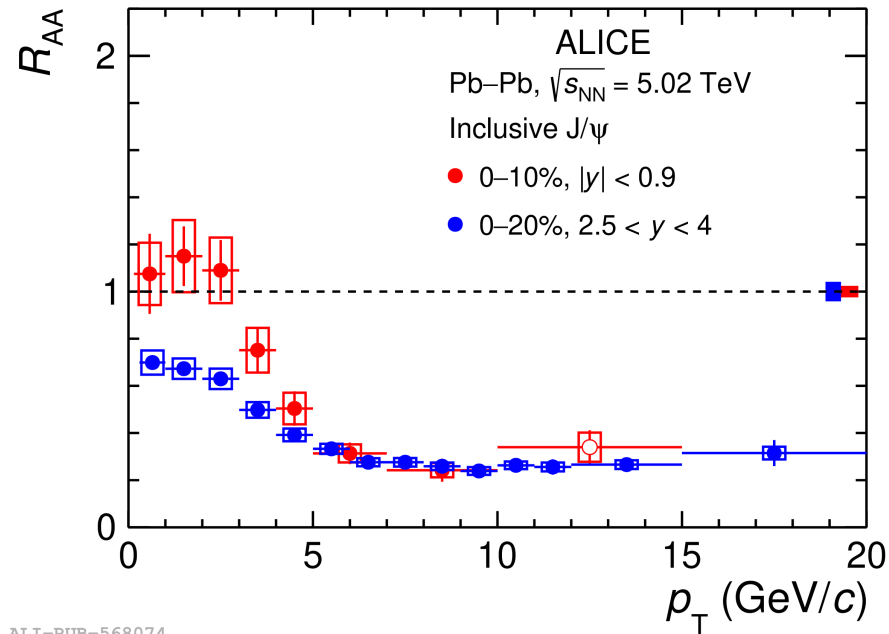


## Particle ratio:

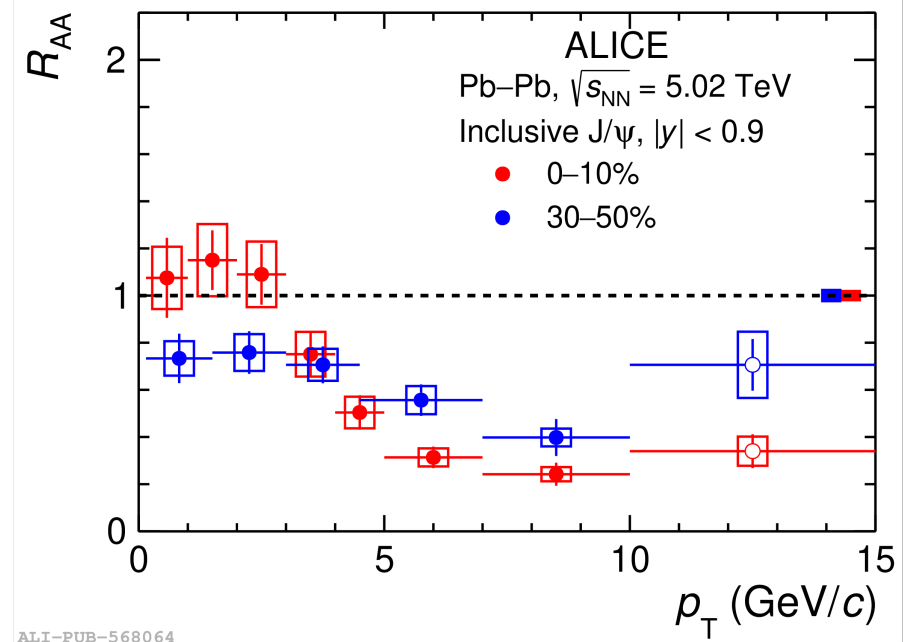
- Increase of J/ $\psi$  yield with respect to  $D^0$  in central collisions

# Differential Measurements at ALICE

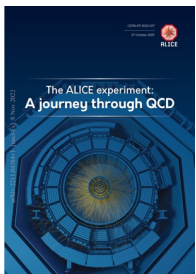
ALICE, PLB 849, 138451 (2024)



- Clear rapidity dependence at low- $p_T$
- Similar suppression at high- $p_T$



- Clear centrality dependence
- Opposite at low and high  $p_T$  region

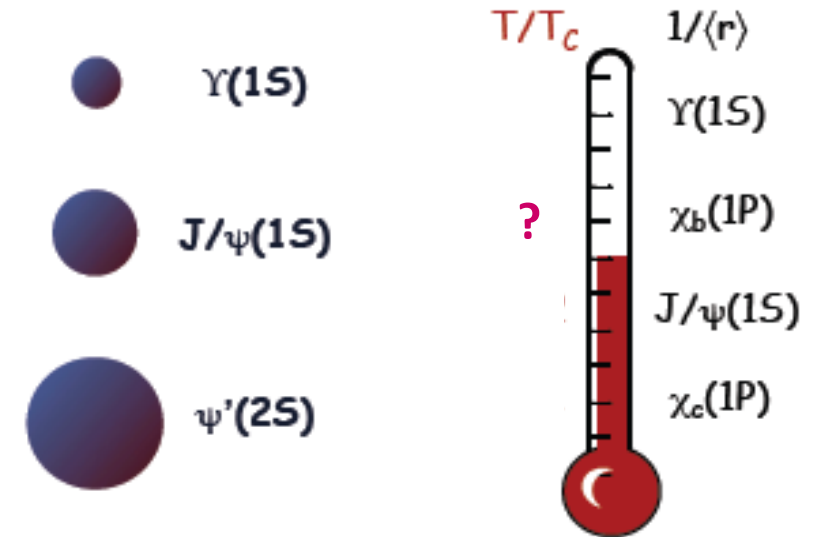
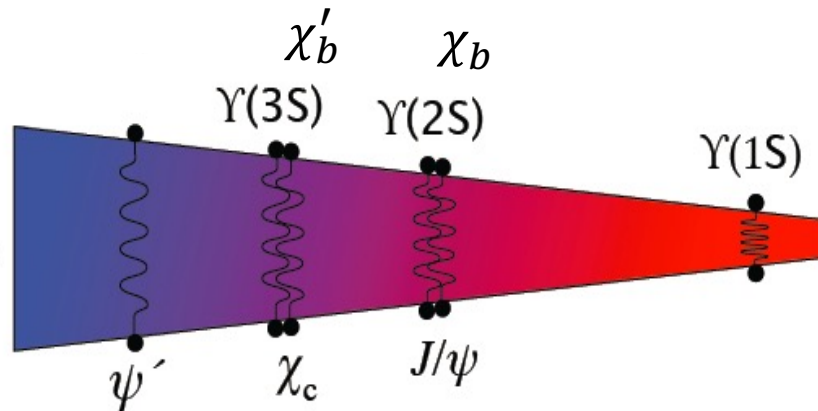


2.5.1 Study of the charmonium ground state: **evidence** for the (re)generation and demonstration of **deconfinement** at LHC energies

