

Analysis Note

Directed flow v_1 of deuterons in the Xe+Cs(I) collisions at 3.8 AGeV (BM@N run8)

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In this note, we present the directed flow v_1 measurements of protons from Xe+Cs(I) collisions at 3.8 AGeV (BM@N run8). We show the datasets, event and track selection cuts, centrality definition, event plane reconstruction and resolution. The v_1 results are presented as function of transverse momentum (p_T) and rapidity (y_{cm}) for 10-30% central Xe+Cs(I) collisions. The systematic uncertainty study will also be presented and discussed. The v_1 measurements are compared with results of JAM transport model calculations and published data from other experiments.

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1 Introduction

Collective motion of particles produced in heavy ion collisions is the result of the evolution of the initial spatial asymmetry of matter created in the overlap region. This evolution occurs by the means of the interaction of the constituents of the matter with each other and with cold spectator fragments. Thus the final-state momentum anisotropy reflects the underlying properties of strongly interacting matter and is sensitive to the baryon densities and temperatures achieved in the collision. This momentum anisotropy is usually quantified with the Fourier-decomposition of the azimuthal angle distribution of the produced in the collision particles with respect to the reaction plane Ψ_{RP} , spanned on the vectors of impact parameter and beam direction [1; 2]:

$$\frac{dN}{d\varphi} \propto \frac{1}{2\pi} \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_{RP}) \right), \quad (1)$$

where φ is the particles azimuthal angle and coefficients v_n calculated as:

$$v_n = \langle \cos n(\varphi - \Psi_{RP}) \rangle, \quad (2)$$

where angle brackets denote the average over all azimuthal and reaction plane angles.

Coefficients v_n , usually referred to as anisotropic flow of n -th harmonic played an important role in discovering the properties of strongly interacting quark-gluon matter (QGM) at top RHIC and LHC energies. Using the data for the elliptic flow v_2 , it was shown that QGM posses the lowest known shear viscosity η/S [3].

At lower beam energies of several GeV per nucleon the passing time of the colliding nuclei is comparable to the time of the evolution of the matter within the overlapping region. Therefore interaction of the produced in the collision matter with cold spectator matter provides the significant contribution the final values of anisotropic flow [4]. The particle emission pattern changes flow in-plane to out-of-plane when going down the collision energy [5]. Additionally, presence of the spectator matter during the evolution of the matter within the overlapping region leads to achieving higher net-baryon densities exceeding by the factor of 3-5 that of regular nuclear matter [6]. Matter under such extreme conditions may be found in

51 the compact stellar objects such as neutron stars [7]. Studying the anisotropic flow
 52 of the produced in the collision particles may shed light on the properties of strongly
 53 interacting matter at these conditions and help to establish the equation of state
 54 (EOS) at the region of high net-baryon densities.

55 The mechanisms of light nuclei production in heavy-ion collisions at the region
 56 of beam energy of several A GeV is not fully understood. However, the recent findings
 57 of the STAR collaboration tend to favor the coalescence mechanism over others [8].
 58 In this model, the light nuclei are produced via binding of nucleons which are close
 59 in momentum space, which leads to the scaling of the directed flow:

$$v_1^p(p_T) \approx v_1^d/2(p_T/2) \approx v_1^t/3(p_T/3). \quad (3)$$

60 The Baryonic Matter at the Nuclotron (BM@N)[9] is a fixed target experiment
 61 at JINR (Dubna). In February 2023, the first physics run of the BMN experiment was
 62 completed with recorded Xe + Cs(I) collision events at $E_{beam} = 3A$ GeV ($\sqrt{s_{NN}} = 3.02$
 63 GeV) and $3.8A$ GeV ($\sqrt{s_{NN}} = 3.26$ GeV). In this analysis note, we present first results
 64 on directed flow (v_1) of protons in 10-30% central Xe + Cs(I) collisions at $E_{beam} =$
 65 $3.8A$ GeV. The analysis presented bellow is based on the analysis performed for the
 66 directed flow of protons. The data selection criteria and symmetry plane estimation
 67 are carried out identically to [10]. While the approach for deuteron selection and
 68 systematic error evaluation is similar to how it was done for protons, there will be
 69 a slight variations. The note is organized as follows. Section 2 is dedicated to event
 70 and track selection. Section 3 briefly describes the methods used for analysis of the
 71 directed flow of deuterons. In section 4 we present the study for the systematic
 72 uncertainty. Finally, Section 5 presents the main results on directed flow v_1 of
 73 deuterons.

74 2 Data, Event and Track Selection

75 In total approximately 500 million events of Xe+Cs(I) collisions at the beam
76 energy of 3.8A GeV were collected by the BM@N experiment in the January of 2023.
77 The event selection procedure is identical to how it was done for proton directed
78 flow analysis. Bellow we briefly remind the selection criteria employed. For further
79 details see [10].

- 80 1. We don't consider runs below RunId=6924 due to unstable operation of the
81 GEM and FSD detectors (BM@N Electronic Logbook).
- 82 2. We removed 74 runs [18M events] based on QA study (see our previous analysis
83 note [10])
- 84 3. We used events from Physical runs and CCT2 trigger [9].
- 85 4. At least 2 tracks in vertex reconstruction
- 86 5. The pileup events were rejected based on the $\pm 3\sigma$ cut on the correlation be-
87 tween the number of FSD digits and the number of charged particles in the
88 tracking system (FSD + GEM).

Table 1: Statistics after applying the selection criteria

Cuts	no. of events	%
def.	530 M	100%
CCT2 trigger	437 M	82%
at least 2 tracks in vertex reconstruction	315 M	59%
Pileup rejection cuts	285 M	53%
QA study	267 M	50%

89 Selection criteria are also imposed on tracks to ensure good tracks for analysis.
90 The selection cuts applied are as follows:

- 91 1. Tracks of charged particles were selected based on the number of stations N_{hits}
92 in the BM@N inner tracking system used for track reconstruction. At least 6
93 were required to satisfy the criteria of a good track: $N_{hits} > 6$.

- 94 2. Only tracks with fit quality $\chi^2/NDF < 5$ were analyzed.
- 95 3. Distance of the closest approach (DCA) of tracks from the primary vertex in
96 the direction perpendicular to the beam: $DCA < 5$ cm

97 Deuterons are identified using the time of flight Δt measured between T0 and
98 the ToF detectors, the length of the trajectory ΔL and the momentum p recon-
99 structed in the BM@N central tracker. Then the squared mass m^2 of a particle
100 is calculated. Figure 1 shows the population of the positively charged particle in
101 the plane spanned by their square mass m^2 and rigidity (p/q) for the TOF400 and
102 TOF700. For each bin in momentum the position $\langle m_d^2 \rangle$ and the width $\sigma_{m_d^2}$ of the
103 deuteron m^2 peak was extracted from the Gaussian fit. The procedure was done
104 separately for TOF-400 and TOF-700 as they have different timing resolution. The
105 proton peak was approximated with gaussian fit as well to calculate the possible con-
106 tamination and to develop the adequate selection criteria to maximize the deuteron
107 purity. Figure 2 shows the m^2 distribution of particles detected in the TOF700 for
108 the particular momenta from the different momentum. Red curves are approxima-
109 tions of the proton and deuteron peaks, and dashed vertical lines indicate the m^2
110 range in which the particles are selected as deuterons.

111 m^2 ranges for different momentum windows used in the analysis are shown in
112 Table 2.

p/q	$m^2 - \langle m^2 \rangle$	
	TOF400	TOF700
0–4 GeV/c	$(-3\sigma, 3\sigma)$	$(-2\sigma, 3\sigma)$
4–6 GeV/c	$(-2\sigma, 3\sigma)$	$(-1\sigma, 3\sigma)$
6–8 GeV/c	$(-1\sigma, 3\sigma)$	$(0, 3\sigma)$

Table 2: Deuteron selection criteria based on the square mass m^2 in different rigidity (p/q) ranges.

