

Colliders and NICA

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Outline

- HEP: Directions, Status and Plans
- Colliders
- Some Basic Concepts of Accelerator Physics
- Present and Future Colliders
- Some Important Accelerator Technologies
- NICA – the First Hadron Collider in Russia

International school of
accelerator physics:
Linear accelerators,
Verbilki, Moscow. Reg.
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HEP: Directions, Status and Plans

Major Questions in High Energy Physics

- The Standard Model doesn't explain:
 - ◆ why the Higgs boson exists.
 - ◆ why the Higgs boson has the mass that it does. The Higgs turned out to be much less massive than predicted (a quadrillion times)
- We did not find a way how to add gravity to the Standard model.
- Where did all the antimatter go after the big bang?
 - ◆ Known CP violation does not look sufficient
- Why lepton number is not conserved in neutrino oscillations?
 - ◆ Neutrino mass = 0 in standard model
- Is neutrino and antineutrino the same particle?
- What are the dark matter and dark energy?
 - ◆ WIMPs, axions
- Does the supersymmetry exist?
- ...
- LHC has been dominating the high energy physics for more than a decade; and it will continue to dominate in feasible future

Major Goals & Directions of High Energy Physics

The goal is to determine the most fundamental building blocks of matter and to understand the interactions between these particles

- Presently the high energy physics extends far beyond the accelerator based HEP
 - ◆ Cosmology & Astrophysics
 - How our universe was created and which laws determine its expansion?
 - High resolution digital map of the universe observed at different wave lengths
 - Gravitational waves
 - ◆ Detection of radiation coming from space
 - ν - Discovery of neutrino mass difference (Nobel prize)
 - ν - Highest observed energy 3×10^{20} eV (far beyond our accelerators)
 - γ - Microwave background radiation
 - γ - Search for dark matter and dark energy
 - e^\pm , p^\pm , nuclei, ...
- The strongest limitation on the neutrino mass came from analysis of the microwave background radiation and universe expansion

Major Directions of Accelerator Based HEP

- LHC is a leader in collider-based physics
 - ◆ Finding physics beyond the standard model
 - Higher energy and luminosity
 - Improvement of detector resolution (space and time)
 - Detailed measurements of Higgs boson
 - FCC is planned as the next step of the CERN program
- Mixing and oscillations of neutrinos
 - ◆ New generation machines come at the end of this decade
 - DUNE (Fermilab) and
Hyper-Kamiokande (Institute for Cosmic Ray Research of University of Tokyo & JPARC & KEK)
 - ◆ The goal: finding CP violation in neutrino sector and neutrino mass hierarchy
- Physics beyond standard model
 - ◆ g-2 experiment (Fermilab) – difference disappears with more data
 - ◆ μ -to-e experiment (Fermilab, the Paul Scherrer Institute (Zurich))
 - lepton number violation is observed in neutrino sector

Colliders

Collision Energy and Luminosity

■ Collision energy

- ◆ Gain in collision energy for ultra-relativistic particles
- ◆ One particle stationary:

$$E_{cm} = \frac{1}{2} E \xrightarrow[\text{case, } E \gg mc^2]{\text{ultra relativistic}} E_{cm} \approx \sqrt{2Emc^2}$$

- ◆ Both particles move:

$$E_{cm} = 2E$$

(120 times gain for the LHC)

■ Luminosity

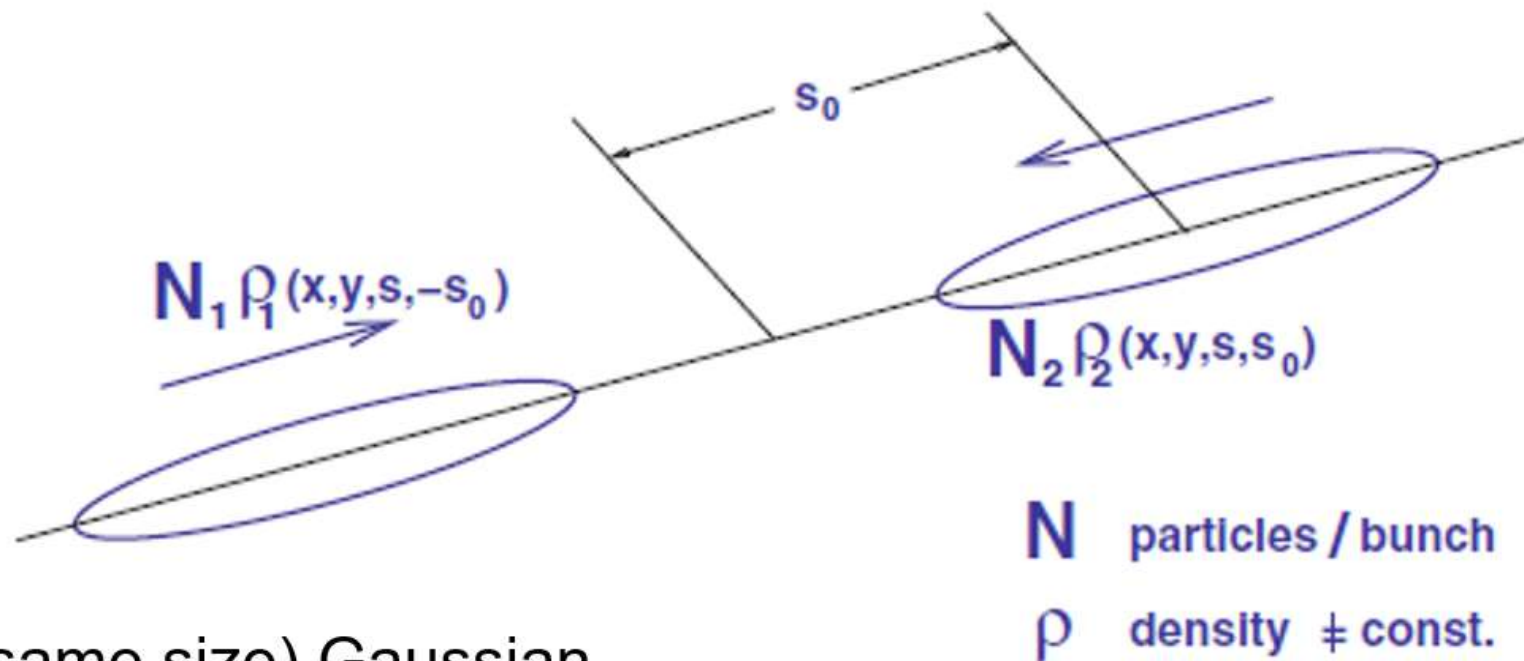
- ◆ Number of events in collisions:

$$\frac{dN}{dt} = L\sigma$$

- The total cross section for Higgs boson production at the LHC operating at **s=13 TeV** is 43 pb = $5 \cdot 10^{-35} \text{ cm}^2$.
⇒ At luminosity of $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ the LHC makes 1 Higgs every 2 s
- ◆ Higgs discovery potential: Tevatron versus LHC: $(E/E)^4(L/L) = 6^4 30 \approx 4 \cdot 10^4$
- ◆ Particle physics detectors want constant luminosity!

Luminosity

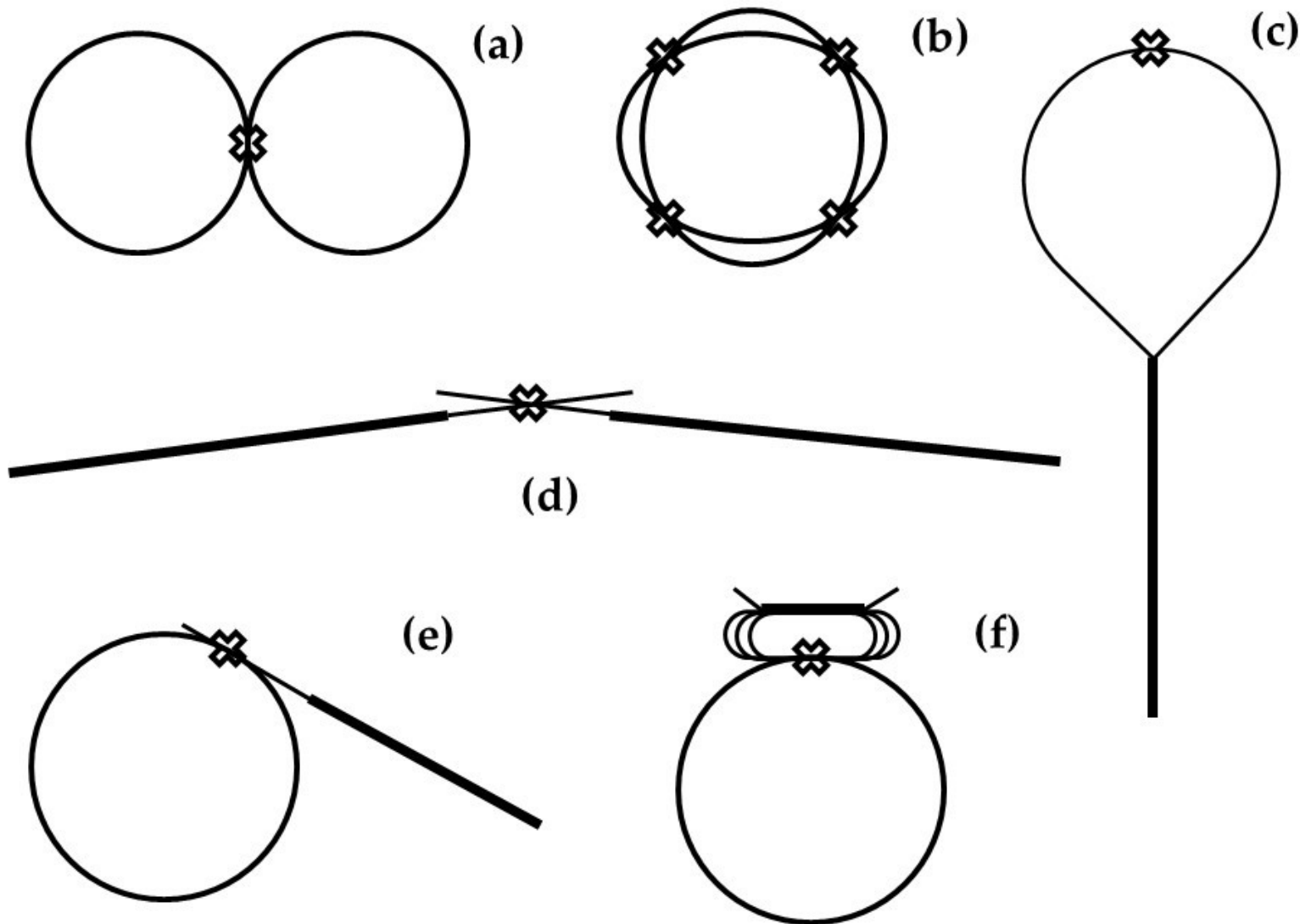
$$N_{\text{exp}} = \sigma_{\text{exp}} \cdot \int \mathcal{L}(t) dt.$$



For (same size) Gaussian bunches:

$$\mathcal{L} = f_{\text{coll}} \frac{N_1 N_2}{4\pi \sigma_x^* \sigma_y^*}$$

Types of Colliding Beams Facilities



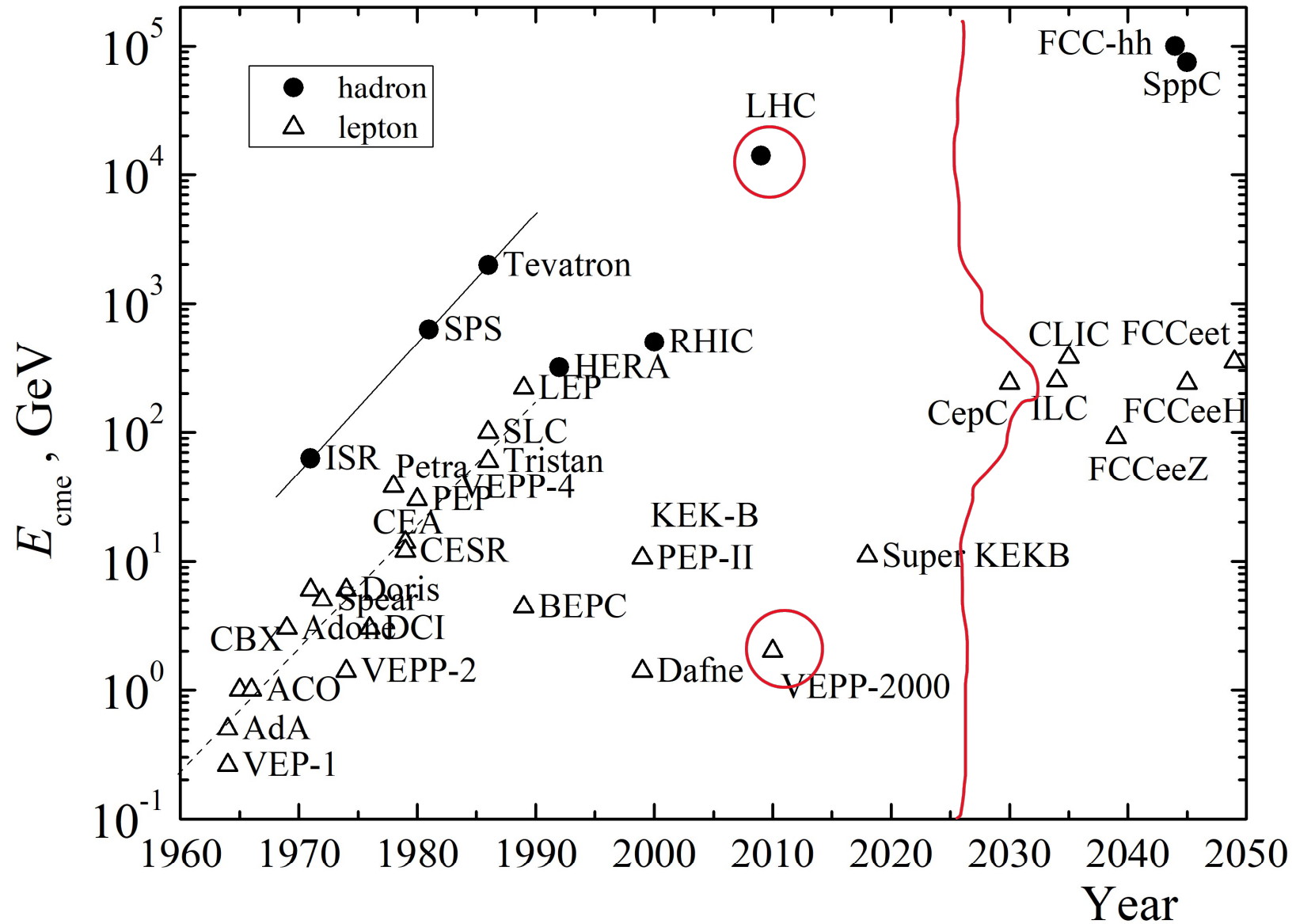
■ Since 60's colliders have been the major instrument in the particle physics

