Observation of a Strange Muonic Atom and Its Antimatter in Heavy-Ion Collisions

Xiaofeng Wang

University of Science and Technology of China, Shandong University

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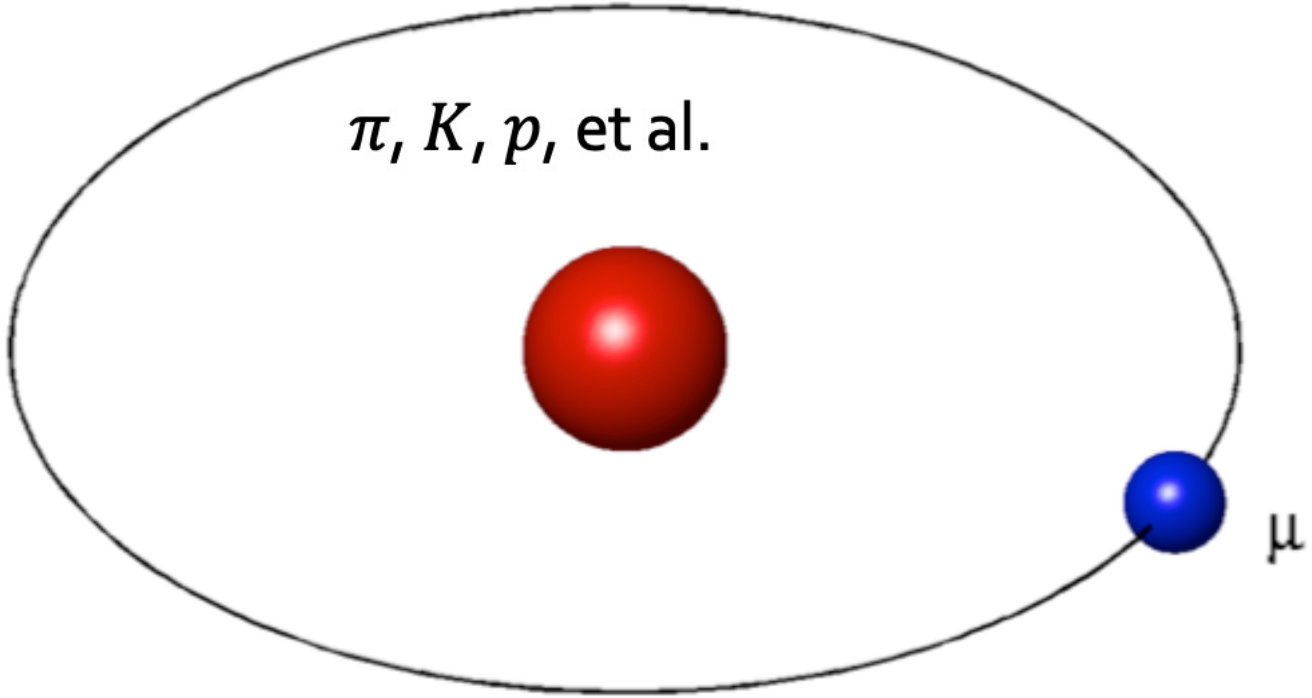
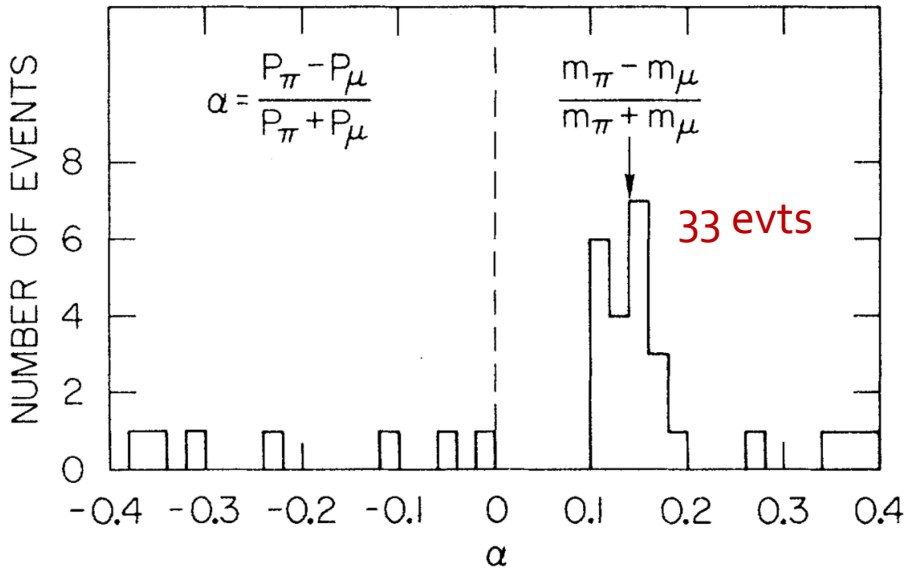
What are Muonic Atoms?



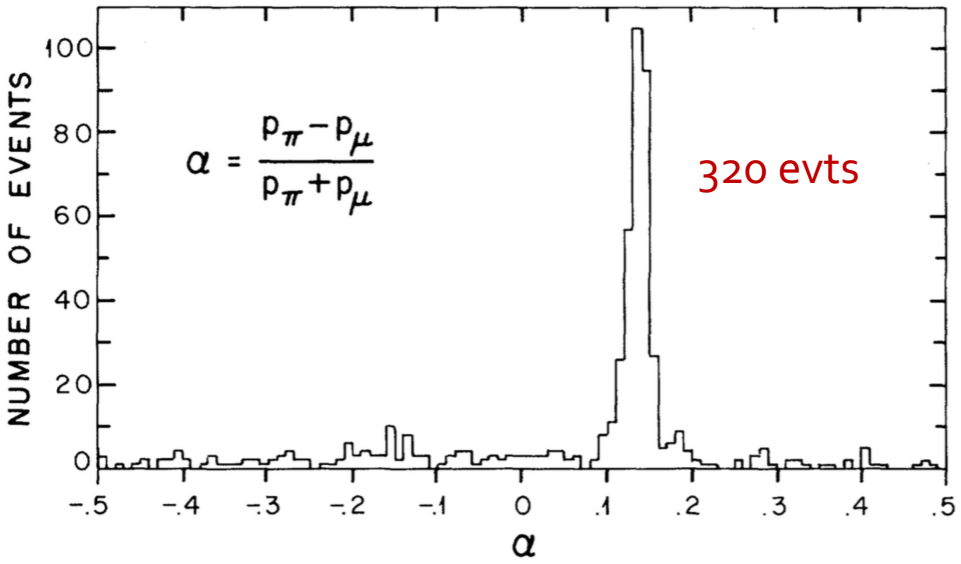
- Muonic atoms consist of a hadron (π, K, p, etc.) and a muon bound by the Coulomb force
- Like an ordinary hydrogen atom, but with a much smaller radius
- Muonic hydrogen dramatically improves proton radius measurement precision

Muonic Atom-$\pi\mu$ Observed from K_L^0 Decay

BNL, Phys.Rev.Lett. 37 (1976) 249-252



FNAL, Phys.Rev.Lett. 48 (1982) 1078-1081



- BNL/Fermilab discovered $\pi\mu$ atoms from intense K_L^0 decays
- Ionized by Al foil with π and μ same velocity
- Sharp peak at expected kinematic variable α
- $\bar{p}\mu^+$ and $K^\pm\mu^\mp$ atoms not yet discovered, need new production mechanism

