INFINUM 2025

Many-body effects of hyperonic interactions in neutron stars

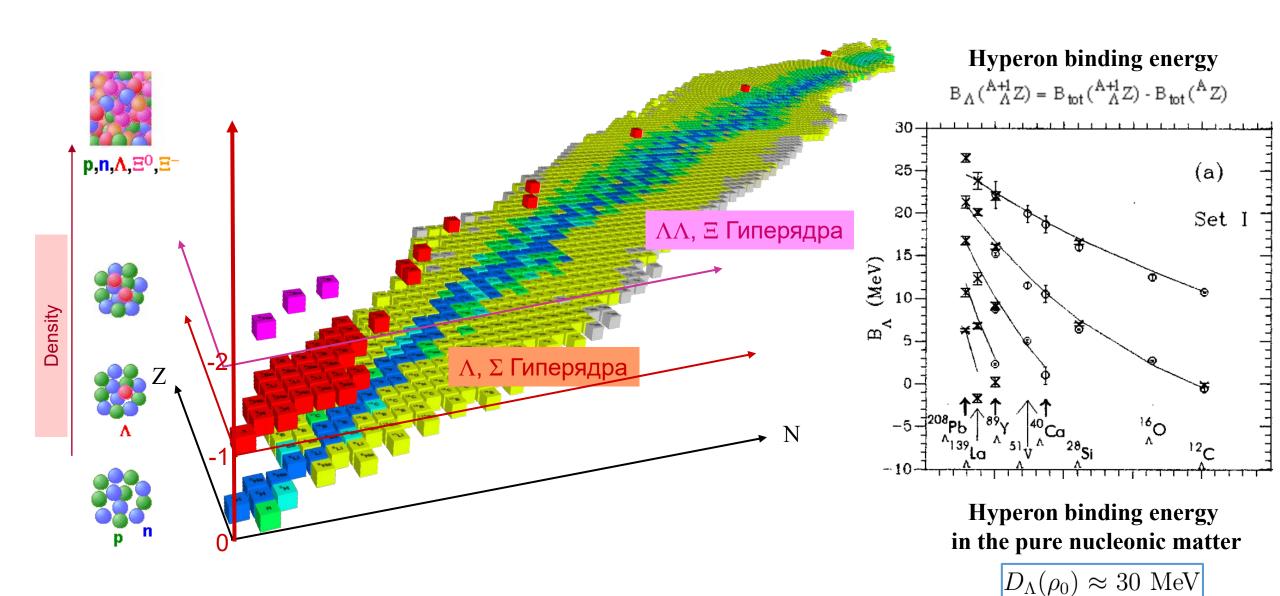
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Hypernuclei and hyperonic interactions



Skyrme interaction

AN-interaction

$$V_{\Lambda N}(\overrightarrow{r_{\Lambda}}, \overrightarrow{r_{N}}) = u_{0}(1 + \xi_{0}P_{\sigma})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})$$

$$+ \frac{1}{2}u_{1}(1 + \xi_{1}P_{\sigma})[\overrightarrow{P}'^{2}\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}}) + \delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\overrightarrow{P}^{2}]$$

$$+ u_{2}(1 + \xi_{2}P_{\sigma})\overrightarrow{P}'\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\overrightarrow{P}$$

$$+ iW_{0}^{\Lambda}\overrightarrow{P}'\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})[\overrightarrow{\sigma} \times \overrightarrow{P}]$$

| Parametrization of AN-interaction | γ |
|-----------------------------------|-----|
| YBZ6 | 1 |
| YBZ2 | 1 |
| SLL4' | 1 |
| LYI | 1/3 |
| YMR | 1/8 |

