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Naturalness Criteria for the Standard Model

Victor T. Kim

**Petersburg Nuclear Physics Institute
NRC «Kurchatov institute»**

Outline:

- **Brout-Englert-Higgs boson mass evolution and naturalness criteria**
- **Brout-Englert-Higgs boson mass renormalization: logarithmic vs. quadratic**
- **Naturalness domain boundaries for the Standard Model**
- **Summary**



Running couplings: α_{QCD} , α_{EW}

Heavy quark (B-quark) running masses

Different mass parameterizations

(different approaches to include higher orders):

- pole (on-shell) mass
- running mass

SM running masses

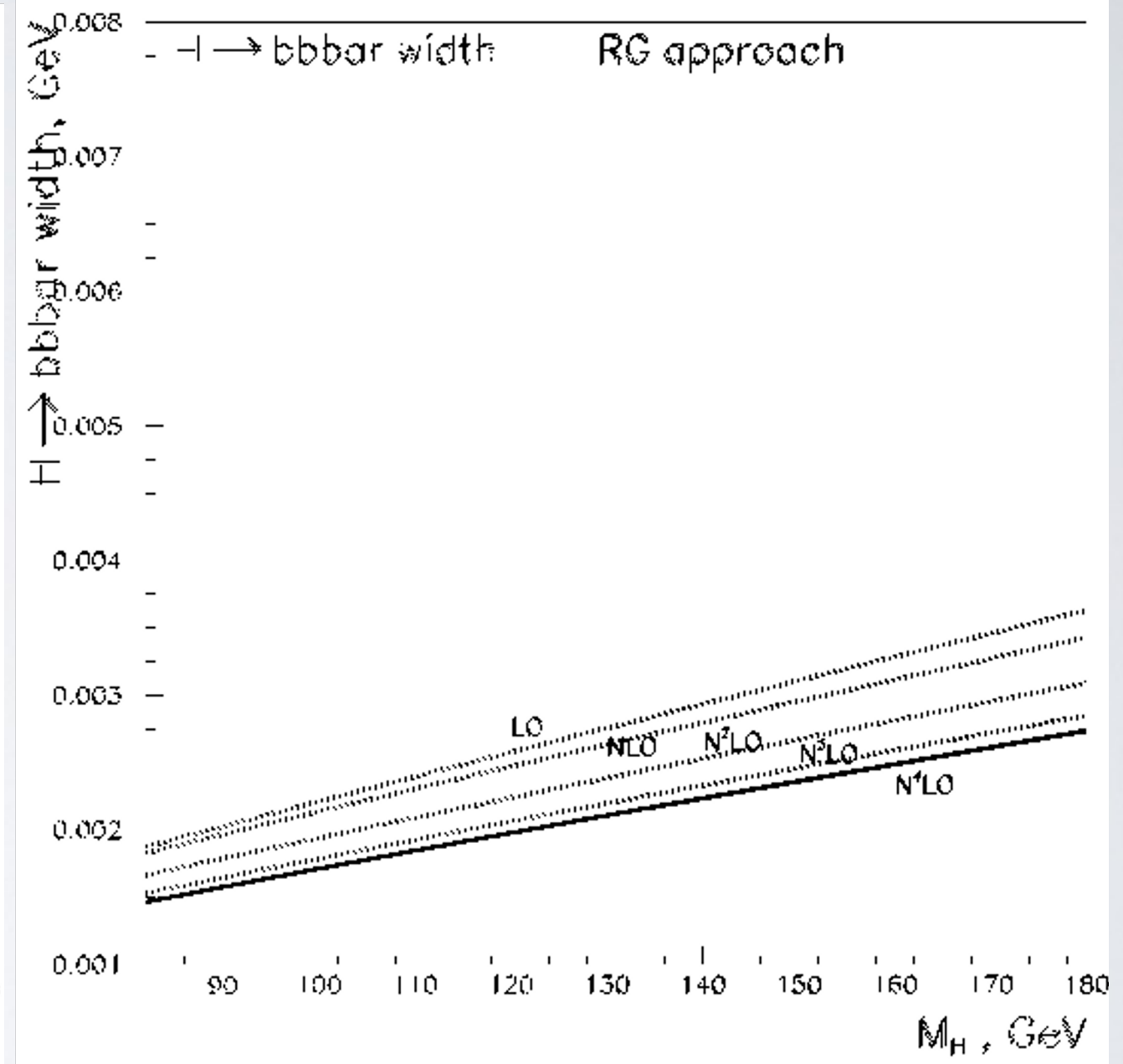
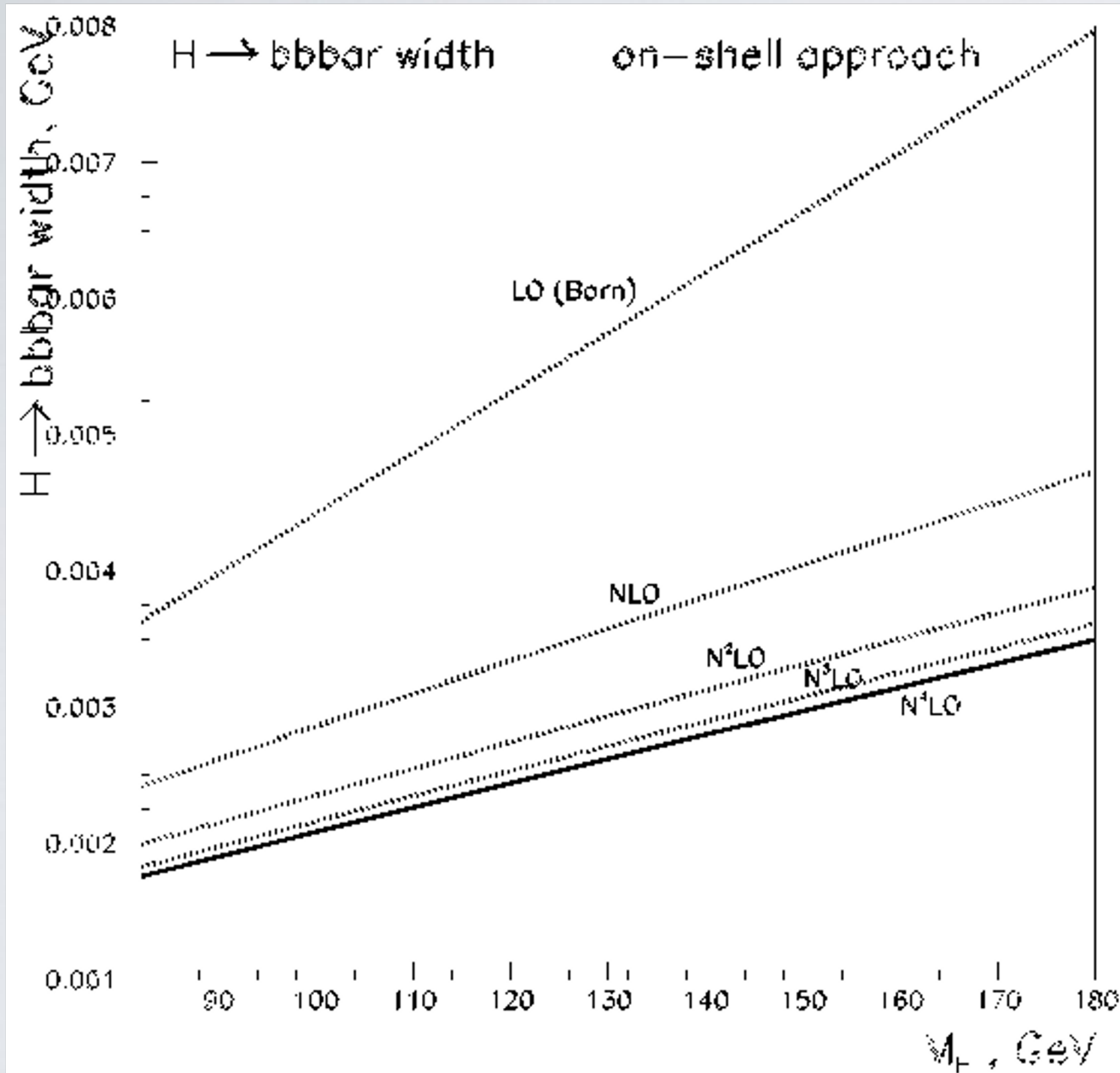
- fermions and vector bosons: logarithmic
- scalar Brout-Englert-Higgs boson: logarithmic
or/and quadratic ?
quadratic -> “non-naturalness”

