

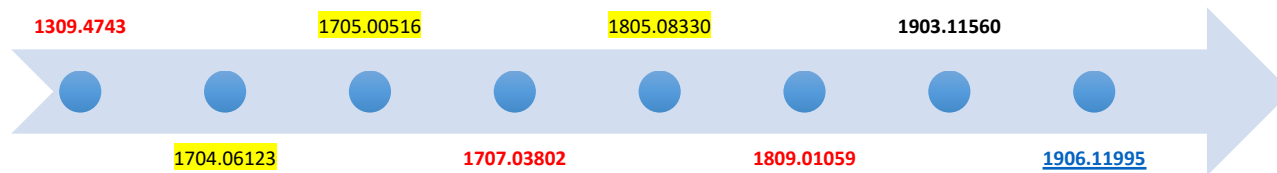


北京航空航天大学
BEIHANG UNIVERSITY



DK interaction as a doorway to manifestly exotic mesons

Lisheng Geng @ Beihang U.



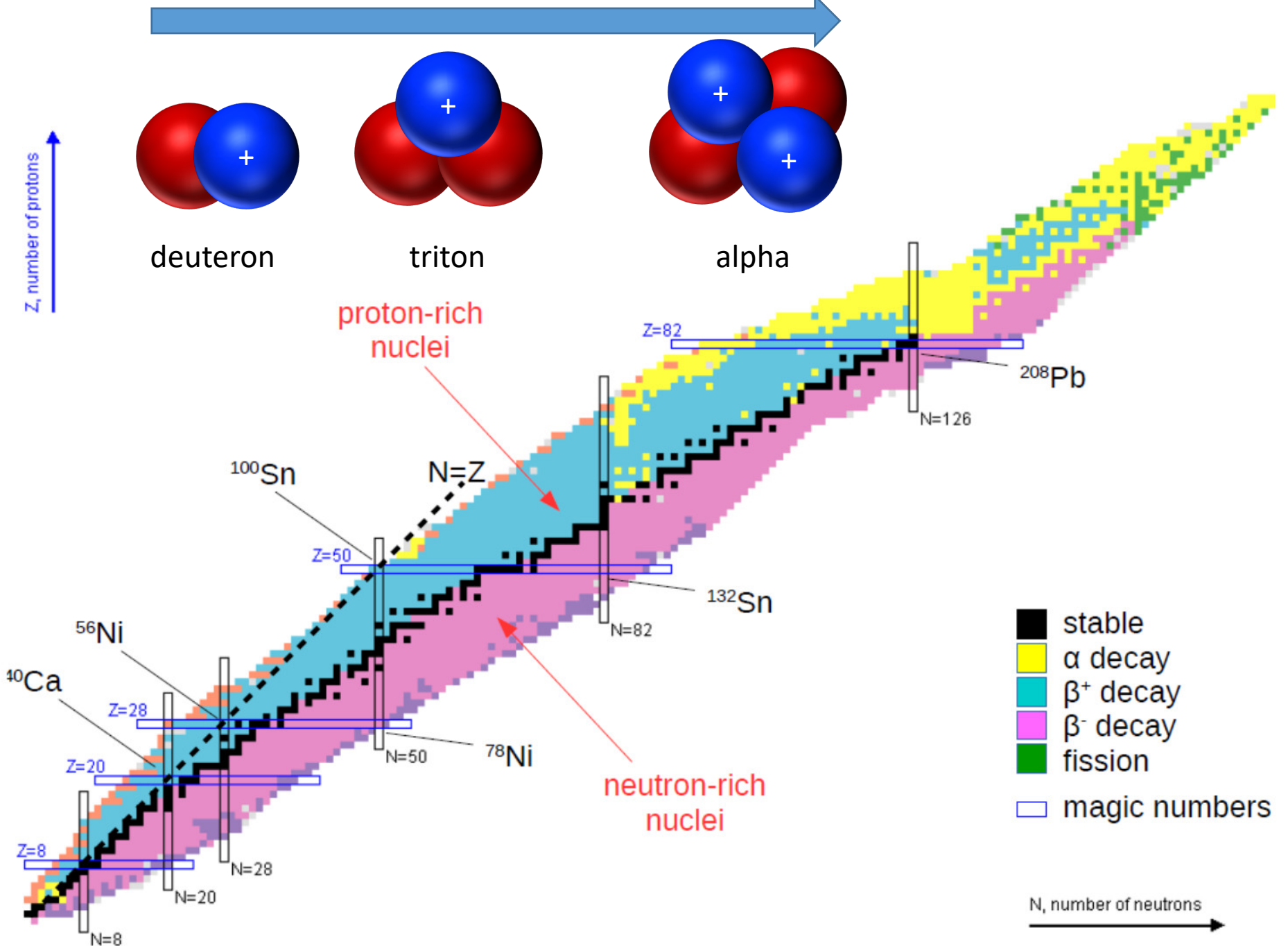
+some preliminary results ₁

Contents

- Motivation: new types of clusters of color singlets in addition to nuclei
- $Ds0^*(2317)$ and $Ds1(2460)$ as DK/D^*K molecules: theory & lattice
- Explicit studies of the DDK & $DDDK$ systems
- A $K^*(4307)$ with hidden charm as $KX(3872)/Z_c(3900)$ molecule
- Summary and outlook

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The existence of triton can be inferred from that of deuteron with reasonable confidence