\*Jet quenching might play an import role at high  $p_T$

# Quarkonium Suppression: QGP Thermometer

## Plasma thermometer



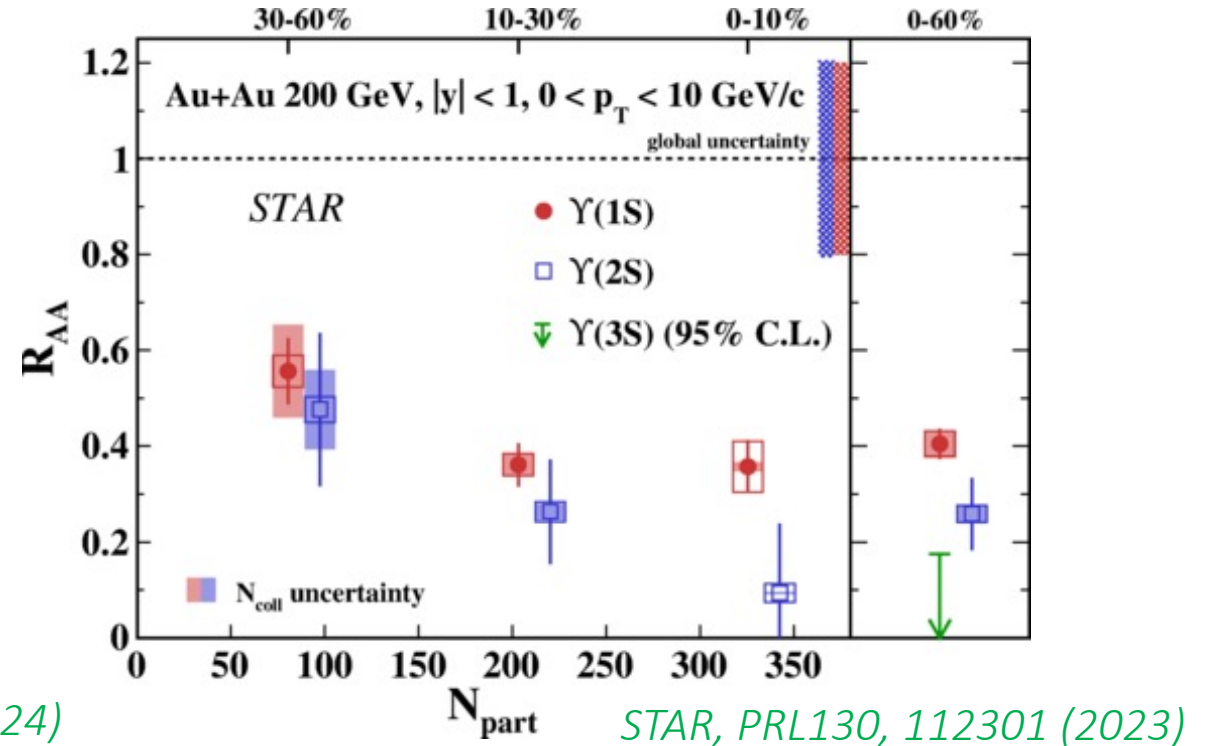
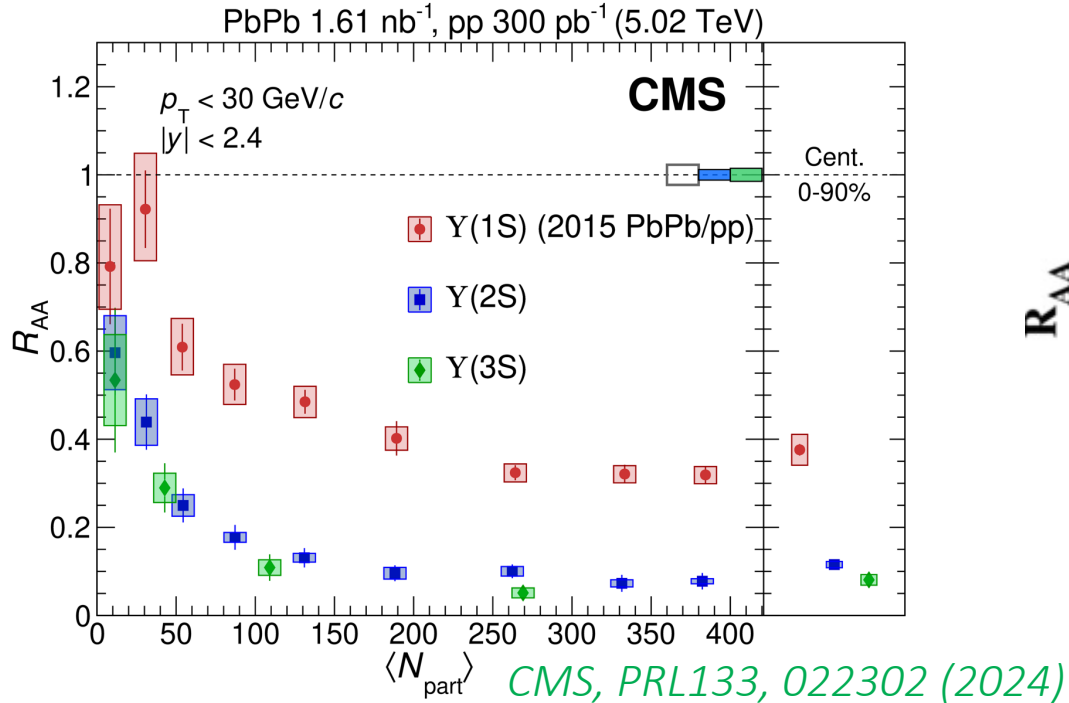
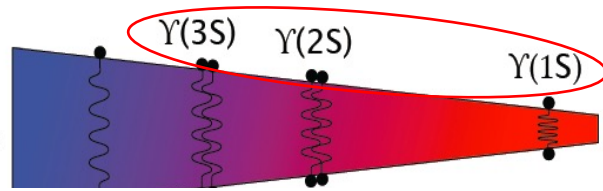
Debye radius is inversely proportional to the temperature of QGP

Different quarkonium states dissociate at different temperature

→ Sequential melting

By measuring sequential melting, one get some information of QGP temperature

# Sequential Melting in Bottom Sector

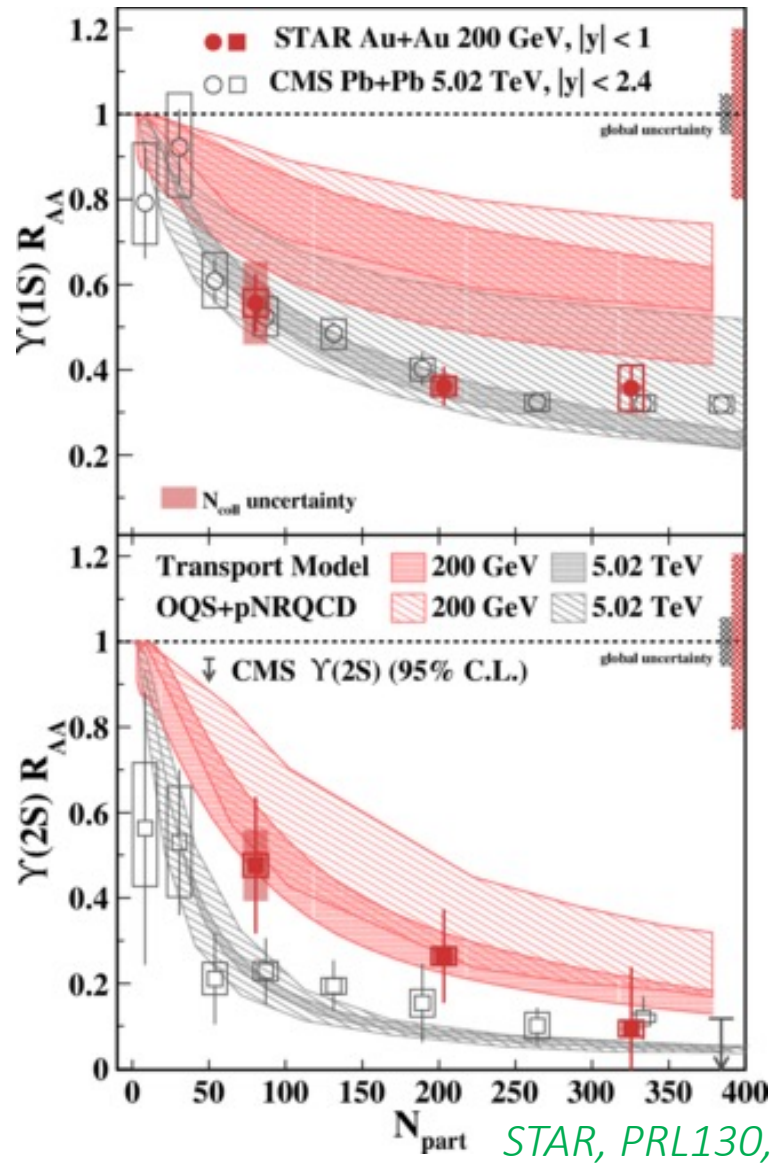


Precise measurement of “sequential melting” at LHC

First observation of “sequential melting” at RHIC



# Sequential Melting in Bottom Sector



## Upsilon(1S):

- Strong suppression, and similar at RHIC and LHC
- Arises mainly from the suppression of excited states feed down to Upsilon(1S) and CNM effects
- Primordial Upsilon(1S) not significantly suppressed

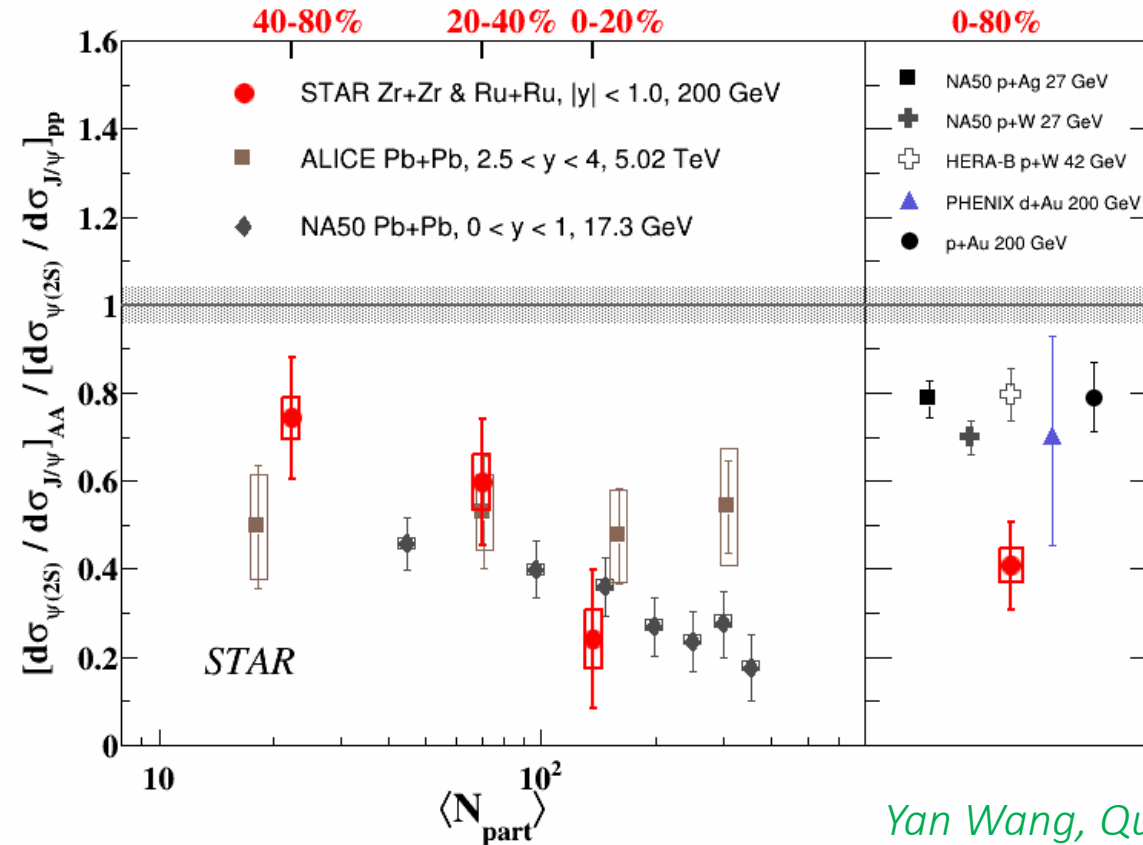
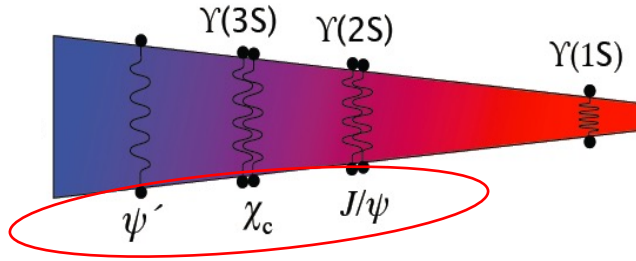
## Upsilon(2S):

- Hints of less suppression at RHIC in peripheral collisions

QGP is formed, and its temperature is high enough to melt excited bottomonium states!!!