113 On Figure 4 shown are the selected deuterons in the m^2 vs. p/q plane for the
114 TOF400 and TOF700 detectors. Figure 5 shows the m^2 distribution of all particles
115 detected in the TOF700 and of the selected deuterons for different y and p_T bins,
116 used in the analysis further.

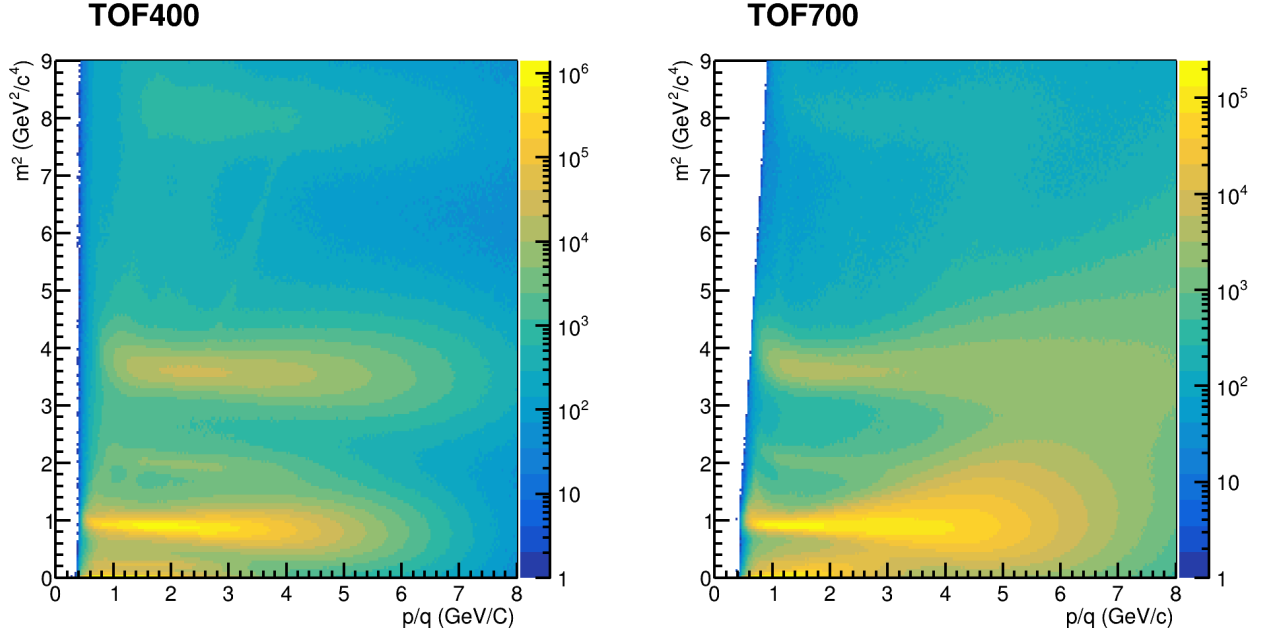


Figure 1: Population of the positively charged particles in the square mass m^2 vs. rigidity (p/q) plane for the TOF400 (left) and TOF700 (right).

117 Figure 6 shows the phase space coverage of identified deuterons as a function
 118 of rapidity y_{cm} and transverse momentum p_T for TOF-400, TOF-700 and for the
 119 combined system. Efficiency of the deuteron reconstruction was calculated using the
 120 realistic Monte-Carlo modelling of the BM@N experiment using GEANT4 transport
 121 code and DCM-QGSM-SMM model [11; 12]. Efficiency of the deuteron reconstruc-
 122 tion with the TOF-detectors acceptance applied is shown in the Figure. 7.

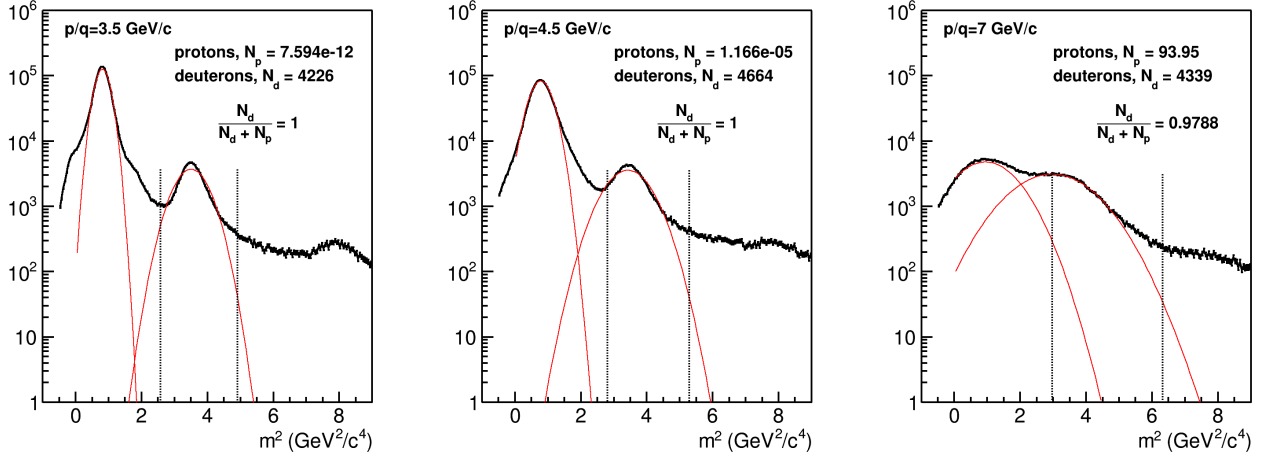


Figure 2: m^2 distribution of particles detected in the TOF700 for several rigidity (p/q) values. Red curves are approximations of the proton and deuteron peaks, and the vertical lines indicate the m^2 ranges for deuteron selection.

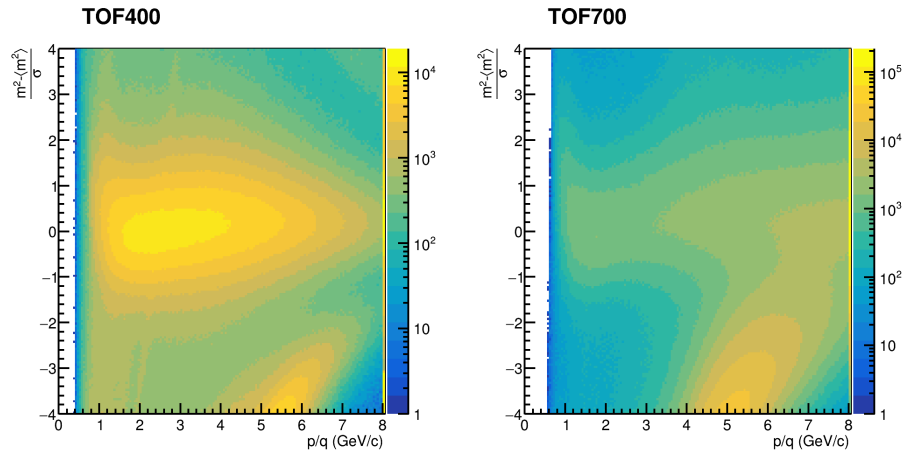


Figure 3: Population of charged particles in the n -sigma (m_p^2) = $(m^2 - \langle m_d^2 \rangle) / \sigma_{m_d^2}$ vs. rigidity (p/q) plane for the TOF-400 (left) and TOF-700 (right) detectors.