“In nature, are there other clusters of color singlet hadrons, similar to atomic nuclei, bound by the residual strong force? ? ? ”

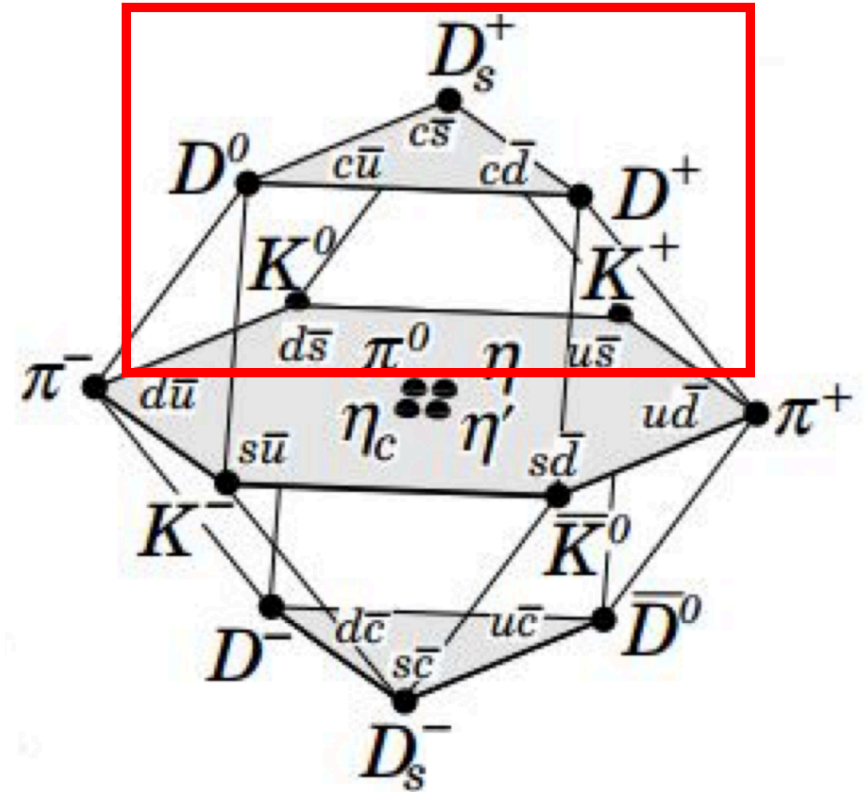


DK/DDK/DDDK molecules???

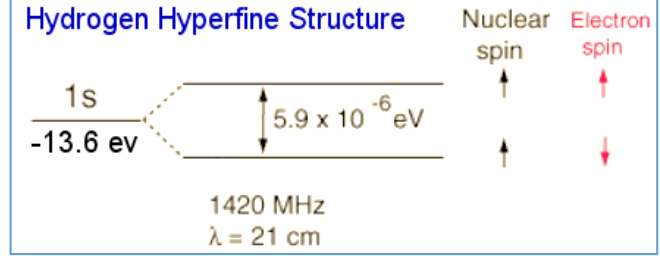
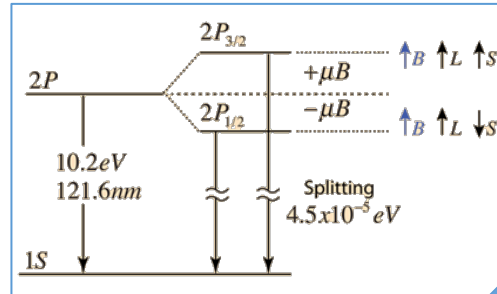
Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III		
mass $\approx 2.2 \text{ MeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ u up	mass $\approx 1.28 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ c charm	mass $\approx 173.1 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ t top	mass 0 charge 0 spin 1 g gluon	mass $\approx 125.09 \text{ GeV}/c^2$ charge 0 spin 0 H higgs
mass $\approx 4.7 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ d down	mass $\approx 96 \text{ MeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ s strange	mass $\approx 4.18 \text{ GeV}/c^2$ charge $-\frac{1}{3}$ spin $\frac{1}{2}$ b bottom	mass 0 charge 0 spin 1 γ photon	
mass $\approx 0.511 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$ e electron	mass $\approx 105.66 \text{ MeV}/c^2$ charge -1 spin $\frac{1}{2}$ μ muon	mass $\approx 1.7768 \text{ GeV}/c^2$ charge -1 spin $\frac{1}{2}$ τ tau	mass $\approx 91.19 \text{ GeV}/c^2$ charge 0 spin 1 Z Z boson	
mass $< 2.2 \text{ eV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_e electron neutrino	mass $< 1.7 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_μ muon neutrino	mass $< 15.5 \text{ MeV}/c^2$ charge 0 spin $\frac{1}{2}$ ν_τ tau neutrino	mass $\approx 80.39 \text{ GeV}/c^2$ charge ± 1 spin 1 W W boson	

QUARKS (left side of fermion table)
LEPTONS (left side of fermion table)
SCALAR BOSONS (right side of boson table)
GAUGE BOSONS VECTOR BOSONS (right side of boson table)



Why spectroscopy—Atomic



$$\frac{1}{\lambda_{\text{vac}}} = R \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