Three-body forces

Density-dependent forces

$$V_{3} = V_{\Lambda NN}(\overrightarrow{r_{\Lambda}}, \overrightarrow{r_{N1}}, \overrightarrow{r_{N2}}) = u_{3}\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N1}})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N2}})$$

$$V_{3} = V_{\Lambda N}(\overrightarrow{r_{\Lambda}}, \overrightarrow{r_{N}}, \rho) = \frac{3}{8}u_{3}(1 + \xi_{3}P_{\sigma})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\rho_{N}^{\gamma}(\frac{\overrightarrow{r_{\Lambda}} + \overrightarrow{r_{N}}}{2})$$

ΛΛ-interaction

$$V_{\Lambda\Lambda}(\overrightarrow{r_1}, \overrightarrow{r_2}) = \lambda_0 \delta(\overrightarrow{r_1} - \overrightarrow{r_2})$$

$$+ \frac{1}{2} \lambda_1 [\overrightarrow{P}^{\prime 2} \delta(\overrightarrow{r_1} - \overrightarrow{r_2}) + \delta(\overrightarrow{r_1} - \overrightarrow{r_2}) \overrightarrow{P}^2]$$

ΛΛ-interaction with density dependence

$$V_{\Lambda\Lambda} = \sum_{1}^{3} (a_i + b_i k_F + c_i k_F^2) e^{-\frac{r^2}{\beta_i^2}}$$

Neutron stars

Chemical equilibrium

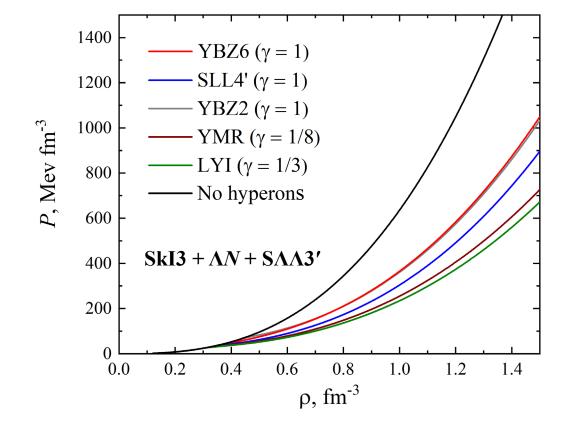
$$\begin{cases} \mu_p + \mu_e = \mu_n \\ \mu_\mu = \mu_e \\ \mu_\Lambda + m_\Lambda = \mu_n + m_n \end{cases}$$

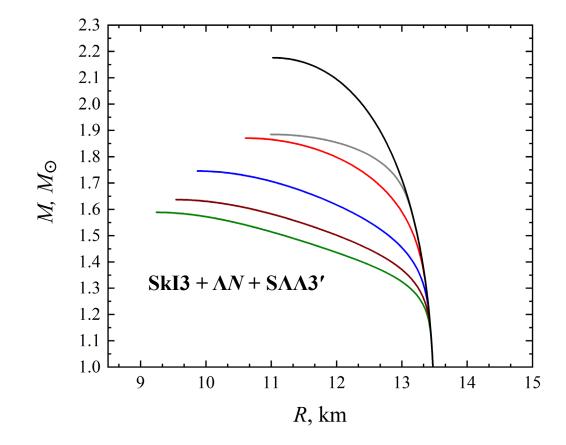
Tolman Oppenheimer Volkov equation

$$\begin{split} \frac{dP}{dr} &= \frac{G}{r^2} \frac{[\rho(r) + P(r)/c^2][m(r) + (4\pi r^3 P(r)/c^2)]}{1 - (2Gm(r)/rc^2)} \\ \frac{dm}{dr} &= 4\pi r^2 \rho(r) \end{split}$$