STAR, PRL130, 112301 (2023)

# Sequential Melting in Charm Sector



- Clearly stronger suppression for  $\psi(2S)$  than  $J/\psi$  from SPS to LHC
- “Sequential melting” in charm sector
- More data is needed to investigate collision energy dependence



# Summary of Quarkonium

Quarkonium is an unique probe of QGP due to its large binding energy

Systematic experimental investigations provide **strong evidence** of

- Color screening in the deconfined medium
- And sequential melting (binding-energy-dependent suppression)
- Regeneration (coalescence) in the deconfined medium

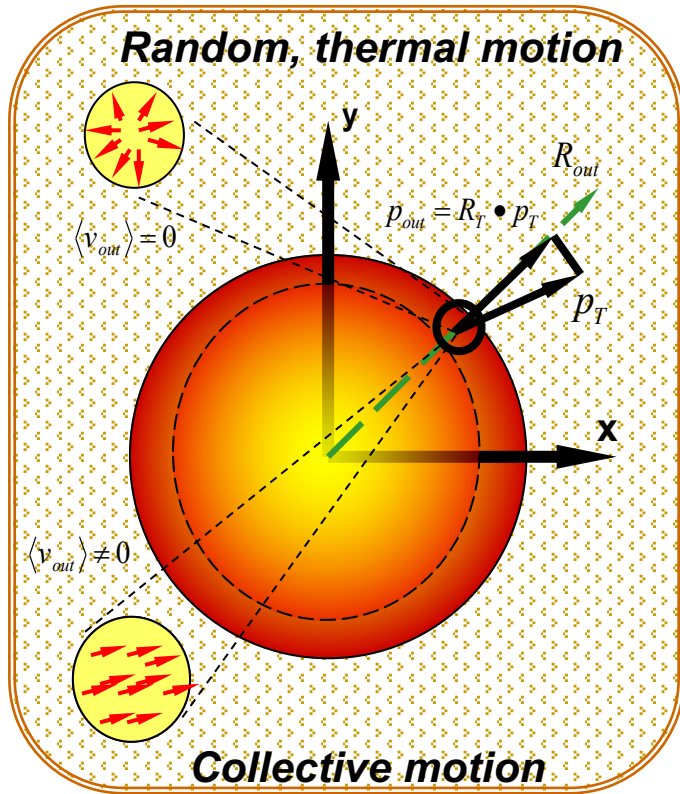
**→ All consistent with the formation of QGP**

**→ Can be (Have been) used to extract QGP properties with precise data and theory**

Deep understanding the inner-working of QGP and quarkonium production mechanism in QGP requires further investigations

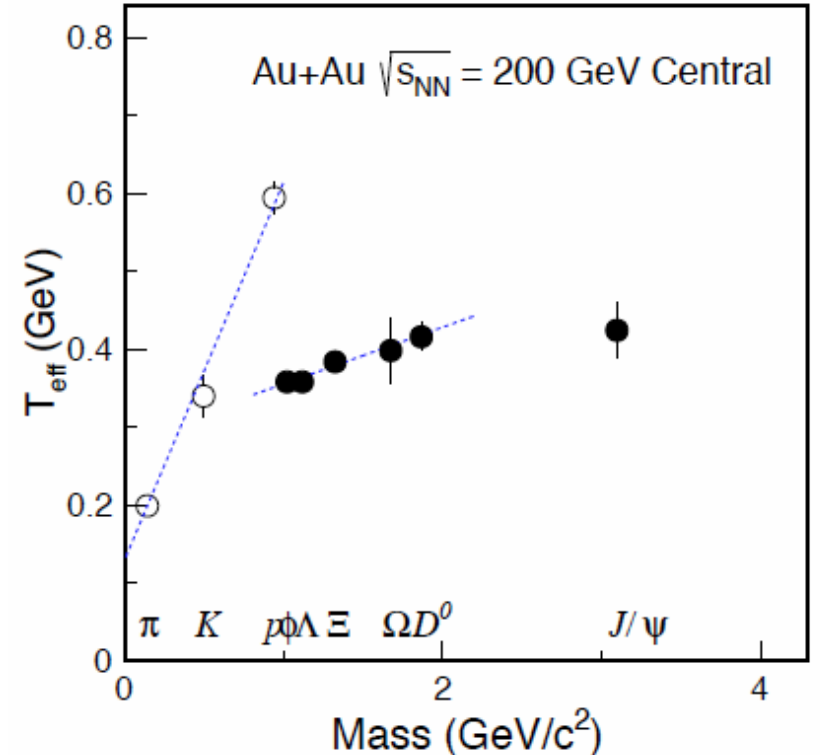
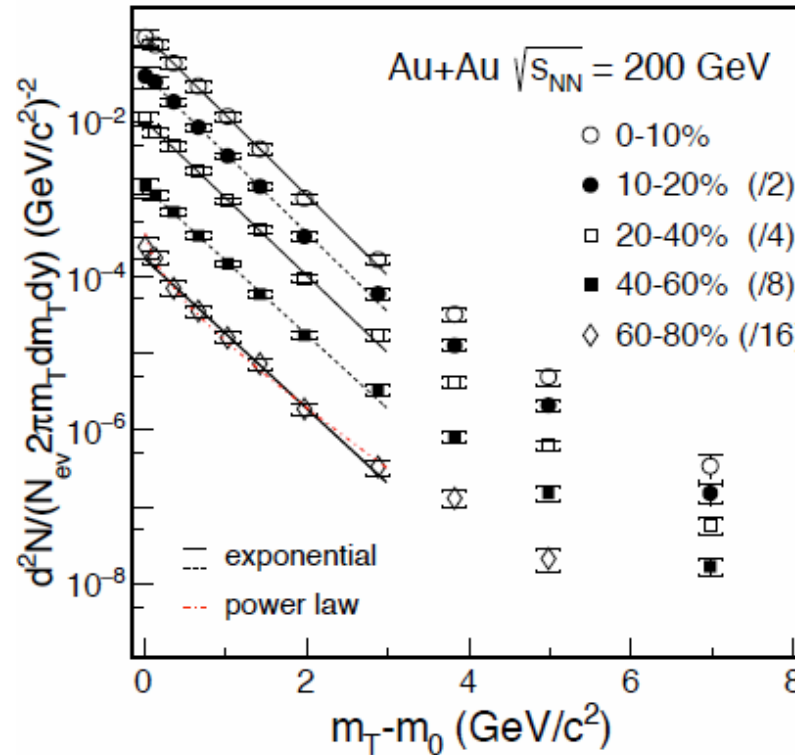