Colliders Landscape

- 61 years since 1st collisions
 - ◆ Spring 1964 AdA and VEP-1
- 31 operated since
- 7 in operations now
 - ◆ S-KEKB, VEPP-2000, VEPP-4M, BEPC, DAFNE
 - ◆ LHC, RHIC
- 1 under construction
 - ◆ NICA (JINR)
- One in a project phase
 - ◆ EIC (BNL)
- Far plans
 - ◆ Higgs/Electroweak factories
 - ILC
 - FCC: e^+e^-
 - ◆ Frontier ($E \gg E_{\text{LHC}}$)
 - FCC: pp

| | Species | E_b , GeV | C , m | \mathcal{L}_{peak}^{max} | Years |
|--------------|------------|-------------|---------|----------------------------|-----------|
| AdA | e^+e^- | 0.25 | 4.1 | 10^{25} | 1964 |
| VEP-1 | e^-e^- | 0.16 | 2.7 | 5×10^{27} | 1964-68 |
| CBX | e^-e^- | 0.5 | 11.8 | 2×10^{28} | 1965-68 |
| VEPP-2 | e^+e^- | 0.67 | 11.5 | 4×10^{28} | 1966-70 |
| ACO | e^+e^- | 0.54 | 22 | 10^{29} | 1967-72 |
| ADONE | e^+e^- | 1.5 | 105 | 6×10^{29} | 1969-93 |
| CEA | e^+e^- | 3.0 | 226 | 0.8×10^{28} | 1971-73 |
| ISR | pp | 31.4 | 943 | 1.4×10^{32} | 1971-80 |
| SPEAR | e^+e^- | 4.2 | 234 | 1.2×10^{31} | 1972-90 |
| DORIS | e^+e^- | 5.6 | 289 | 3.3×10^{31} | 1973-93 |
| VEPP-2M | e^+e^- | 0.7 | 18 | 5×10^{30} | 1974-2000 |
| VEPP-3 | e^+e^- | 1.55 | 74 | 2×10^{27} | 1974-75 |
| DCI | e^+e^- | 1.8 | 94.6 | 2×10^{30} | 1977-84 |
| PETRA | e^+e^- | 23.4 | 2304 | 2.4×10^{31} | 1978-86 |
| CESR | e^+e^- | 6 | 768 | 1.3×10^{33} | 1979-2008 |
| PEP | e^+e^- | 15 | 2200 | 6×10^{31} | 1980-90 |
| $Sp\bar{p}S$ | $p\bar{p}$ | 455 | 6911 | 6×10^{30} | 1981-90 |
| TRISTAN | e^+e^- | 32 | 3018 | 4×10^{31} | 1987-95 |
| Tevatron | $p\bar{p}$ | 980 | 6283 | 4.3×10^{32} | 1987-2011 |
| SLC | e^+e^- | 50 | 2920 | 2.5×10^{30} | 1989-98 |
| LEP | e^+e^- | 104.6 | 26659 | 10^{32} | 1989-2000 |
| HERA | ep | 30+920 | 6336 | 7.5×10^{31} | 1992-2007 |
| PEP-II | e^+e^- | 3.1+9 | 2200 | 1.2×10^{34} | 1999-2008 |
| KEKB | e^+e^- | 3.5+8.0 | 3016 | 2.1×10^{34} | 1999-2010 |
| VEPP-4M | e^+e^- | 6 | 366 | 2×10^{31} | 1979- |
| BEPC-I/II | e^+e^- | 2.3 | 238 | 10^{33} | 1989- |
| DAΦNE | e^+e^- | 0.51 | 98 | 4.5×10^{32} | 1997- |
| RHIC | p, i | 255 | 3834 | 2.5×10^{32} | 2000- |
| LHC | p, i | 6500 | 26659 | 2.1×10^{34} | 2009- |
| VEPP2000 | e^+e^- | 1.0 | 24 | 4×10^{31} | 2010- |
| S-KEKB | e^+e^- | 7+4 | 3016 | $8 \times 10^{35} *$ | 2018- |

Colliders: Energy



Colliders: Luminosity

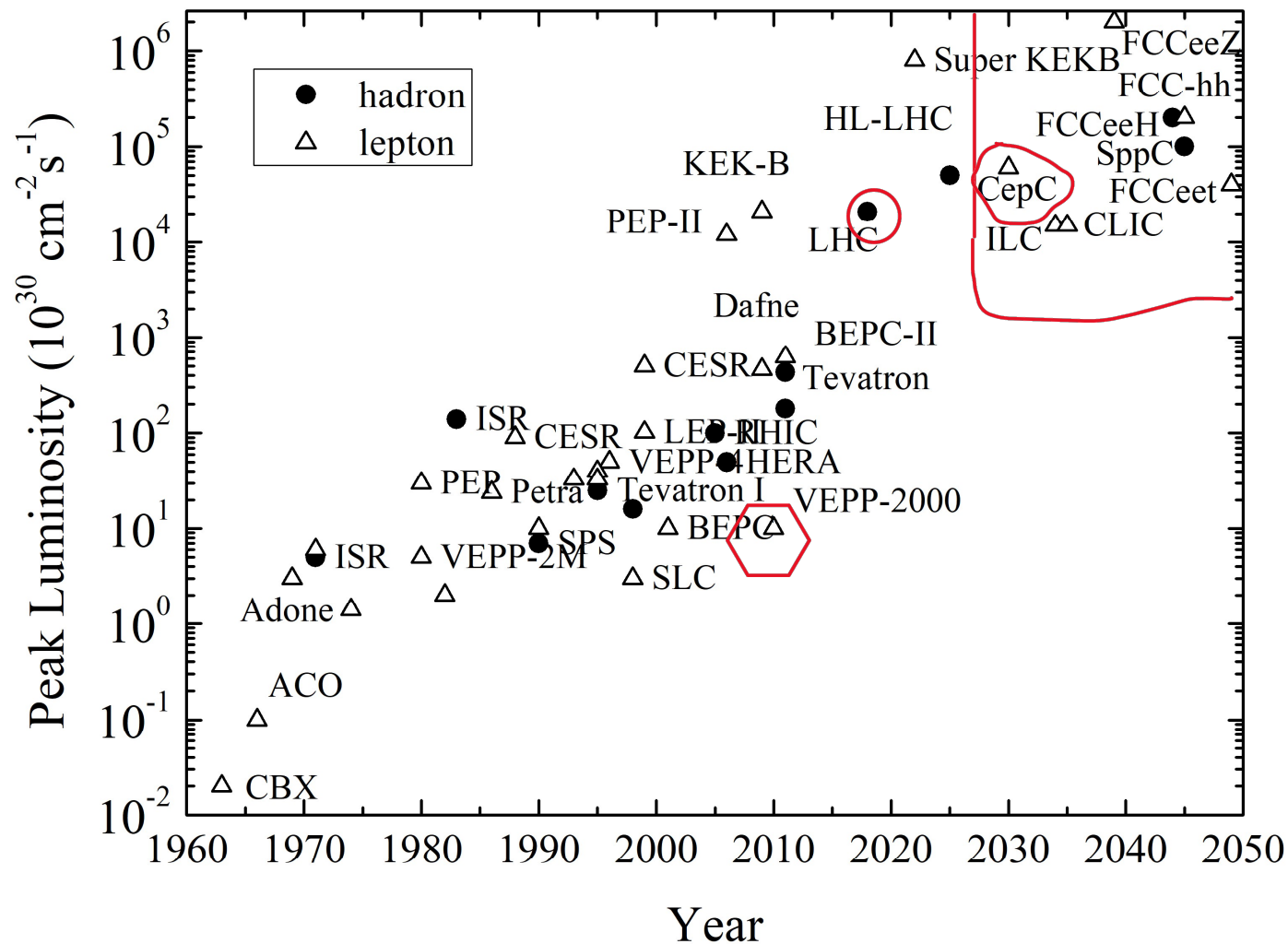
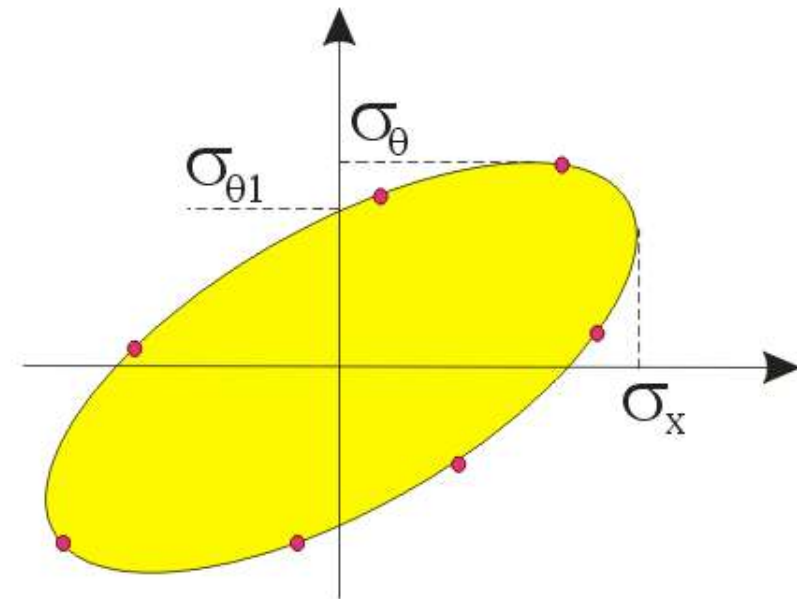


FIG. 3. Luminosities of particle colliders (triangles are lepton colliders and full circles are hadron colliders, adapted from [37]). Values are per collision point.

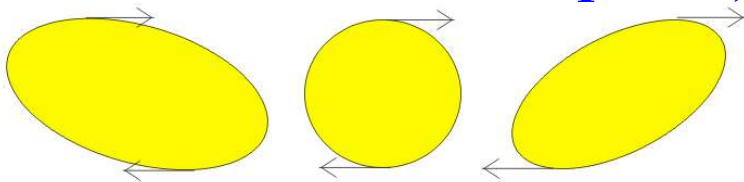
Some Basic Concepts of Accelerator Physics

Emittance

- Two sides of the emittance concept
 - ◆ Liouville theorem
 - ◆ Action - Single particle emittance
- As a particle returns to the same point on subsequent revolutions, it will map out an ellipse in the phase space



- Emittance = $\sigma_x \sigma_\theta$
(σ_θ – local momentum spread)



- Normalized emittance:
 $\varepsilon_n = \varepsilon \gamma \beta$ - adiabatic invariant

- Beam size:

$$\sigma_{x,y} = \sqrt{\frac{\varepsilon_n \cdot \beta_{x,y}}{\gamma}}$$

- Luminosity $\sim 1/\varepsilon$

$$\varepsilon = \sigma_x \sigma_{\theta 1}$$

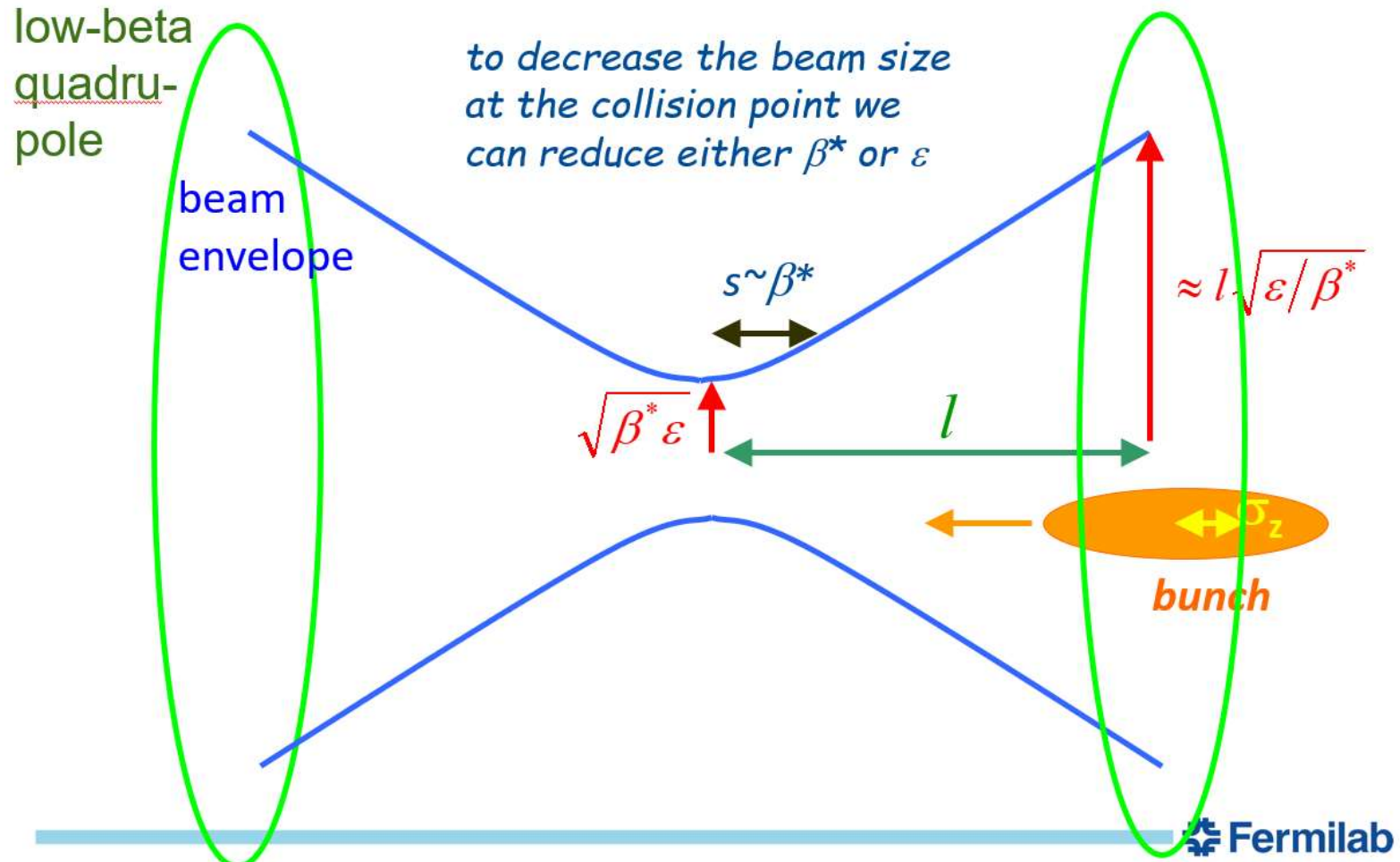
$$\sigma_x = \sqrt{\varepsilon \beta_x}$$