Brout-Englert-Higgs boson decay width

Width of BEH boson decay into b-quarks (up to N⁴LO)

P. Baikov, K. Chetyrkin, J. Kuhn (2006)

A.P.Bakulev, S.V.Mikhailov, N.G.Stefanis(2007,2008), A.L. Kataev, VK (2007,2008)



b-quark mass running:

- 4.5 GeV Upsilon mass
- 2.8 GeV BEH boson mass



Brout-Englert-Higgs boson discovery of CMS and ATLAS in 2012 is most important physics result at LHC upto now

ATLAS, Phys. Lett. B 716 (2012) 1

CMS, Phys. Lett. B 716 (2012) 30

Brout-Englert-Higgs-Guralnik-Hagen-Kibble mechanism of spontaneous symmetry breaking

R. Brout, F. Englert, Phys. Rev. Lett. 13 (1964) 321

P.W. Higgs, Phys. Rev. Lett. 13 (1964) 508

G.S. Guralnik, C.R. Hagen, T.W.B. Kibble, Phys. Rev. Lett. 13 (1964) 585

P. Higgs & F. Englert: Nobel Prize (2013)

Higgs boson is only scalar elementary particle known up to now



**Brout-Englert-Higgs boson:
if only logarithmic mass evolution**

**Brout-Englert-Higgs boson defines electroweak vacuum density
(meta)stable vacuum up to Planck scales**

**F. Bezrukov, M. Kalmykov, B. Kiehl & M. Shaposhnikov, JHEP 10 (2012) 140
within $\overline{\text{MS}}$ -scheme**

**One may conclude: M. Shaposhnikov et al.
(Almost) no need for a New Physics up to Planck scales**

Only needs:

- (~ 1 GeV) BSM neutral leptons to explain Dark Matter
- strong CP-problem
- neutrino masses
- baryon-antibaryon asymmetry

...