# Radial Flow of D-meson



Courtesy of Nu Xu

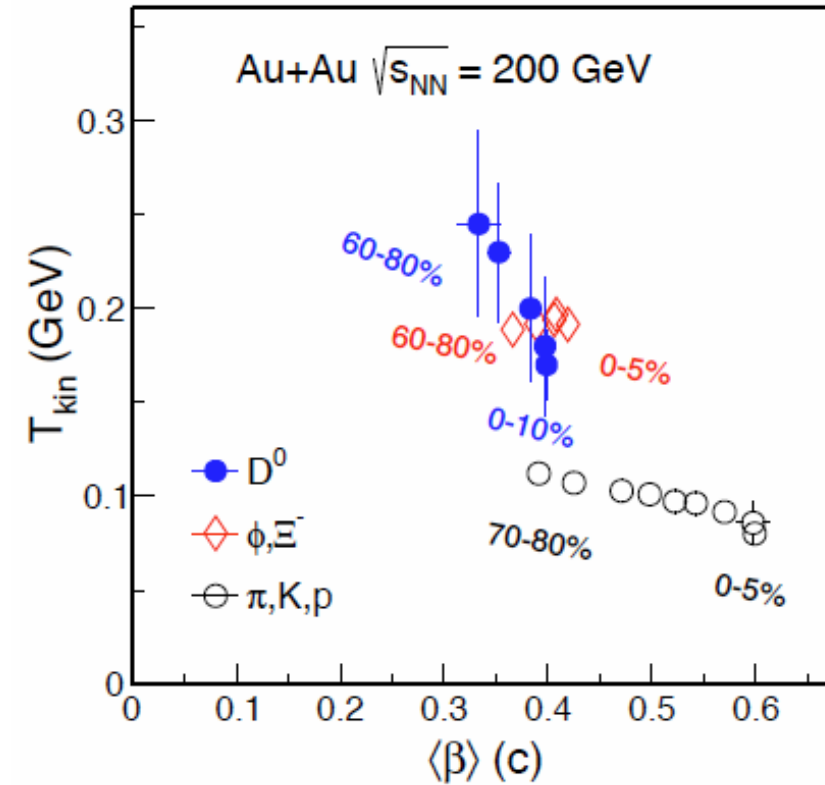
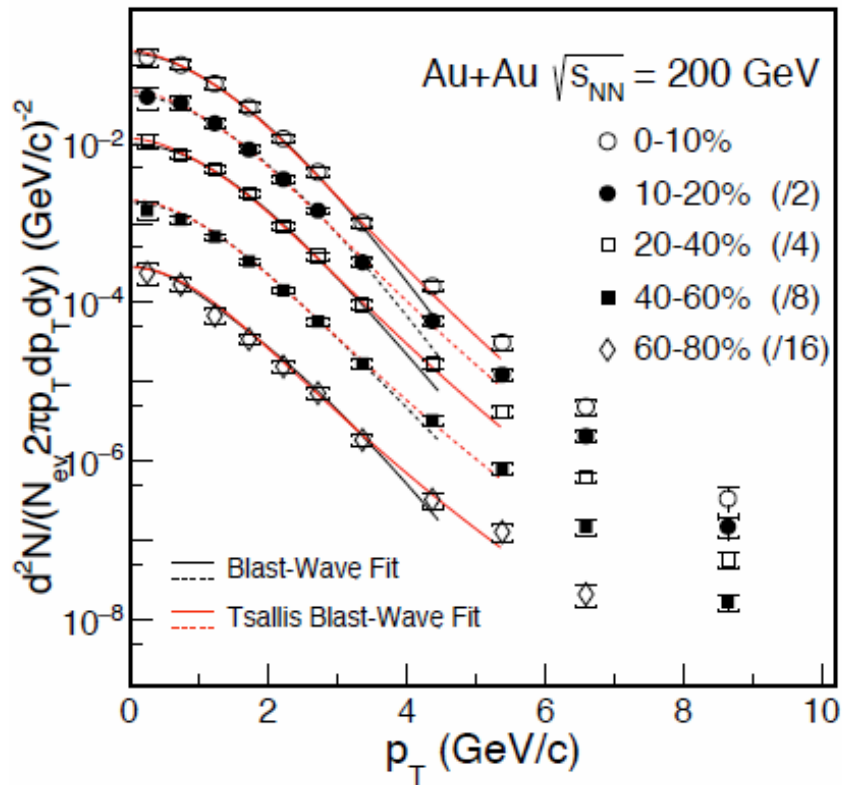
Charm should have blue shift if it has enough interactions with the **expanding** medium



$$T_{eff} \propto T_{thermal} + m \times \langle v_T \rangle^2$$

- $D^0$  follow the trend of strange particles
- Non-zero slop, but smaller than light flavor

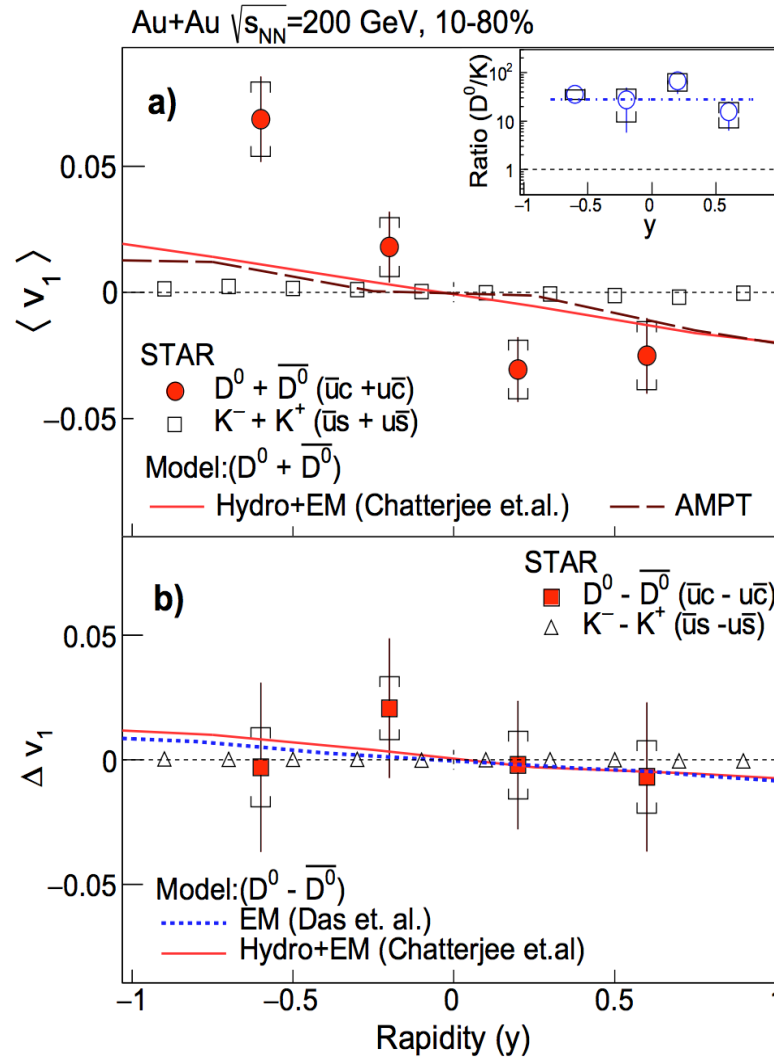
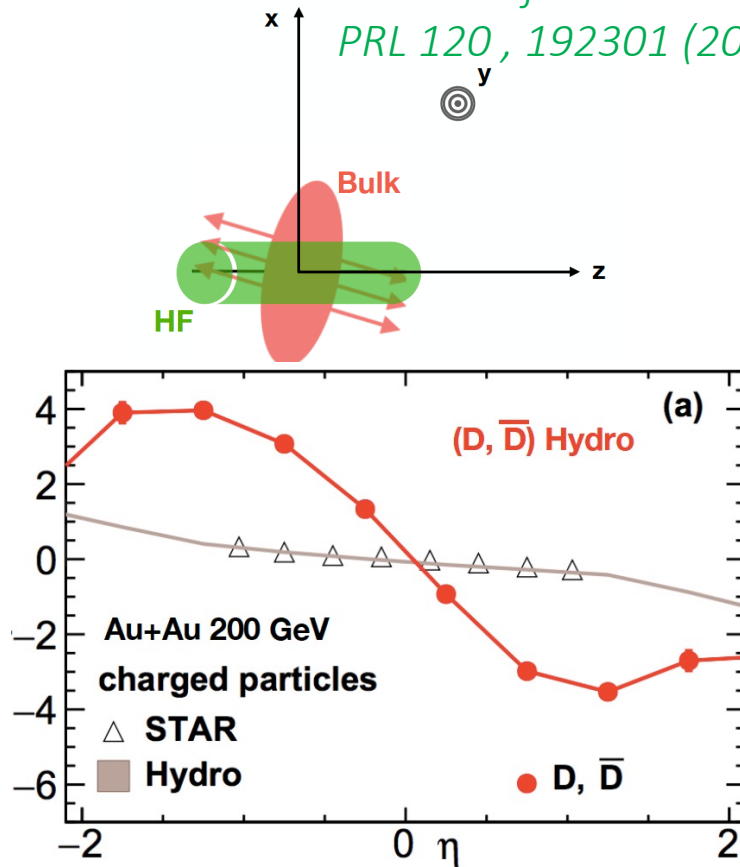
# Radial Flow of D-meson



- Kinetic freezeout parameters can be extracted with (Tsallis) Blast-Wave model
- $D^0$  has similar parameters as strange particles

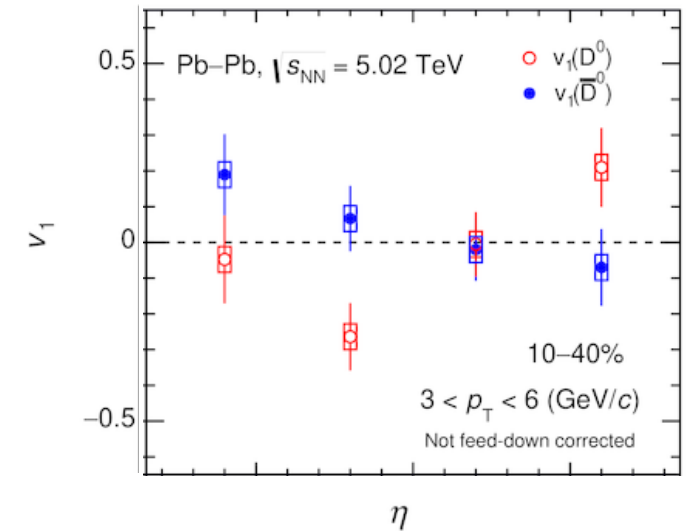
# Directed Flow of $D^0$

*S. Chatterjee and P. Božek,  
PRL 120, 192301 (2018)*



*STAR: PRL123, 162301 (2019).*

*ALICE, PRL125, 022301 (2020)*

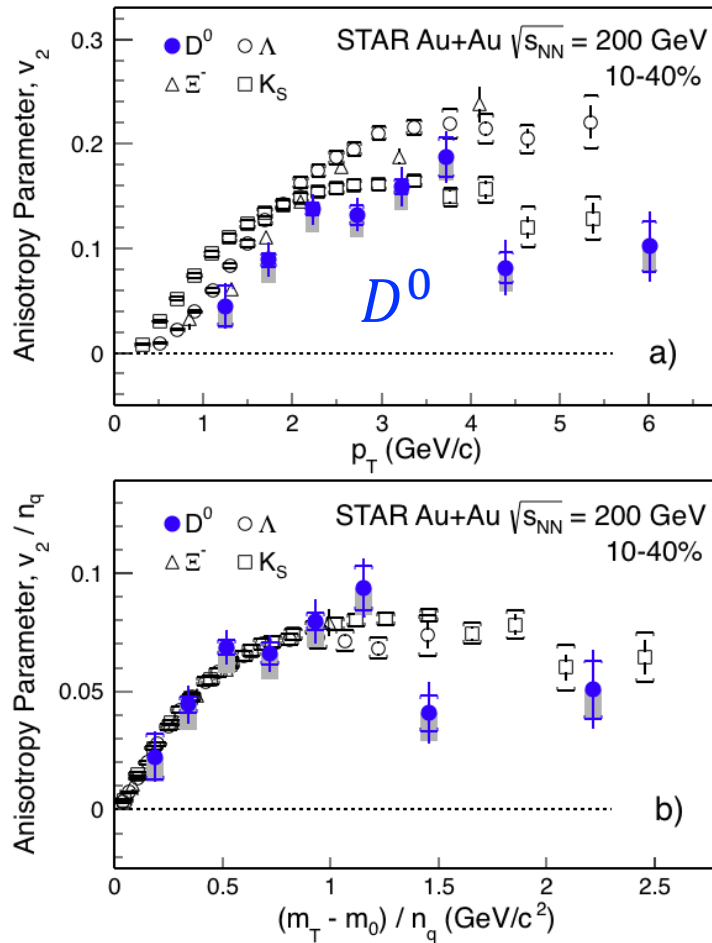


- Huge  $v_1$  slop of D-meson is observed
- Need much more statistics to prob early electromagnetic field