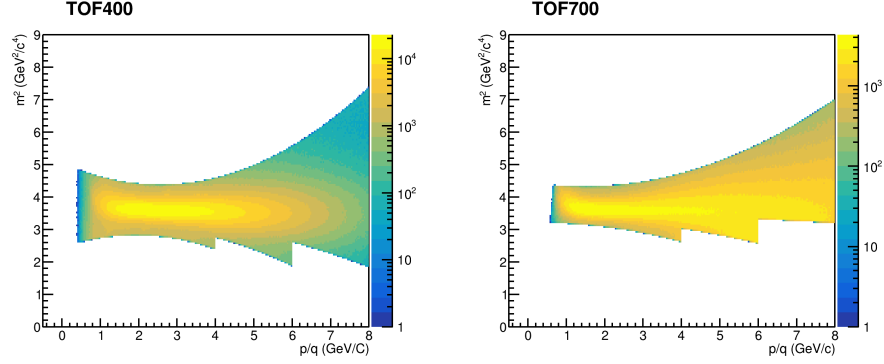


Figure 4: Population of selected deuterons in the m^2 vs. rigidity (p/q) plane for the TOF-400 (left) and TOF-700 (right) detectors. The deuteron selection criteria are shown in Table 2.

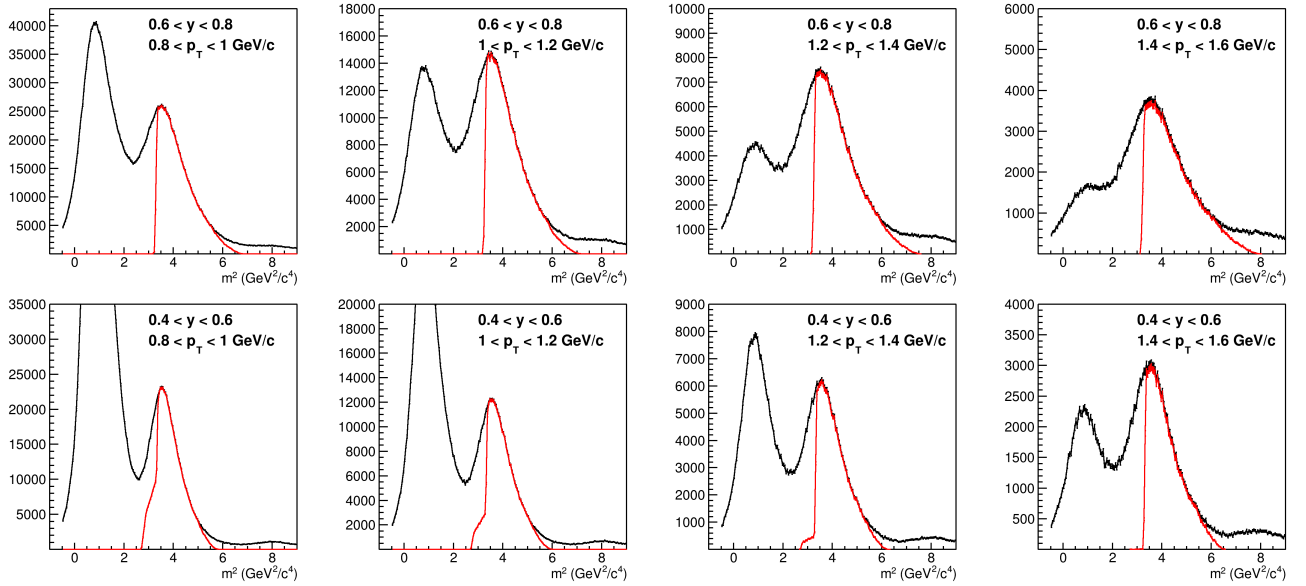


Figure 5: m^2 distribution of all particles (black line) detected in the TOF700 and of deuterons selected according to the n -sigma cuts (red line) for different y and p_T bins.

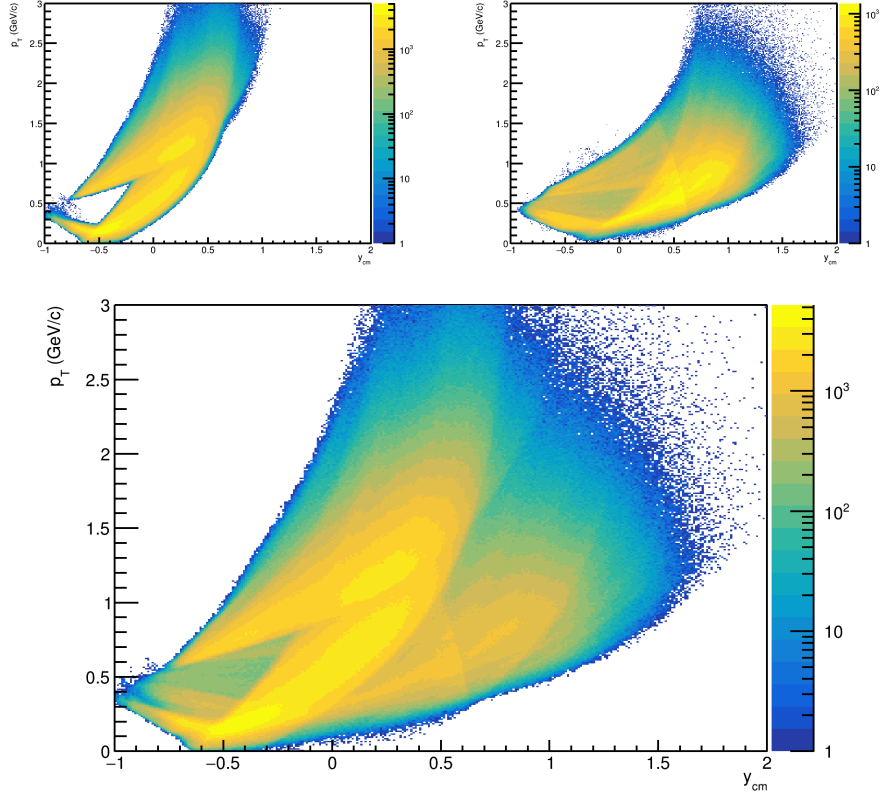


Figure 6: The phase space coverage of identified deuterons as a function of the centre-of-mass rapidity y_{cm} and transverse momentum p_T .

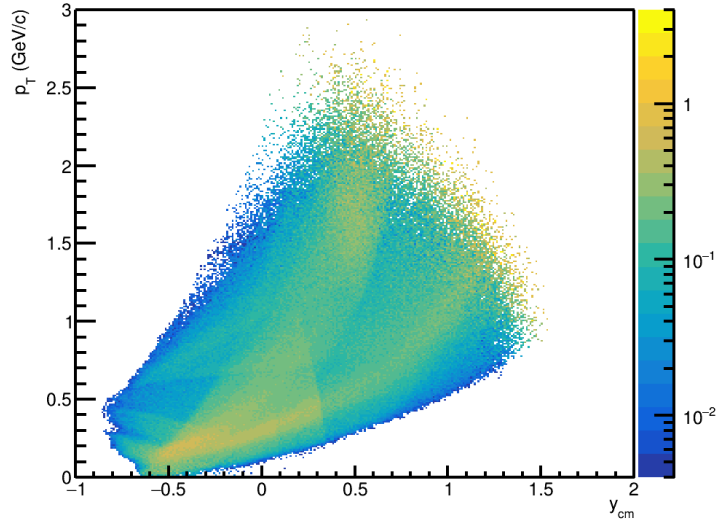


Figure 7: Efficiency of the deuteron reconstruction in the phase space of rapidity y_{cm} and transverse momentum p_T .

123 3 The analysis of v_1 of deuterons from BM@N run8 124 data

125 In this section, we discuss the details of analysis of directed flow v_1 of deuterons
126 in Xe+Cs(I) collisions at 3.8 AGeV using the BM@N run8 data. The symmetry
127 plane estimation and resolution correction factor calculation was performed identical
128 to how it was done for the analysis of directed flow of protons, so for details on
129 the u_n - Q_n -vector formalism, corrections on the azimuthal non-uniformity, BM@N
130 performance for the flow measurements, please consult the [10].