Fine structure

Hyperfine structure



Johannes Rydberg

Rydberg formula



Niels Henrik David Bohr



Edward W. Morley



Arnold Sommerfeld



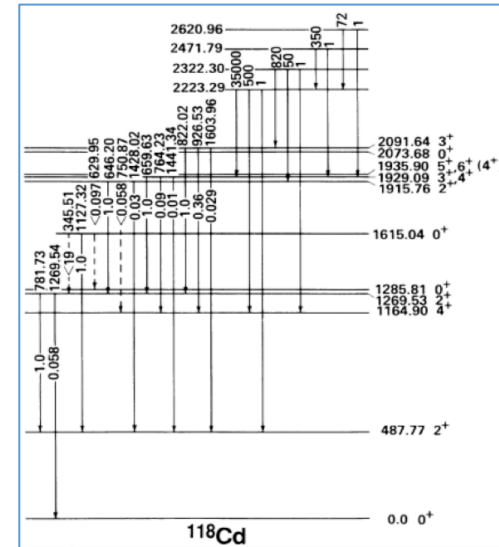
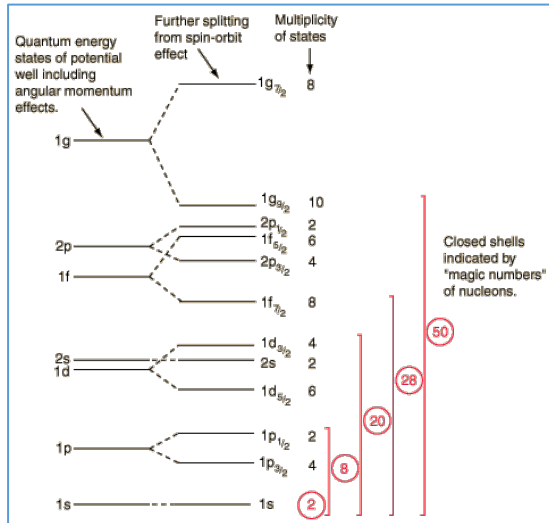
Albert Abraham Michelson



Wolfgang Pauli

Why spectroscopy—Nuclear

Single
particle
motion



Collective
motion



Maria Goeppert Mayer J. Hans D. Jensen



Aage Niels Bohr, Ben Roy Mottelson Leo James Rainwater

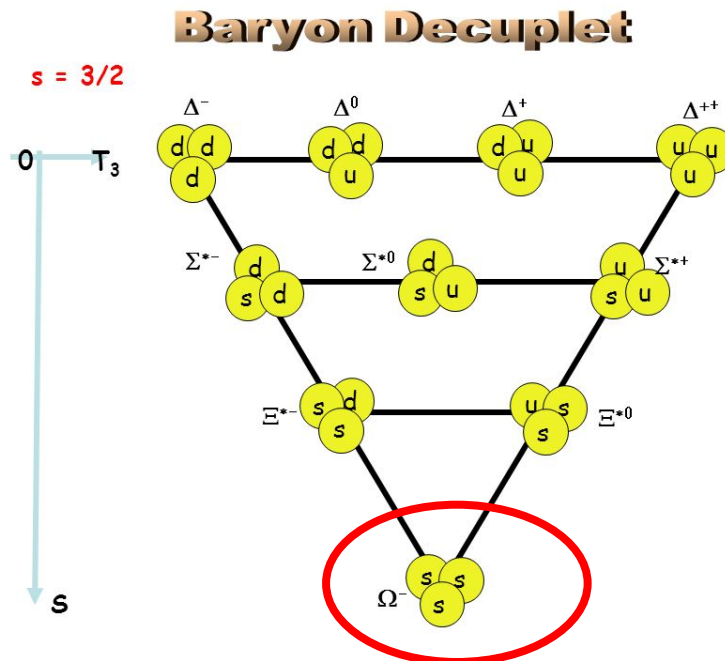
Why spectroscopy—particle/hadron



Murray Gell-Mann



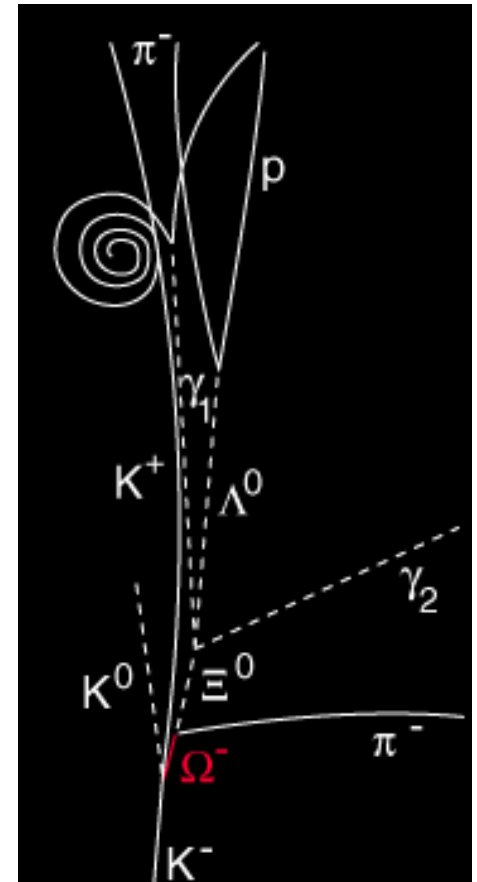
Yuval Ne'eman.