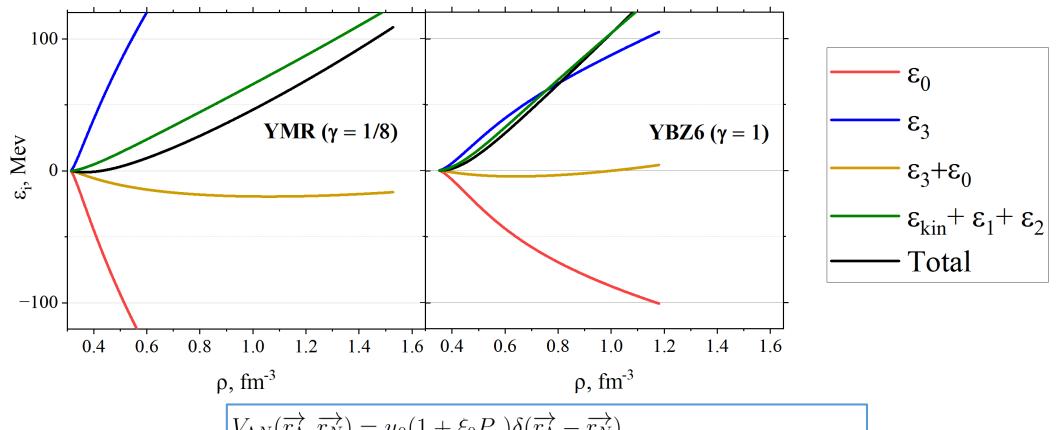
Hyperon puzzle

PSR J0740+6620, $M = 2.08 \pm 0.07 M_{\odot}$ PSR J0952-0607, $M = 2.35 \pm 0.17 M_{\odot}$





Contributions of various terms in energy per baryon



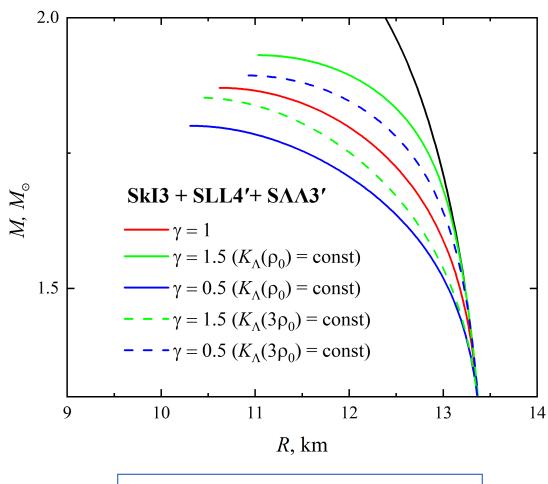
$$V_{\Lambda N}(\overrightarrow{r_{\Lambda}},\overrightarrow{r_{N}}) = u_{0}(1 + \xi_{0}P_{\sigma})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})$$

$$+ \frac{1}{2}u_{1}(1 + \xi_{1}P_{\sigma})[\overrightarrow{P}'^{2}\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}}) + \delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\overrightarrow{P}^{2}]$$

$$+ u_{2}(1 + \xi_{2}P_{\sigma})\overrightarrow{P}'\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\overrightarrow{P}$$

$$+ \frac{3}{8}u_{3}(1 + \xi_{3}P_{\sigma})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\rho_{N}^{\gamma}(\frac{\overrightarrow{r_{\Lambda}} + \overrightarrow{r_{N}}}{2})$$