Muonic Atoms in Ultrarelativistic Heavy-Ion Collisions (HICs)

- Small volume, copious hadrons and μ,
  enhance formation of hydrogen-like systems: $\pi\mu$, $K\mu$ and $p\mu$

  G. Baym, et al. Phys. Rev. D 48, R3957–R3959 (1993)
- Only way to discover some of those exotic atoms
- Test QED and meson form factors
- Muonic atom yield proportional to muon and hadron yields

  G. Baym, et al. Phys. Rev. D 48, R3957–R3959 (1993)
- Muonic atom yields give access to primordial muon production

  Joseph Kapusta and Agnes Mocsy, Phys. Rev. C 59, 2937–2940 (1999)
  - ✓ The soft leptons produced in the early evolution are difficult to measure (mixed with decay muon)
  - ✓ Only thermal muons or muons from short-lived resonances can form muonic atoms
  - ✓ Muons from weak or electromagnetic decays are produced too far away from the collision zone
  - ✓ The K_L^0 yields and the fraction of its decay make negligible contribution to the total $\pi\mu$ yields

Theoretical Prediction of Muonic Atoms in HICs (1)

- Muonic atom yield proportional to muon and hadron yields

G. Baym, et al. Phys. Rev. D 48, R3957–R3959 (1993)

$$\frac{dN_{\text{atom}}}{dyd^2p_{T,\text{atom}}} = 8\pi^2\zeta(3)\alpha^3 m_{\text{red}}^2 \frac{dN_h}{dyd^2p_{T,h}} \frac{dNl}{dyd^2p_{T,l}}$$

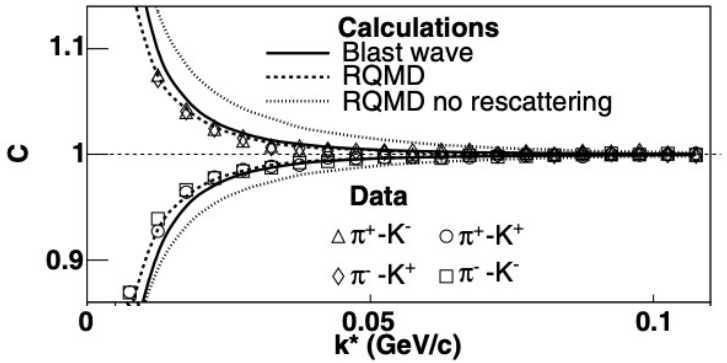
- ✓ $m_{red}^{-1} = m_h^{-1} + m_\mu^{-1}$
- ✓ All particles at same velocity

- ✓ Did not include the Coulomb potential at freeze-out in relativistic heavy-ion collisions

- ✓ Relative velocity of the two constituents at freeze-out < Bohr velocity
  → effectively assumes no correlations between the particles

Theoretical Prediction of Muonic Atoms in HICs (2)

J. Adams, et al., Phys. Rev. Lett. 91 262302 (2003)
X. Wang et al. Phys. Lett. B 861,139242(2025)



- The Coulomb interaction at freeze-out strongly enhances the two-particle correlation function for charged particles

$$\frac{q_0^2}{2m_{\mathrm{red}}} \leq \frac{1}{4\pi\varepsilon_0}\frac{e^2}{r_0}$$

- Relative momentum at the fireball freeze-out is not enough to overcome their Coulomb potential

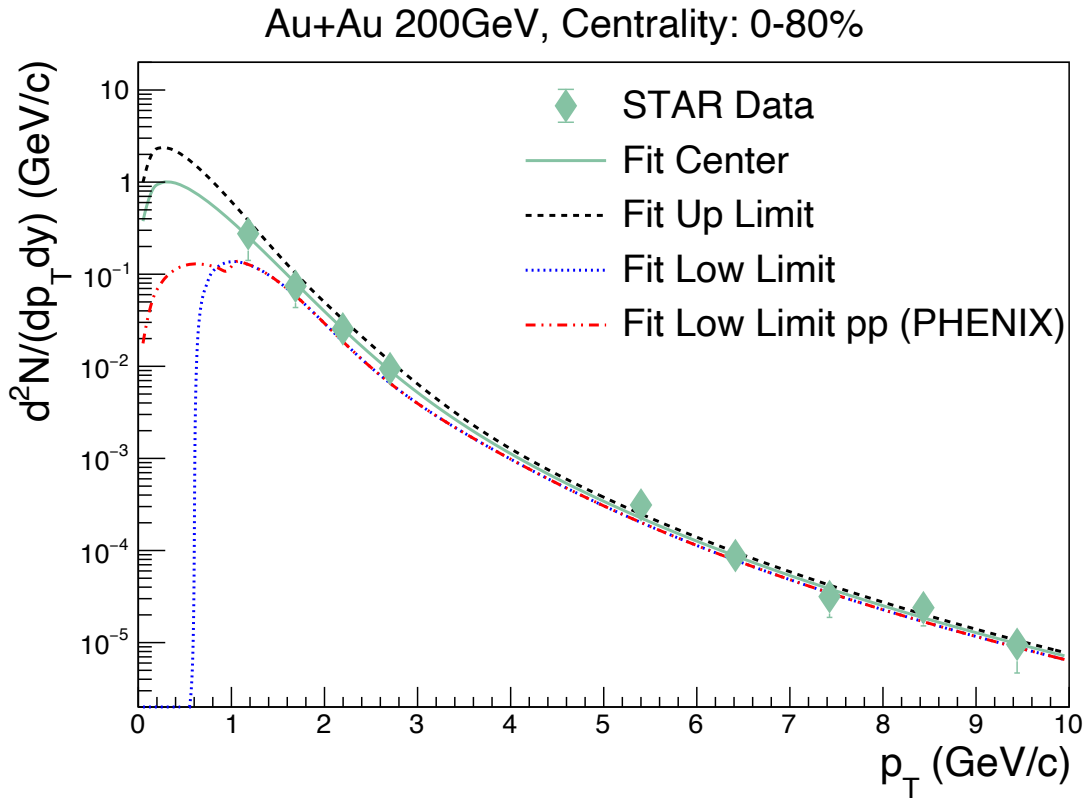
$$\frac{dN_{\mathrm{atom}}}{dyd^2p_{T,\mathrm{atom}}} = \frac{4\pi}{3}(2\alpha\hbar c/r_0)^{3/2}m_{\mathrm{red}}^{1/2}\frac{dN_h}{dyd^2p_{T,h}}\frac{dN_l}{dyd^2p_{T,l}}$$

- ✓ r_0: homogeneous coalescence radius
- ✓ $m_{red}^{-1} = m_h^{-1} + m_\mu^{-1}$
- ✓ All particles at same velocity

$$\chi = \frac{1}{6\pi\zeta(3)}\left(\frac{2\hbar c}{m_{\mathrm{red}}r_0\alpha}\right)^{3/2}$$

- $r_0 = 10$ fm, enhancement factor for $K\mu$ atom is $\chi \simeq 21.6$

Primordial μ ↔ Direct γ



X. Wang et al., Arxiv: 2603.09327
(STAR Collaboration), Phys. Lett. B 770, 451 (2017)

- Dileptons and direct virtual photons: key probes of thermal radiation
- Low-p_T thermal region NO direct data → extrapolate from high-p_T fits → order-of-magnitude uncertainties
- Need complementary observables sensitive to primordial low-p_T muons