131 **1)** To address the effects of the non-uniform acceptance we applied the cor-
132 rections for both u_1 and Q_1 vectors :recentering, twist and rescaling. The QnTools
133 framework [13] was used for corrections of u_1 and Q_1 vectors and flow analysis. For
134 u_1 -vector corrections were employed multi-differentially on transverse p_T , rapidity
135 y and centrality. For the Q_1 -vectors corrections were applied only differentially on
136 centrality.

137 **2)** The detailed performance study [10] shows that due to magnetic field acting
138 along the y -axis and deflecting charged particles along the x axis, we can measure
139 the directed flow v_1 of protons using only the y components of flow vectors:

$$v_1 = 2 \frac{\langle y_1 Y_1^a \rangle}{R_1^y \{a\}}, \quad (4)$$

140 where the resolution correction factor is calculated using the method of three sub-
141 events:

$$R_1^y \{a(b, c)\} = \sqrt{\frac{\langle Y_1^a Y_1^b \rangle \langle Y_1^a Y_1^c \rangle}{\langle Y_1^b Y_1^c \rangle}}, \quad (5)$$

142 or by the four sub-event method:

$$R_1^y \{a(d)(b, c)\} = \langle Y_1^a Y_1^d \rangle \sqrt{\frac{\langle Y_1^d Y_1^b \rangle \langle Y_1^d Y_1^c \rangle}{\langle Y_1^b Y_1^c \rangle}}, \quad (6)$$

143 **3)** The Y_1 component of the Q_1 vector for symmetry planes in the FHCAL have
 144 been obtained using the formula:

$$Y_1 = \frac{\sum_{k=1}^M E_k \sin(\phi_k)}{\sum_{k=1}^M E_k} \quad (7)$$

145 , where E_k and ϕ_k is the energy deposition and azimuthal angle of the k -th module
 146 and M is the number of modules used.

147 Modules of the FHCAL were divided into three groups (sub-events): F1, F2, F3
 148 as it is shown in the Figure. 8. According to the simulations, due to charge splitting
 149 in the dipole analyzing magnet SP-41, F1 sub-event primarily registers the spectator
 protons, F2 — spectator fragments and F3 — neutrons.

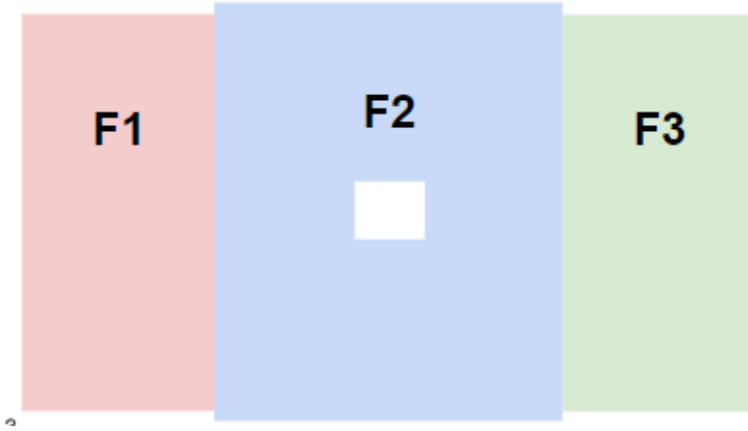


Figure 8: Layout of the FHCAL modules division into three groups (sub-events): F1, F2, F3.

150
 151 Two additional sub-events were introduced from the tracks of the charged
 152 particles in the inner tracking system of BM@N. All negatively charged particles with
 153 pseudorapidity $1.5 < \eta < 3$ and transverse momentum $p_T > 0.2$ GeV/c comprise the
 154 T- sub-event. The T+ sub-event consists of positively charged particles in following
 155 kinematic region: $2 < \eta < 3$ and $p_T > 0.2$ GeV/c.

156 Resolution correction factor was calculated for 3 spectator symmetry planes F1,
 157 F2, and F3 using the three sub-event method (for F2 four sub-event technique was
 158 employed as well) using the Equation 5 (and for four sub-events Equation 6). Figure 9
 159 shows the centrality dependence of the resolution correction factors R_1 for sub-event
 160 symmetry planes F1, F2 and F3 from left to right. For each symmetry plane R_1

161 was estimated using 3 combinations of sub-events (as indicated in the figure). One
 162 can observe that all three estimations for each symmetry plane are in reasonable
 163 agreement. This fact may suggest that the contribution of non-flow correlations in
 164 the final values of R_1 is very small. The R_1 for all three symmetry planes calculated
 165 using T - sub-event is off in the last centrality bin (30-40%) due to the fact that
 166 directed flow of negatively charged pions crosses zero at this centrality. We employ
 167 this estimation purely for the verification purpose and this outlier does not affect
 168 the estimated values for v_1 . For correction on the symmetry plane resolution we use
 $F1(F2, F3)$, $F2(F1, F3)$ and $F3(F1, F2)$ combinations.

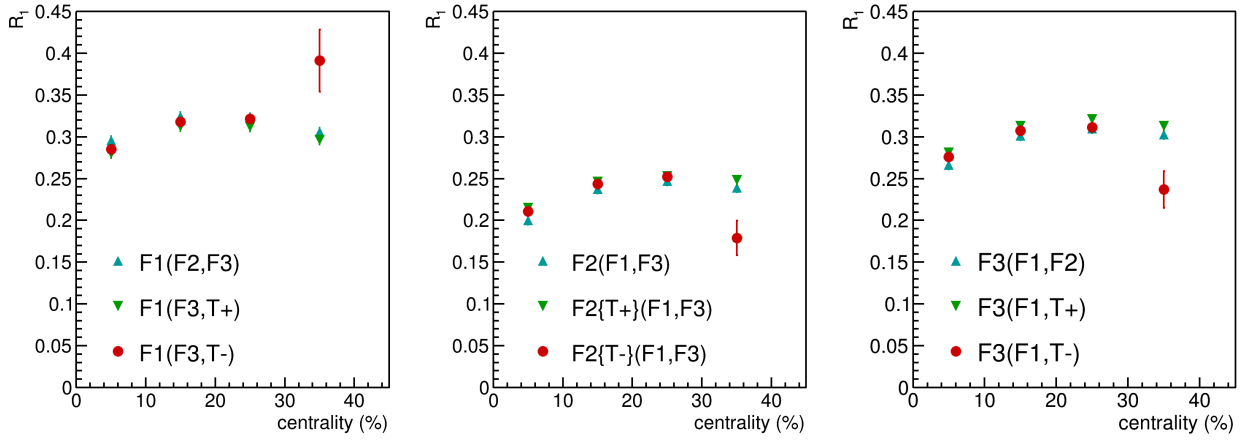


Figure 9: Resolution correction factor R_1 calculated using different combinations as a function of centrality for sub-event symmetry planes F1, F2 and F3 from left to right.

169
 170 4) Figure 10 shows the rapidity dependence y_{cm} of directed flow v_1 of deuterons
 171 in in 10-30% central Xe+Cs(I) collisions at 3.8 A GeV. The measurements have been
 172 performed with respect to the F1, F2, F3 and combined(F2+F3) symmetry planes.
 173 The resulting v_1 values of deuterons are in a good agreement for the measurements
 174 with respect to F2, F3 and combined(F2+F3) symmetry planes. The small difference
 175 in resulting v_1 values for the measurements with respect to the F1 plane, can be
 176 explained by the small contribution of non-flow effects. In order to get the final
 177 results, the measurements of directed flow v_1 have been performed with respect to
 178 the combined (F2+F3) symmetry plane, see Figure 17.