Masses and radii of neutron star for different values of γ



Changed parameters: γ, u₀, u₃

The binding energy of Λ -hyperon in the pure nucleonic matter

$$D_{\Lambda} = -\mu_{\Lambda}$$

$$D_{\Lambda}(\rho_0) \approx 30 \text{ MeV}$$

Compression power of AN-interaction

$$K_{\Lambda}=3
horac{dD_{\Lambda}(
ho)}{d
ho}$$

$$K_{\Lambda}(
ho_0)=const$$

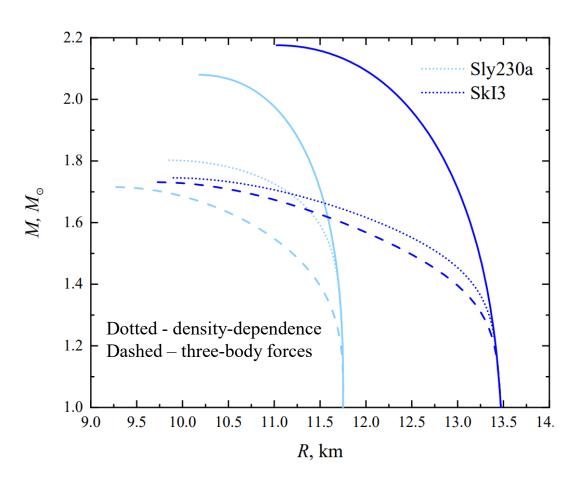
$$K_{\Lambda}(3
ho_0)=const$$

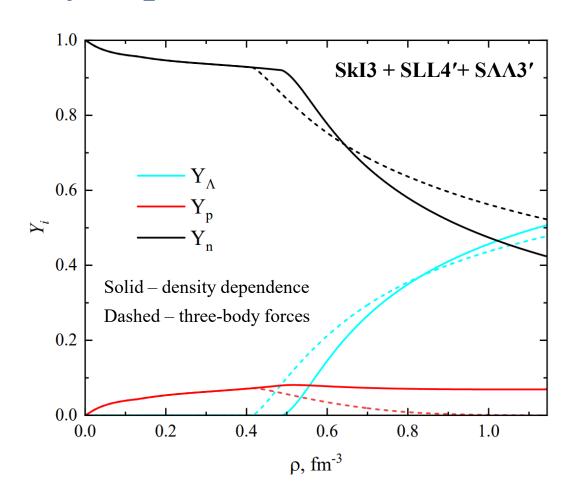
$$K_{\Lambda}(3
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Three-body and density-dependent forces





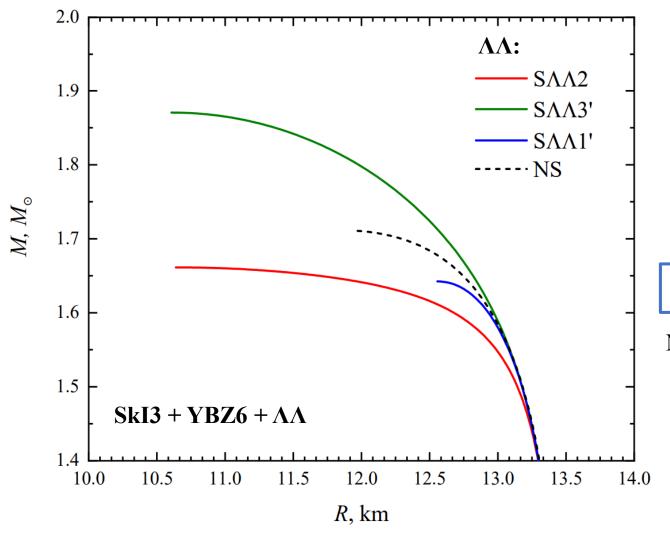
Three-body forces

Density-dependent forces

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$$V_{3} = V_{\Lambda N}(\overrightarrow{r_{\Lambda}}, \overrightarrow{r_{N}}, \rho) = \frac{3}{8}u_{3}(1 + \xi_{3}P_{\sigma})\delta(\overrightarrow{r_{\Lambda}} - \overrightarrow{r_{N}})\rho_{N}^{\gamma}(\frac{\overrightarrow{r_{\Lambda}} + \overrightarrow{r_{N}}}{2})$$

AA-interaction



D. E. Lanskoy, 1998, Minato F., 2011

| ΛΛ-interaction | Radius of interaction |
|----------------|-----------------------|
| SΛΛ1' | Small |
| SΛΛ2 | Medium |
| SAA3' | Large |

NS – density dependent $\Lambda\Lambda$ -interaction

NSC89 (Nijmegen) \longrightarrow NS (Gauss) \longrightarrow NS (Skyrme)

Conclusion

- The equation of state of the neutron star matter is determined by a complex interplay of the contributions of the different terms of the potential of ΛN -interaction and various values of γ in the density dependence can be acceptable.
- The three-body ΛNN forces in neutron stars lead to a softer equation of state than density-dependent ΛN forces and can even considerably affect the chemical composition of the star leading to disappearance of protons at a certain density.
- Density dependence in $\Lambda\Lambda$ -interaction does not have such a large effect on the characteristics of neutron stars as in ΛN -interaction interactions and this issue requires further study.