**- and still needs to explain:
why there is naturalness (New Physics?!)**

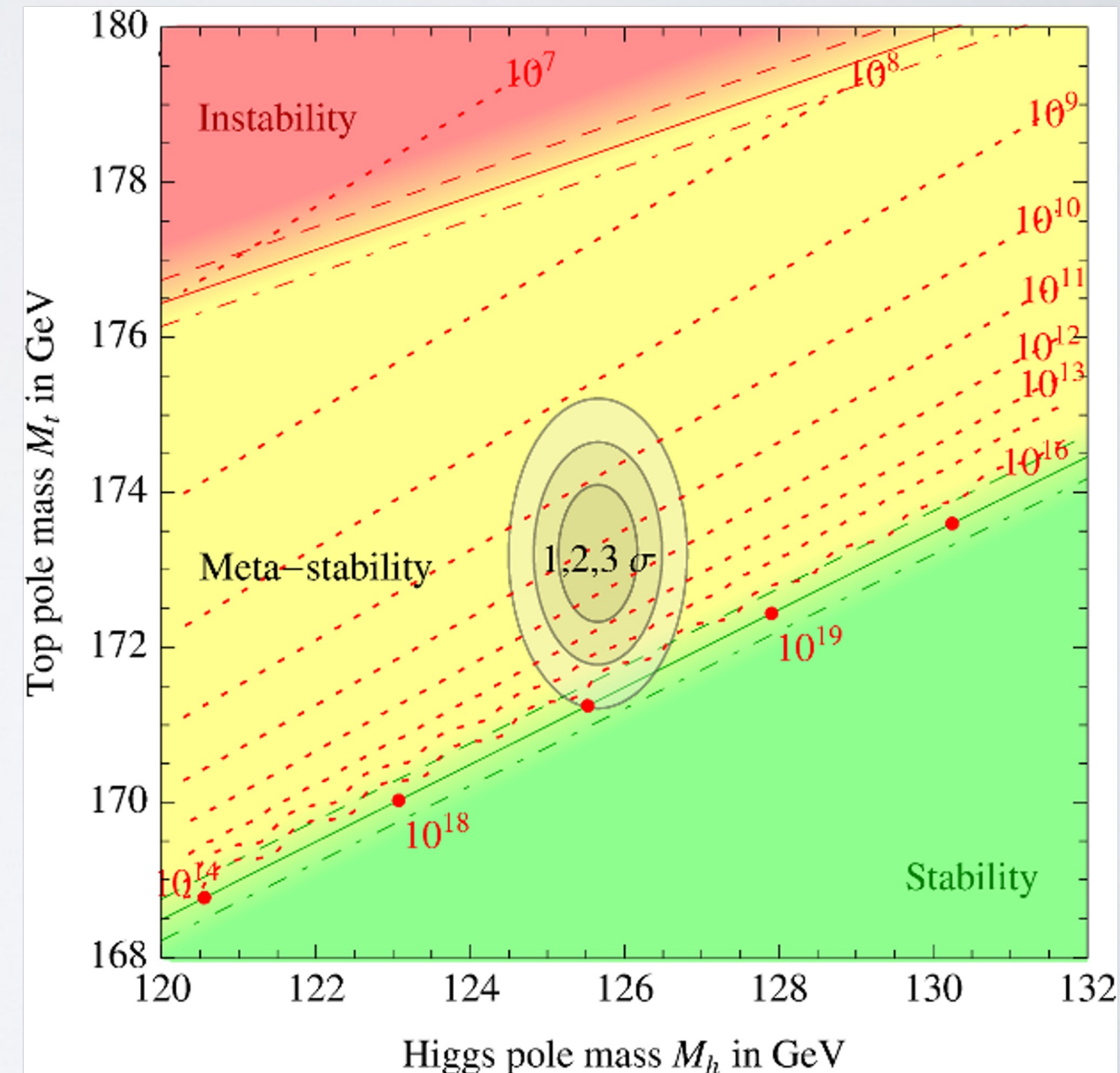
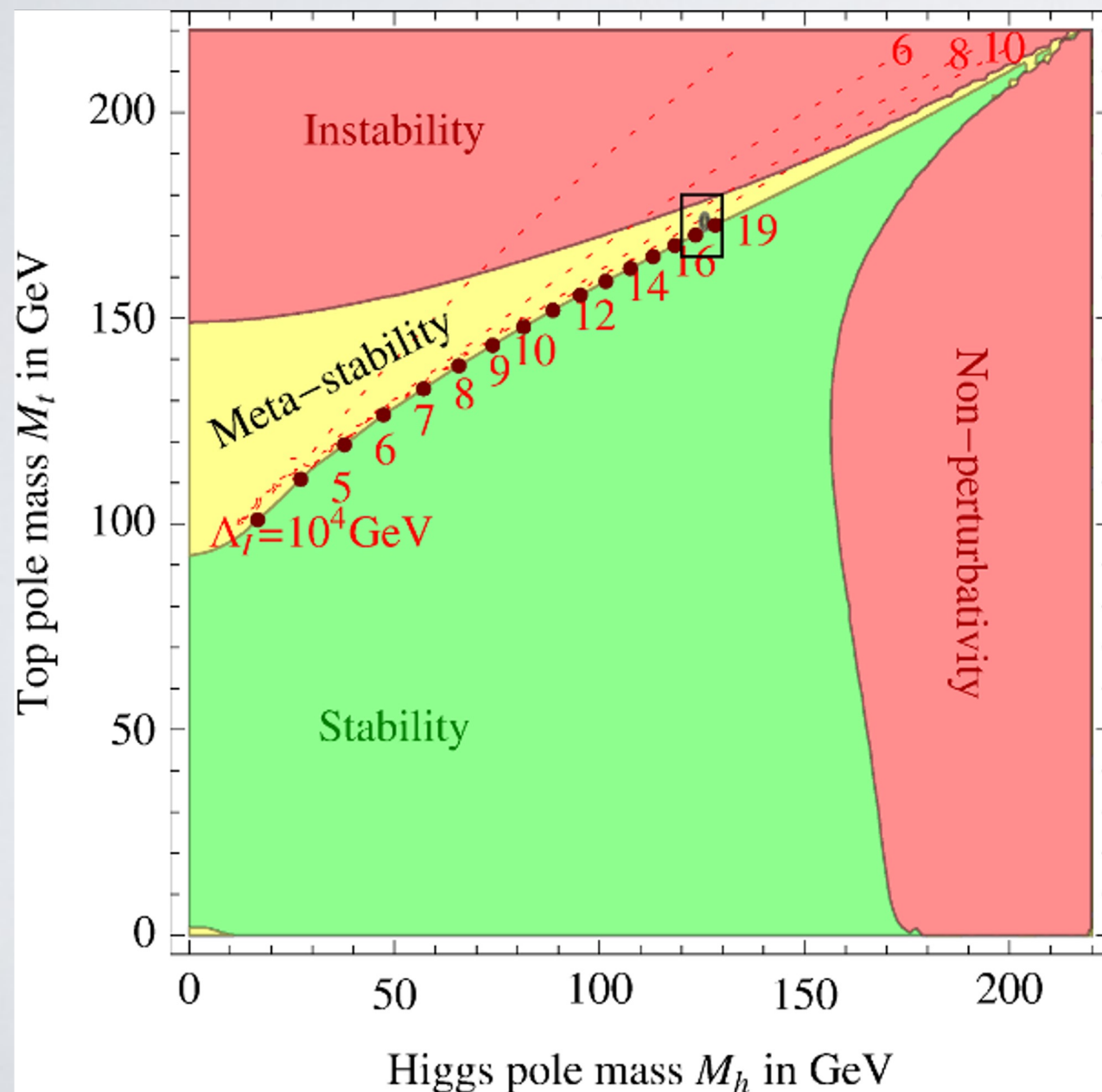
The Standard Model with 125 GeV Higgs boson