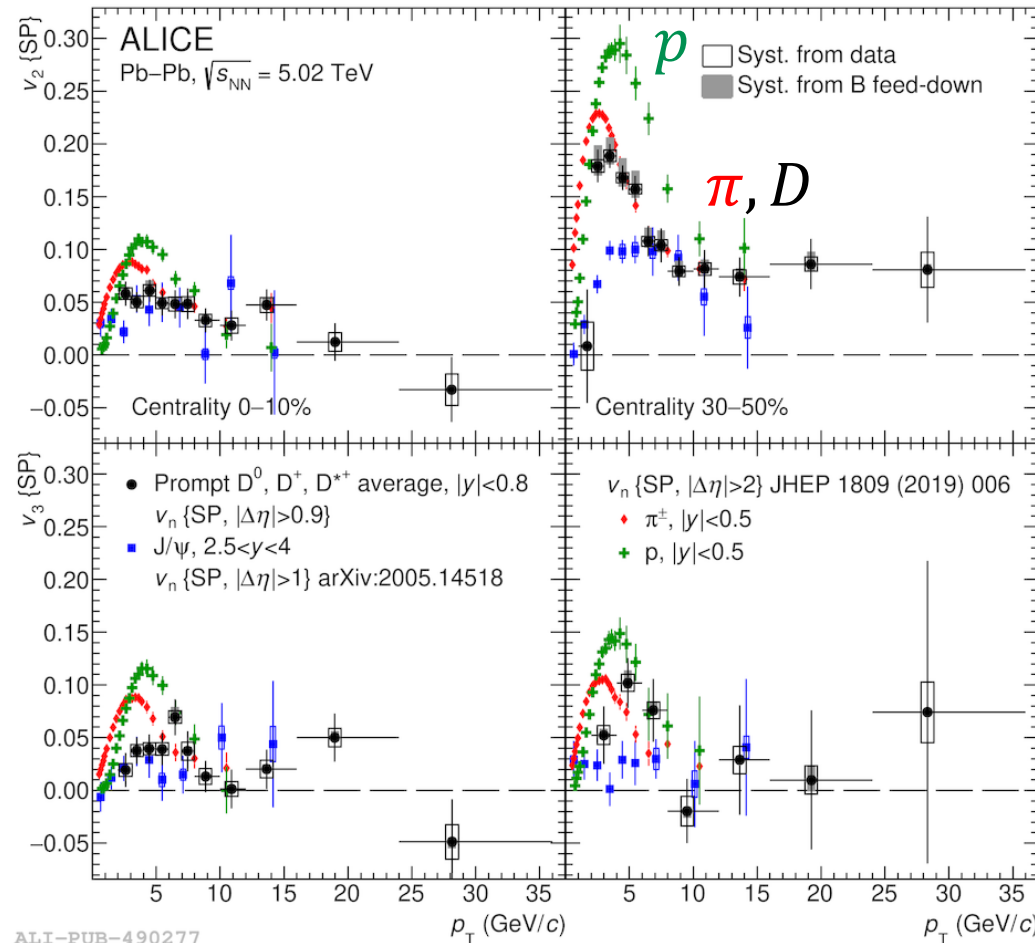
Charm should gain significant directed flow if it has enough interactions with the **tilted** medium



# Elliptic/Triangular Flow of D-meson



STAR, PRL  
118, 212301  
(2017)



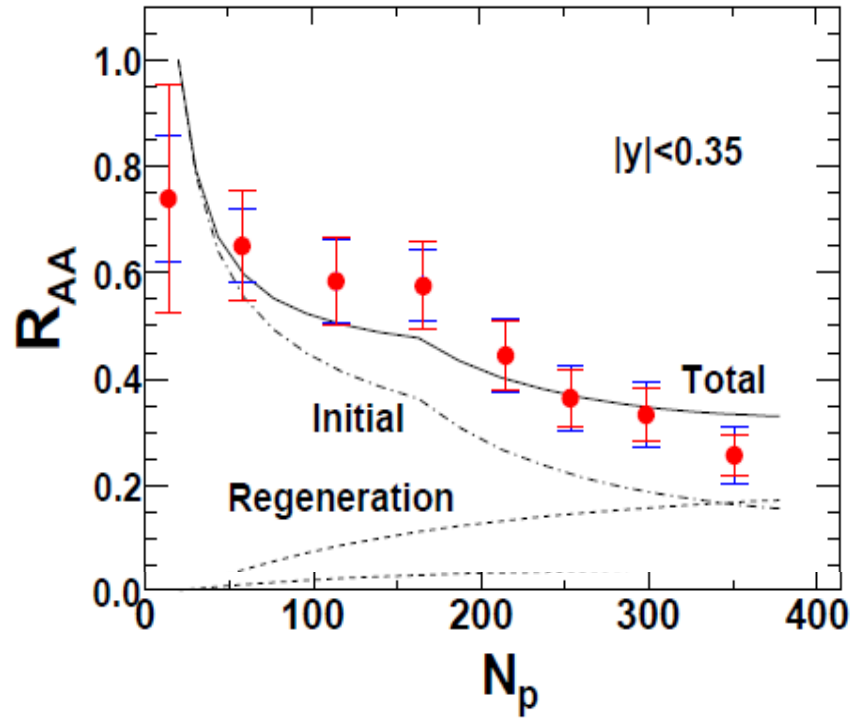
**QUESTION:**  
Why large  
flow at high  
 $p_T$ ?

ALICE, PLB813,  
136054 (2021)

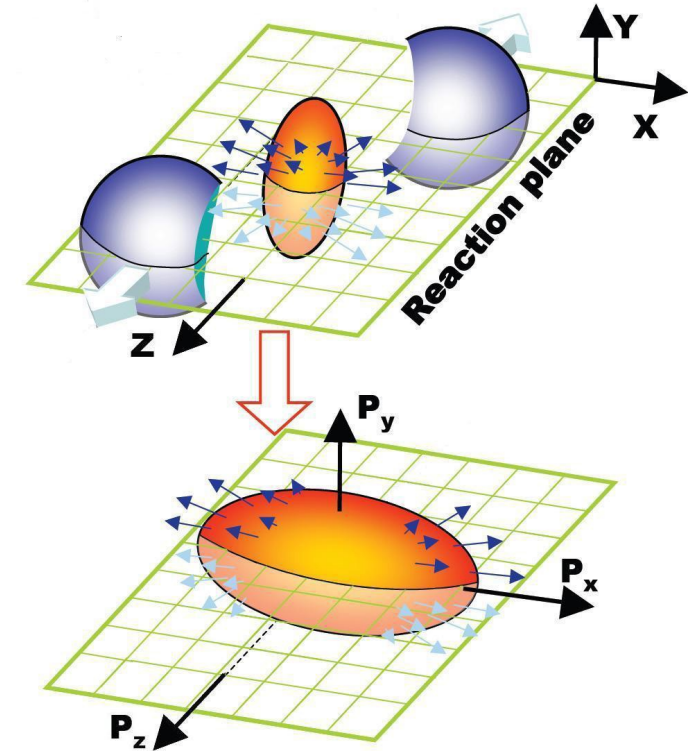
Significant  $v_2$  and  $v_3$  observed for  $D^0$ -meson

- Mass ordering at low  $p_T$  – **hydrodynamics** behavior
- Follow mesons' flow at intermediated  $p_T$  – quark **coalescence**

# Test $J/\psi$ Regeneration with Elliptic Flow



Theory predicts comparable contribution of color screening and regeneration in central collisions at RHIC