Parameters in Coalescence Model

X. Wang et al., Arxiv: 2603.09327
(STAR Collaboration), Phys. Lett. B 770, 451 (2017)
J. Adams, et al., Phys. Rev. Lett. 91 262302 (2003)
X. Wang et al. Phys. Lett. B 861,139242(2025)

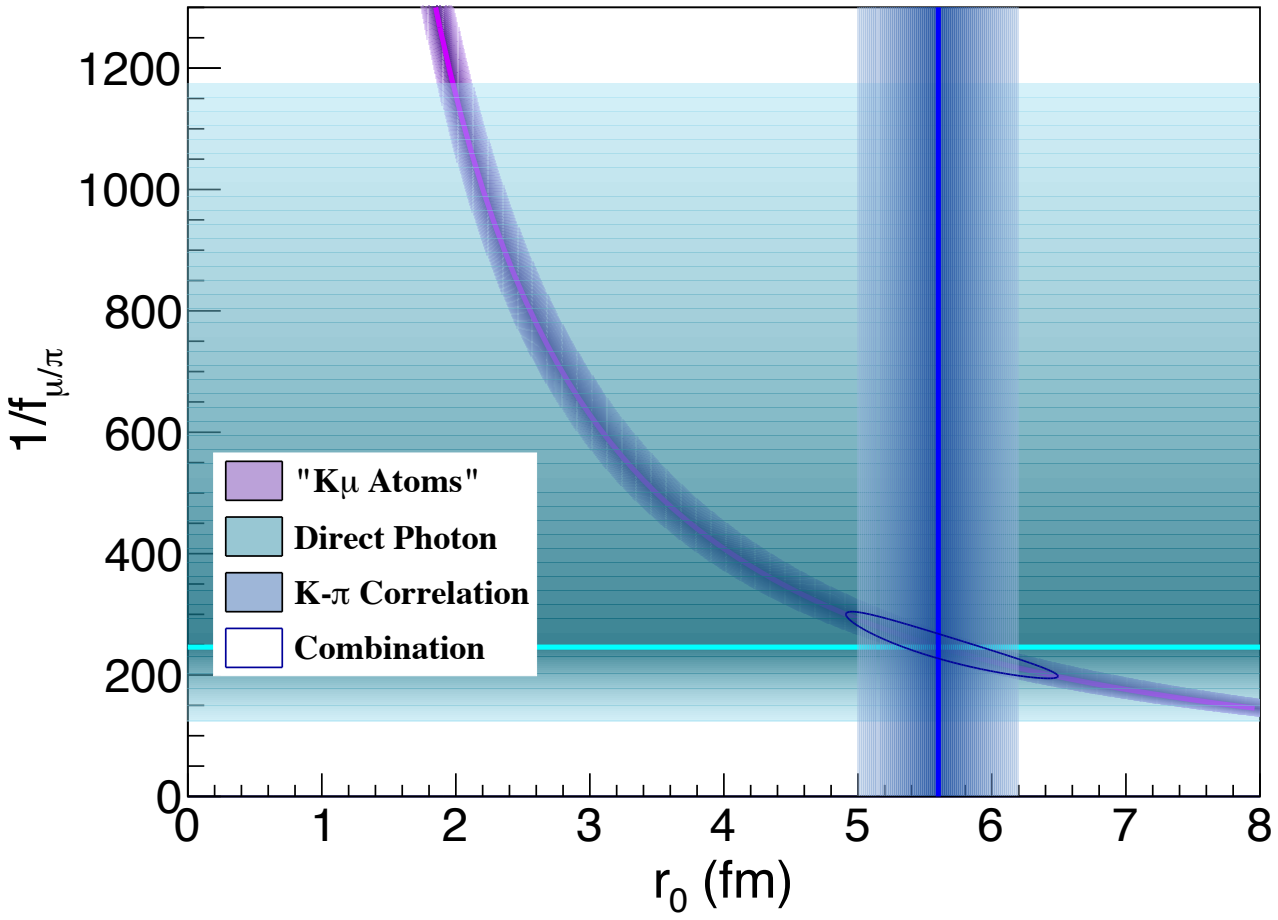
- $\boldsymbol{1/f_{\mu/\pi} = 246 + 929 - 124}$ **:**
  - ✓ Primordial muon spectrum obtained via Monte Carlo simulation of direct-photon decays
  - ✓ Pion yield taken from Tsallis–Blast-Wave Model

- $\boldsymbol{r_0 = 5.6 \pm 0.6}$ **fm**
  - ✓ From $\boldsymbol{K-\pi}$ correlation in Au+Au collisions at 130 GeV

- **In the coalescence model, assume $\boldsymbol{1/f_{\mu/\pi} = 246}$, $\boldsymbol{r_0 = 5.6}$ for the central value of K$\boldsymbol{\mu}$ atom yield, with a 10% uncertainty, yielding a joint constraint on $\boldsymbol{1/f_{\mu/\pi}}$ and $\boldsymbol{r_0}$**

$$\frac{dN_{\text{atom}}}{dyd^2p_{T,\text{atom}}} = \frac{4\pi}{3}(2\alpha\hbar c/r_0)^{3/2} m_{\text{red}}^{1/2} \frac{dN_h}{dyd^2p_{T,h}} \frac{dN_l}{dyd^2p_{T,l}}$$