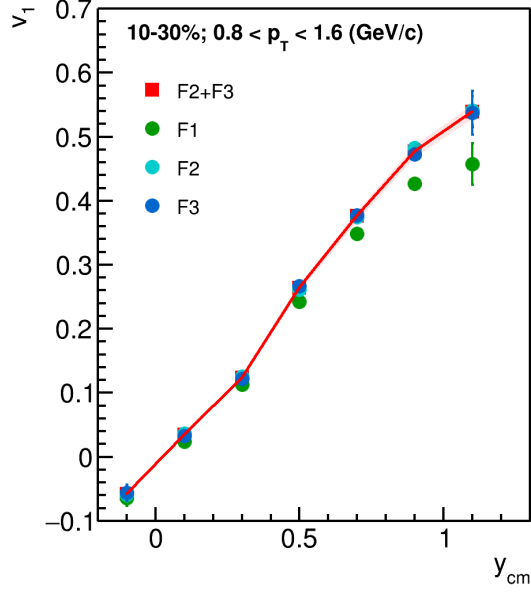


Figure 10: Directed flow v_1 of deuterons as a function of rapidity y_{cm} measured with respect to different spectator symmetry planes: F1, F2, F3 and combined (F2+F3), see text for the details.

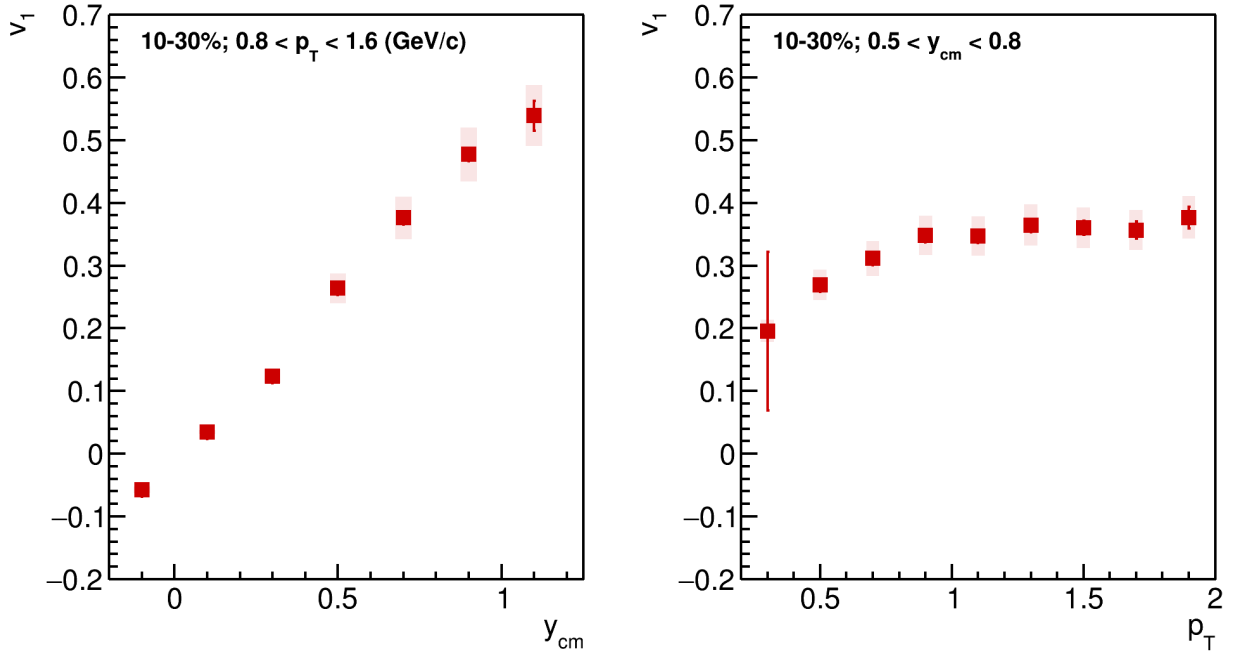


Figure 11: Directed flow v_1 of deuterons in 10-30% central Xe+Cs(I) collisions at 3.8 A GeV as a function of rapidity y_{cm} (left panel) and transverse momentum p_T (right panel).

179 4 Systematic uncertainties of v_1 measurements

180 In order to estimate systematic uncertainties of v_1 measurements , the following
181 sources were considered:

- 182 • Uncertainty in deuteron momentum reconstruction. We varied the number of
183 stations N_{hits} in inner tracking system used for track reconstruction as well
184 as the values of track χ^2/NDF quality, see results in Figure 12. The overall
185 systematic uncertainty is found to be below 5%.

- 186 • Contribution from the secondary particles. We studied the difference in v_1
187 results for tracks with different Distance of the Closes Approach (DCA) to the
188 primary vertex, see the left panel of Figure 13 for results. It is found that
189 deuteron v_1 values obtained with different DCA cut are in agreement within
190 1-2%.

- 191 • Contamination from the different particle species. We used stricter criteria for
192 deuteron selection in TOF700 , which are: $-2\sigma < (m^2 - \langle m^2 \rangle) < 2.5\sigma$ for
193 $0 < p/q < 4$ GeV/c; $-0.5\sigma < (m^2 - \langle m^2 \rangle) < 2.5\sigma$ for $4 < p/q < 6$ GeV/c;
194 $0.5\sigma < (m^2 - \langle m^2 \rangle) < 2.5\sigma$ for $6 < p/q < 8$ GeV/c. See the right panel of
195 Figure 13 for results. Observed is the systematic uncertainty below 5%

- 196 • Contribution due to off-target collisions. We divided the events based on the
197 azimuthal angle of the vertex position and compared the v_1 of deuterons in each
198 group of events, see results in Figure 14. Systematic variation stays below 3%.

- 199 • Acceptance and efficiency. We perform the v_1 flow measurements for deuterons
200 identified with TOF-400 and TOF-700 separately. We perform the measure-
201 ments with and without the applying the efficiency correction for deuterons
202 based on MC simulations for run8, see Figure 15 for results. The results are
203 in a good agreement and we can conclude that the mean value of transverse
204 momentum p_T is not shifted in this rapidity range.

- 205 • Run-by-run systematics was estimated dividing the events into several run
206 periods and comparing the results in each group, see the left panel of Figure 16

207 for results. The systematic uncertainty is less than 5% and found to be less
 208 than statistical.

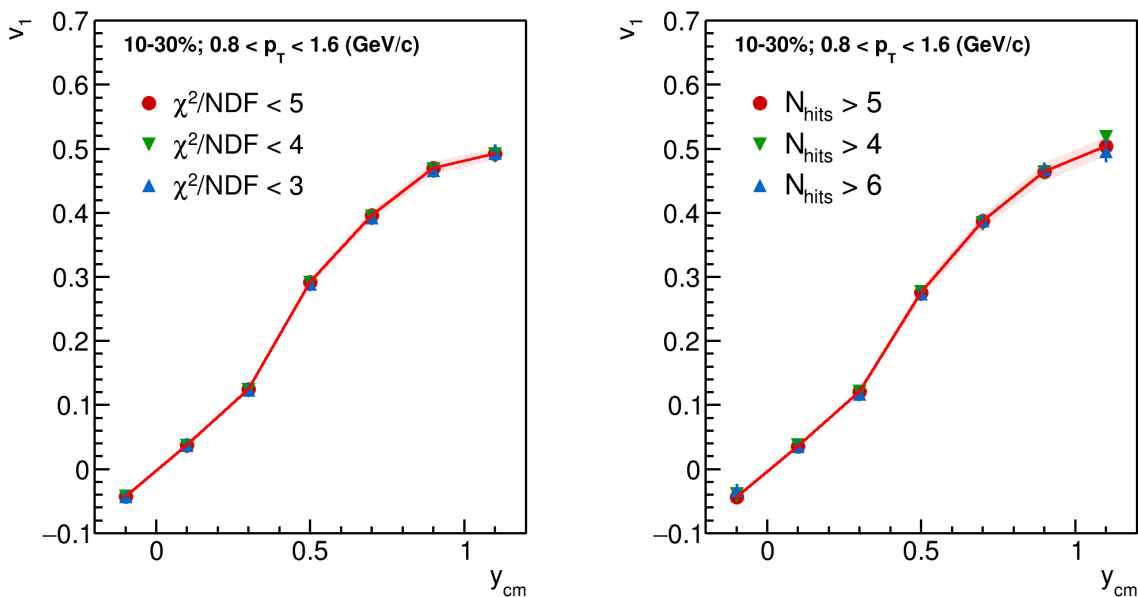


Figure 12: Directed flow v_1 of deuterons as a function of rapidity y_{cm} measured for different values of the track χ^2/NDF quality (left) and the number of stations used for track reconstruction N_{hits} (right).