Brout-Englert-Higgs boson mass defines
electroweak vacuum density: meta-stable vacuum

G. Degrassi et al., JHEP 08 (2012) 098

D. Butazzo et al., JHEP 12 (2013) 089

A.V.Bednyakov, B.A.Kniehl, A.F.Pikelner, O.L.Veretin, PRL 115 (2015) 201802





**Logarithmic evolution of theory parameters:
weak dependence between low and very large scales
-> concept of "Naturalness"**

- Scalar field is simple, but “non-natural”:
scalar mass evolution is quadratic, not only logarithmic**

K. Wilson, Phys. Rev. D3 (1971) 1818

L. Susskind, Phys. Rev. D20 (1979) 2619

- Scalar field is not protected by a symmetry,
while fermions are protected by chiral symmetry**

G. 't Hooft, Proc. Cargese Summer Inst. (1980) 135-157

for reviews see, e.g., G.F. Giudice (2008), L. Alvarez-Gaume, M.A.Vazquez-Mozo (2012)

Quantifying Naturalness: Criteria



G. 't Hooft, Proc. Gargese Inst. (1980) pp.135-157

P.A.M. Dirac 139 (1937) 323, Proc. Roy. Soc. A165 (1948) 199

Naturalness Criterion by G. 't Hooft:
physics parameter can be small at any scale only
if setting this parameter to zero would increase
the symmetry of the system

E.g., naturalness for fermions: $m=0$ leads to chiral symmetry

However, nothing like that for scalar particles ...

dilaton symmetry W. Bardeen (1995)