$$\sigma_{\theta 1} = \sqrt{\frac{\varepsilon}{\beta_x}}$$

$$\sigma_\theta = \sqrt{\frac{\varepsilon (1 + \alpha_x^2)}{\beta_x}}$$

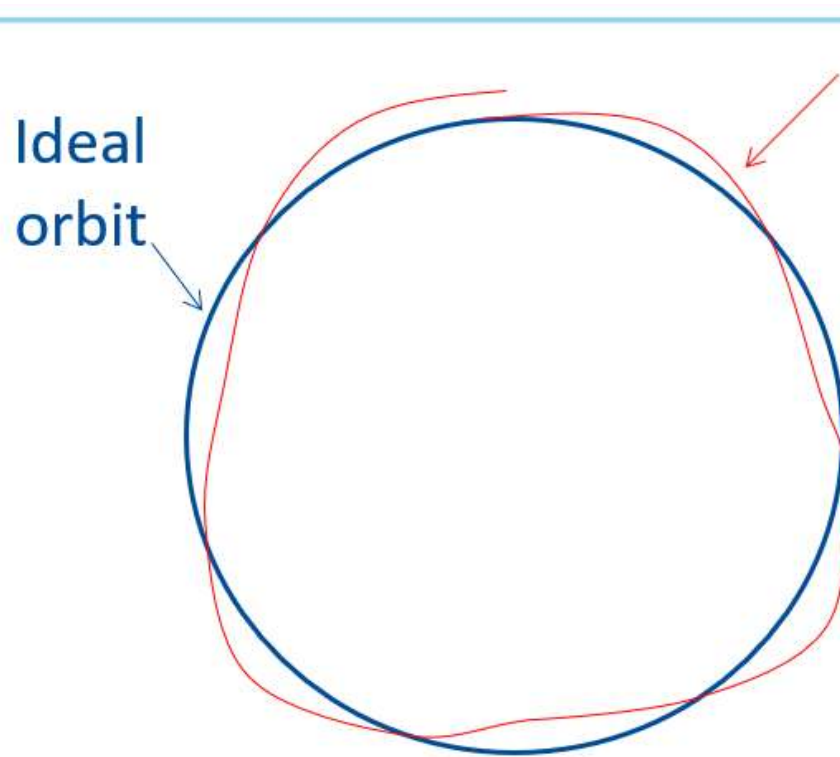
$$\alpha_x = -\frac{1}{2} \frac{d\beta}{ds}$$

Collider Spot Size



- β^* must be equal or larger than σ_z ('hourglass effect')
 - ◆ with exception of crab-waist (e⁺e⁻ colliders)
- Quadrupole aperture must be respected

Betatron Oscillations, Tune



Particle trajectory

- As particles go around a ring, they will undergo a number of betatron oscillations ν (sometimes Q) given by

$$\nu = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$

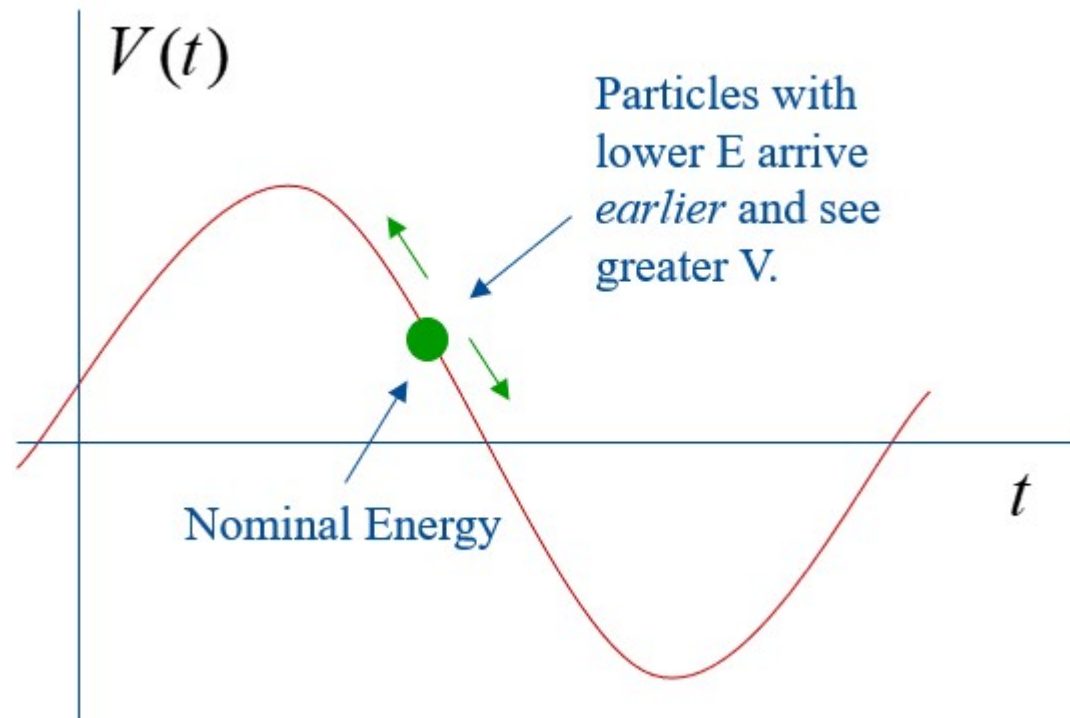
- This is referred to as the “tune”

- We can generally think of the tune in two parts:

Integer : magnet/aperture optimization → **64.31** ← Fraction: Beam Stability

Longitudinal Motion: Phase Stability

- Particles are typically accelerated by radiofrequency (“RF”) structures.
- Stability depends on particle arrival time relative to the RF phase.
 - ◆ Time of arrival depends mostly on the energy deviation relative to “the reference (central) particle”



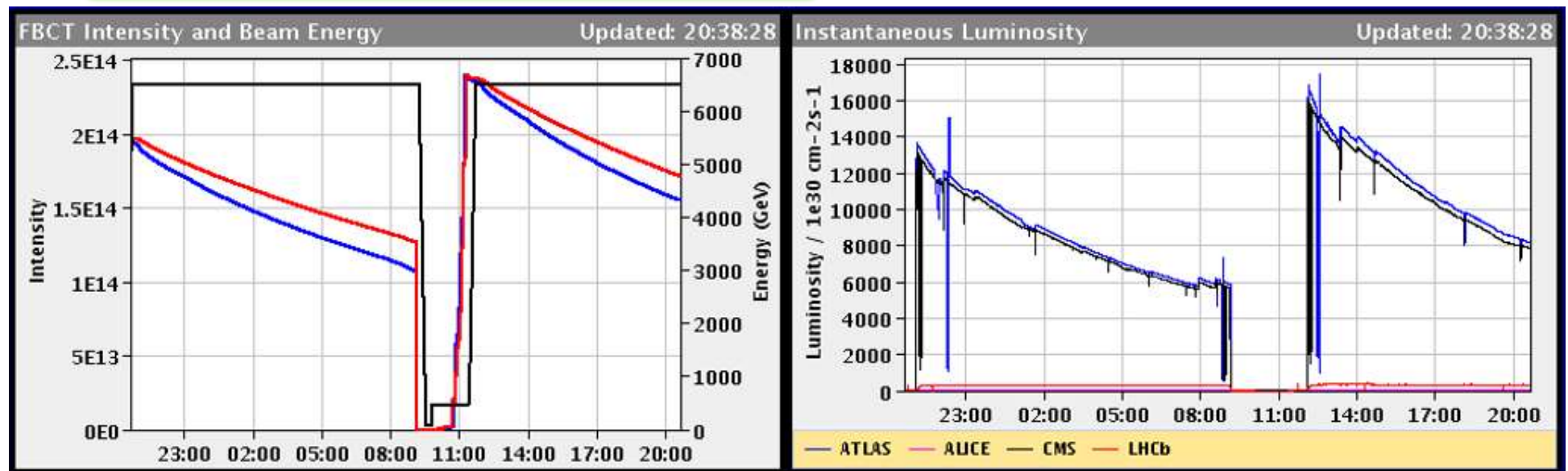
Luminosity Evolution

$$L = \gamma f_B \frac{N_1 N_2}{4\pi\beta^* \varepsilon} H(\sigma_s / \beta^*)$$

■ Factors change in time: $L(t) = C \frac{N_1(t)N_2(t)}{\varepsilon(t)} H(t)$

■ Therefore, in the absence of cooling the lifetime

$$\tau_L^{-1} = \frac{dL(t)}{L(t)dt} = \tau_{N1}^{-1} + \tau_{N2}^{-1} - \tau_{\varepsilon}^{-1} + \tau_H^{-1}$$



LHC luminosity plot

Electrons versus protons

■ Electrons

- ◆ (+) Point-like objects
=> the entire energy may go to creation of a particle-of-interest
- ◆ (+) Well-determined energy
=> better resolution; in particular, for narrow resonances
- ◆ (+) Smaller backgrounds
=> Easier to separate events from backgrounds => less expensive detector
- ◆ (-) Energy is limited by SR ($dE/dt \propto E^4$)
 - In LEP (LHC tunnel, $C=26.7$ km) operating at $E=104$ GeV the beam was losing 3% of its energy per turn

■ Protons

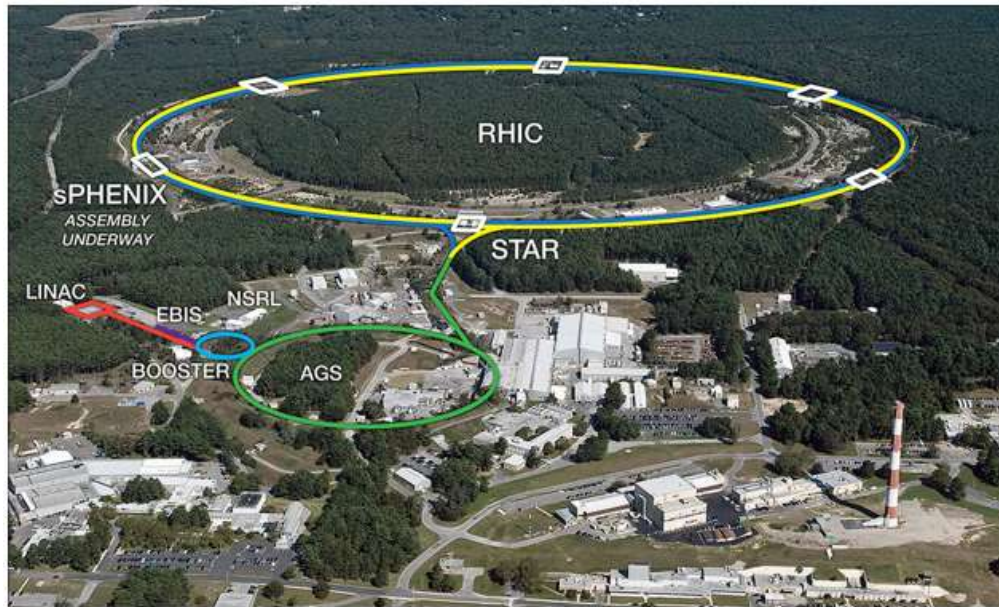
- ◆ (-) Large nuclear cross sections => large background
- ◆ (-) Quarks carry out a fraction of energy
=> effective energy = $\sim 1/6$ of total (LHC may create particles with $\sqrt{s} \leq 2$ GeV)
- ◆ (-) Wide PDF (parton distribution function) => poor knowledge of initial energy of collisions
- ◆ (++) May operate at very high energy: LHC - $E_{\max}(\text{protons})=6.8$ TeV
- ◆ (+) Much larger cross sections for creation of hadrons. For creation of B-mesons the cross section in LHCb is ~ 4 order of magnitude higher than in KEKB

Electrons versus protons (continue)

- Development of detector technology in the last ~50 years proved that in a proton collider a modern detector can deal with backgrounds even at luminosity few units of $10^{34} \text{ cm}^{-2}\text{s}^{-1}$
- If built, all other types of colliders have to be competitive to the proton colliders (i.e. to the LHC) in luminosity and/or energy
- ILC – International Linear collider
 - ◆ Very expensive, looks like cannot not reach the design luminosity for reasonable time. SLC did not get the design luminosity (4 orders lower ILC) after 10 years of commissioning
 - Looks like is coming too late with too little E & L, compared to LHC
 - ◆ FCC e⁺e⁻ looks much better choice – CERN's choice
 - requires ~doubling of CERN budget
- Muon collider looked as a very promising choice
 - ◆ Point-like particles, small SR
 - but muons survive only about 1000 turns => cannot compete with LHC in the luminosity
 - ◆ Some people still believe that we can built one

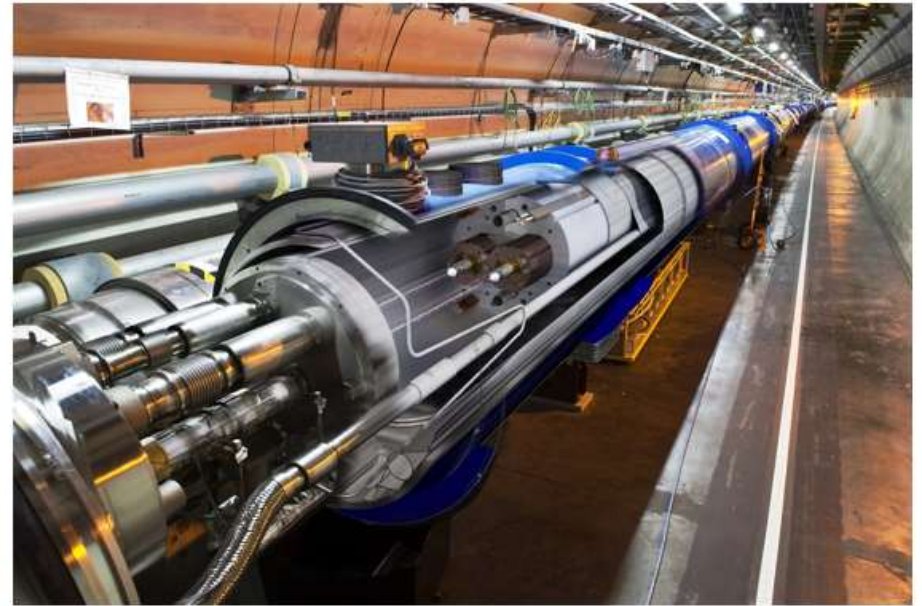
Present and Future Colliders

Present Hadron Colliders



RHIC (BNL, Brookhaven)