209 Systematic uncertainties were calculated by the square root of quadratic sum
 210 of uncertainties from each source.

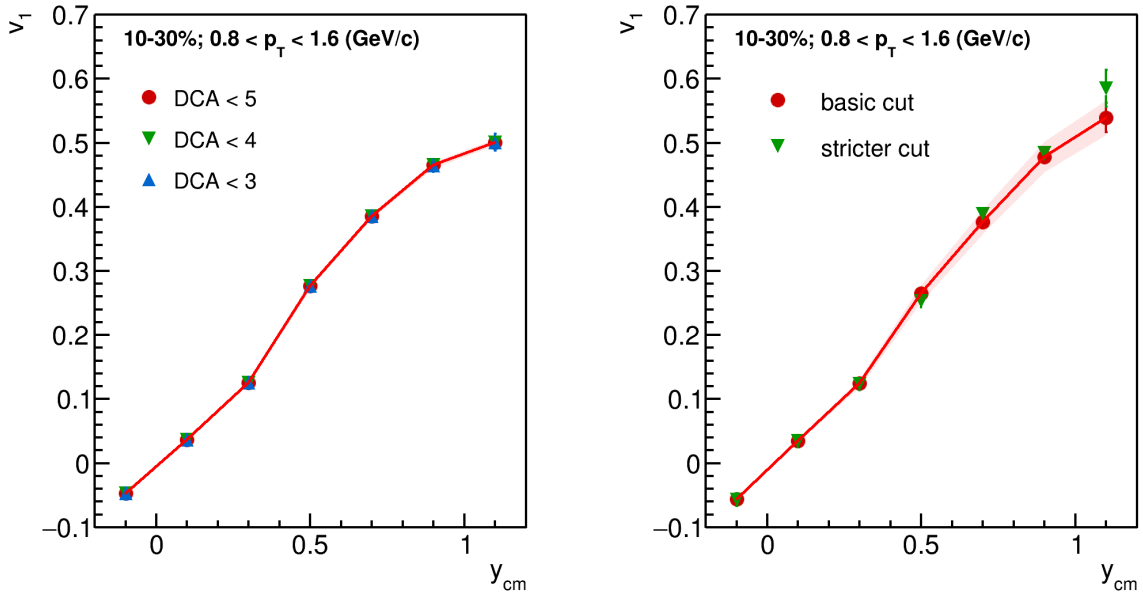


Figure 13: Directed flow v_1 of deuterons as a function of rapidity y_{cm} measured for different values of the DCA cut and different $n\text{-}\sigma$ PID cuts for the deuteron identification: $(m^2 - \langle m_p^2 \rangle) < 1, 2, 3 \sigma_{m_p^2}$ cut (right).

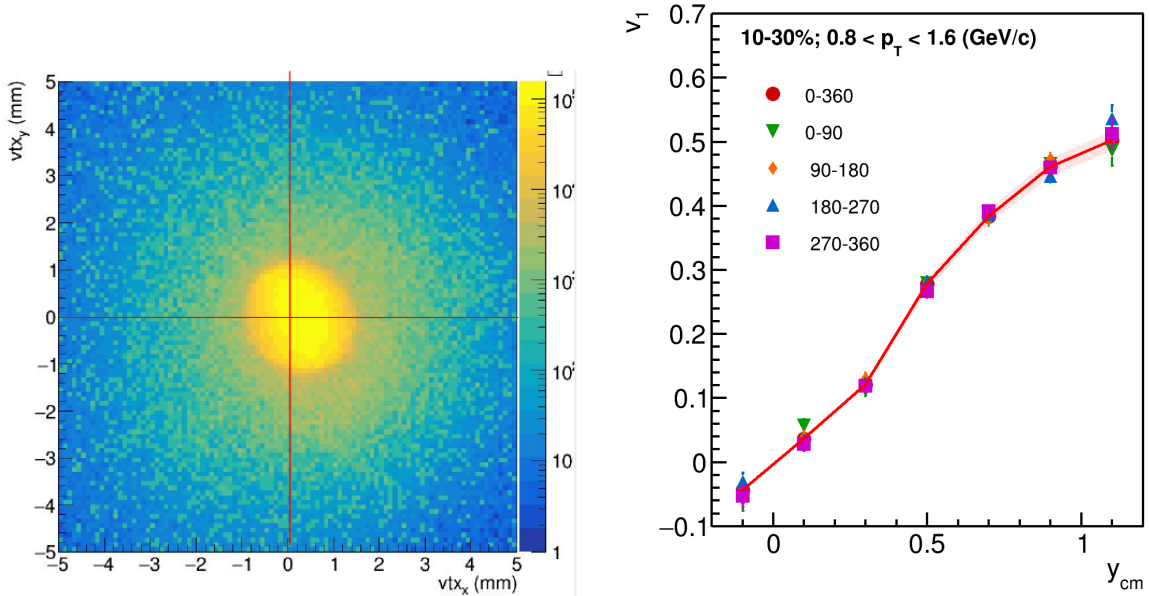


Figure 14: Left: the distribution of the primary vertex in X-Y plane. Right: Directed flow v_1 of deuterons as a function of rapidity y_{cm} calculated with varying the reconstructed primary vertex position of the collision.

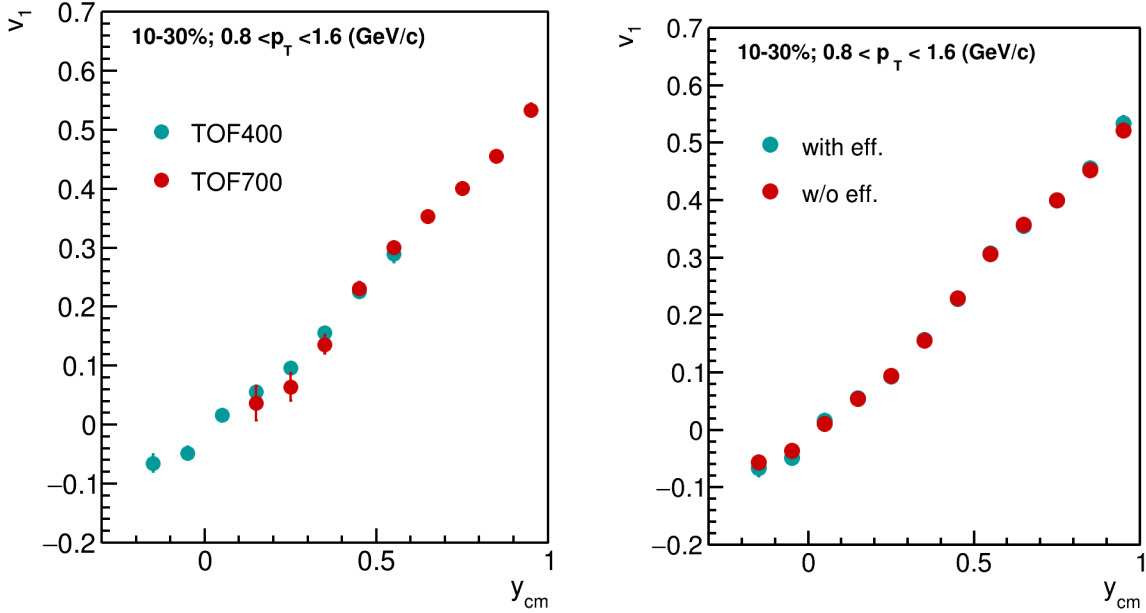


Figure 15: Directed flow v_1 of deuterons as a function of rapidity y_{cm} measured for deuterons identified using different TOF-systems (left) and deuterons weighted and not weighted with efficiency based on MC simulations for run8 (right).