Multi-dimension Constraints on μ, Kμ and r_0

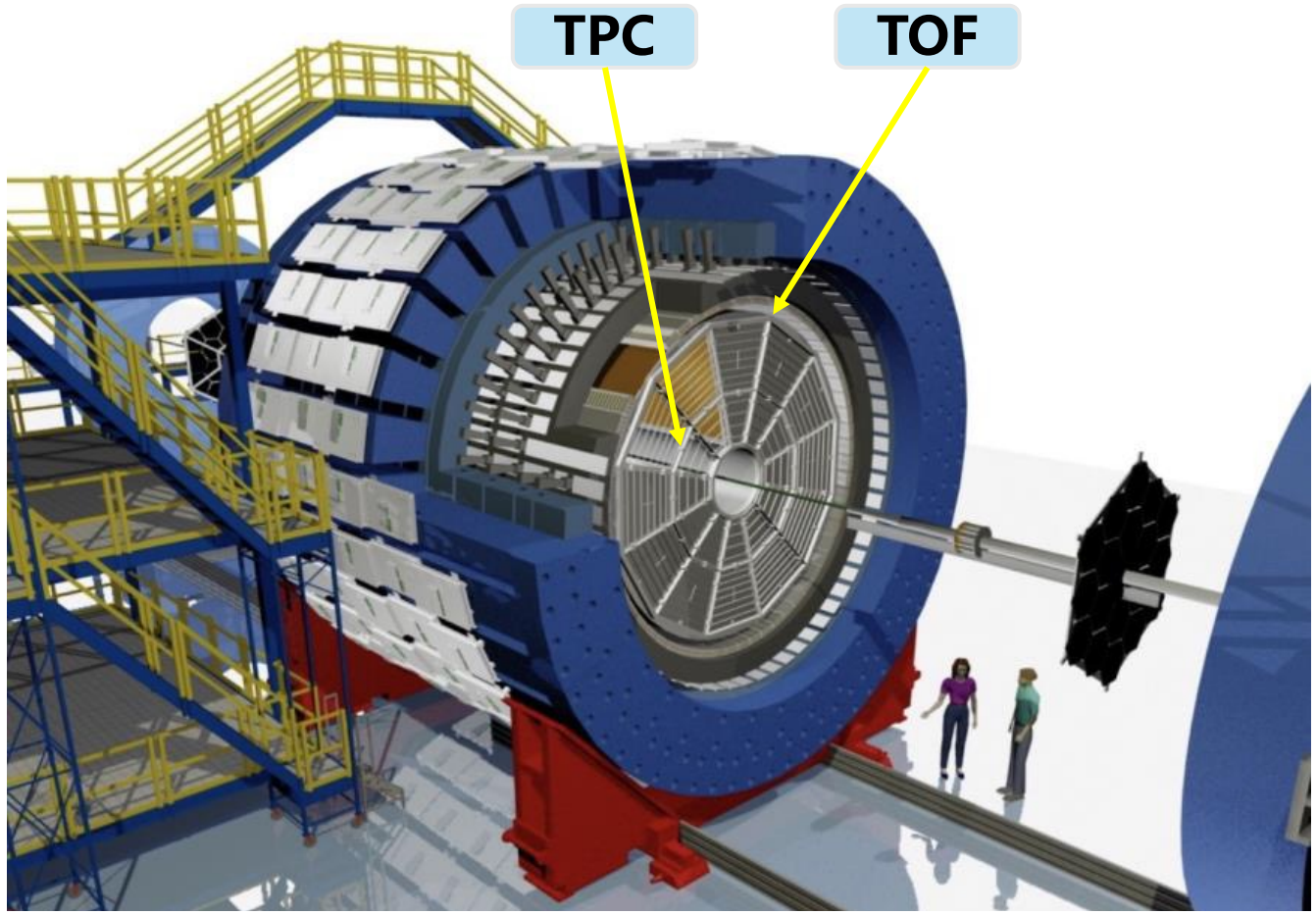


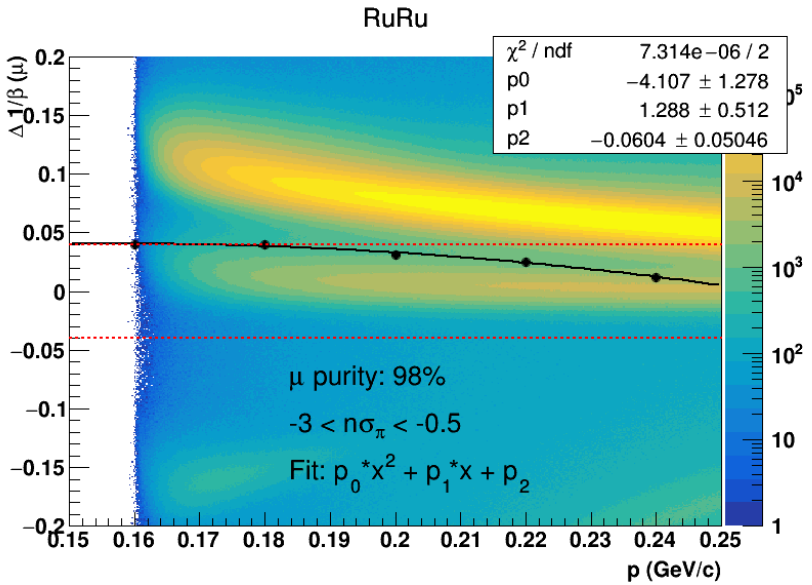
X. Wang et al., Arxiv: 2603.09327
(STAR Collaboration), Phys. Rev. Lett. 91, 262302 (2003)
(STAR Collaboration), Phys. Lett. B 770, 451 (2017)

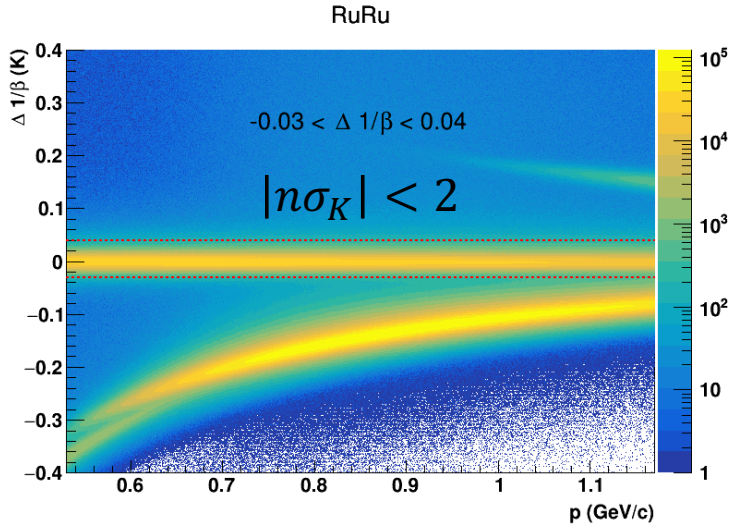
- Direct $\gamma \rightarrow 1/f_{\mu/\pi}$
- K-π Correlation $\rightarrow r_0$
- Kμ atoms: projected measurement