$C=3.84$ km,
 $E_{\text{max}}(\text{protons})=255$ GeV
■ RHIC is our main competitor



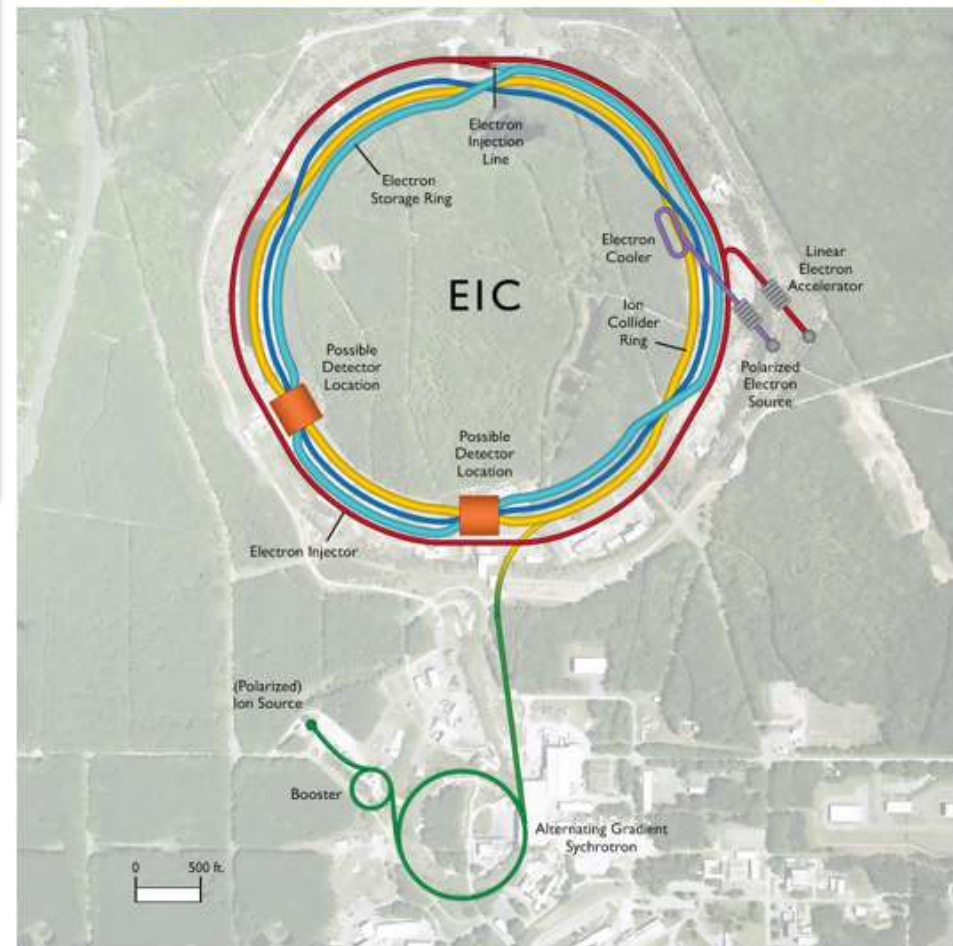
LHC (CERN)

$C=26.7$ km,
 $E_{\text{max}}(\text{protons})=6.8$ TeV

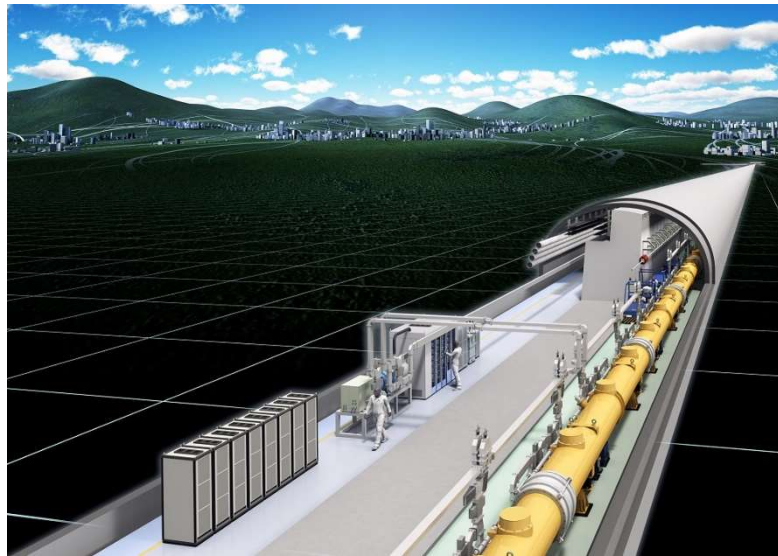
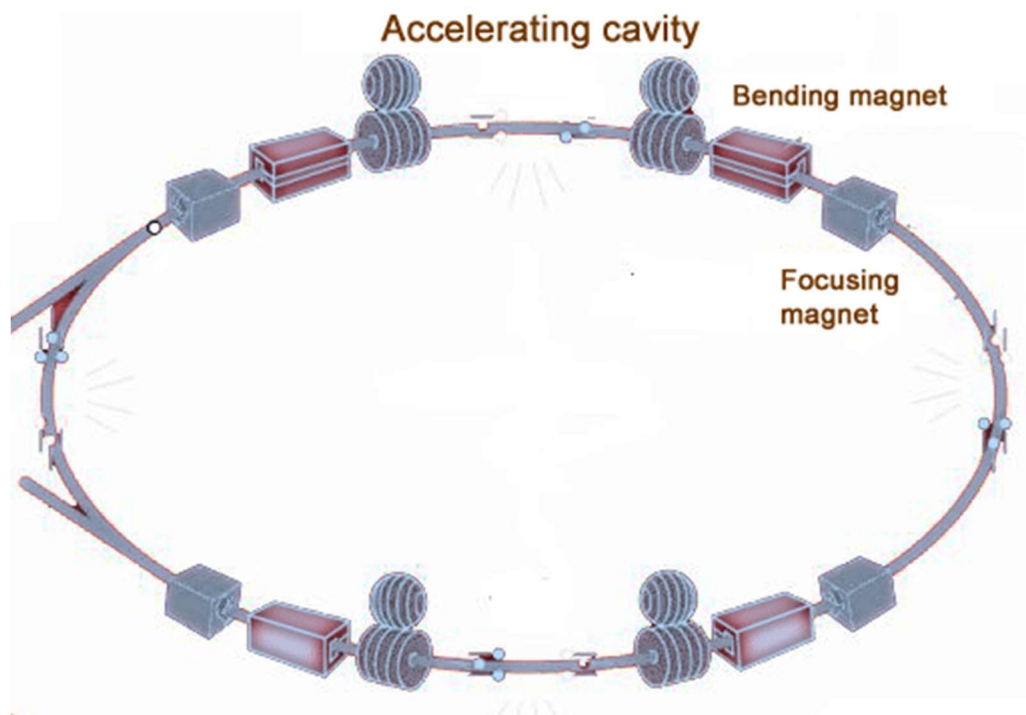
Colliders That Will Be



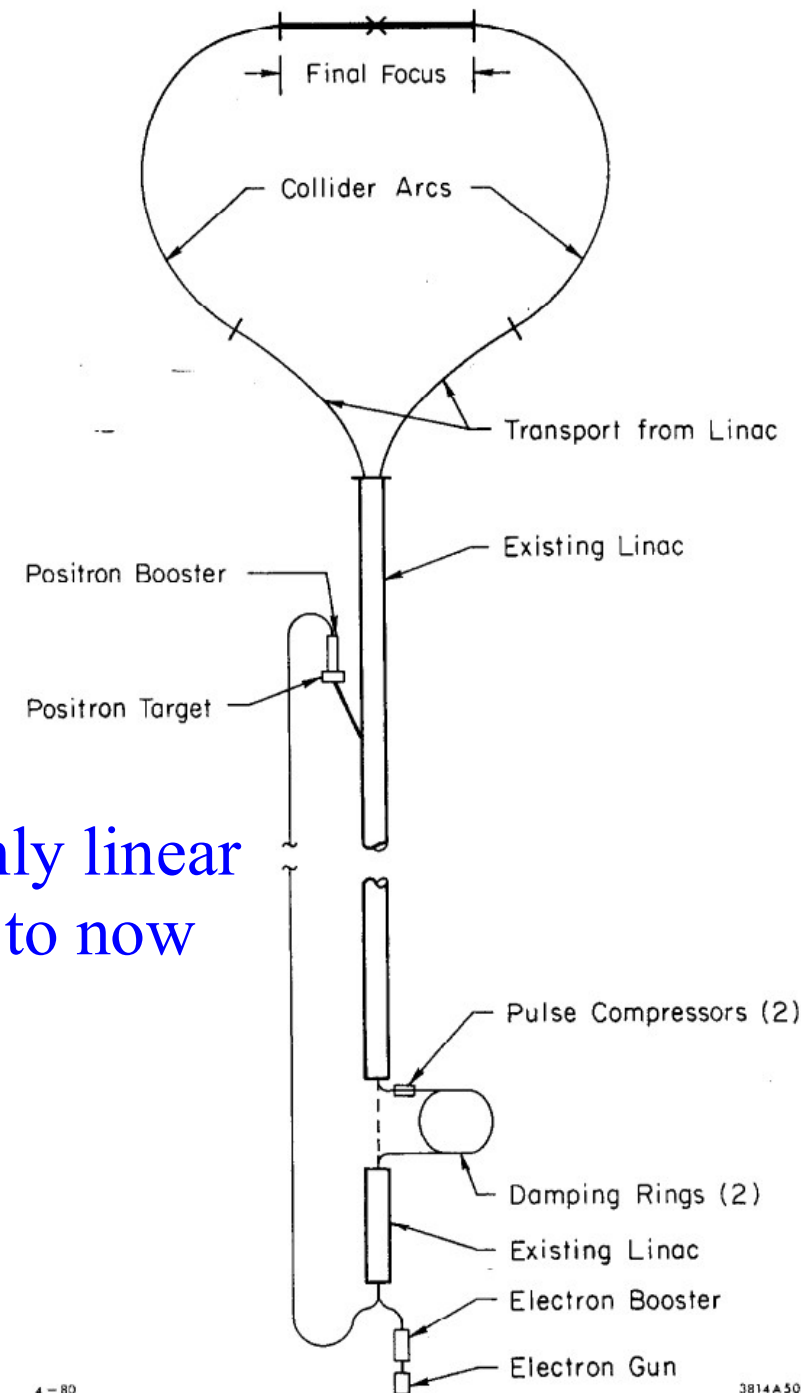
EIC (BNL, Brookhaven)



e^+e^- : Rings vs Linacs



■ SLC – the only linear collider built up to now

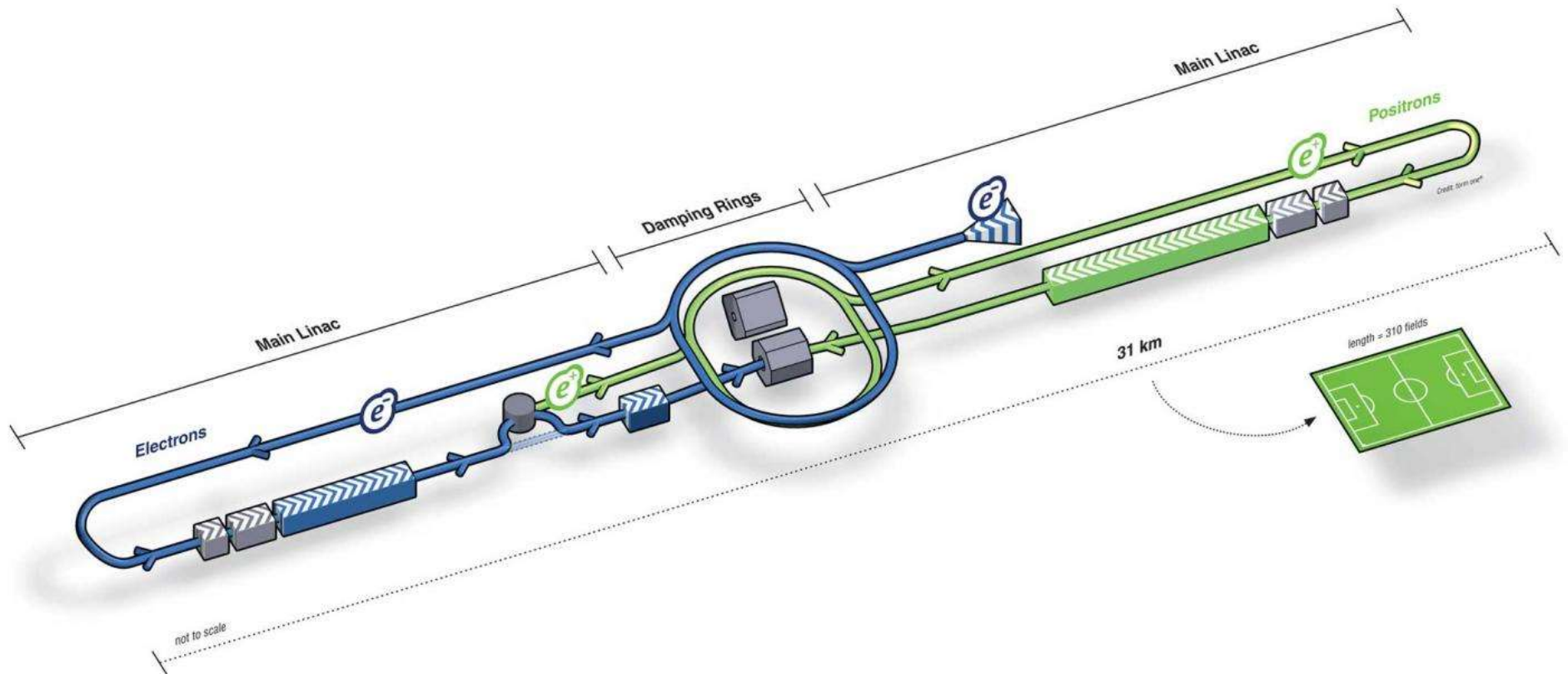


ILC (e^+e^-)