Eightfold
way

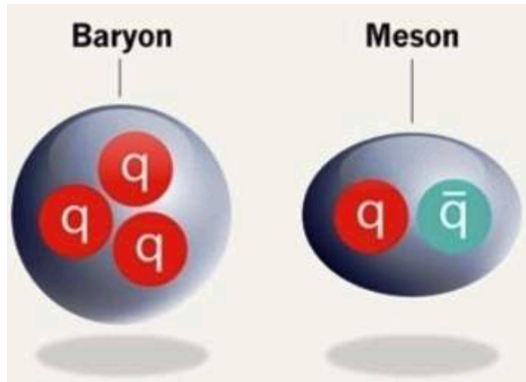
Quark
model

QCD/SM



V. E. Barnes et al., Phys. Rev. Lett. 12, 204 (1964)

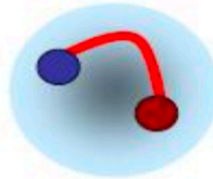
Why spectroscopy—particle/hadron



In the **naïve quark model**

In principle,
QCD **allows**

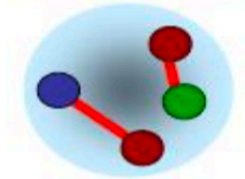
Hybrid



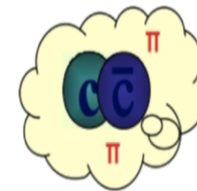
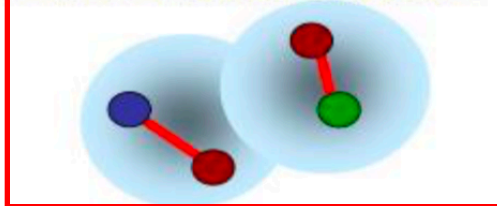
Glueball



Tetraquark

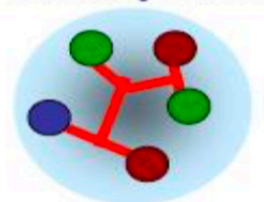


Hadronic molecule



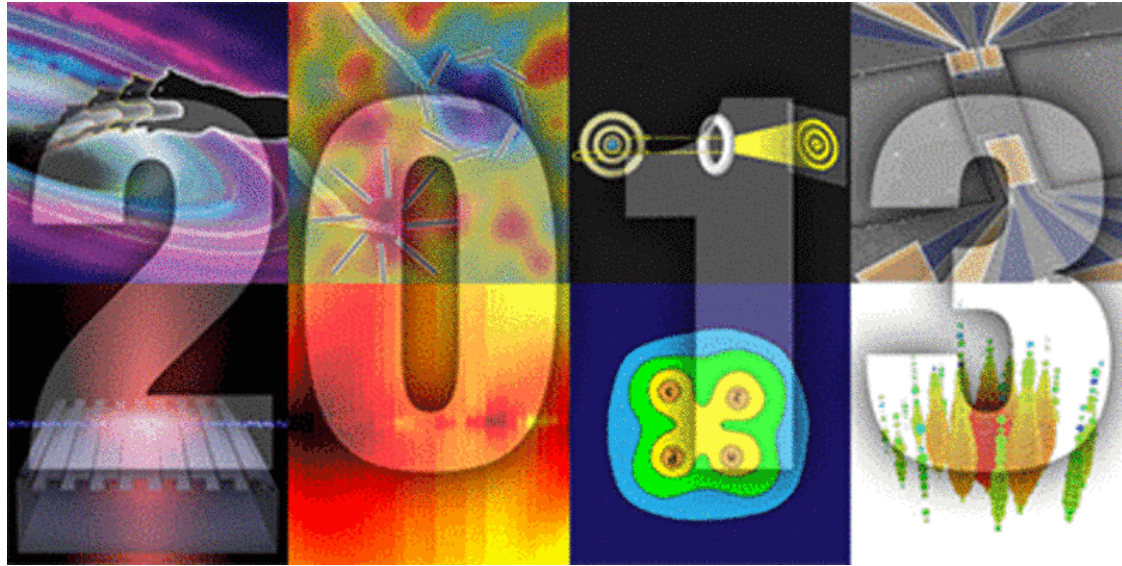
Hadro-
quarkonium

Pentaquark



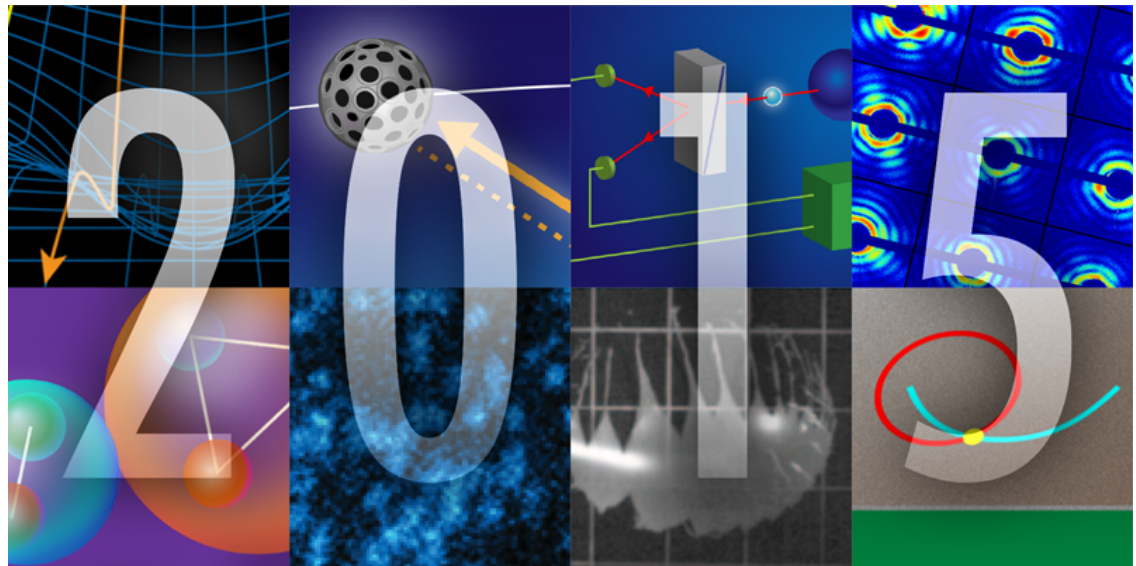
Highlights of the year

the research covered in Physics that **really made waves in and beyond the physics community.**