A powerful probe of low-p_T thermal dilepton and direct-photon production

The Solenoid Tracker At RHIC (STAR) and PID

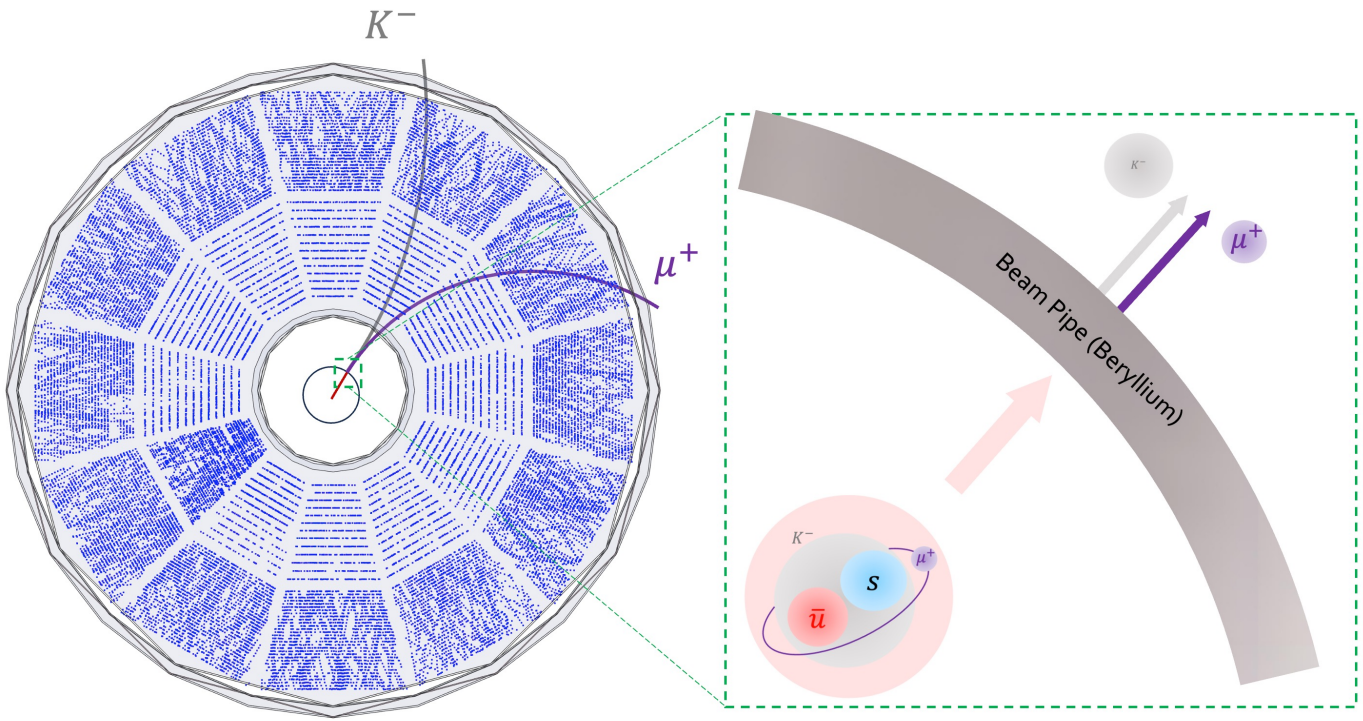




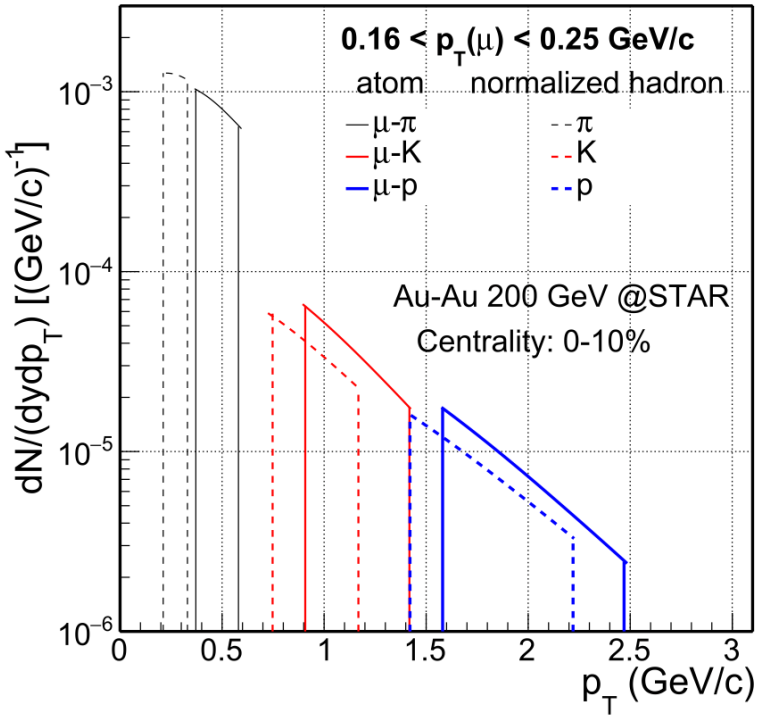


- Time Projection Chamber (**TPC**): momentum and energy loss
- Time Of Flight (**TOF**): velocity

Measurement of muonic atoms at STAR



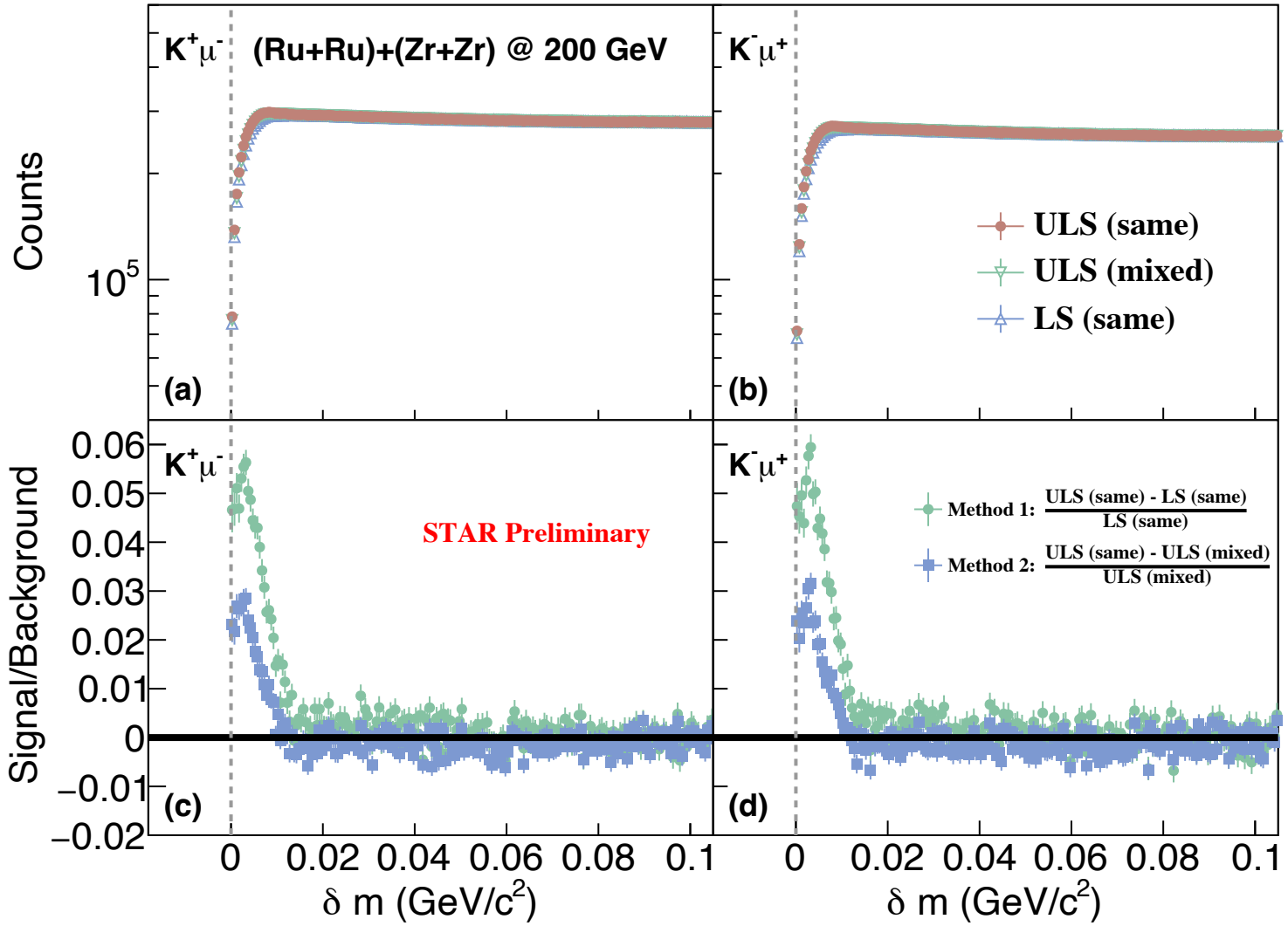
X. Wang et al. Phys. Lett. B 861,139242(2025)



- Example of a $K^-\mu^+$ event produced and detected by STAR detector
- Possible for $\pi\mu$、$K\mu$、and $p\mu$ (anti-)atoms