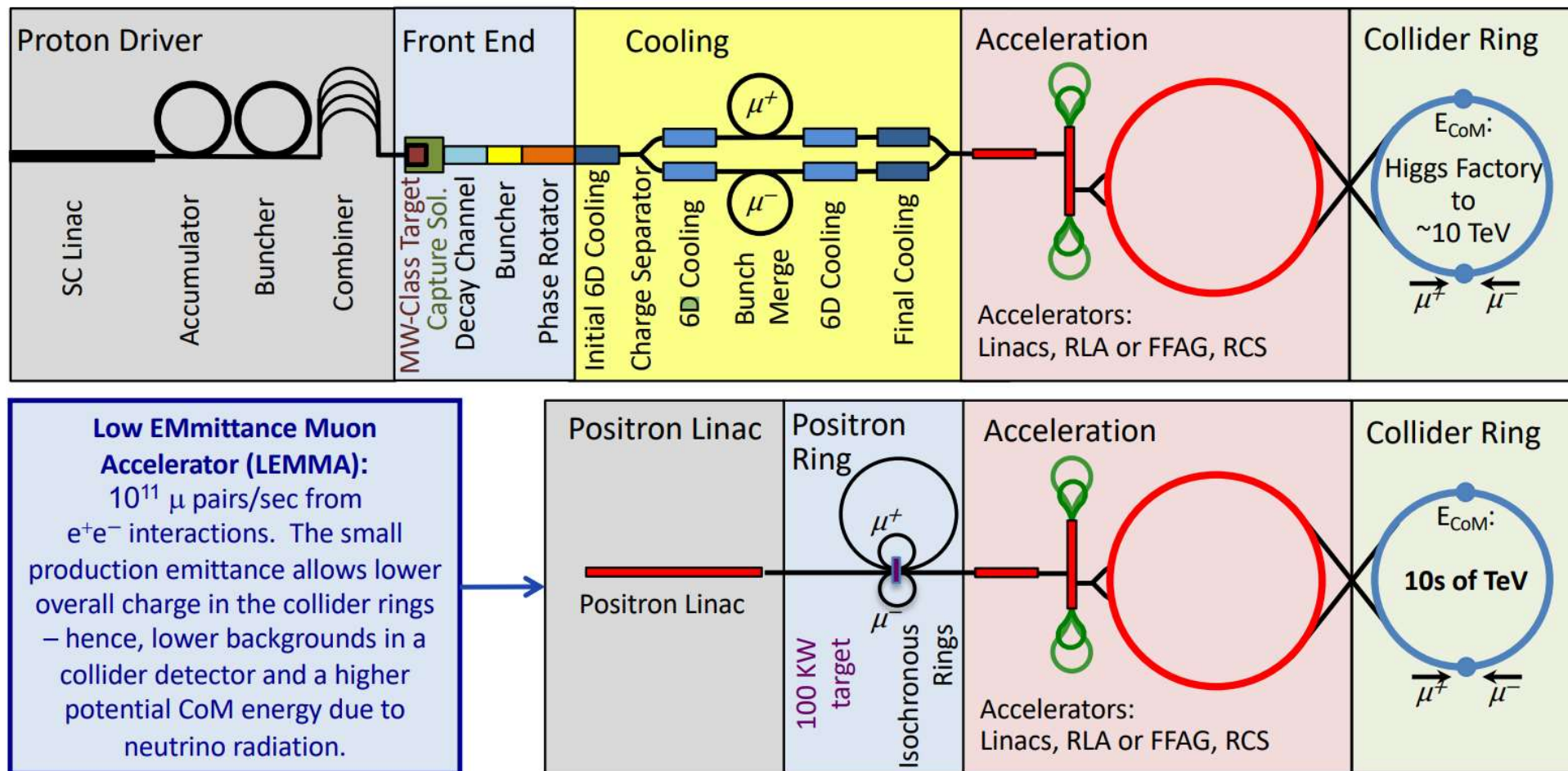
$\sqrt{s} = 500 \rightarrow 1000 \text{ GeV}$

$C = 30 \rightarrow 50 \text{ km}$

Polarized electrons and positrons



Muon Collider



■ Great challenge in accumulation and cooling muons

- ◆ Multimegawatt proton driver
- ◆ Ionization cooling

■ Higgs factory: $L \sim 10^{29} \text{ cm}^{-2}\text{s}^{-1}$ doable, $L \sim 10^{32} \text{ cm}^{-2}\text{s}^{-1}$ needed,

- ◆ Expected Higgs width 4.1 MeV, $W/M \approx 3 \cdot 10^{-5}$ ($\sqrt{s} \leq 125 \text{ GeV}$, s-channel)

Some Important Accelerator Technologies

Highest Energy = Highest Field SC Magnets

4.5T

Tevatron,
6 m, 76 mm
774 dipoles



4.5 K He, NbTi
+ warm iron
small He-plant

5.3T

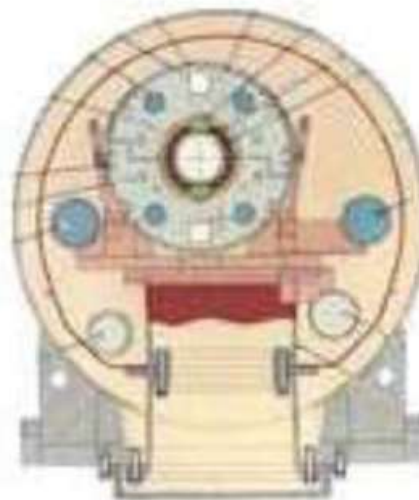
HERA,
9 m, 75 mm
416 dipoles



NbTi cable
cold iron
Al collar

3.5T

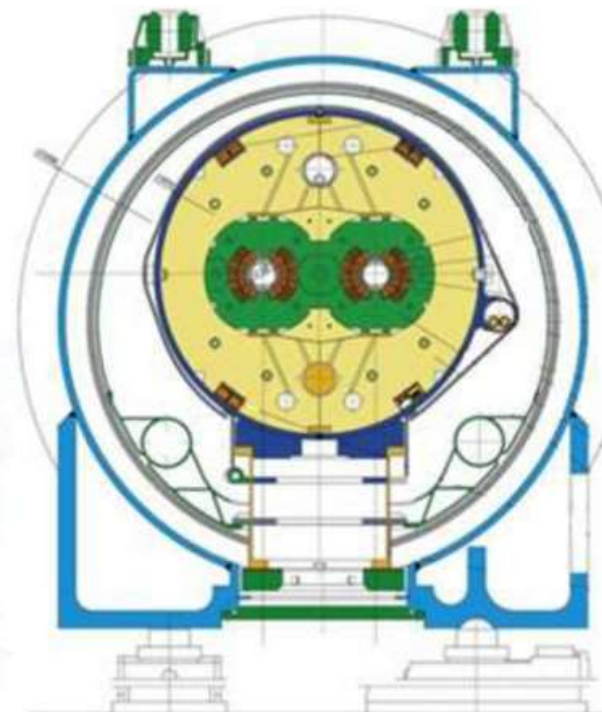
RHIC,
9 m, 80 mm
264 dipoles



NbTi cable
simple &
cheap

8.3T

LHC,
15 m, 56 mm
1276 dipoles

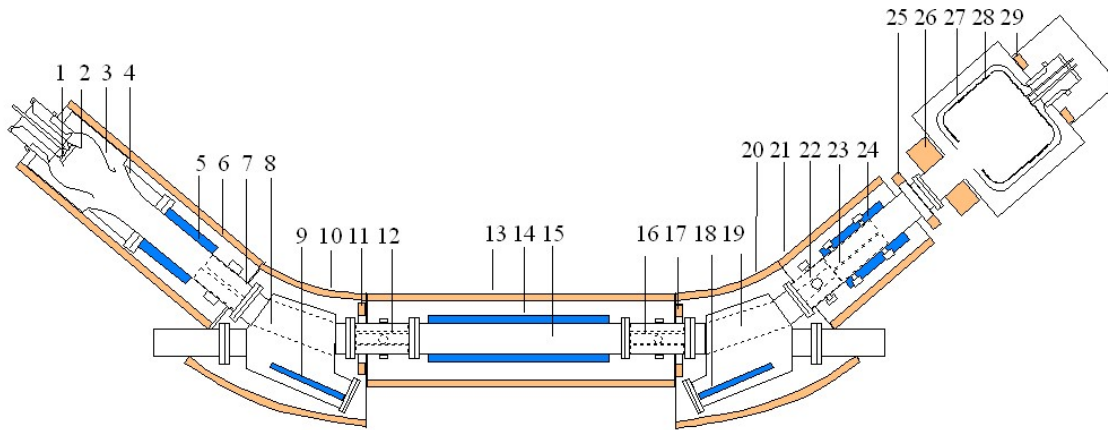


NbTi cable
2K He
two bores

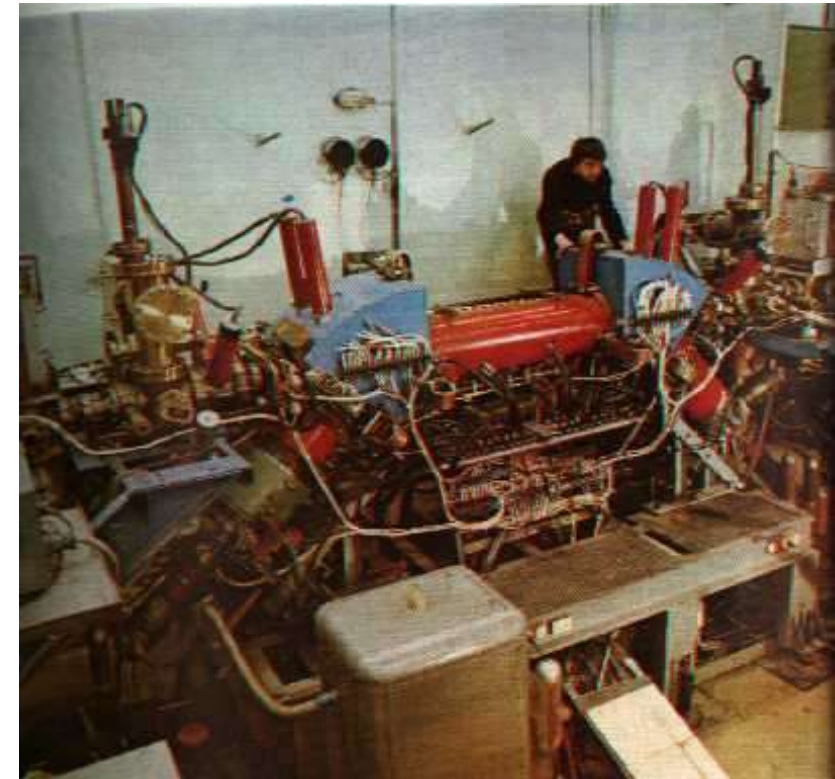
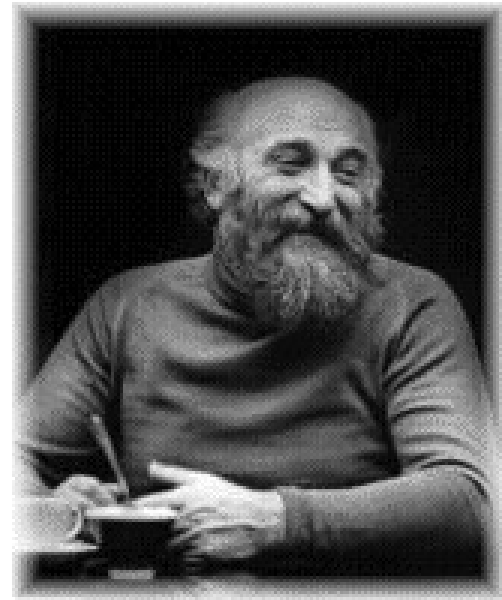


Electron cooling

- Invented in 1966 by A. M. Budker
 - ◆ In the beam frame - heavy particles come into equilibrium with electron gas
- Tested experimentally in BINP, Novosibirsk, in 1974-79 at NAP-M
 - ◆ 35 MeV electron beam (65 MeV protons)
 - ◆ Magnetized electron cooling

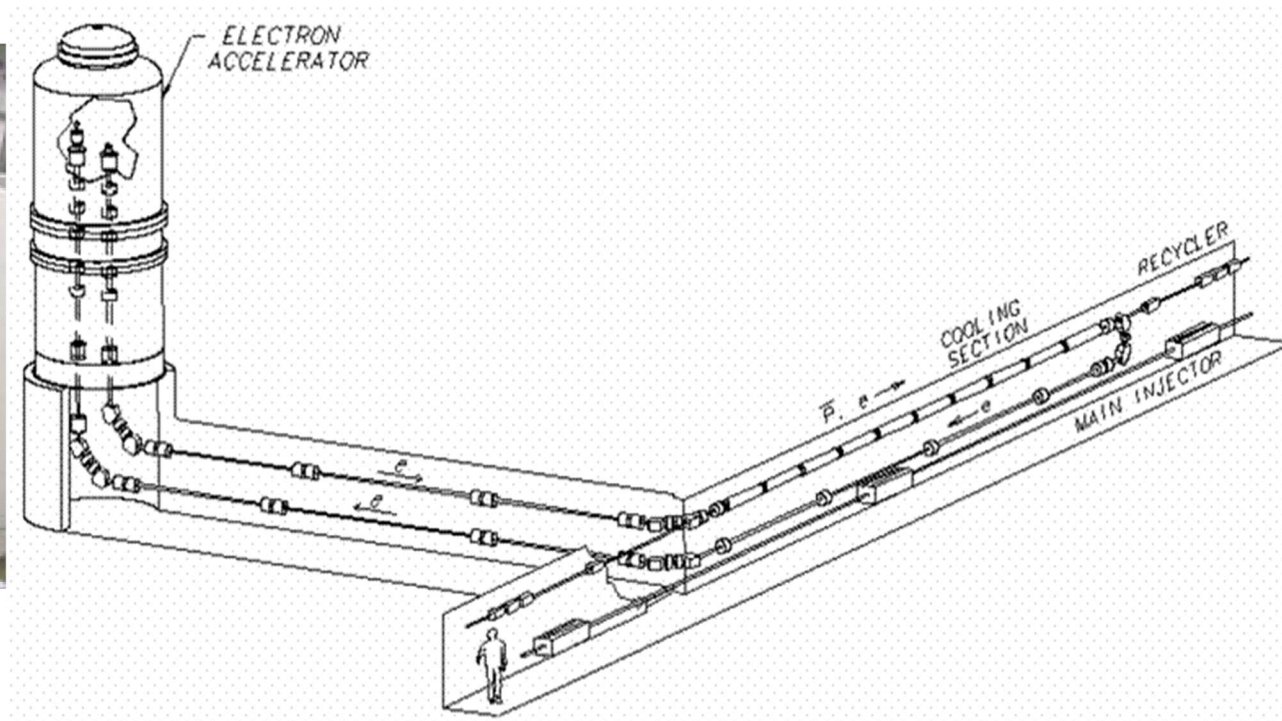


- Many installations since then, up to 300 kV electron beam (GSI, Darmstadt)
- FNAL 4.3 MeV cooler – next step in technology