Four-Quark Matter

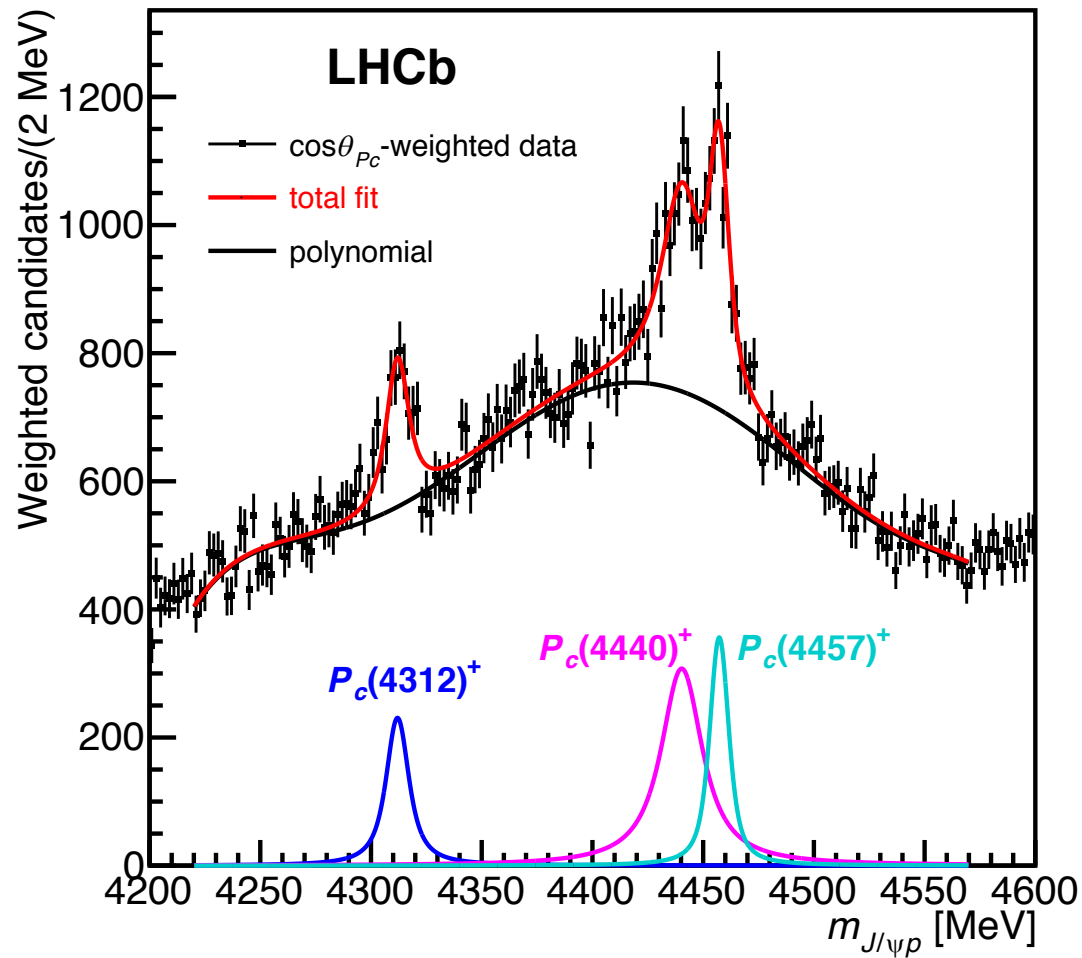
Particle High Five



Latest LHCb discovery:

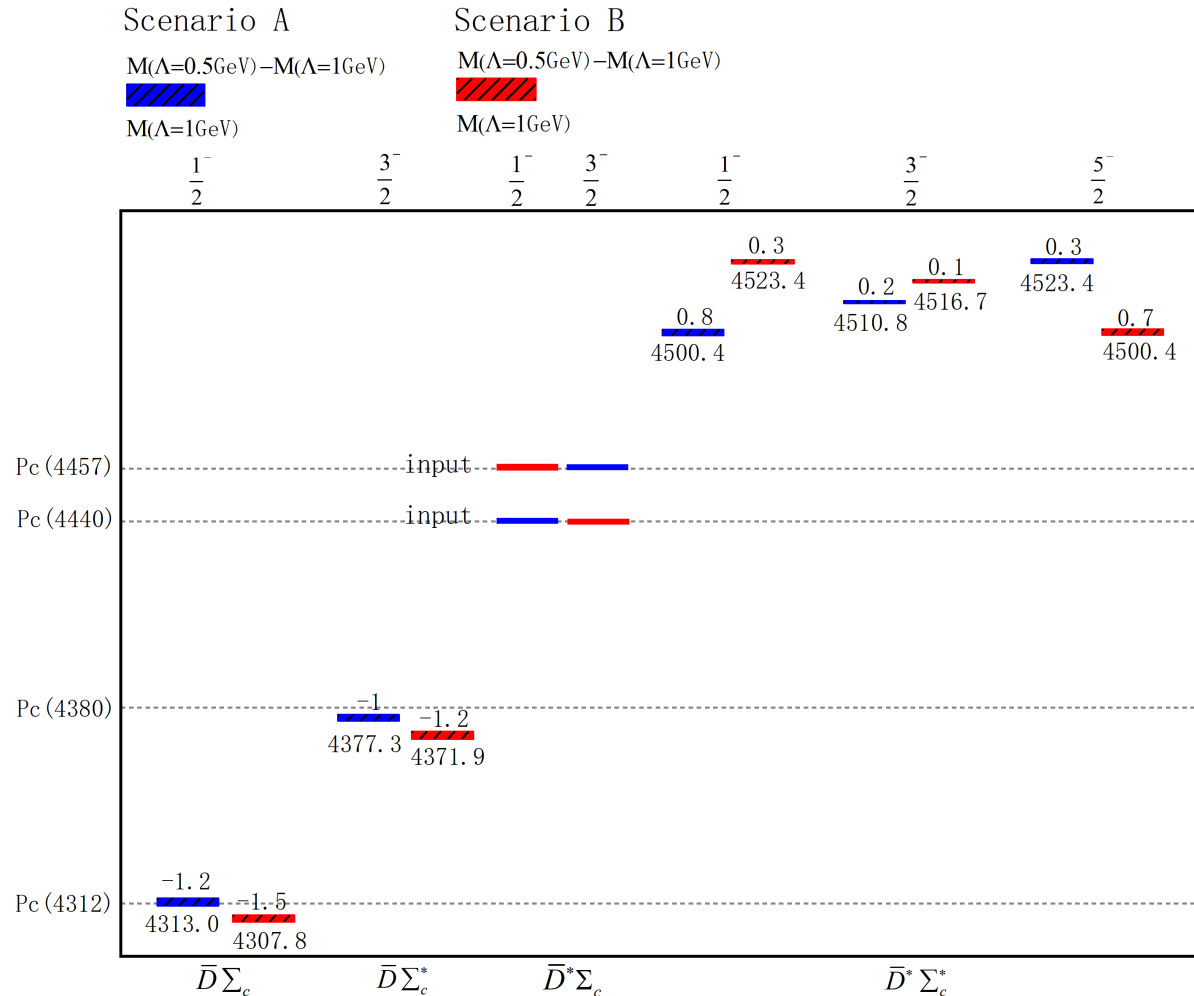


Fine structure observed by LHCb, 1904.03947