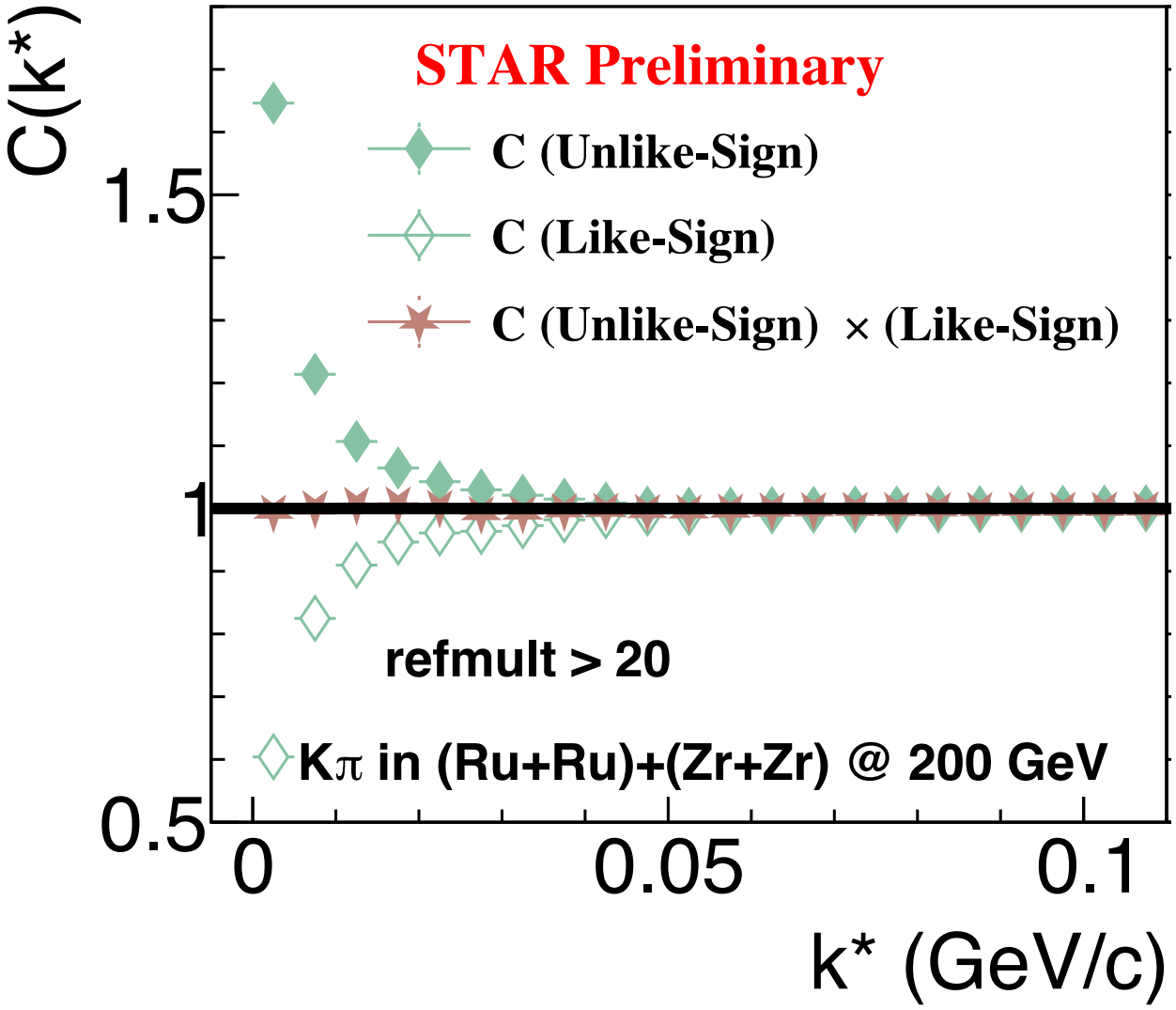
δm Distribution (primary track)

Primary tracks: tracks with DCA<3cm to the primary vertex, whose momenta are obtained from a fit assuming they originate from the primary vertex



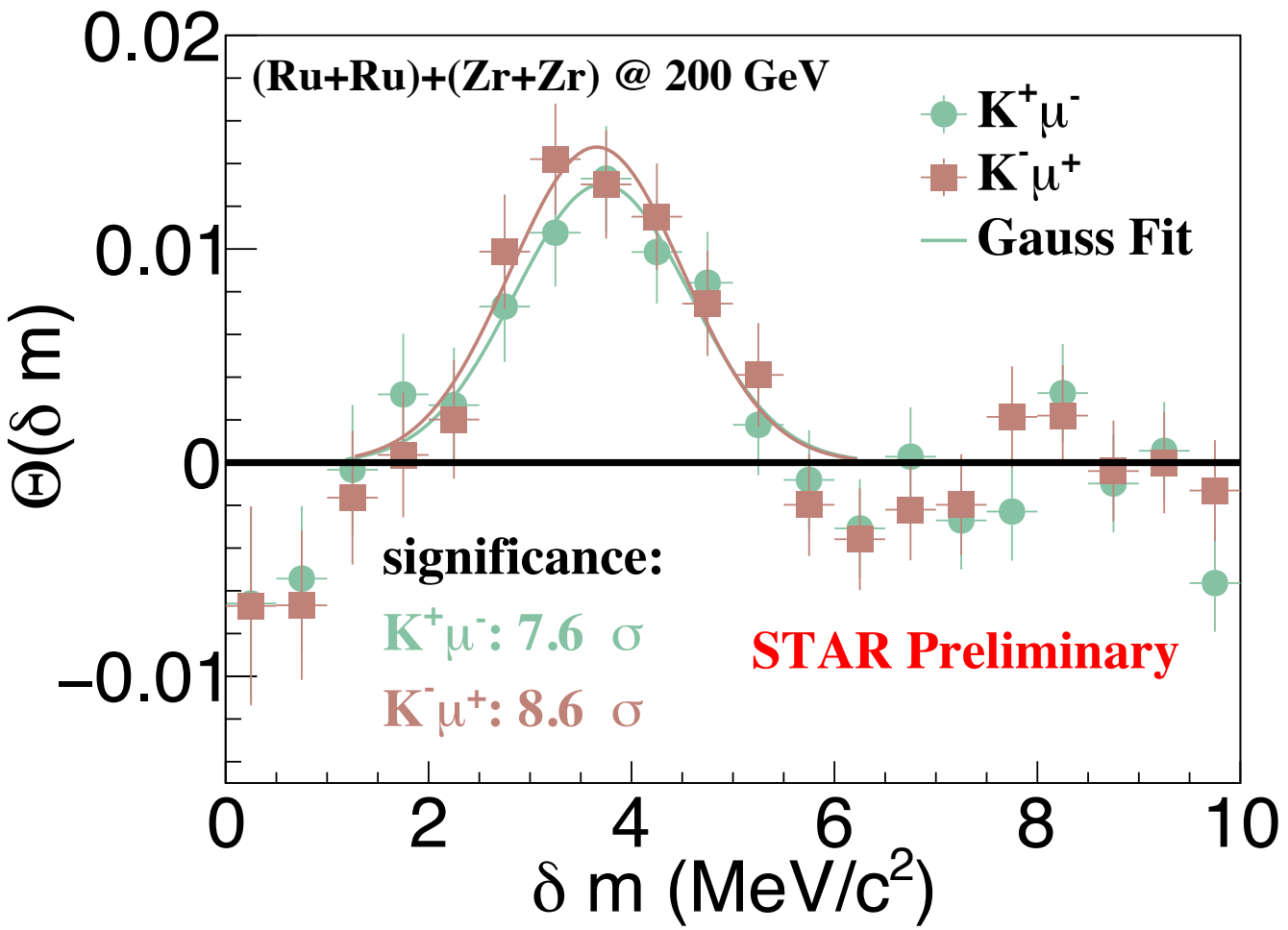
- $\delta m \equiv M_{inv} - m_{\mu} - m_{hadron}$
- A clear yield hierarchy: (same event unlike-sign) > (mixed-event unlike-sign) > (same event like-sign)
- Observed **sharp peaks** for pairs at expected $\delta m \sim 0$ from both background subtraction methods
- Flat at higher mass (0.02~0.1 GeV/c²)
- Like-Sign background has repulsive Coulomb contribution, and thus underestimates the background, leading to a higher "signal" than Mixed-Event

