

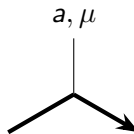
# Vogel universality and diagrammatic technique

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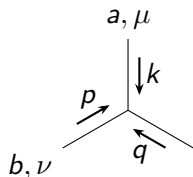
Problems of Modern Mathematical Physics, February 2025

# Yang-Mills diagrammatic technique



A diagrammatic representation of a vertex. A vertical line enters from the top, labeled  $a, \mu$ . Two lines exit from the bottom-left and bottom-right, representing outgoing particles.

$$= ig\gamma^\mu T^a$$



A diagrammatic representation of a three-gluon vertex. A vertical line enters from the top, labeled  $a, \mu$  and has a downward arrow labeled  $k$ . Two lines exit from the bottom-left and bottom-right, labeled  $b, \nu$  and  $c, \rho$  respectively. The bottom-left line has an upward arrow labeled  $p$ , and the bottom-right line has a downward arrow labeled  $q$ .

$$= g f^{abc} [g^{\mu\nu}(k-p)^\rho + g^{\nu\rho}(p-q)^\mu + g^{\rho\mu}(q-k)^\nu]$$

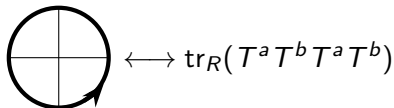
# Chern-Simons diagrammatic technique

In Chern-Simons 3D topological theory observables are the Wilson loops:

$$\langle W_R(K) \rangle = \langle \text{tr}_R \text{P exp} \left( \oint_K A \right) \rangle$$

Perturbative expansion of Wilson loop vacuum expectation value yields:

$$\langle W_R(K) \rangle = \sum_m h^m \oint dx_1 \dots dx_m \langle A^{a_1}(x_1) \dots A^{a_m}(x_m) \rangle \text{tr}_R (T^{a_1} \dots T^{a_m})$$


$$\text{Diagram} \longleftrightarrow \text{tr}_R(T^a T^b T^a T^b)$$

# STU, AS and IHX relations

$$f^{ab}_c T^c = [T^a, T^b] = T^a T^b - T^b T^a$$

$$f^{ad}_e f^{bc}_d + f^{bd}_e f^{ca}_d + f^{cd}_e f^{ab}_d = 0$$

$$f^{ab}_c = -f^{ba}_c$$

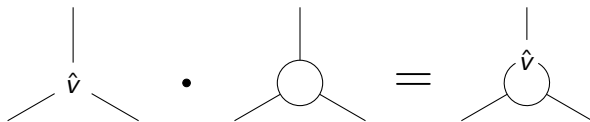
Vogel universality is an idea that some Lie algebra properties can be expressed in a uniform way as a function of the so-called Vogel parameters. Currently there are several features that are known to admit Vogel universalization:

- Decompositions of adjoint representation powers [Vogel, 1999; Isaev, Krivonos, 2024 & others]
- Higher Casimir eigenvalues [Mkrtchyan, Sergeev, Veselov, 2012]
- Quantum dimension [Mkrtchyan, Veselov, 2012]
- Quantum knot invariants [Mironov, Mkrtchyan, Morozov, 2016]

## Definition

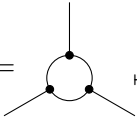
$\Lambda$  is an algebra over  $\mathbb{Q}$  generated by 3-legged diagrams modulo AS and IHX relations antisymmetric with respect to permutations of legs.

Multiplication in  $\Lambda$  is given by insertion of one diagram into any vertex of the other diagram.



# Characters on $\Lambda$

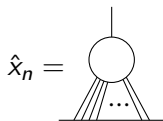
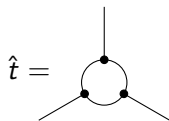
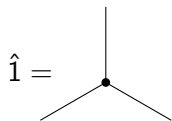
Trivalent diagrams correspond to contractions of structure constants. In case of  $\Lambda$  algebra three free indices remain, therefore the result is proportional to the structure constant.

$$\hat{t} = \text{trivalent diagram} \mapsto f^{aik} f^{bjl} f^{ckj} = \chi_L(\hat{t}) f^{abc}$$


The coefficient of proportionality is called a *character* of this element of  $\Lambda$ . Each Lie algebra induces its own character on  $\Lambda$ . One can say that Lie algebras can be parameterized by these characters.

# Structure of $\Lambda$

In  $\Lambda$  there exist the following diagrams that happen to be multiplicative generators of the studied sector of  $\Lambda$ :



However, there are many relations between the  $x_n$  diagrams, that allow to rewrite them in terms of polynomials in three parameters  $t$ ,  $\sigma$  and  $\omega$  of the form  $\mathbb{Q}[t] \oplus \omega\mathbb{Q}[t, \sigma, \omega]$

The most interesting part about the  $\Lambda$  algebra is an existence of zero divisors. They effectively restrict possible values of the Vogel parameters to those that satisfy:

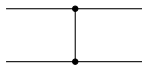
$$t\omega(2t\sigma - \omega - 2t^3)(27\omega - 45t\sigma + 40t^3) \times \\ \times (27\omega^2 - 72t\sigma\omega + 40t^3\omega + 4\sigma^3 + 29t^2\sigma^2 - 24t^4\sigma) = 0$$