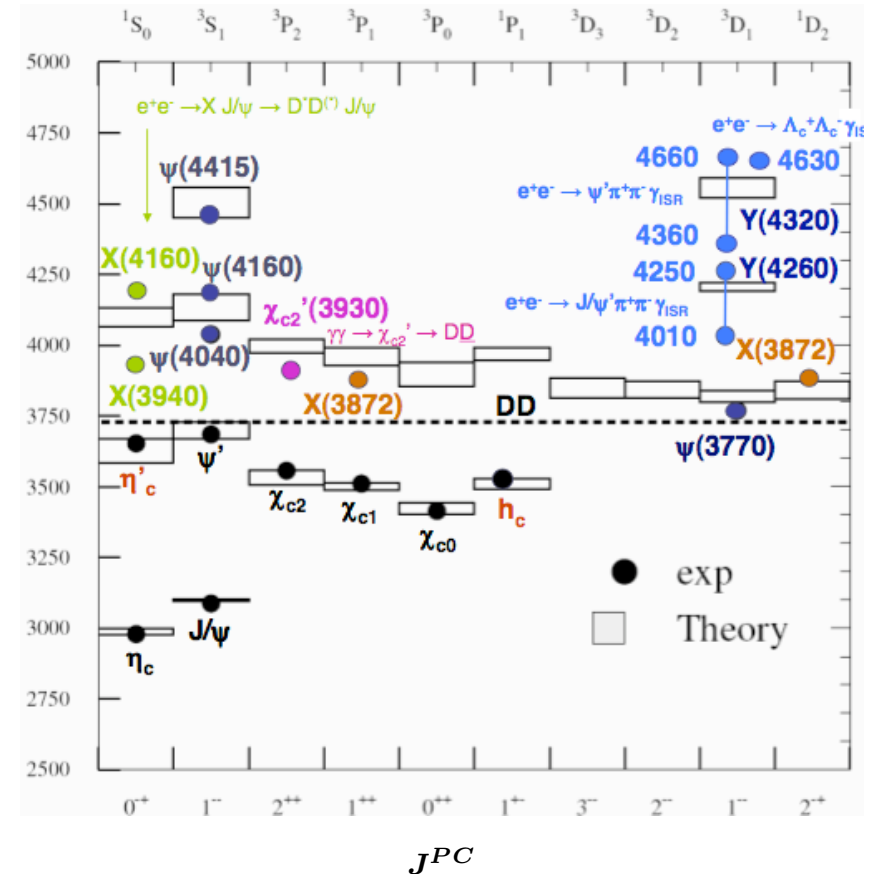
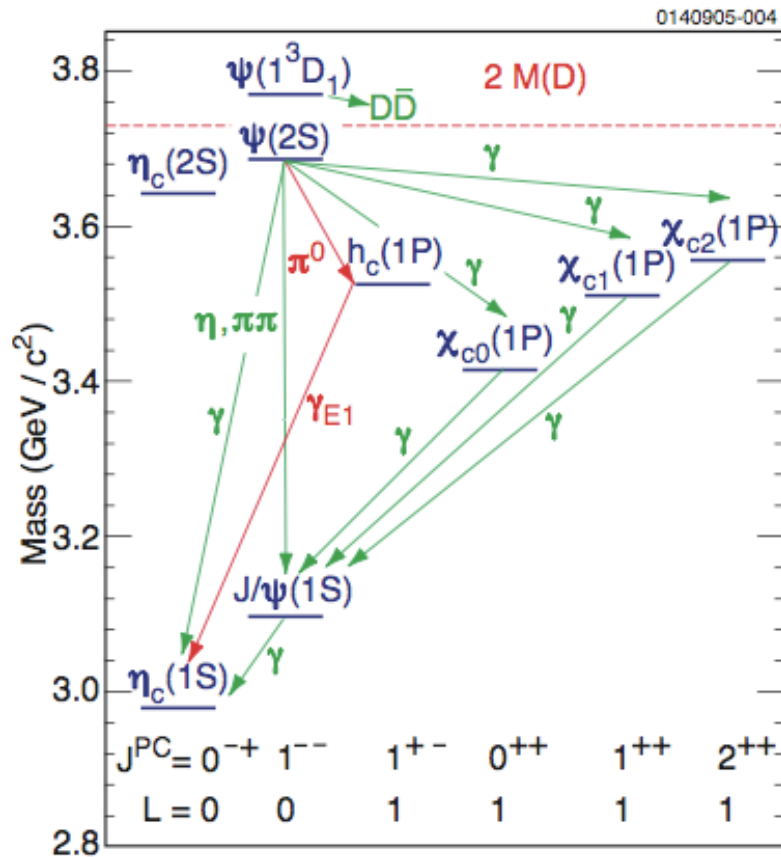