K- π Correlation Function (primary track)



- $\mathrm{C}(k^*) = \frac{k^* \text{ in same event}}{k^* \text{ in mixed event}}$, k^*: relative momentum of K-π
- Unlike-sign: Coulomb attraction → enhanced correlation
- Like-sign: Coulomb repulsion → opposite correlation
- Unity product: dominant Coulomb correlation removed in this method

Signal-to-Background Ratio (Θ, primary track)



- $\Theta \equiv \frac{(\mathrm{ULS}\times\mathrm{LS})_{same\ event}}{(\mathrm{ULS}\times\mathrm{LS})_{mixed\ event}} - 1$
- Assume $K\pi$ atom does not exist
- Subtract the $K\pi$ $\Theta(\delta m)$ distribution from the $K\mu$ $\Theta(\delta m)$ distribution to account for residual Coulomb effects
- Significant $K^+\mu^-(7.6\sigma)$ and $K^-\mu^+(8.6\sigma)$ atomic peaks observed
- Using primary tracks shift the mass peak away from zero because breakup happens away from primary collision vertex

$K\mu$ Atom Dissociation Fraction

S. Mrowczynski, Phys. Rev. A 33, 1549–1555 (1986)
X. Wang et al., Arxiv: 2603.09327

Dissociation fraction: $f = 1 - e^{-n\sigma l}$

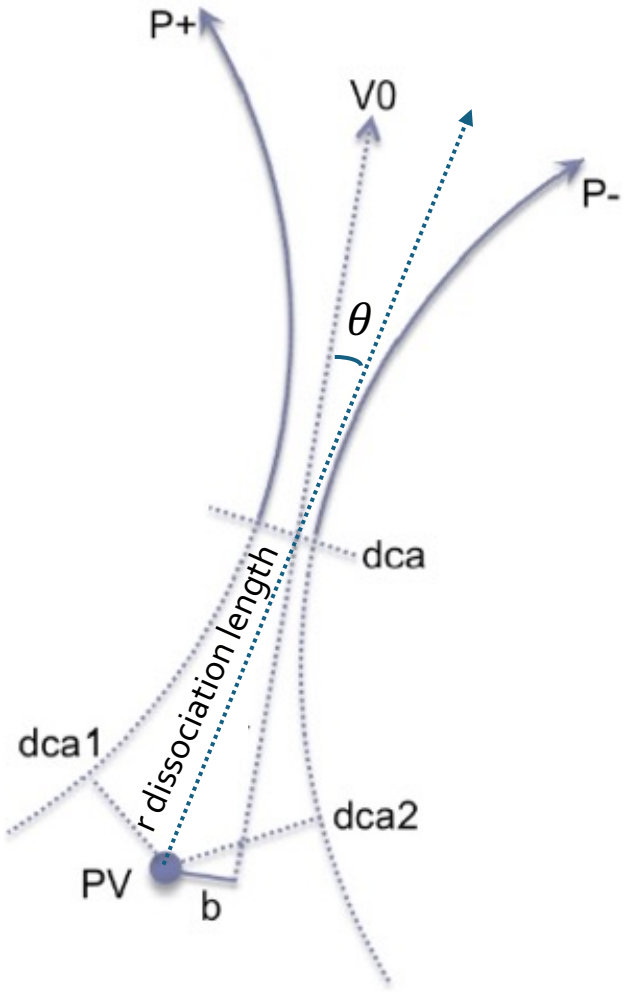
- n: material number density,
- l: material thickness
- σ: atom–material interaction cross section
- Using $k\pi$ dissociation cross sections:

$$1.4\times10^{-22}\ \mathrm{cm}^2$$

	RHIC-STAR
Beampipe (Z, n, l)	Be(4, 1.236×10^{23} atoms/cm^3, 0.1 cm)
Dissociation at beampipe	53.5%
Air (Z, n, l)	(7.2, 2.54×10^{19} atoms/cm^3, 46.19 cm)
Dissociation in air	$(1 - 53.5\%) \times 20.9\%$
First-layer detector (Z, n, l)	Inner field cage (Kapton, $0.6\%X_0$)
Total dissociation before deeper tracking	$\simeq 100\%$

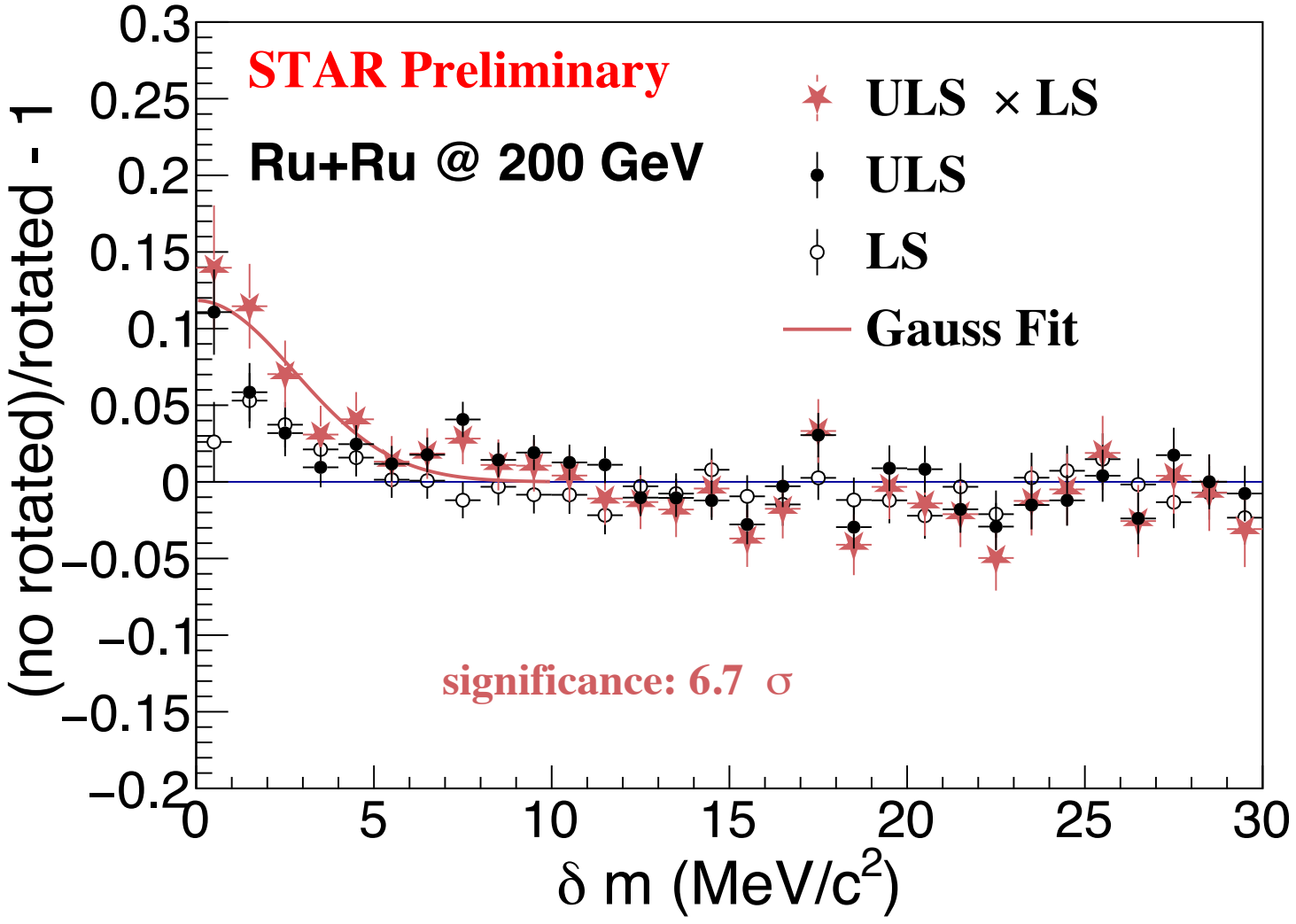
- Muonic Atom interaction with material calculated
- Breakup at several locations away from collisions
- **New signature: breakup at secondary vertex**
- Global Tracks reconstruct the breakup vertex with correct momenta

