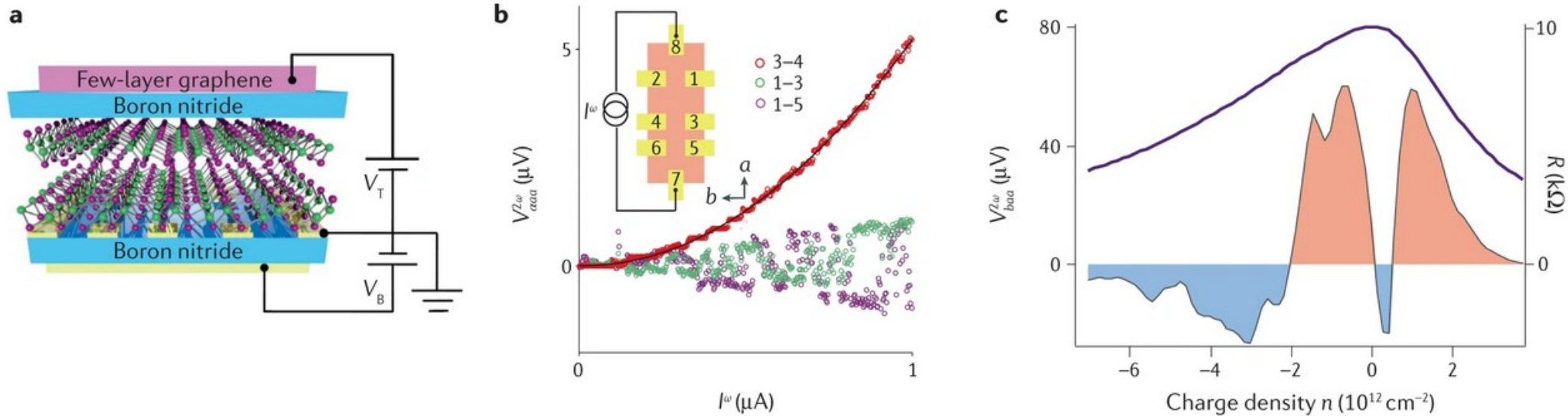


Non-Linear Hall effect

- The nonlinear Hall effect shows a tranverse voltage of $V_{2\omega}$ for an applied current of $I_{2\omega}$
- Unlike usual (linear) Hall effects time reversal symmetry is not broken in this case.



Orbital Magnetization

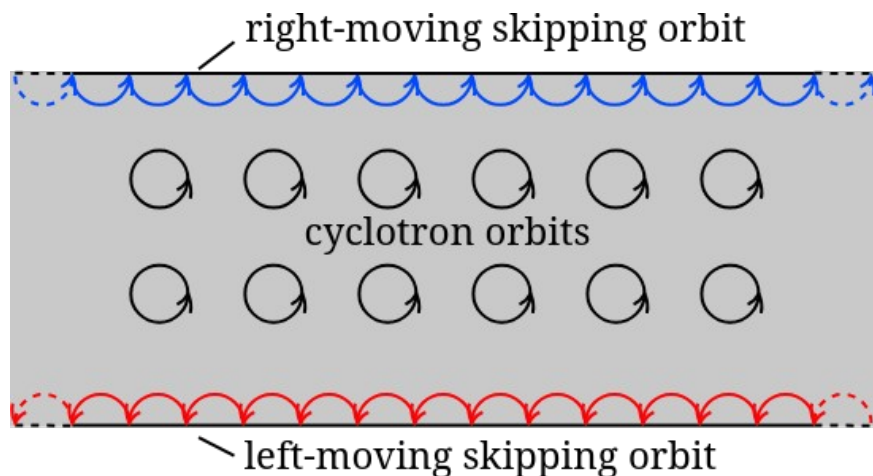
- Orbital magnetization is part of the magnetic moment due to orbital motion of electrons.
- In condensed matter theory it is correlated to the Berry curvatures.

$$\mathbf{M}_{\text{orb}} = \frac{e}{2\hbar c} \text{Im} \sum_n \int \frac{d^3k}{(2\pi)^3} f_{n\mathbf{k}} \langle \partial_{\mathbf{k}} u_{n\mathbf{k}} | \times (H_{\mathbf{k}} + E_{n\mathbf{k}} - 2\mu) | \partial_{\mathbf{k}} u_{n\mathbf{k}} \rangle ,$$

Orbital Magnetization due to Berry curvature

- Orbital magnetization due to chiral edge states

$$M = \frac{e}{\hbar} \sum_a \int_{\mathbf{k}} \Omega_{a,\mathbf{k}} \frac{1}{\beta} \log \left[1 + e^{-\beta(E_{a,\mathbf{k}} - \mu)} \right],$$



$$\Omega_n(\mathbf{k}) = -\text{Im} \langle \nabla_{\mathbf{k}} u_{n\mathbf{k}} | \times | \nabla_{\mathbf{k}} u_{n\mathbf{k}} \rangle$$

Relation between orbital magnetization and Hall current density

- The Hall current density is defined as

$$J_{H,i} = \epsilon_{ijk} \frac{\partial M_k}{\partial r_j} = \epsilon_{ijk} \frac{\partial \mu}{\partial r_j} \frac{\partial M_k}{\partial \mu}.$$

- Electric field is the gradient of chemical potential

$$\mathbf{E} = -\frac{1}{e} \partial_{\mathbf{r}} \mu,$$

Relation between orbital magnetization and Hall current density

- The Hall current density is:

$$J_{H,i} = \sigma_{ij}^H E_j$$

- The conductivity is:

$$\sigma_{ij}^H = -\frac{e^2}{\hbar} \sum_a \int_{\mathbf{k}} \epsilon_{ijk} \Omega_{a,\mathbf{k}}^k f(E_{a,\mathbf{k}}) d\mathbf{k},$$

Orbital magnetization due to perturbation

- For the case of magnetoelectric case, the perturbation of magnetization gives

$$M_i = \chi_{ij} E_j.$$

$$\chi_{ij} = -\frac{e^2}{\hbar} \sum_a \int_{\mathbf{k}} \left(\frac{\tau}{\hbar} \Omega_{a,\mathbf{k}}^i v_{a,\mathbf{k}}^j + 2\epsilon_{ikq} G_{a,\mathbf{k}}^{jk} v_{a,\mathbf{k}}^q \right) f(E_{a,\mathbf{k}})$$

- Substituting perturbed magnetization back into hall conductivity we get

$$J_{H,i} = \sigma_{ijk}^H E_j E_k,$$

$$\begin{aligned} \sigma_{ijk}^H = & -\frac{e^3}{\hbar} \sum_a \int_{\mathbf{k}} \left[\frac{\tau}{2\hbar} \epsilon_{iqk} \partial_{k_j} \Omega_{a,\mathbf{k}}^q + \left(\partial_{k_i} G_{a,\mathbf{k}}^{jk} - \partial_{k_j} G_{a,\mathbf{k}}^{ik} \right) \right. \\ & \left. + (j \leftrightarrow k) \right] f(E_{a,\mathbf{k}}). \end{aligned} \quad (5)$$