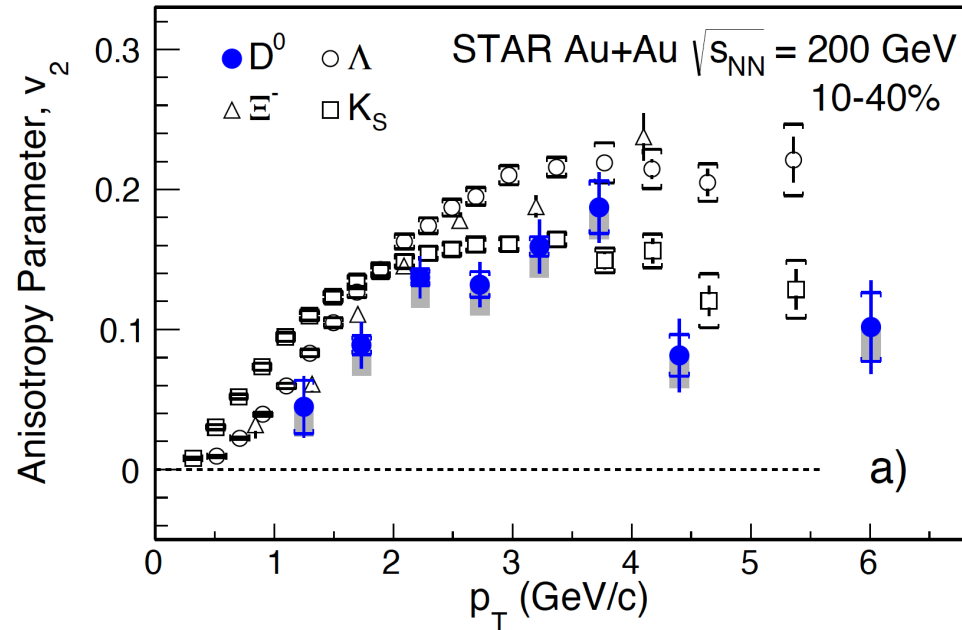
May study via azimuthal anisotropy (elliptic flow)

If charm quark have elliptic flow, regenerated  $J/\psi$  should inherit it

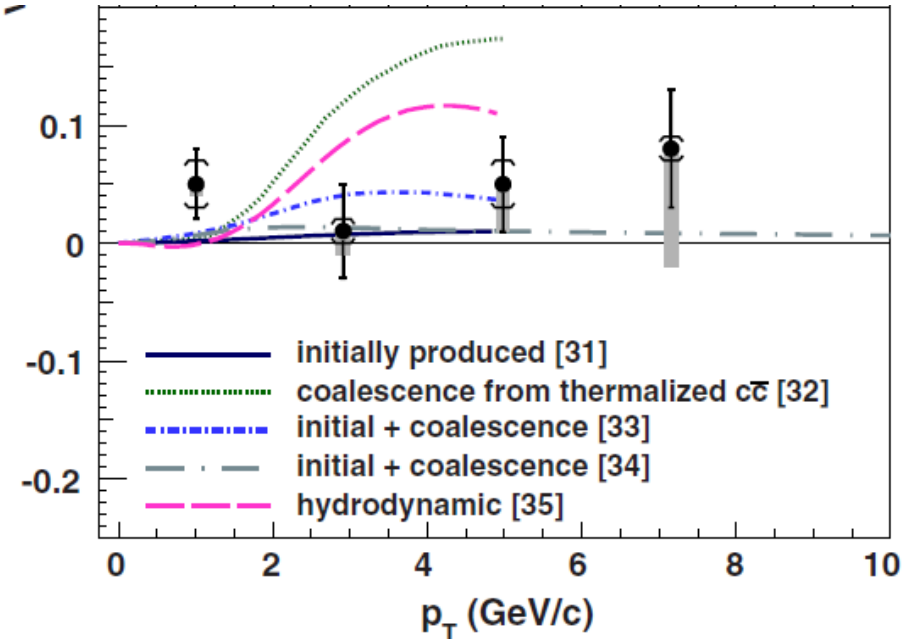


# Elliptic Flow J/ $\psi$ Mesons at RHIC

STAR, PRL118, 212301 (2017)



STAR, PRL 111, 052301 (2013)



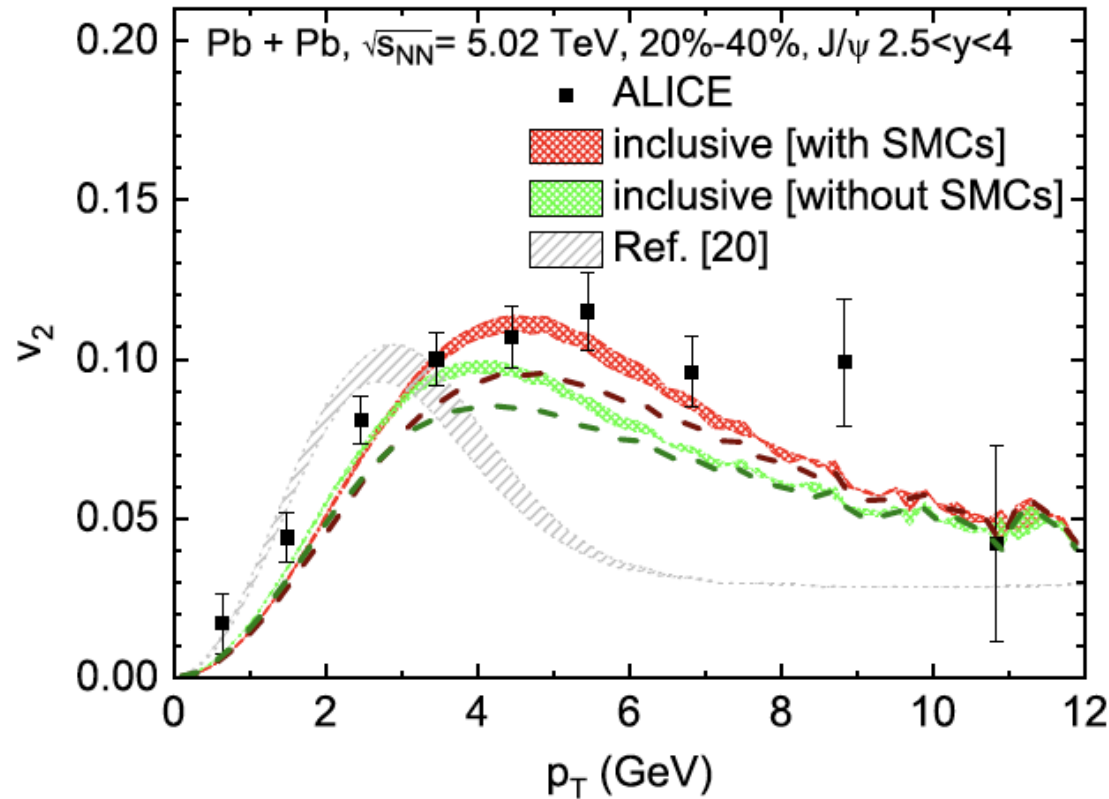
Significant elliptic flow for D mesons  $\rightarrow$  thermalization of charm quark

At  $p_T > 2$  GeV/c, J/ $\psi$  elliptic flow consistent with 0

$\rightarrow$  **Disfavor** the case of **dominantly** produced by **thermalized** charm quarks **coalescence**



# J/ $\psi$ Elliptic Flow at LHC



*ALICE, JHEP10, 141 (2021)*

*Transport model:*

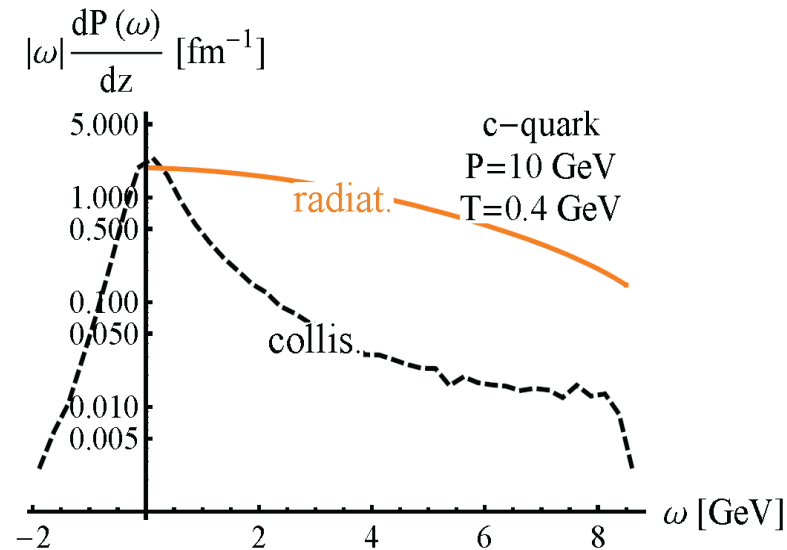
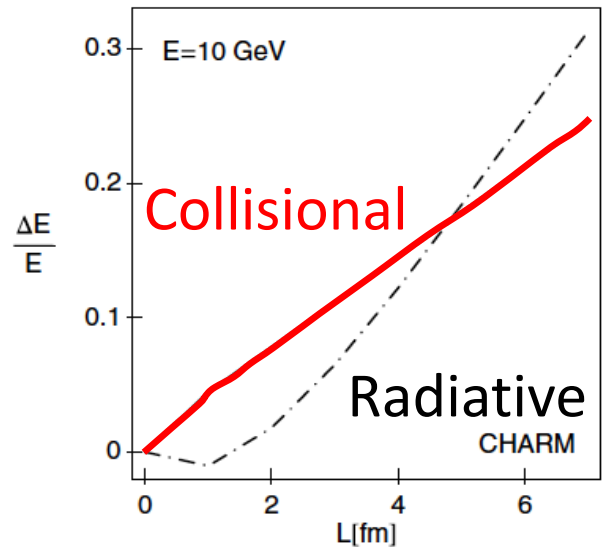
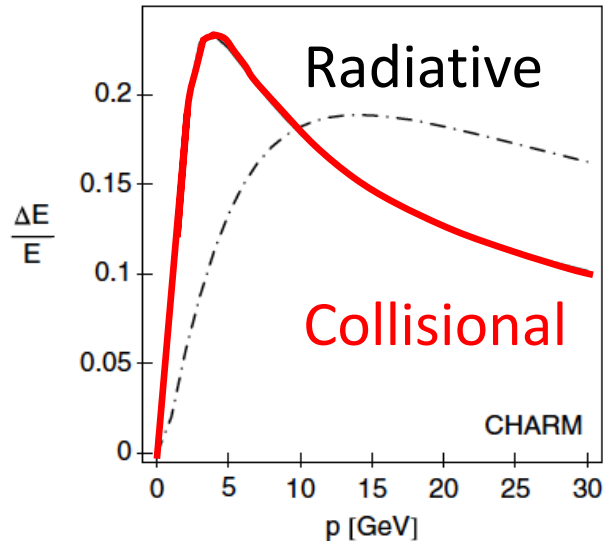
*X. Du and R. Rapp, NPA943, 147 (2015)*

*M. He, B. Wu and R. Rapp, PRL128, 162301 (2022)*

- Significant J/ $\psi$   $v_2$  observed at forward-rapidity (via dimuon trigger)
- Primordial only can not explain the large  $v_2 \rightarrow$  (Re)generation
- The role of jet fragmentation at high  $p_T$ ?