Quantifying Naturalness: Criteria



K. Wilson, Phys. Rev. D3 (1971) 1818

L. Susskind, Phys. Rev. D20 (1979) 2619

G. 't Hooft, Proc. Gargese Inst. (1980) pp.135-157

G.F. Giudice (2008) [arXiv:0801.2562 [hep-ph]]

L. Alvarez-Gaume, M.A. Vazquez-Mozo, Invitation to Quant. Field Theory (2012)

E.g., naturalness for fermions: $m=0$ leads to chiral symmetry

However, nothing like that for scalar particles ...

dilaton symmetry W. Bardeen (1995)

The Standard Model: Brout-Englert-Higgs boson roles



The two main roles of BEH boson in the SM:

- **providing mass** to the SM particles by Brout-Englert-Higgs mechanism
 - **restoring unitarity** for EW vector boson scattering: Brout-Englert-Higgs boson cancels quadratic growth of longitudinal components for EW vector bosons with collision energy
-

- **if Higgs boson could be very light** -> no noticeable growth with collision energy
- **if Higgs boson could be very heavy** -> strong growth of EW vector boson interaction -> new SM dynamics: nonperturbative strong EW interaction can lead to heavy EW resonances



The two main regularization methods:

- **Dimensional regularization**

“Geometrical”: regulator depends on geometry of space-time

- **Cut-off regularization**

Covariant Pauli-Villars

“Physical”



Собственная энергия (пропагатор):

$$-i\Sigma(p^2) = -\frac{\lambda}{2} \int \frac{d^4 k}{2\pi^4} \frac{i}{k^2 - \mu^2 + i\epsilon} = \frac{-i\lambda\pi^{\frac{n}{2}}\Gamma\left(1 - \frac{n}{2}\right)}{32\pi^4(\mu^2)^{1-\frac{n}{2}}}$$

Регуляризованная собственная энергия:

$$\begin{aligned} -i\Sigma(p^2) &= \frac{\lambda}{2} \int \frac{d^4 k}{2\pi^4} \frac{\Lambda^4}{(k^2 - \mu^2 + i\epsilon)(k^2 - \Lambda^2 + i\epsilon)^2} \\ &= \frac{-i\lambda\Lambda^4}{32\pi^2} \int_0^1 \frac{\alpha d\alpha}{\alpha\Lambda^2 + (1-\alpha)\mu^2} \\ &= \frac{-i\lambda}{32\pi^2} \left[\Lambda^2 - \mu^2 \log\left(\frac{\Lambda^2}{\mu^2}\right) \right] \end{aligned}$$



Размерность пространства $D = n - \varepsilon < 4$, $\varepsilon > 0$

$$-i\Sigma(p^2) = \frac{\lambda}{2} \int \frac{d^n k}{2\pi^4} \frac{1}{k^2 - \mu^2 + i\epsilon} = \frac{-i\lambda\pi^{\frac{n}{2}}\Gamma(1 - \frac{n}{2})}{32\pi^4(\mu^2)^{1-\frac{n}{2}}}$$

$$\Gamma\left(1 - \frac{n}{2}\right) = \frac{\Gamma(3 - \frac{n}{2})}{(1 - \frac{n}{2})(2 - \frac{n}{2})}$$

Полюс: $n = 4$ (и при $n = 2$)

При $n \rightarrow 4$ только логарифмическая расходимость (!):

$$-i\Sigma(0) = \frac{i\lambda\mu^2}{16\pi^2} \frac{1}{(4 - n)}$$

Н.В. Вклад $n \rightarrow 2$ дает квадратичную расходимость (!)

Но зависит от числа петель L : $D = 4 - 2/L$

При $L \rightarrow \infty$ $D = n \rightarrow 4$

M. Veltman, Acta Phys. Pol. B12 (1981) 437

Naturalness of the Standard Model at 1-loop



M. Veltman, Acta Phys. Pol. B12 (1981) 437

renormalization scheme dependence for scalar particles:

$$m_H^2 = m_{0H}^2 + C_L(\lambda_i, m_i) \cdot \log\left(\frac{\Lambda_{UV}^2}{m^2}\right) + C_X(\lambda_i, m_i) \cdot \Lambda_{UV}^2$$