THANK YOU FOR ATTENTION

Back up

NS approximation

$$V_{\Lambda\Lambda} = \sum_{1}^{3} (a_i + b_i k_F + c_i k_F^2) e^{-\frac{r^2}{\beta_i^2}}$$

$$\begin{split} V_{\Lambda\Lambda}(\mathbf{r}_{1}, \mathbf{r}_{2}) &= \lambda_{0}\delta(\mathbf{r}_{1} - \mathbf{r}_{2}) \\ &+ \frac{1}{2}\lambda_{1}[\mathbf{P}^{'2}\delta(\mathbf{r}_{1} - \mathbf{r}_{2}) + \delta(\mathbf{r}_{1} - \mathbf{r}_{2})\mathbf{P}^{2}] \\ &+ \lambda_{2}\mathbf{P}^{'}\delta(\mathbf{r}_{1} - \mathbf{r}_{2})\mathbf{P} \\ &+ \sum_{i}\lambda_{3}^{i}\delta(\mathbf{r}_{1} - \mathbf{r}_{2})\rho_{N}^{\delta_{i}} \\ &+ \frac{1}{2}\sum_{i}\lambda_{4}^{i}[\mathbf{P}^{'2}\delta(\mathbf{r}_{1} - \mathbf{r}_{2}) + \delta(\mathbf{r}_{1} - \mathbf{r}_{2})\mathbf{P}^{2}]\rho_{N}^{\delta_{i}}. \end{split}$$

CSB in three-body forces

$$\varepsilon_3 = \frac{1}{4}u_3Y_{\Lambda}(\rho_N^2 + 2\rho_p\rho_n) = \frac{1}{8}u_3Y_{\Lambda}(3\rho_N^2 - \rho_-^2).$$

NS parameters

$$\lambda_0 = \pi^{3/2} \sum_{1}^{3} a_i \beta_i^3,$$

$$\lambda_1 = -\frac{1}{2} \pi^{3/2} \sum_{1}^{3} a_i \beta_i^5,$$

$$\lambda_3^1 = \left(\frac{3\pi^2}{2}\right)^{1/3} \pi^{3/2} \sum_{1}^{3} b_i \beta_i^3,$$

$$\lambda_3^2 = \left(\frac{3\pi^2}{2}\right)^{2/3} \pi^{3/2} \sum_{1}^{3} c_i \beta_i^3,$$

$$\lambda_4^1 = -\frac{1}{2} \left(\frac{3\pi^2}{2}\right)^{1/3} \pi^{3/2} \sum_{1}^{3} b_i \beta_i^5,$$

$$\lambda_4^2 = -\frac{1}{2} \left(\frac{3\pi^2}{2}\right)^{2/3} \pi^{3/2} \sum_{1}^{3} c_i \beta_i^5.$$

| Table 1. Parameters λ_0 , λ_1 , λ_3^i and λ_4^i (9) of NS parametrization in Skyrme form | | | | | | |
|--|-------------------------------|-------------------------------------|-------------------------------------|-------------------------------------|-------------------------------------|--|
| $\lambda_0, \mathrm{MeV} \mathrm{fm}^3$ | $\lambda_1,{ m MeV}{ m fm}^5$ | λ_3^1 , MeV fm ⁴ | λ_3^2 , MeV fm ⁵ | λ_4^1 , MeV fm ⁶ | λ_4^2 , MeV fm ⁷ | |
| -833.1 | 646.4 | 1268.8 | -960.5 | -735.4 | 625.0 | |