# Elastic vs. Inelastic Collisions in QGP



## Elastic collision:

- Dominant at low- $p_T$
- Proportional to  $L$
- Not always energy “loss”
- Responsible for heavy quark collectivity

## Inelastic collision:

- More important at high- $p_T$
- Proportional to  $L^2$
- Main contributor of energy loss (jet quenching)

Both have much less effect on bottom than charm

*M. Djordjevic, PRC74, 064907 (2006)*

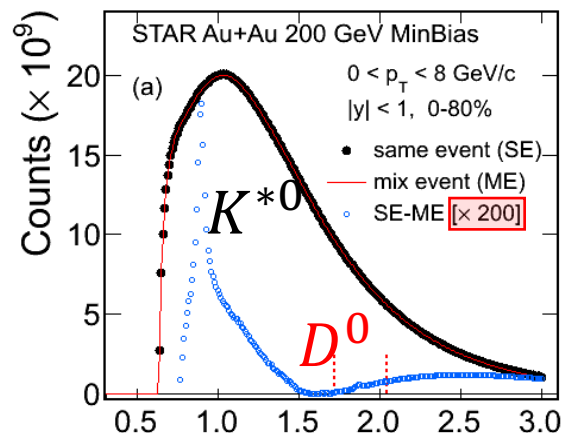


# D<sup>0</sup> Signal with Topological Cuts

2010+2011

820M events

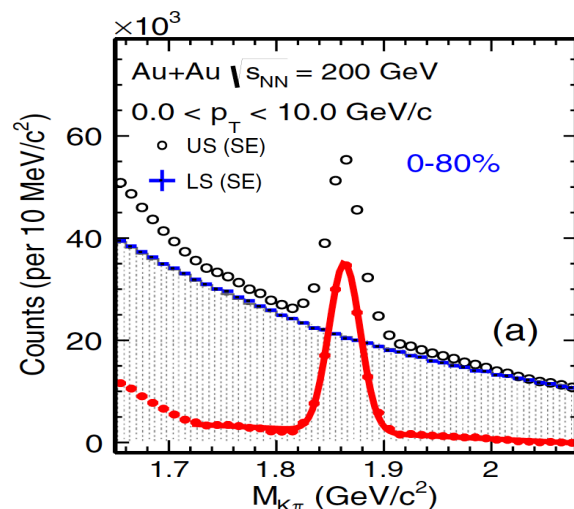
w/o HFT



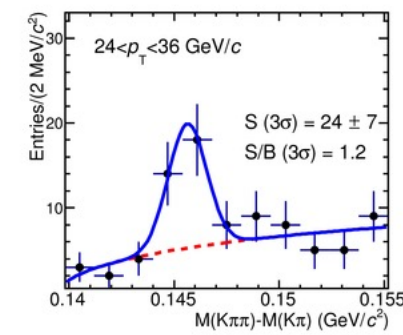
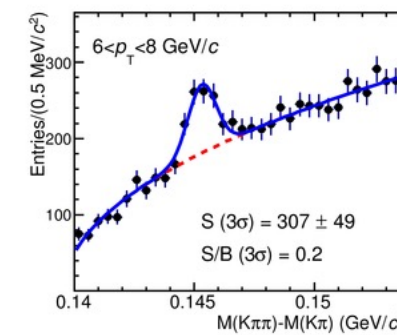
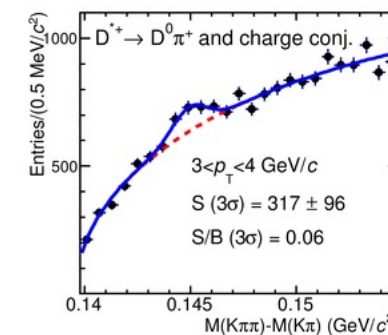
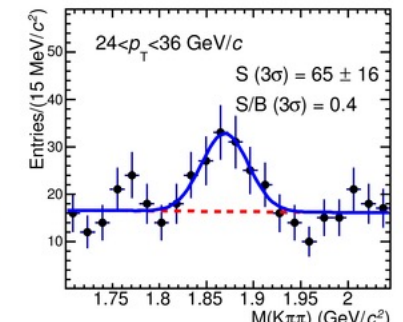
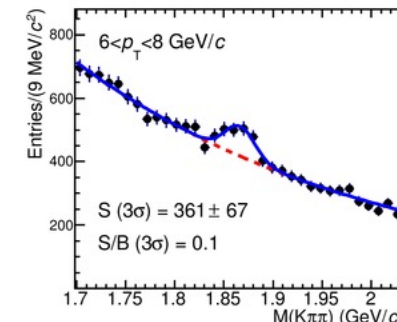
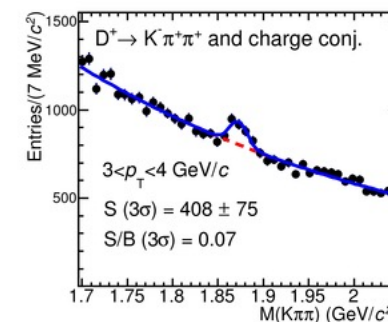
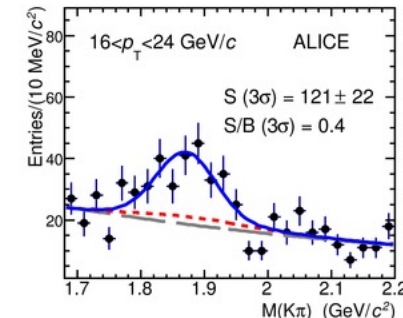
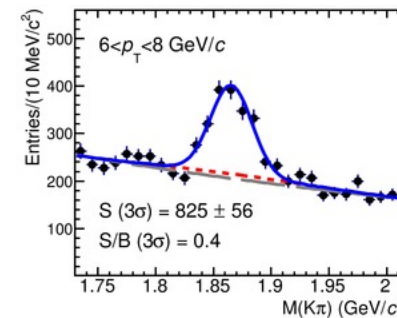
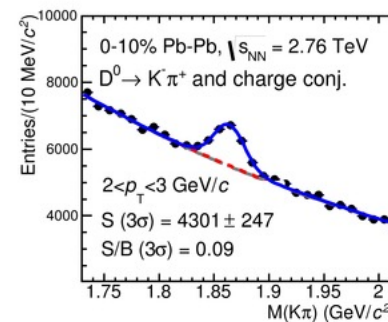
2014

900M events

w/ HFT



S/B improved by O(10<sup>4</sup>)



ALICE, JHEP03, 081 (2016)



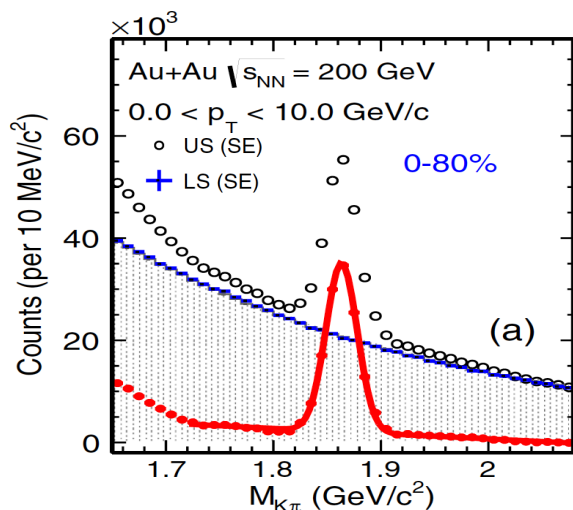
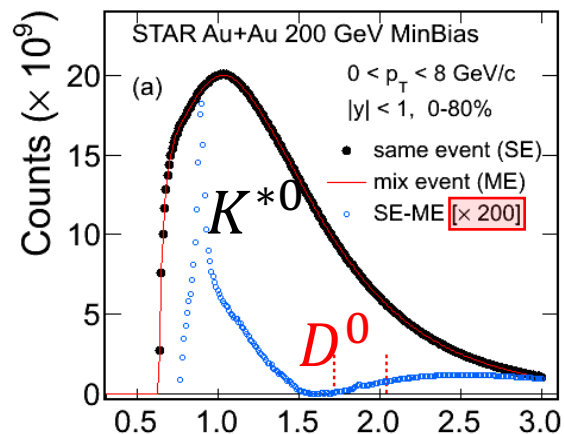