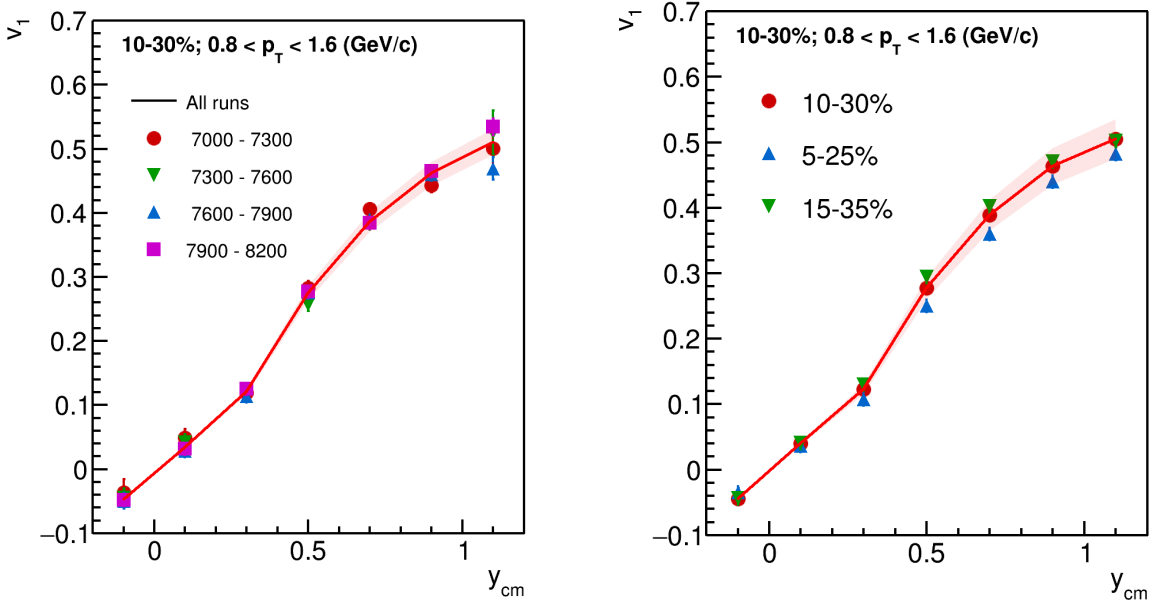


Figure 16: Directed flow v_1 of deuterons as a function of rapidity y_{cm} measured in the different run periods (left) and for different bins in collision centrality (right).

211 5 Results of the directed flow measurements

212 Directed flow v_1 of deuterons was measured in 10-30% central Xe+Cs(I) col-
 213 lisions at 3.8 A GeV as a function of rapidity y_{cm} and transverse momentum p_T ,
 214 see Figure. 17. Rapidity-dependence and transverse-momentum-dependence of v_1 of
 215 deuterons from the experimental data has been compared with proton v_1 obtained
 216 for the same data [10]. The directed flow of deuterons is greater that that of protons
 as was expected.

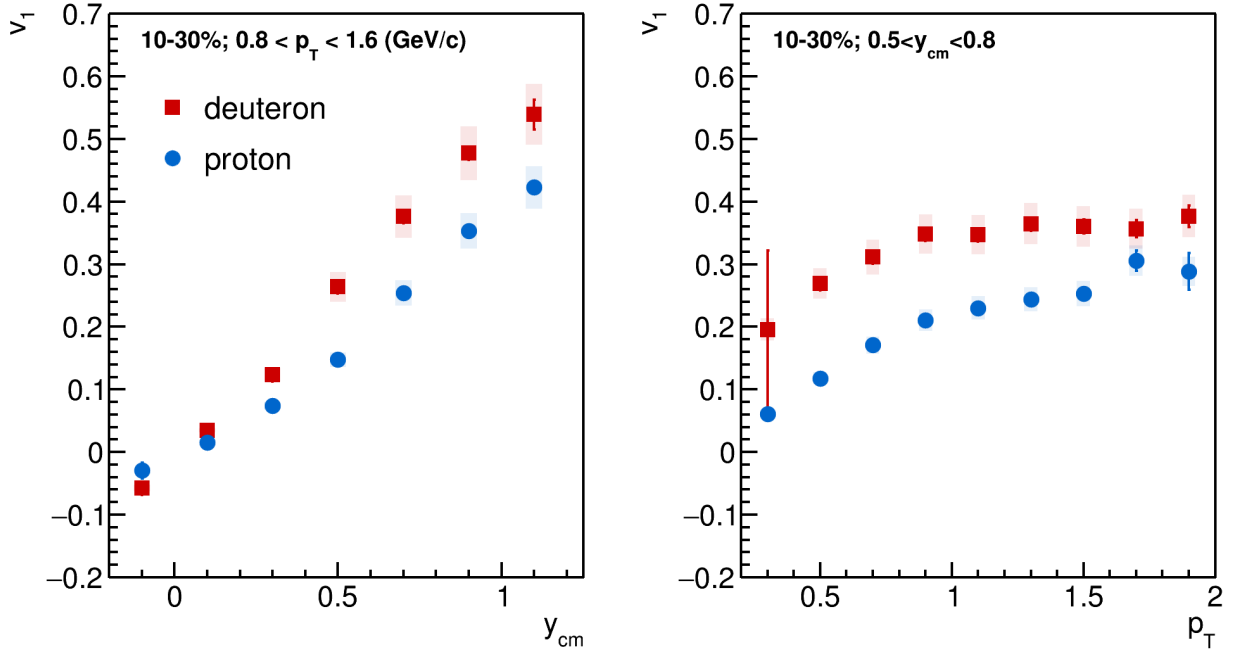


Figure 17: Directed flow v_1 of deuterons in 10-30% central Xe+Cs(I) collisions at 3.8 A GeV as a function of rapidity y_{cm} (left panel) and transverse momentum p_T (right panel).

217 The coalescence scaling was studied similar to how it was done by the STAR
 218 experiment [8]. The directed flow of light nuclei is divided over a respective mass
 219 number A together with the transverse momentum. The underlying assumption is
 220 that nucleons with similar momenta can strongly bind producing the corresponding
 221 light nuclei such as deuterons, tritons, helium-3 and etc. Directed flow of protons
 222 and deuterons before (upper plots) and after (lower plots) applying the scaling in
 223 different rapidity bins is shown in fig. 18. We observe that v_1 of light nuclei follow
 224 the scaling in selected rapidity ranges.
 225

226 The slope of the directed flow v_1 at midrapidity $dv_1/dy_{cm}|_{y_{cm}=0}$ is extracted
227 by fitting the $v_1(y_{cm})$ with polynomial function $v_1 = a + by_{cm} + cy_{cm}^3$ as it was
228 done in other experiments [14–17]. The slope of v_1 of deuterons at midrapidity
229 $dv_1/dy_{cm}|_{y_{cm}=0}$ as a function of collision energy is presented in the fig. 19. The results
230 for the BM@N experiment are compared with existing data from other experiments
231 [15–17]. Directed flow slope at midrapidity $dv_1/dy_{cm}|_{y_{cm}=0}$ are found to be in a
232 reasonable agreement with the existing measurements.