Electron Cooling at FNAL

- Fermilab made next step in the electron cooling technology
- Main Parameters
 - ◆ 4.34 MeV pelletron
 - ◆ 0.5 A DC electron beam with radius of 6 mm
 - ◆ Magnetic field in the cooling section - 100 G
 - ◆ Interaction length – 20 m (out of 3319 m of Recycler circumference)



Stochastic Cooling

- Invented in 1969 by Simon van der Meer

- Naïve cooling model

- ◆ 90 deg. between pickup and kicker

$$\delta\theta = -g\theta$$

Averaging over betatron oscillations yields

$$\overline{\delta\theta^2} = -\frac{1}{2}2g\overline{\theta^2} \equiv -g\overline{\theta^2}$$

- Adding noise of other particles yields

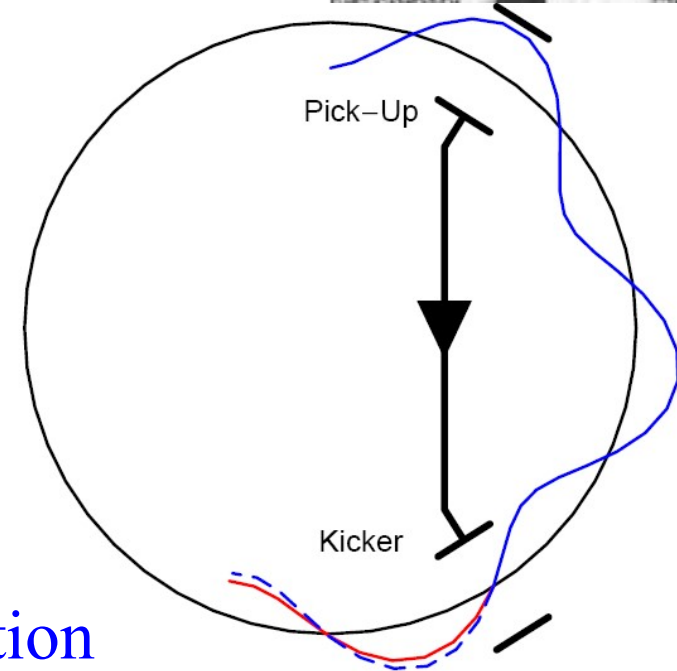
$$\overline{\delta\theta^2} = -g\overline{\theta^2} + N_{\text{sample}}g^2\overline{\theta^2} \equiv -(g - N_{\text{sample}}g^2)\overline{\theta^2}$$

- That yields

$$\overline{\delta\theta^2} = -\frac{1}{2}g_{\text{opt}}\overline{\theta^2}, \quad g_{\text{opt}} = \frac{1}{2N_{\text{sample}}}, \quad N_{\text{sample}} \approx N \frac{f_0}{W}$$

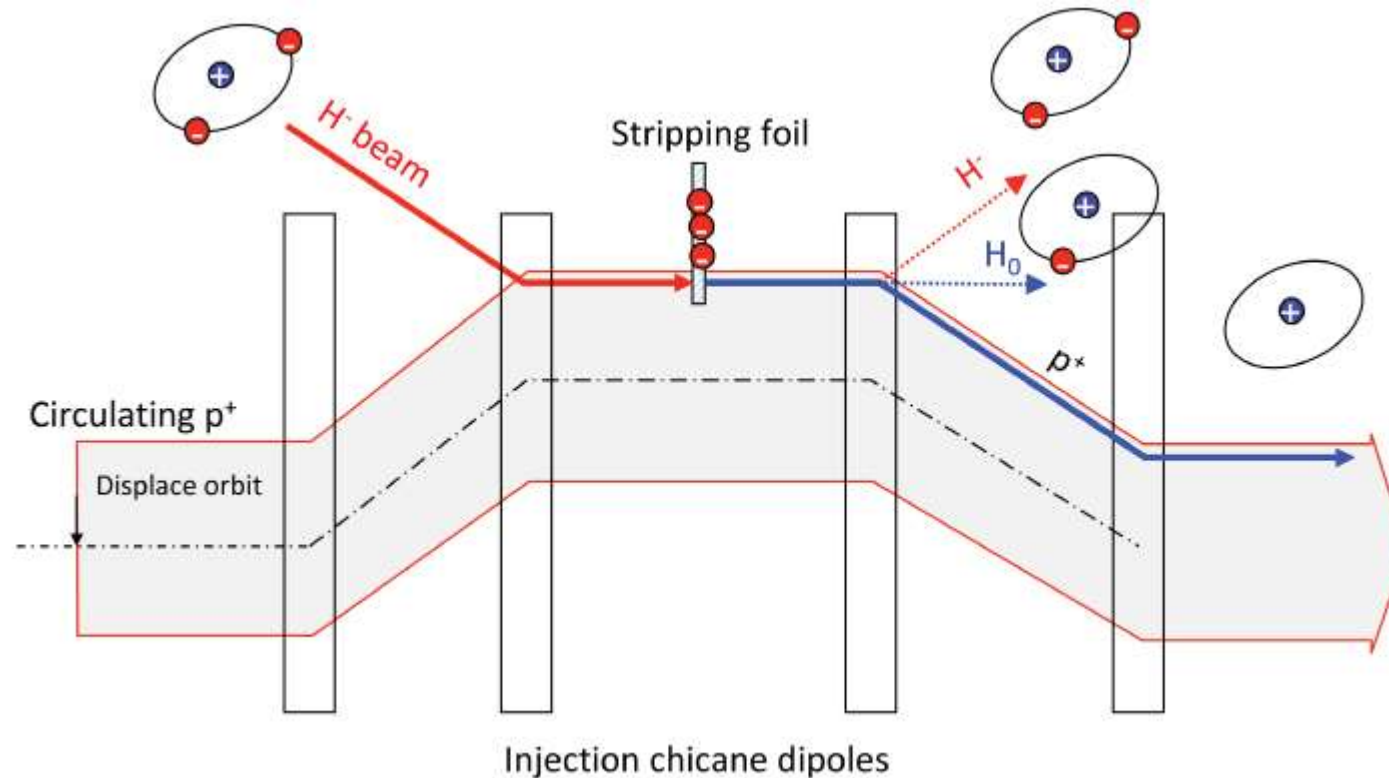
- In accurate analytical theory the cooling process is described by Fokker-Planck equation

- ◆ The theory is built on the same principle as plasma theory – which is a perturbation theory (large number of particles in the Debye sphere versus large number of particles in the sample)



Strip Injection

- Invented by Budker, First implemented in INP (Novosibirsk)
- Used in many labs: Fermilab, CERN, Oakridge NL, JPARC, ...

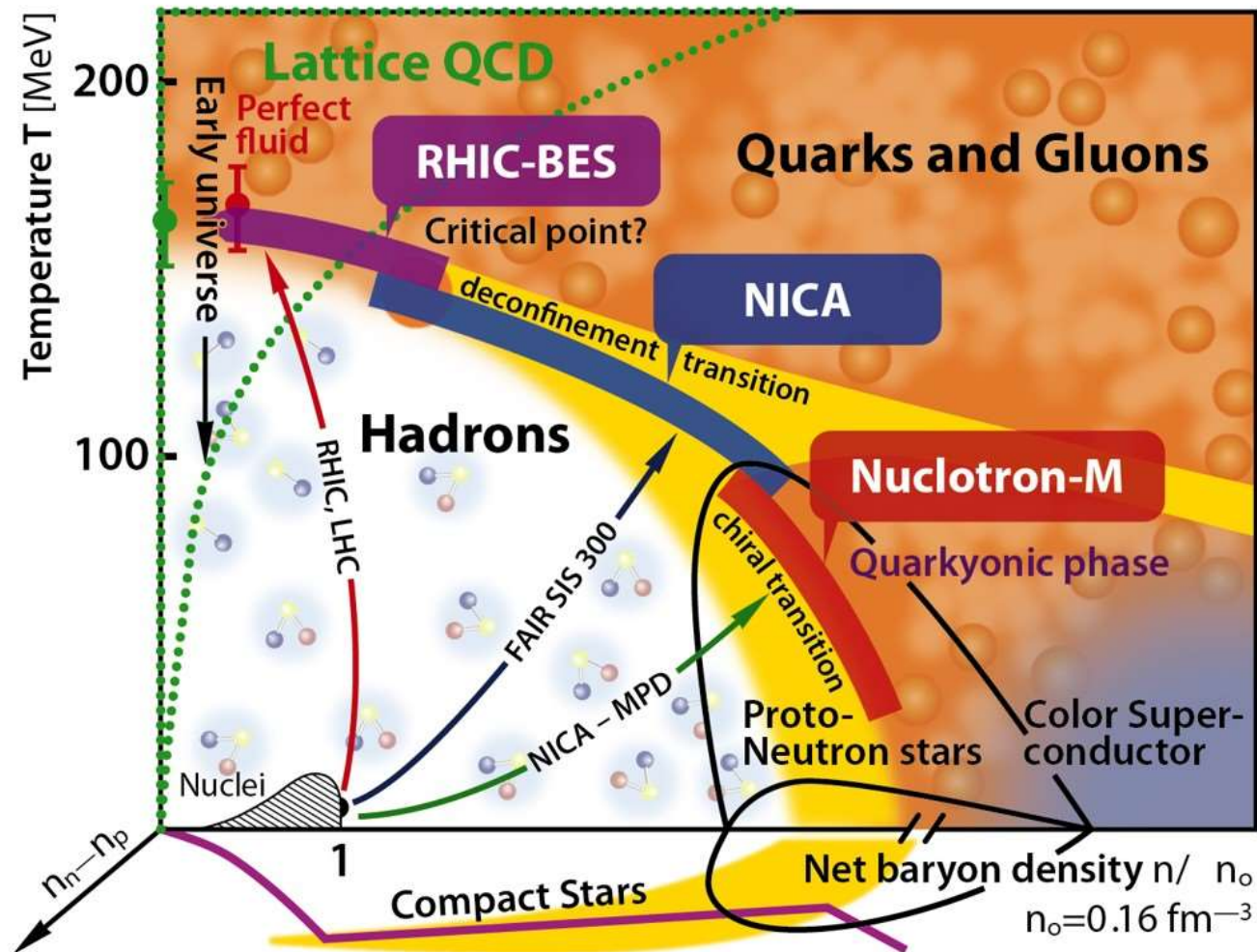


- Modern reincarnations (suggested in SNS in Oakridge):
 - ◆ Painting
 - ◆ Laser stripping

NICA – the First Hadron Collider in Russia

Major Questions in Nuclear Physics

- How do quarks and gluons give rise to the properties of strongly interacting particles?
- How does the structure of nuclei emerge from nuclear forces?
- What physics lies beyond the Standard Model?
- What are the phases of strongly interacting matter, and what roles do they play in the cosmos? (MPD)
- Spin structure of the proton/deuteron (g-factor). (SPD)



Why NICA?

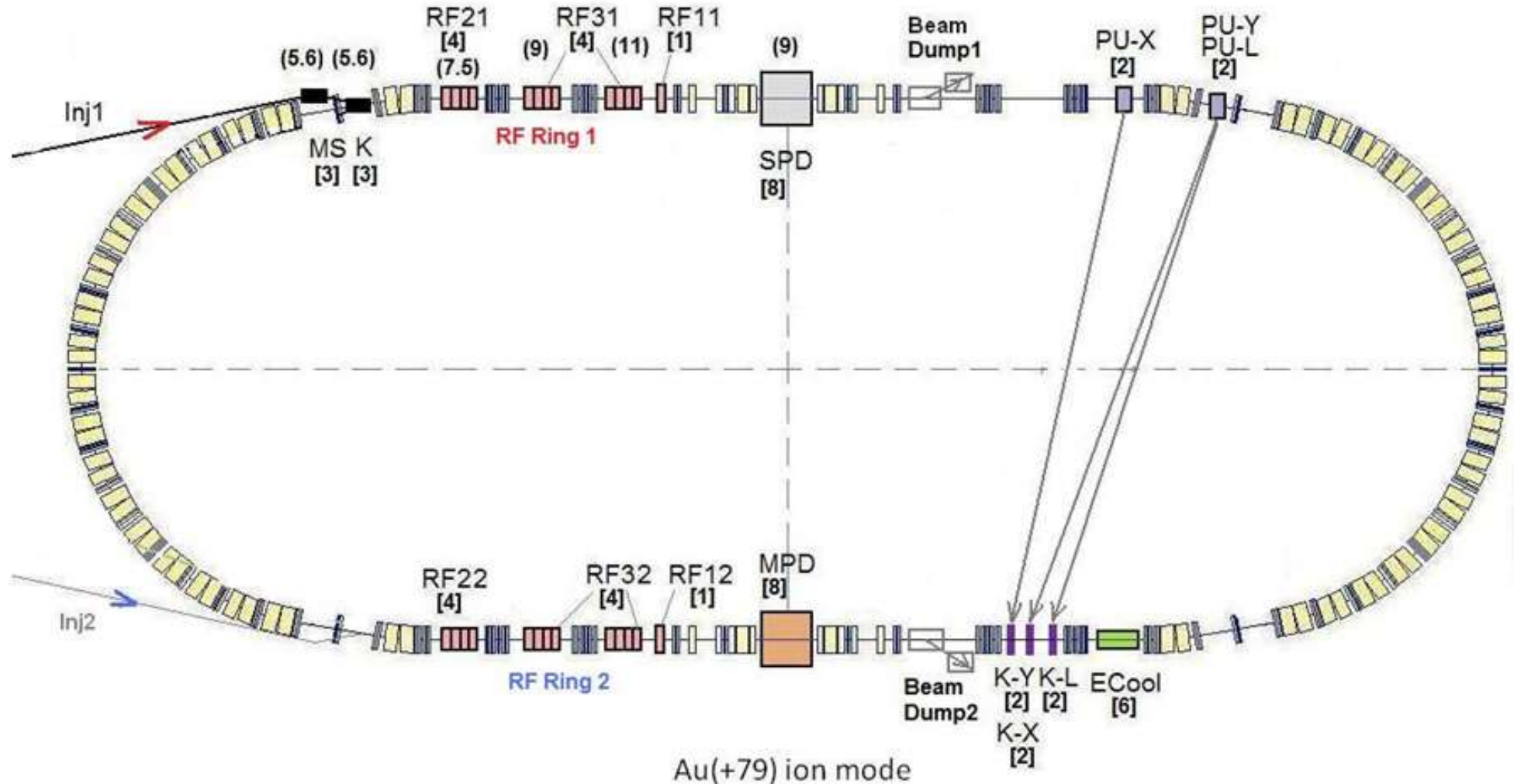
- NICA is built to answer the last 2 questions
- Unique niche
 - ◆ Two major competitors (LHC & RHIC) have too large energy to get to the ultimate luminosity in the interesting region of low energy of few GeV/n
- From accelerator physics point of view, NICA has complete set of problems/technologies present in modern hadron colliders
 - ◆ Ultrahigh vacuum
 - ◆ Superconducting (superferric) magnets
 - ◆ Large beam current results in beam instabilities
 - ⇒ Feedback systems for suppression of instabilities
 - ◆ Low-beta optics brings dynamic aperture limitations
 - Careful design of machine optics, optical measurements and correction
 - ◆ Electron and stochastic cooling at collisions
 - ◆ Instrumentation and controls required for modern colliders
 - ◆ ...