"physical" schemes $\rightarrow C_X \neq 0$

schemes with dim. regularization ($\overline{\text{MS}}$, ...) $\rightarrow C_X = 0$

$\overline{\text{MS}}$ reproduces quadratic divergence at $D = 2$, $L=1$

$$m_H^2 = m_{H0}^2 + \delta m_H^2$$

$$v = 246 \text{ GeV}$$

$$\delta m_H^2 \approx \frac{\Lambda^2}{16\pi^2} (24y_t^2 - 6(2y_W^2 + y_Z^2 + y_H^2)) \sim 8.2 \frac{\Lambda^2}{16\pi^2}$$

$$y_i \equiv \frac{m_i}{v}$$

**Non-naturalness of BEH boson at $\Lambda > 550 \text{ GeV}$
(Veltman criterion):**

$$\delta m_H^2 \approx m_H^2 \quad (\Lambda = 550 \text{ GeV}, \quad m_H = 125 \text{ GeV})$$

The Standard Model: Higgs boson mass evolution



M. Veltman, Acta Phys. Pol. B12 (1981) 437

quadratic mass divergences within MSbar renormalization:

Dim = 4 – 2/L

$$m_R^2 = m_B^2 + \xi \Lambda^2,$$

where $\xi = \xi(m_H, m_t, m_W, m_Z)$

Veltman's condition for absence of quadratic mass divergences:

$$\xi(m_H, m_t, m_W, m_Z) = 0$$

Veltman's condition holds up to 3-loops:

but in higher orders it cannot be hold in self-consistent way

M.S. Al-sarhi, I. Jack, D.R.T. Jones, Zeit fur Physik Pol. C55 (1992) 283

Veltman's condition and Higgs effective potential

M.B. Einhorn, D.R.T. Jones, Phys. Rev. D42 (1992) 5206

Quadratic divergence cancellation: Veltman's condition



Масса бозона Браута-Энглера-Хиггса:

$$m^2 = m_0^2 + \Delta m^2 = m_0^2 + \xi(\lambda_0, g) \Lambda^2$$

Условие сокращения квадратичной расходимости Вельтмана:

$$\xi(\lambda_0, g) = 3(3g_2^2 + g_1^2 + 2\lambda_0 - 4y_t^2)/(32\pi^2) = 0$$

g_1 , g_2 , y_t - константы СМ: групп $U(1)$, $SU(2)$, юкавская константа t -кварка

M. Veltman, Acta Phys. Pol. B12 (1981) 437

Нарушается с 4-х петель $L \geq 4$

M.S. Al-sarhi, I. Jack, D.R.T. Jones, Z.Phys. C55 (1992) 283



M. Veltman, Acta Phys. Pol. B12 (1981) 437

$$m_H^2 = m_{H0}^2 + \delta m_H^2$$

$$\delta m_H^2 \approx \frac{\Lambda^2}{16\pi^2} (24y_t^2 - 6(2y_W^2 + y_Z^2 + y_H^2)) \sim 8.2 \frac{\Lambda^2}{16\pi^2}$$

$$y_i \equiv \frac{m_i}{v} \quad v = 246 \text{ GeV}$$

Veltman's criterion:

$$\left| \frac{m^2 - m_0^2}{m_0^2} \right| = \left| \frac{\Delta m^2}{m_0^2} \right| \leq q = 1$$

Non-naturalness of BEH boson at $\Lambda > 550 \text{ GeV}$:

$$\delta m_H^2 \approx m_H^2 \quad (\Lambda = 550 \text{ GeV}, \quad m_H = 125 \text{ GeV})$$

Log-Derivative criterion:

$$\left| \frac{\lambda_0}{m^2} \frac{\partial m^2}{\partial \lambda_0} \right| \leq q \simeq 10$$

R. Barbieri, G.F. Giudice, Nucl. Phys. B306 (1988) 63

J.R.Ellis, K.Enquist, D.V.Nanopoulos, F.Zwirner, Mod.Phys.Lett. A1(1986)57

$$\begin{aligned} m_0^2 &= m^2 - \xi(\lambda, g)\Lambda^2 \\ \lambda_0 &= \lambda + \frac{\beta(\lambda, g)}{2} \log\left(\frac{\Lambda^2}{m^2}\right) \end{aligned}$$