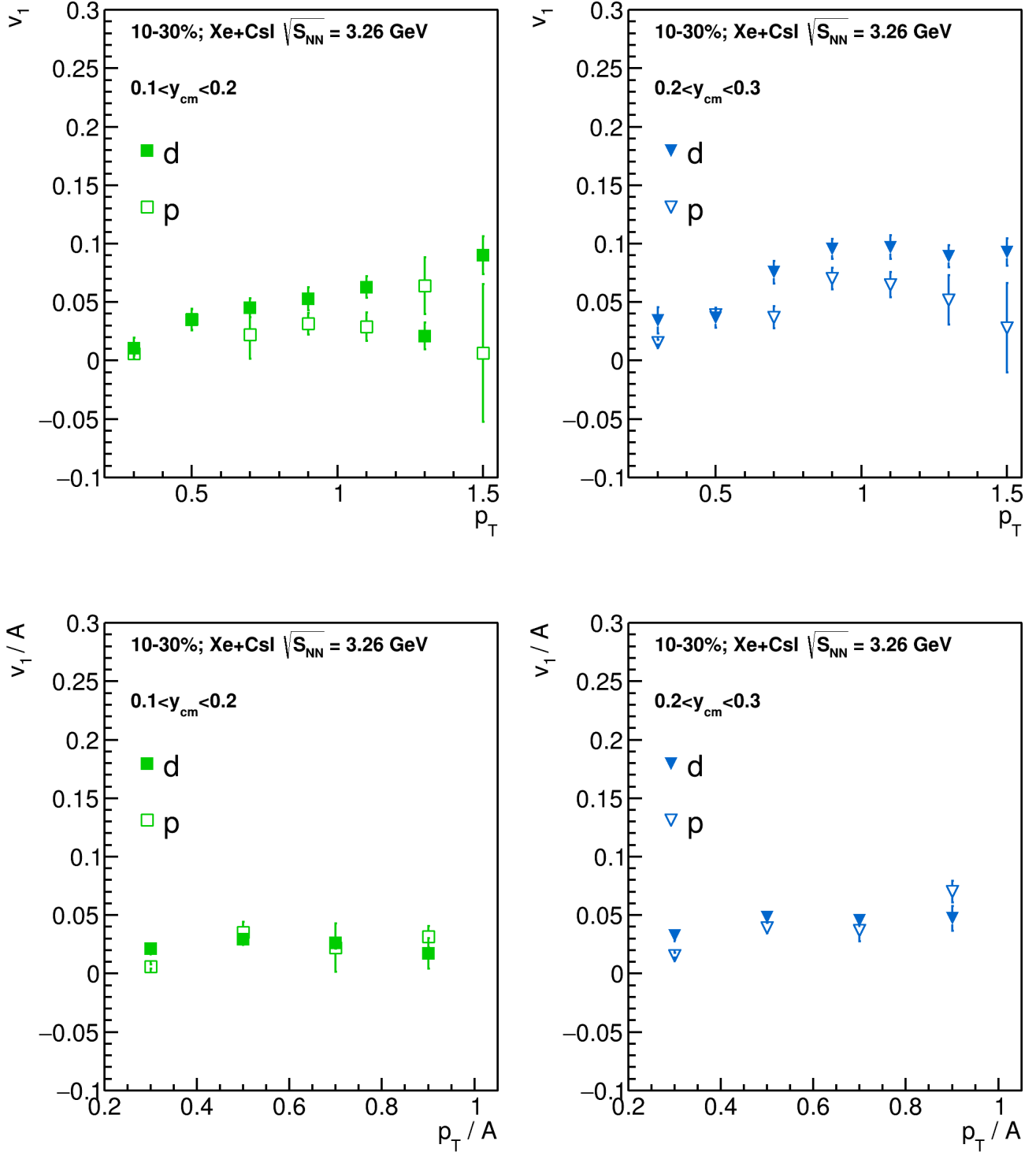


Figure 18: Upper plots: directed flow v_1 of protons and deuterons in 10-30% central Xe+Cs(I) collisions at 3.8 A GeV as a function of transverse momentum p_T . Lower plots: directed flow scaled on the corresponding mass number v_1/A for protons and deuterons in 10-30% central Xe+Cs(I) collisions at 3.8 A GeV as a function of scaled transverse momentum p_T/A .

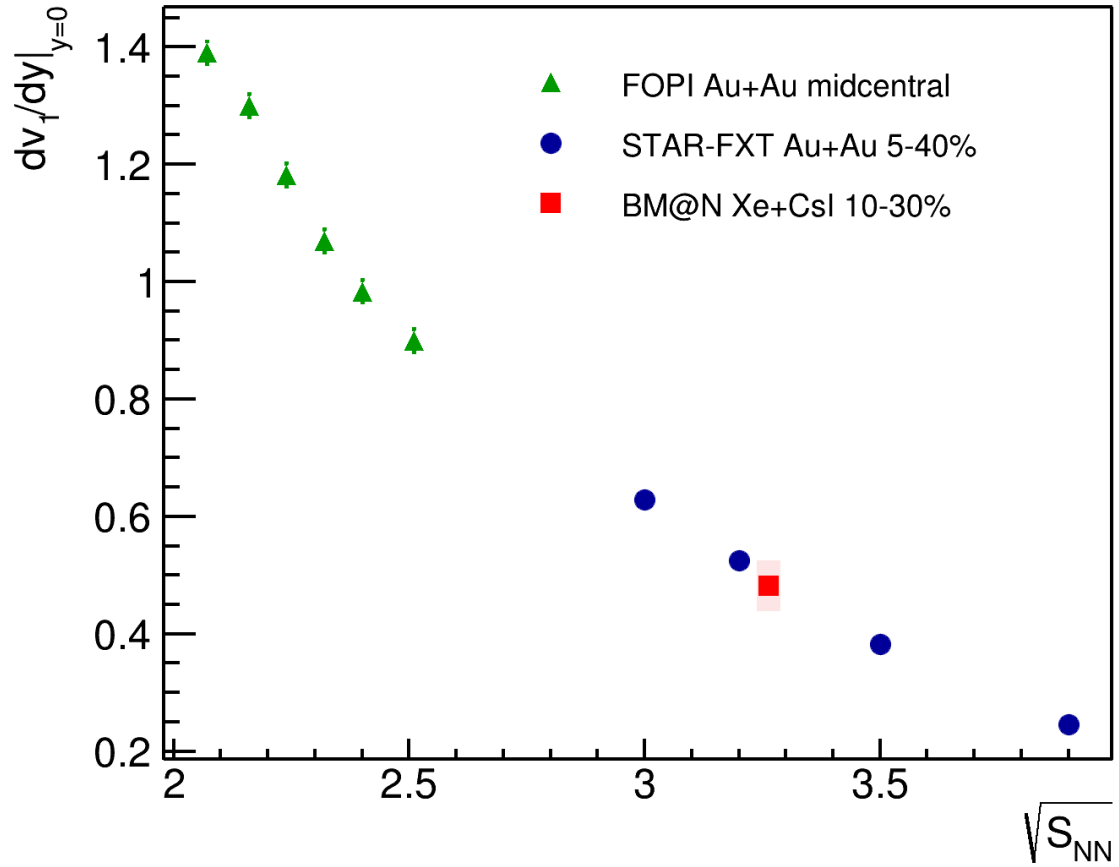


Figure 19: The slope of v_1 of deuterons at midrapidity $dv_1/dy_{cm}|_{y_{cm}=0}$ as a function of collision energy. The obtained BM@N results were compared with existing data from other experiments [15–17].

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