# Adjoint square decomposition

On the tensor square of adjoint representation there exists an operator  $\Psi$ ,

also known as the split Casimir operator:



It was shown by Vogel that this diagram satisfies a cubic relation on the symmetric square, factorised by a one-dimensional representation:

$$\Psi^3 - t\Psi^2 + (\sigma - 2t^2)\Psi - (\omega - t\sigma)\hat{1} = 0$$

This implies the following decomposition:

$$S^2\text{adj} = \Omega \oplus Y_\alpha \oplus Y_\beta \oplus Y_\gamma$$

Where  $\alpha$ ,  $\beta$  and  $\gamma$  are the eigenvalues of  $\Psi$ . They are related to other Vogel parameters as follows:

$$\alpha + \beta + \gamma = t, \quad \alpha\beta + \beta\gamma + \alpha\gamma = \sigma - 2t^2, \quad \alpha\beta\gamma = \omega - t\sigma.$$

# Universal dimension

It Vogel's paper there is following diagrammatic equation:

$$\sigma \left( \begin{array}{c} \diagup \quad \diagdown \\ | \quad | \\ \diagdown \quad \diagup \end{array} - t \left( \begin{array}{c} \text{---} \quad \diagdown \\ \text{---} \quad \diagup \end{array} \right) \right) = \left( \begin{array}{c} \diagup \quad \diagdown \\ | \quad | \\ \bullet \\ 0 \\ \diagdown \quad \diagup \end{array} - \omega \left( \begin{array}{c} \text{---} \quad \diagdown \\ \text{---} \quad \diagup \end{array} \right) \right)$$

Taking trace over it yields:

$$2t\sigma \dim - t\sigma(\dim - 1) = -\omega(\dim - 1)$$

Hence, the dimension of Lie algebra can be expressed as:

$$\dim = \frac{\omega - 3t\sigma}{\omega - t\sigma} = \frac{(\alpha - 2t)(\beta - 2t)(\gamma - 2t)}{\alpha\beta\gamma}$$

# Universal Casimir eigenvalues

Consider the following Casimir operators:

$$C_p = \text{Tr}_{ad}(X^{\mu_1} X^{\mu_2} \dots X^{\mu_p}) X_{\mu_1} X_{\mu_2} \dots X_{\mu_p}$$

The correspond to diagrams of the following form:

$$\hat{u}_p = \text{Diagram 1} = \text{Diagram 2} = \text{Diagram 3} = 2t\hat{x}_{p-1} \text{ ---}$$

$C_p = (-1)^p u_p = (-1)^p \cdot 2t x_{p-1}$ .  $x_n$  is related to  $\alpha$ ,  $\beta$  and  $\gamma$  as follows [Kneissler, 2001]:

$$x_n = -\frac{\alpha\beta\gamma(2t)^n}{(2t-\alpha)(2t-\beta)(2t-\gamma)} + \frac{t^n}{2} + c_\alpha\alpha^n + c_\beta\beta^n + c_\gamma\gamma^n$$

Same result has been achieved by Mkrtchyan, Sergeev and Veselov for classical Lie algebras, later generalized by Isaev and Provorov for Lie superalgebras.

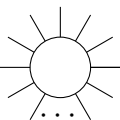
# Universal quantum dimension

The formula for the quantum dimension [Mkrtchyan, Veselov, 2012]:

$$r + \sum_{\mu \in R} e^{x(\mu, \rho)} = \frac{\sinh((\alpha - 2t)x/4) \sinh((\beta - 2t)x/4) \sinh((\gamma - 2t)x/4)}{\sinh(\alpha x/4) \sinh(\beta x/4) \sinh(\gamma x/4)}$$

This formula corresponds to the following diagram expansion, known as the Kontsevich integral of unknot:

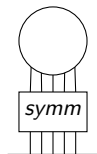
$$\text{qdim} = I(\bigcirc) = \exp\left(\sum_{n=1}^{\infty} b_{2n} w_{2n}\right), \text{ where } b_{2n} = \frac{(-1)^{n+1} \zeta(2n)}{2n(2\pi)^{2n}} \text{ and}$$

$$w_{2n} = \text{Diagram of a circle with } 2n \text{ radial lines extending outwards, representing a } 2n\text{-th order term in the expansion.}$$


# Universal symmetrized Casimir eigenvalues

$$C_n^{\text{symm}} = \text{tr}_{\text{adj}} (T^{i_1} T^{i_2} T^{i_3} \dots T^{i_n}) \frac{1}{n!} \sum_{\sigma \in S_n} T_{i_{\sigma(1)}} T_{i_{\sigma(2)}} T_{i_{\sigma(3)}} \dots T_{i_{\sigma(n)}}$$

Action of these operators corresponds to the following diagram:



It is unknown how to express this diagram through regular Vogel parameters, but in can be done at low orders:

$$C_4^{\text{symm}} = \frac{20}{3} t^4 - 3t\omega \quad C_6^{\text{symm}} = \frac{28}{3} t^6 - \frac{42}{5} t^3 \omega + \frac{15}{8} t\sigma\omega$$