NICA Layout



- Initial operation (MPD): Xe-Xe collisions \rightarrow Bi-Bi
- The second stage (5-10 years later)(SPD): collisions of polarized protons/deuterons (spin structure)

Scheme of the Collider Ring

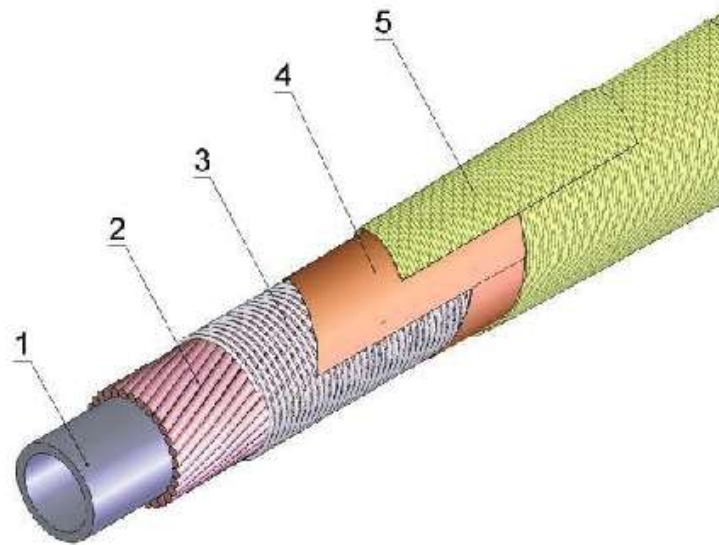
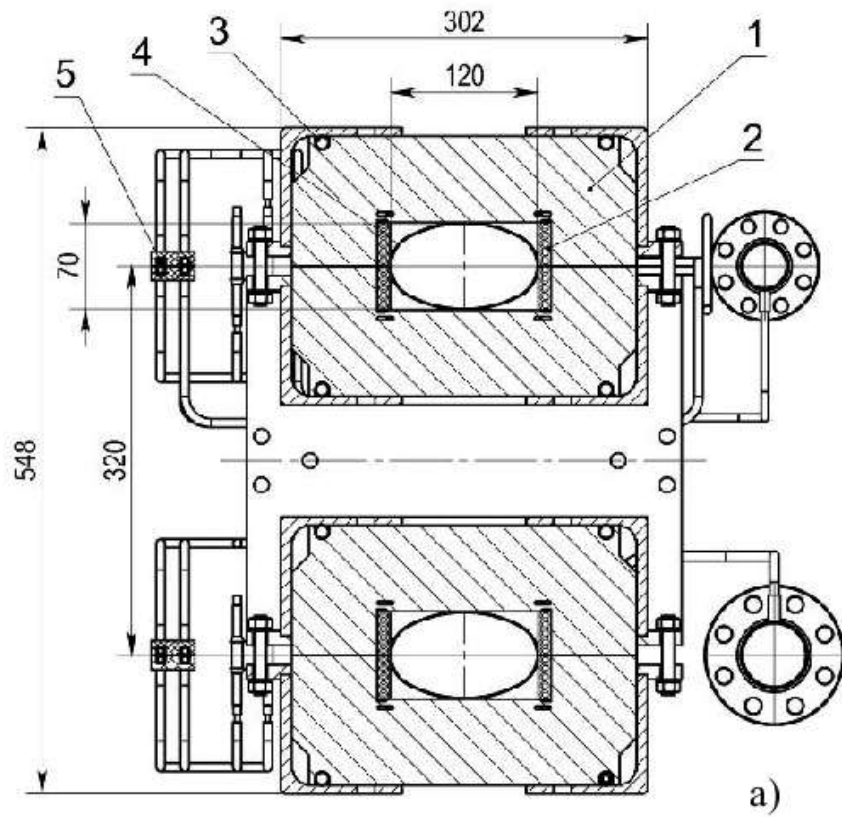


Two rings: one above another, 503 m circumference

Collision energy in the heavy ion mode: $\sqrt{s} = 2 \cdot (2.5 \div 5.5) \text{ GeV/n}$

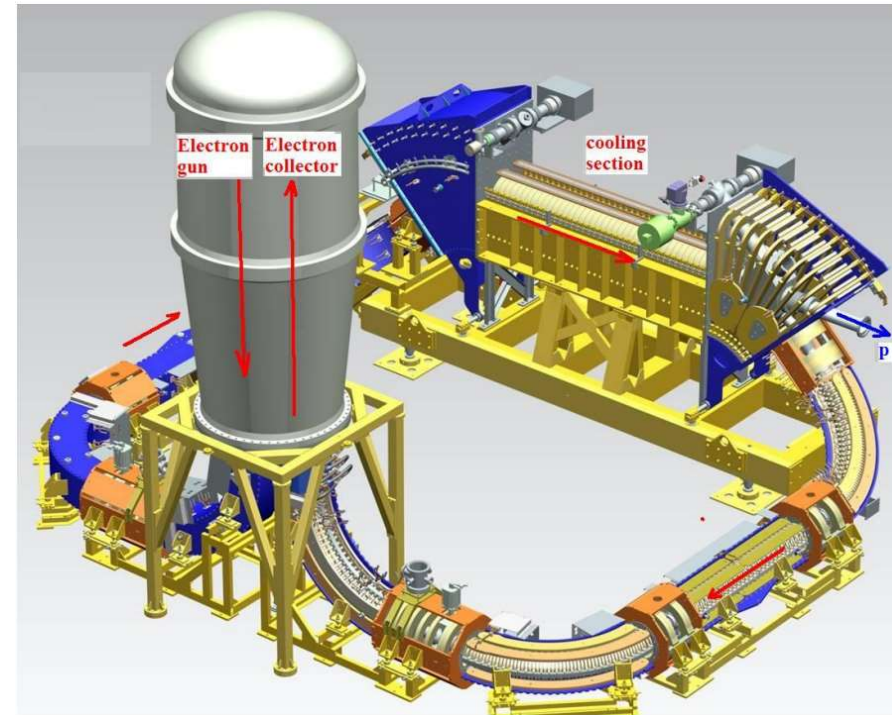
1.5 – 4.5 GeV kinetic energy

NICA dipoles

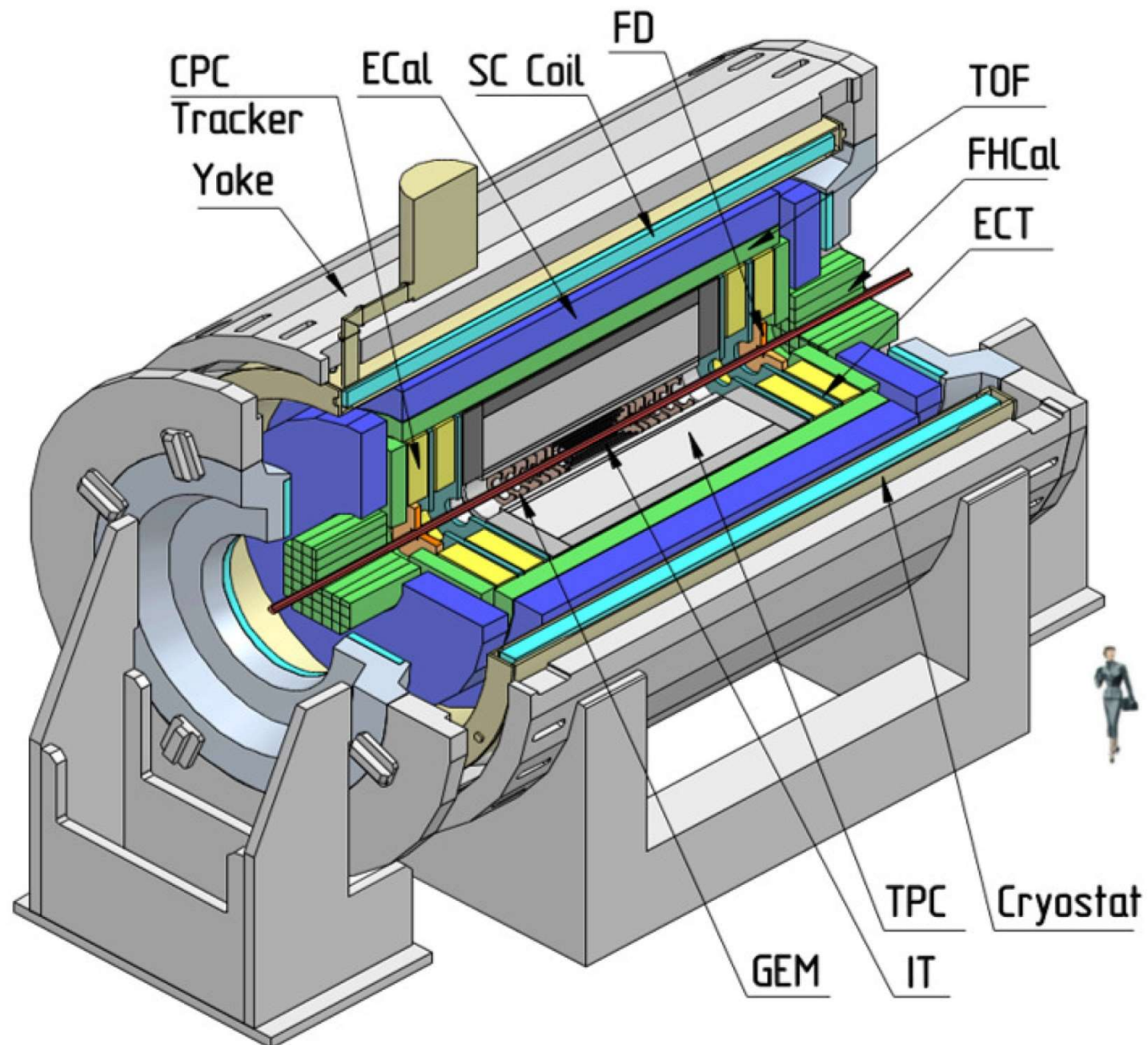


Beam Cooling

- Two systems of beam cooling will be present in NICA: electron cooling and stochastic cooling
- They are complimentary
- Stochastic cooling
 - ◆ Initially was expected to be as the main and only cooling system
 - ◆ Lack of expertise strongly delayed its development
 - ◆ Still, we plan it be ready in ~ 2 years
 - ◆ Quite challenging system to cool a bunched beam. Very little margin for errors for cooling at the collisions. Poor performance below 2.5 GeV
- Electron cooling
 - ◆ Good expertise accumulated in Novosibirsk for high energy cooling
 - 2 MeV system was supplied to COSY, Julich, Germany
 - ◆ Very good cooling of small amplitudes. Much slower cooling at high amplitudes where help from stochastic cooling would be valuable
 - ◆ Poor beam lifetime due to capture of electrons



Detector MPD



NICA: Most Important Topics/Effects

- Engineering of magnets, RF, Power supplies, vacuum, particle sources, targets, diagnostics, collimators, cryogenics, *etc.*
- Beam physics (incomplete list)
 - ◆ One particle: beam optics, long-term stability, resonances, losses, noises, diffusion/emittance growth, *etc.*
 - ◆ One beam: instabilities, beam-induced radiation deposition, intrabeam scattering, cooling, space-charge effects and compensation
 - ◆ Two-beams: beam-beam effects and compensation, instabilities in two-beam system, machine-detector interface, *etc.*
 - ◆ Beam cooling (electron, ionization, stochastic)
- Construction
 - ◆ Schedules, costs, deliveries of components
- Operations
 - ◆ Transition to whole year operation with 2-3 months shutdown
 - ◆ New operations department to drastically reduce staff required for operations

Instead of Conclusions

- At the end of this year we plan to inject beams into collider
- At the year beginning we started operations of injection complex (KRION ion source, heavy ion linac, Booster & Nuclotron)
 - ◆ The goal is an increase of particle flux by at least an order of magnitude relative to the last Run carried out 2.5 years ago
 - ◆ The means: beam accumulation with electron cooling in Booster and loss reduction in the accelerator string
- In about 2 years we plan completion of all collider systems including high voltage electron cooling, stochastic cooling, feedbacks, all 3 RF systems of each ring and MPD detector
- The program with polarized protons and deuterons will be aimed at operation with the slow beam extraction to target(s)
 - ◆ SPD detector will follow later
- Although relatively small the NICA collider will be at the front line of modern accelerator and nuclear physics
 - ◆ We are looking for young and enthusiastic people

Backup slides

Betatron Tune Shift due to Beam Space Charge

- Dependence of betatron tunes on the betatron amplitude results in that the tunes of some particles stay at non-linear resonances
 - ◆ Consequently, particle amplitudes grow resulting in the beam loss
 - ◆ SC effect is diminishing fast with beam energy

$$\begin{bmatrix} \delta\nu_{SCX} \\ \delta\nu_{SCY} \end{bmatrix} = \frac{r_p Z^2 N_i}{2\pi A \beta^2 \gamma^3} \frac{C}{\sqrt{2\pi} \sigma_s} \left\langle \frac{1}{(\sigma_x + \sigma_y)} \begin{bmatrix} \beta_x / \sigma_x \\ \beta_y / \sigma_y \end{bmatrix} \right\rangle_s, \quad \sigma_{x,y} = \sqrt{\beta_{x,y} \epsilon_{x,y} + (D_{x,y} \sigma_p)^2}$$

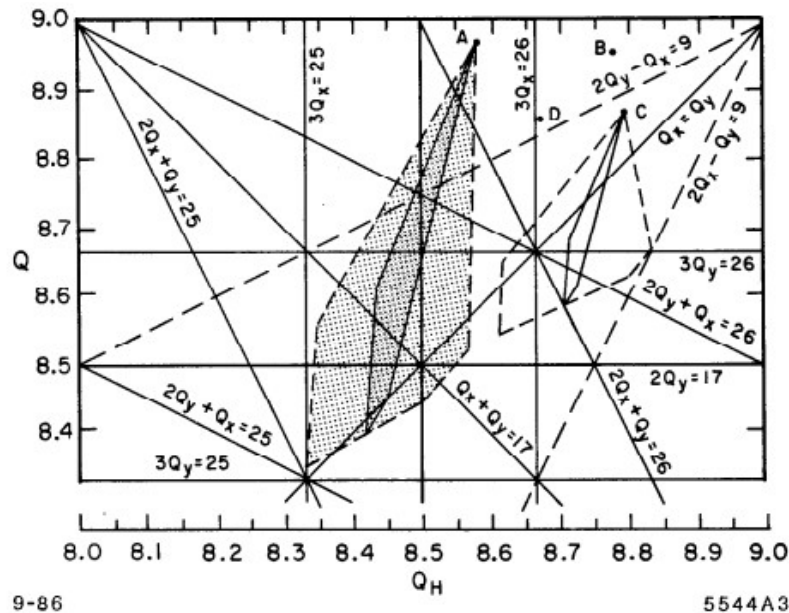


Fig. 3. Space charge tune shift of the AGS.