Topological Reconstruction (global track)



- Global track: tracks fitted without being constrained to the primary vertex
- Topological cuts reduce combinatorial background:
  - ✓ dca2(μ)>1.1 cm;
  - ✓ dca1(K) > 0.75 cm;
  - ✓ dca<1 cm;
  - ✓ Cos(θ)>0.99;
  - ✓ b(V0 dca) < 3cm;
  - ✓ 10 cm < r dissociation length < 50 cm

Signal-to-Background Ratio (global track)



- δm peaks at zero
- Rotated kaon track 180^0, background
- Next steps: optimize topological cuts to suppress background and improve signal extraction

Summary and Outlook

Summary

- First significant $K\mu$ atom signals observed in isobar collisions
  - ✓ Significant mass peaks shifted away from zero (primary tracks)
  - ✓ Significant mass peak at zero (global tracks and secondary vertex)

- Consistent with $K\mu$ atoms formation and breakup in detector material

Outlook 1:

- Topological cut optimization needed to improve background suppression
- Extract muonic atom and primordial muon yields
- More detailed studies are ongoing to improve the understanding and quantification of the background

Outlook 2: $K\mu$ Atom- QGP Coalescence vs. D^0 Decay

X. Wang et al., Arxiv: 2603.09327

TABLE I. Total muonic atom yields at different facilities. RHIC Au+Au collisions correspond to STAR/sPHENIX running. LHC p+p assumes an effective trigger strategy for charm production. STCF assumes the integrated luminosity accumulated over 20 years of operation. The QGP-coalescence yields are taken from Ref. [14].

	RHIC (Au+Au)	LHC (p+p)	LHC (A+A)	STCF
Luminosity	20–50 billion events	3000 fb^{-1}	8550 μb^{-1}	1 ab^{-1}
D^0 counts	10^{10}	2×10^{15}	4×10^9	8×10^{10}
$(K^{\pm}\mu^{\mp})_{\text{atom}}$ from D^0 decay	~ 10	5×10^5	1	~ 20
$(K^{\pm}\mu^{\mp})_{\text{atom}}$ from QGP coalescence	10^5	N/A	10^5	N/A

$$D^0 \rightarrow (K\mu)_{\text{atom}}\,\nu_\mu$$

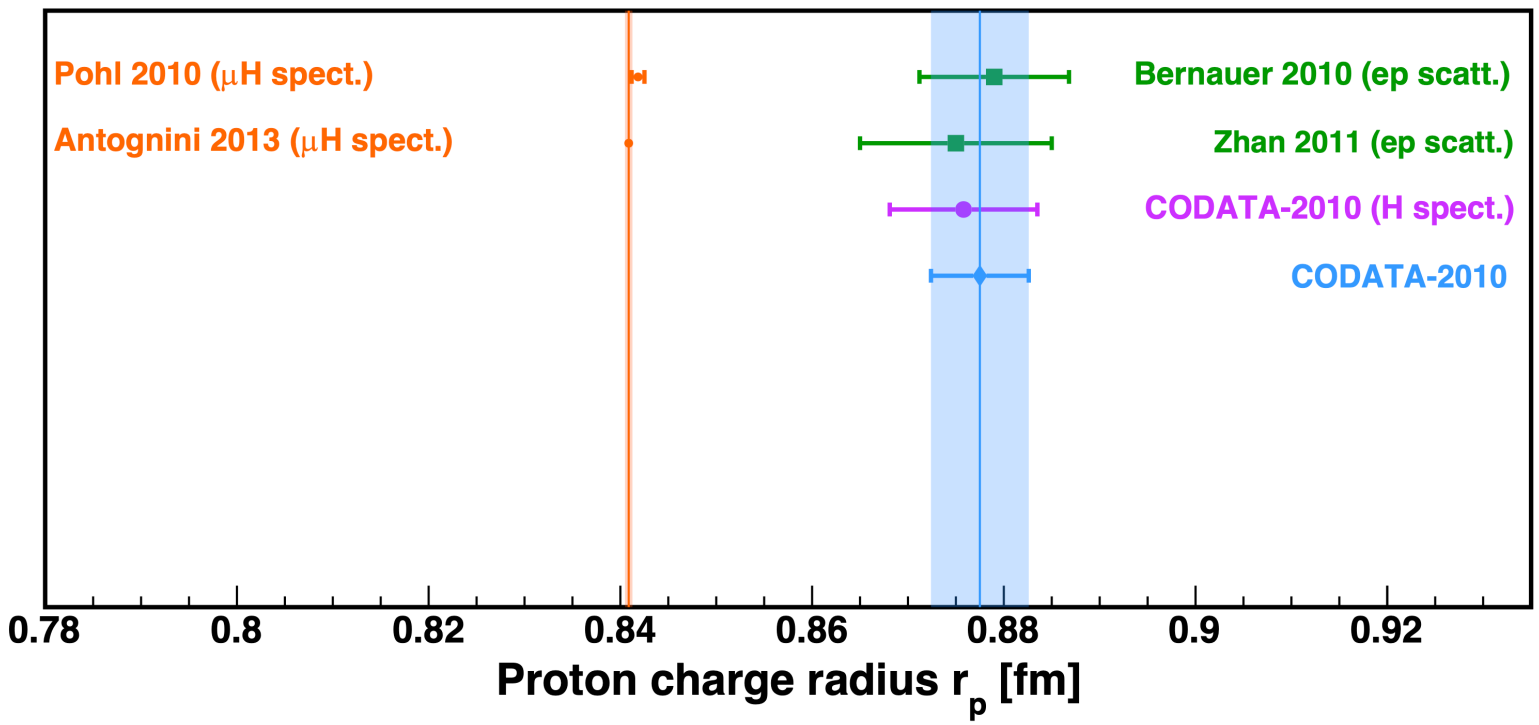
$$\text{BR}\left(D^0 \rightarrow (K\mu)_{nS\,\text{all}}\,\nu_\mu\right) = 2.29 \times 10^{-10}$$

- Based the effective weak Hamiltonian, a helicity-based treatment of the leptonic current, and a nonrelativistic bound-state projection

Thanks for your attention!

Backup

Proton charge radius puzzle





- CODATA world avg.
- H spectroscopy

- ep scattering
- μH spectroscopy

- Muonic hydrogen enhances sensitivity
- Testbed for Fundamental Physics
  - The Rydberg constant?
  - New physics beyond the standard model?