# D<sup>0</sup> Signal with Topological Cuts

2010+2011

820M events

w/o HFT

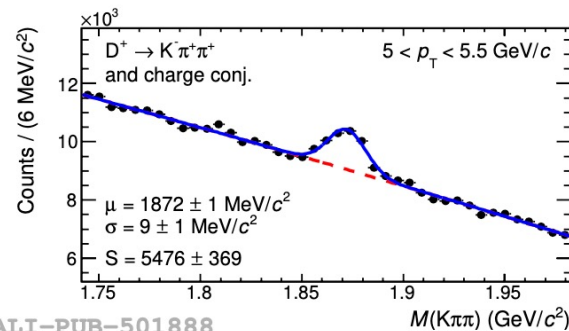
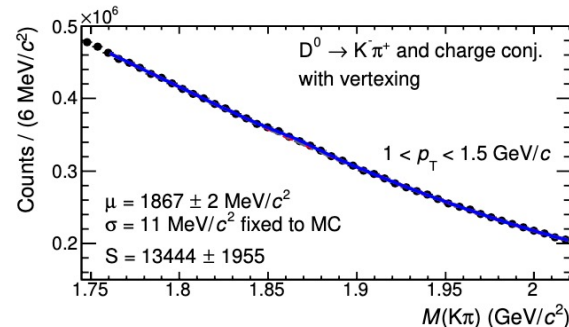
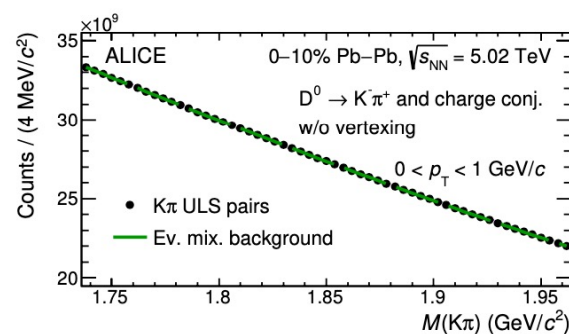


S/B improved by  $O(10^4)$

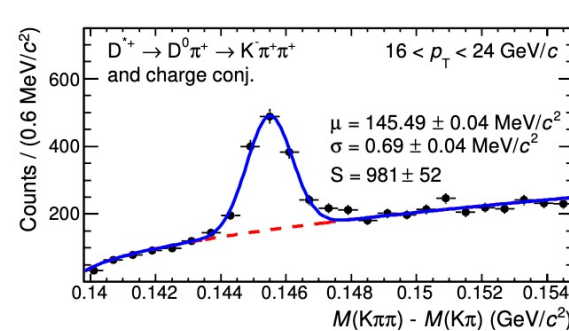
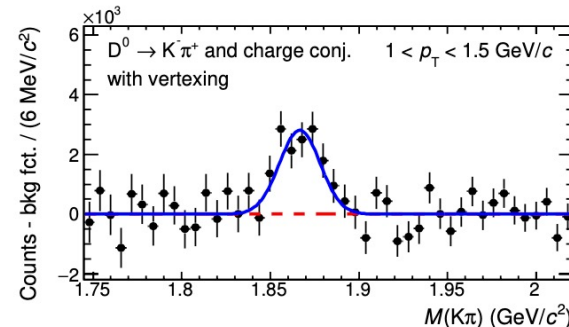
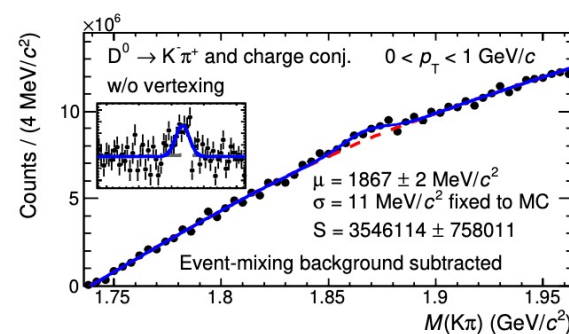
2014

900M events

w/ HFT



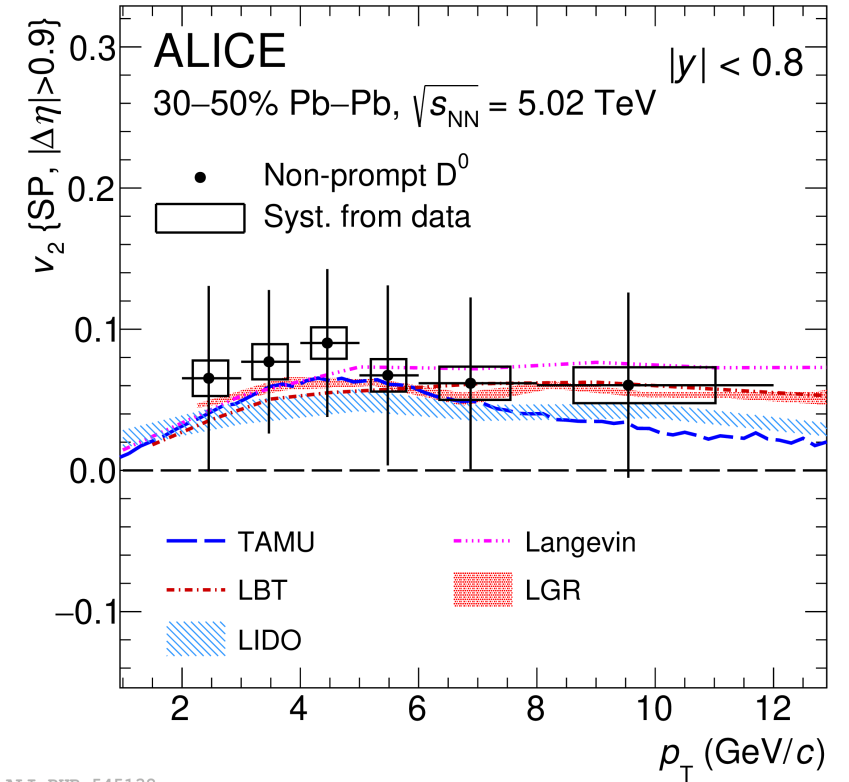
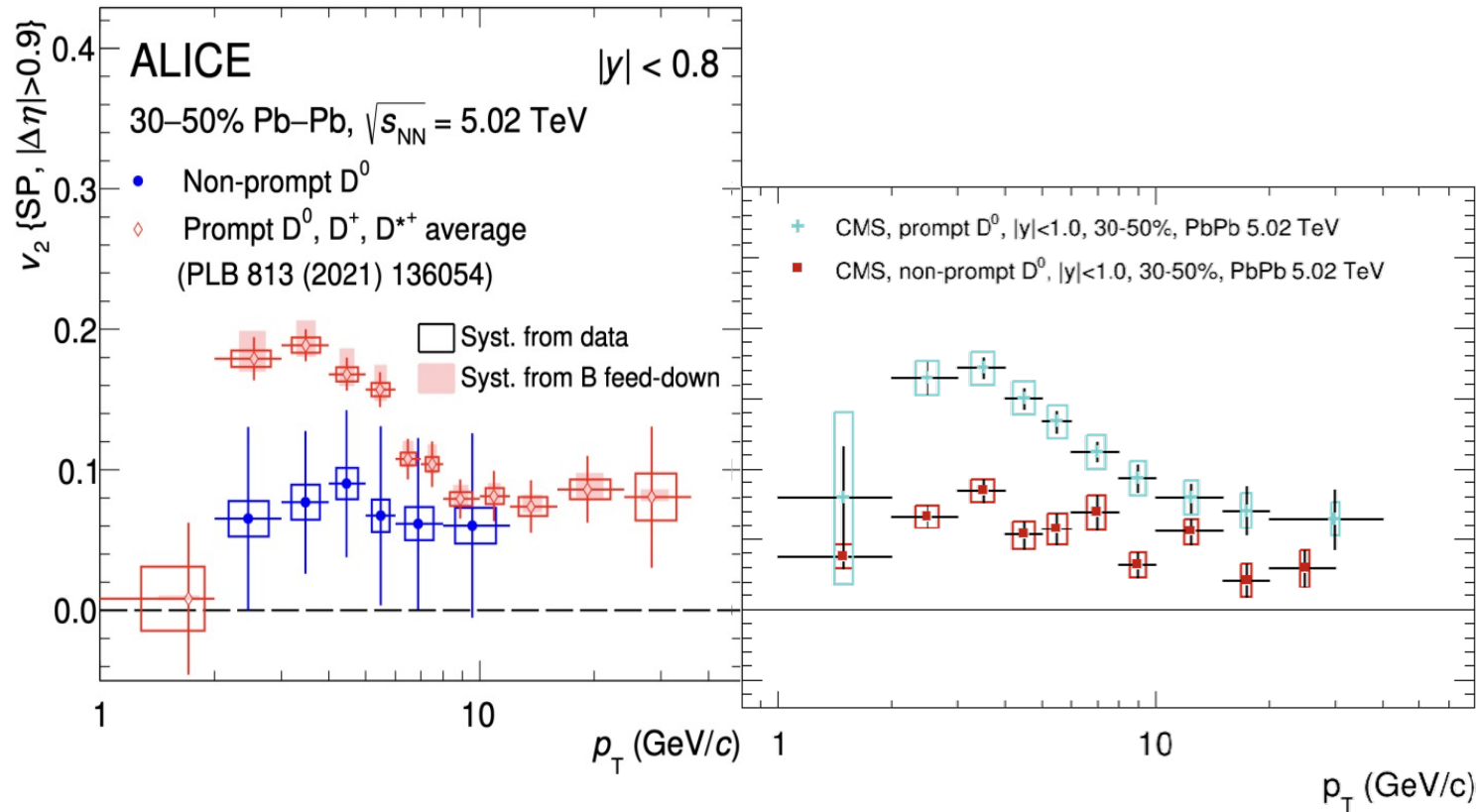
ALI-PUB-501888



ALICE, JHEP01, 174 (2022)



# Collectivity of Bottom



- Significant flow of D from B decay
- Much smaller than prompt D
- Different degree of thermalization of charm and bottom