- Beam magnetic field $\sim \beta^2$, partially compensates electric field, $1 - \beta^2 = 1/\gamma^2$

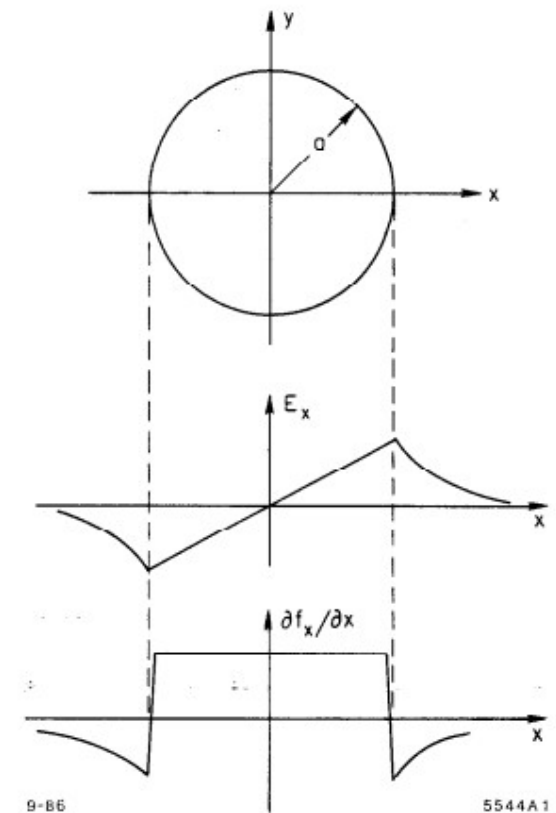


Fig. 1. Space Charge force of a uniform cylindrical beam.

Beam-beam Effects

- The beam-beam tune shift is similar to the space charge tune shift but is engaged in the IPs only. The tune shift per IP:

$$\begin{bmatrix} \delta\nu_{BBx} \\ \delta\nu_{BBy} \end{bmatrix} = \frac{r_p Z^2 N_i}{4\pi A \beta^2 \gamma} \frac{1 + \beta^2}{(\sigma_x + \sigma_y)} \begin{bmatrix} \beta_x^* / \sigma_x \\ \beta_y^* / \sigma_y \end{bmatrix}, \quad \sigma_{x,y} = \sqrt{\beta_{x,y}^* \varepsilon_{x,y} + (D_{x,y}^* \sigma_p)^2}$$

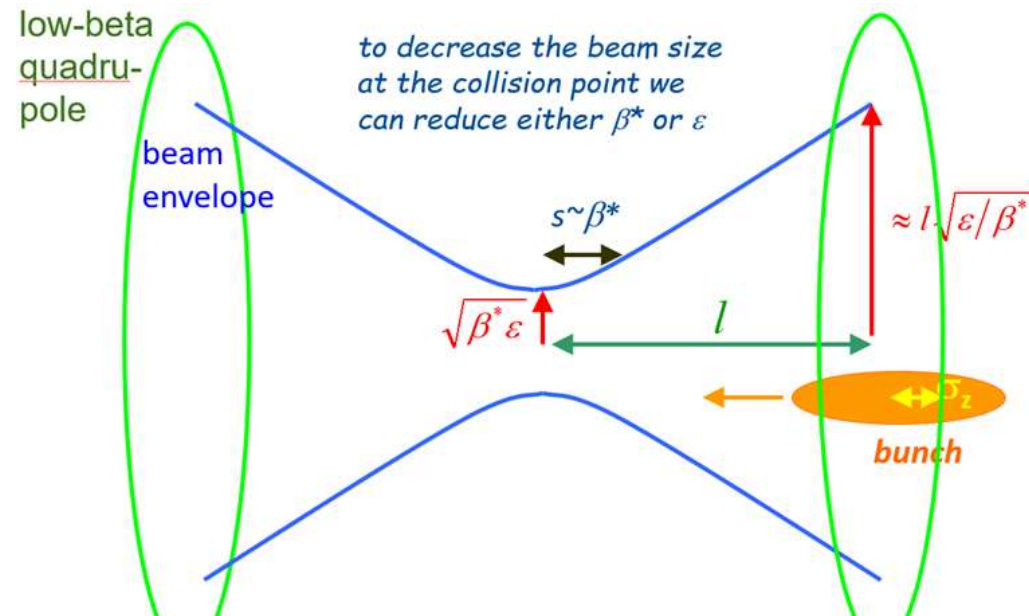
For round beam

$$\delta\nu_{SCx} = \frac{r_p Z^2 N_i}{8\pi A \beta^2 \gamma} \frac{1 + \beta^2}{\varepsilon}$$

- ◆ Magnetic field of counter rotating beam almost doubles force, $1 + \beta^2$
- ◆ Note that for large synchrotron amplitude the tune shift increase due to larger beta-function with longitudinal displacement is compensated by decrease of space charge field
=> no dependance on bunch length

- Smaller β^* yields larger β -function and beam size in quads

$$\beta(s) = \beta^* + s^2 / \beta^*$$



Possible Values of Tune Shifts

■ Achieved values of tune shifts

◆ Space charge

- NAPM ~ 0.15 (strong el. cooling, 200000 turns)
- Fermilab Booster ~ 0.3 (only ~ 2000 turns at low energy)
- JPARC, PS Booster $\sim 0.5-0.6$ (high accuracy of super-periodicity)

◆ Beam-beam

- VEPP-2 ~ 0.2 (round beams)
- Typical e^+e^- ~ 0.05 (fast SR damping)
- Typical hadron beams (Tevatron, LHC) $\sim 0.01-0.015$ per IP
- Low energy RHIC ~ 0.1 (bad life time)

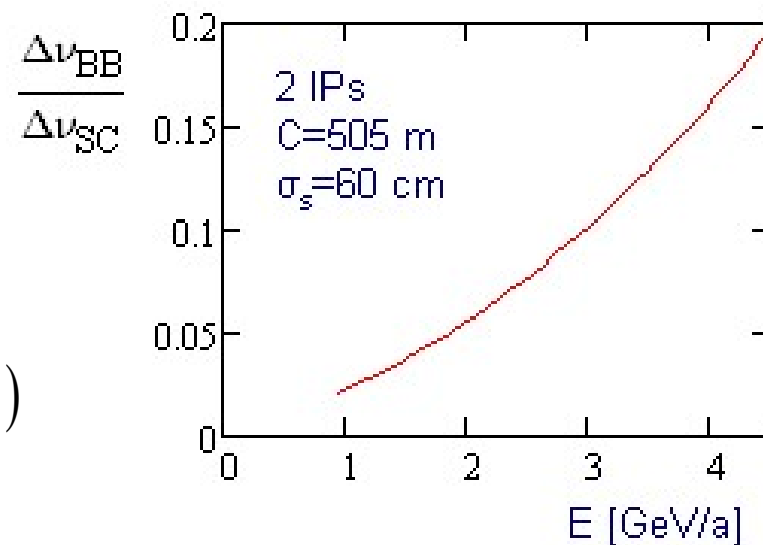
■ Ratio of tune shifts:
$$\frac{\delta\nu_{BB}}{\delta\nu_{SC}} = N_{IPs} \sqrt{\frac{\pi}{2}} \frac{\sigma_s}{C} \gamma^2 (1 + \beta^2)$$

■ For the present NICA parameters, the beam-beam tune shifts are much smaller than the space charge ones and, in the first approximation, can be neglected

■ Note that for the same tune shift the beam-beam effect is more destructive than the space charge due to kick concentration near IPs

■ For NICA we choose total $\Delta\nu = \Delta\nu_{SC} + 2\Delta\nu_{BB} \sim 0.05$

◆ Cooling helps, still quite optimistic



Luminosity Limitation due to Beam Space Charge

- Luminosity of round beams ($\beta_x^* = \beta_y^*$ & head-on collisions)

$$L = \frac{f_0 n_b N_i^2}{4\pi\beta^* \varepsilon} H_L(\sigma_s / \beta^*), \quad H_L(x) = \frac{2}{\sqrt{\pi}} \int_0^\infty \frac{\exp(-y^2)}{1+x^2 y^2} dy$$

- SC tune shift: round beam, smooth focusing & $D=0$

$$\delta\nu_{SC} \approx \frac{r_p Z^2 N_i}{4\pi A \beta^2 \gamma^3 \varepsilon} \frac{C}{\sqrt{2\pi} \sigma_s}$$

- ◆ Weak dependence of SC tune shifts on optics

- SC limits the beam longitudinal density, N_i / σ_s

- Combining the above equations, one obtains a luminosity limitation

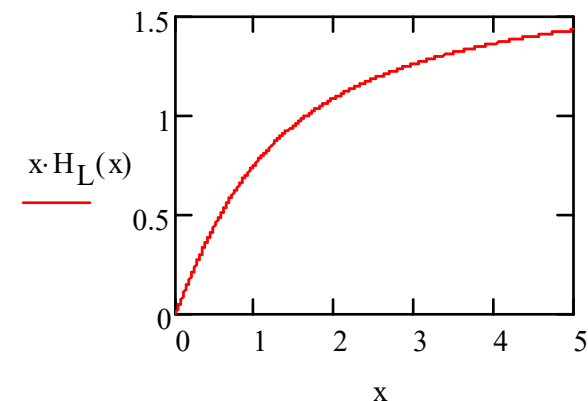
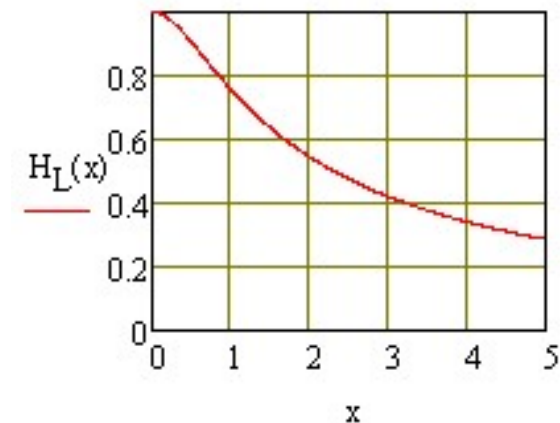
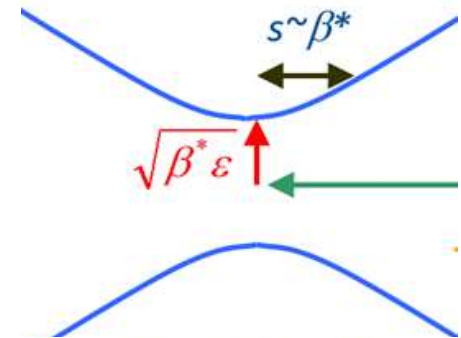
$$L = \frac{\sqrt{2\pi} A \beta^2 \gamma^3}{r_p Z^2} \frac{f_0 N_i}{(C/n_b)} \left(\frac{\sigma_s}{\beta^*} H\left(\frac{\sigma_s}{\beta^*}\right) \right) \delta\nu_{SC}$$

- ◆ Strong dependence of L on the beam energy

- ◆ Longer bunch => larger luminosity

- Still collisions must be within detector
- Luminosity distribution along IP has
the rms length of $\sigma_s / \sqrt{2} \sim 42$ cm

- ◆ $\varepsilon \propto N_i \Rightarrow$ larger luminosity \rightarrow larger acceptance



Intrabeam Scattering

■ Intrabeam scattering is determined by two major mechanisms

- ◆ Temperature exchange between degrees of freedom
 - Landau collision integral describes the temperature exchange:

$$\frac{\partial f}{\partial t} = -\frac{2\pi e^4 n L_c}{m^2} \frac{\partial}{\partial \mathbf{v}_i} \int \left(f \frac{\partial f'}{\partial \mathbf{v}'_j} - f' \frac{\partial f}{\partial \mathbf{v}_j} \right) \frac{u^2 \delta_{ij} - u_i u_j}{u^3} d^3 \mathbf{v}'$$

$$\mathbf{u} = \mathbf{v} - \mathbf{v}', \quad \int f d^3 \mathbf{v} = 1$$

- ◆ Additional heating related to non-zero dispersion
 - Scattering with particle momentum change results in additional betatron oscillations due to instant change of reference orbit

$$\Delta x = D \frac{\Delta p}{p} \xrightarrow[\text{approximation}]{\text{smooth lattice}} \Delta \varepsilon_x = \frac{1}{2} \frac{\Delta x^2}{\beta_x} = \frac{D^2}{2\beta_x} \left(\frac{\Delta p}{p} \right)^2$$

■ Relatively simple equations in the smooth lattice approximation

- ◆ Below transition there is an equilibrium state where no emittance growth
- ◆ Particle mass changes “its sign” above the transition. That yields unlimited emittance growth (energy is taken from the beam energy)

■ In heavy ion mode NICA operates in the regime of quasi-equilibrium (all 3 temperatures are approximately equal)