**Logarithmic sensitivity
Transformation matrix**

RG mixing -> matrix non-degeneracy -> matrix

V.K., G. Pivovarov, Phys. Rev. D78 (2008) 016001

V.K., Phys. Part. Nucl. 55 (2024) 156

$$\begin{pmatrix} m_0^2 \\ \lambda_0 \end{pmatrix} = A \times \begin{pmatrix} m^2 \\ \lambda \end{pmatrix}$$

$$A = \begin{pmatrix} \frac{\partial \lambda}{\partial \lambda_0} & \frac{\partial \lambda}{\partial m_0^2} \\ \frac{\partial m^2}{\partial \lambda_0} & \frac{\partial m^2}{\partial m_0^2} \end{pmatrix}$$

$$A = \frac{1}{\det(A^{-1})} \begin{pmatrix} 1 & \frac{\beta(\lambda, g)}{2m^2} \\ \xi'(\lambda, g)\Lambda^2 & 1 + \frac{\beta'(\lambda, g)}{2} \log\left(\frac{\Lambda^2}{m^2}\right) \end{pmatrix}$$

$$\det(A^{-1}) = 1 - \xi'(\lambda, g) \frac{\beta(\lambda, g)}{2} \frac{\Lambda^2}{m^2} + \frac{\beta'(\lambda, g)}{2} \log\left(\frac{\Lambda^2}{m^2}\right)$$

BEH-boson mass evolution and naturalness criterion



Logarithmic sensitivity

Transformation matrix

RG mixing -> matrix non-degeneracy -> matrix

V.K., G. Pivovarov, Phys. Rev. D78 (2008) 016001

V.K., Phys. Part. Nucl. 55 (2024) 156

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**Non-naturalness from BEH Higgs boson mass
at $\Lambda \sim 0$ (10 TeV)**

Previously without logs: $\Lambda \sim 0$ (1 TeV)



Quantifying Naturalness

Naturalness log-derivative criterion:

weak sensitivity physical parameters for small variation of bare ones

J.R.Ellis, K.Enquist, D.V.Nanopoulos, F.Zwirner, Mod.Phys.Lett. A1(1986)57

R. Barbieri, G.F. Giudice, Nucl. Phys. B306 (1988) 63

Using BG condition with both quadratic and logarithmic contributions leads to extension of Naturalness domain of SM:

up $\sim O(10 \text{ TeV})$ instead of $\sim O(1 \text{ TeV})$

V.K., G. Pivovarov, Phys. Rev. D78 (2008) 016001

V.K., Phys. Part. Nucl. 55 (2024) 156

Regular way for scalar boson mass evolution with quadratic mass divergences

G. Pivovarov, Phys. Rev. D81 (2010) 076077

K. Fujikawa, Int. Mod. Phys. A (2016)

The SM with the “non-natural” Brout-Englert-Higgs boson



**Proper physical consideration with quadratic evolution
for Brout-Englert-Higgs boson mass:**

**Higgs boson observables (mass, self-coupling, EW vacuum density)
gets critical values at later scales
than in popular “standard” naturalness treatments**

- > at those scales $\sim O(10 \text{ TeV})$ one should expect
new physics manifestations:**
 - new strong EW dynamics**
 - or/and New Physics beyond Standard Model**

Summary



- **The Standard Model without quadratic evolution for Brout-Englert-Higgs boson mass requires (!) New Physics to have Naturalness**
- **Naturalness domain of the Standard Model with quadratic evolution for Brout-Englert-Higgs boson mass may be larger than generally accepted: instead of $\sim O(1 \text{ TeV})$ it is up $\sim O(10 \text{ TeV})$**
- **Present LHC physics: new physics is unavoidable either as a new dynamics of SM or/and a New Physics. Besides search direct search of New Physics it requires ‘non-naturalness’ studies**