Emergence of a complete heavy-quark spin symmetry multiplet

1903.11560

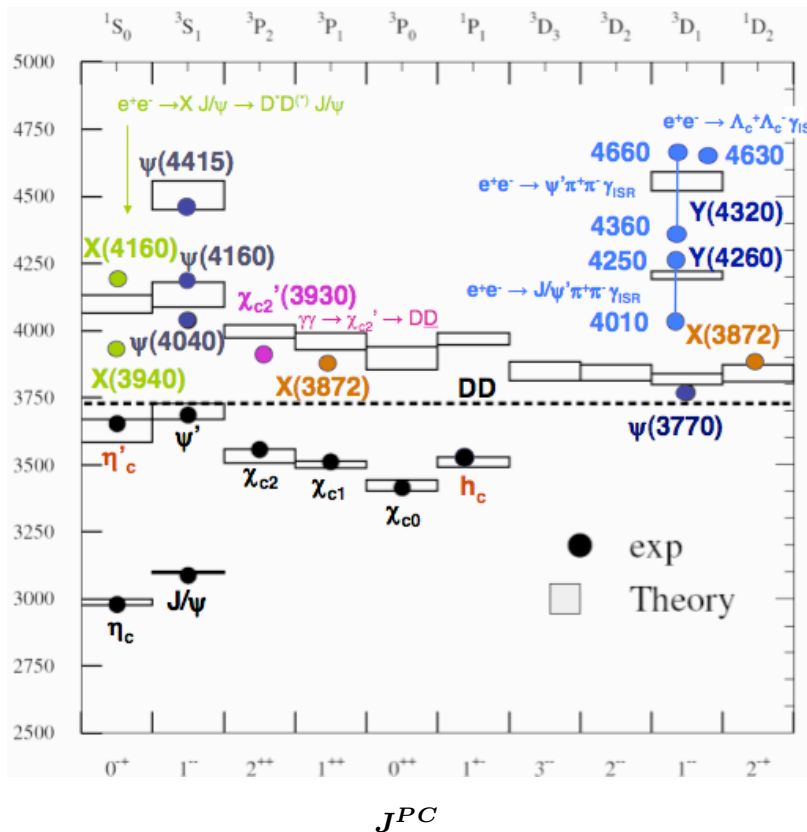


Charmonium spectroscopy before/after B factories



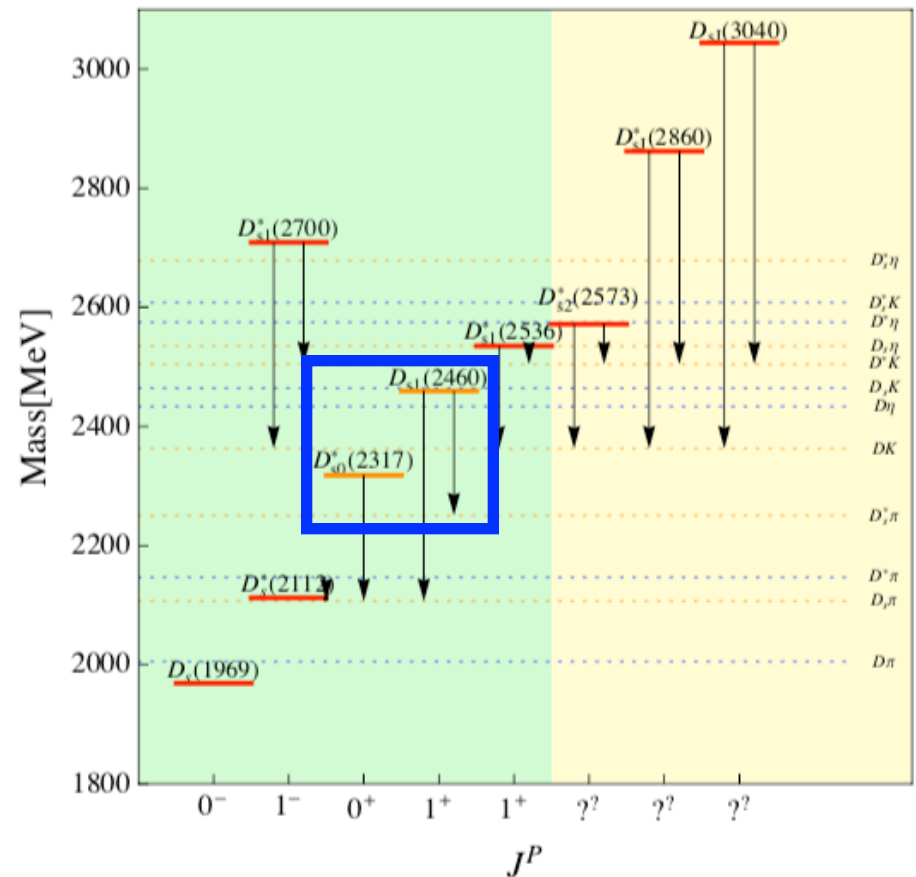
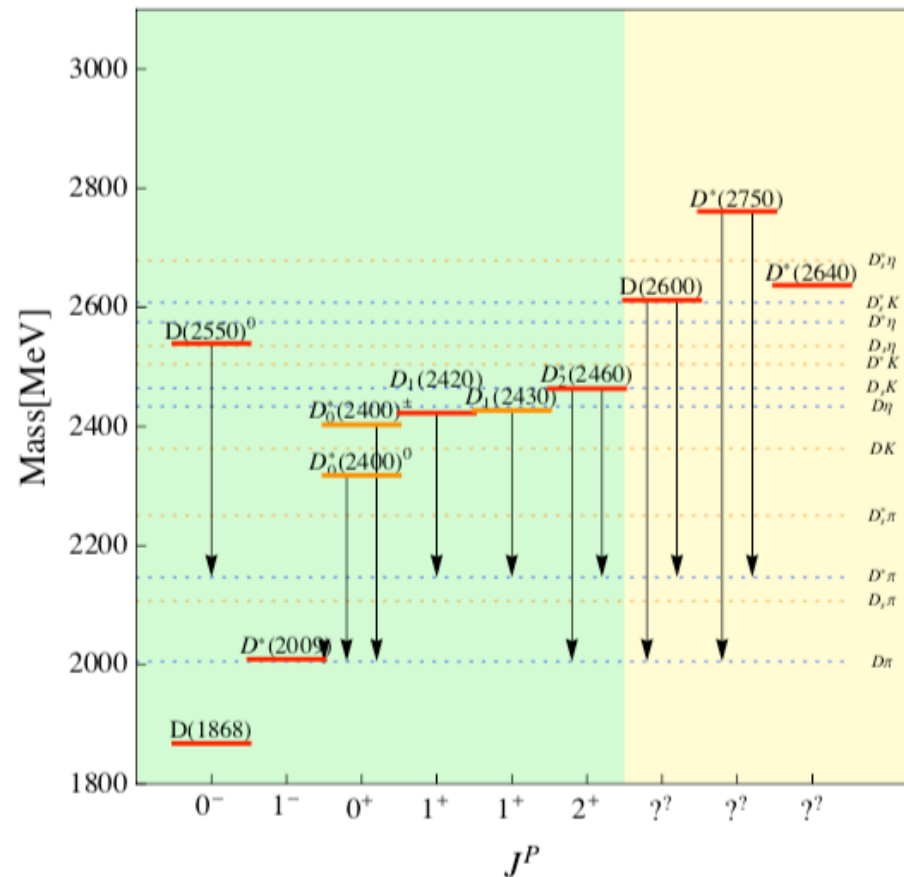
- A lot of new states, but the nature of some states are not well understood, e.g. **Y(4260)**, **X(3872)**, **Z_c(3900)**
- In complete contrast to the before-B-factory period, when potential models worked quite well.

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Channel	Threshold Energy (MeV)
$D^0 \bar{D}^0$	3729.4
$D^+ D^-$	3738.8
$D^0 \bar{D}^{*0}$ or D^*	3871.5
$\rho^0 J/\psi$	3872.7
$D^\pm D^{*\mp}$	3879.5
$\omega^0 J/\psi$	3879.6
$D_s^+ D_s^-$	3936.2
$D^{*0} \bar{D}^{*0}$	4013.6
$D^{*+} D^{*-}$	4020.2
$\eta' J/\psi$	4054.7
$f^0 J/\psi$	≈ 4077
$D_s^+ \bar{D}_s^{*-}$ or D_s^*	4080.0
$a^0 J/\psi$	4081.6
$\varphi^0 J/\psi$	4116.4
$D_s^{*+} D_s^{*-}$	4223.8

Open charm systems

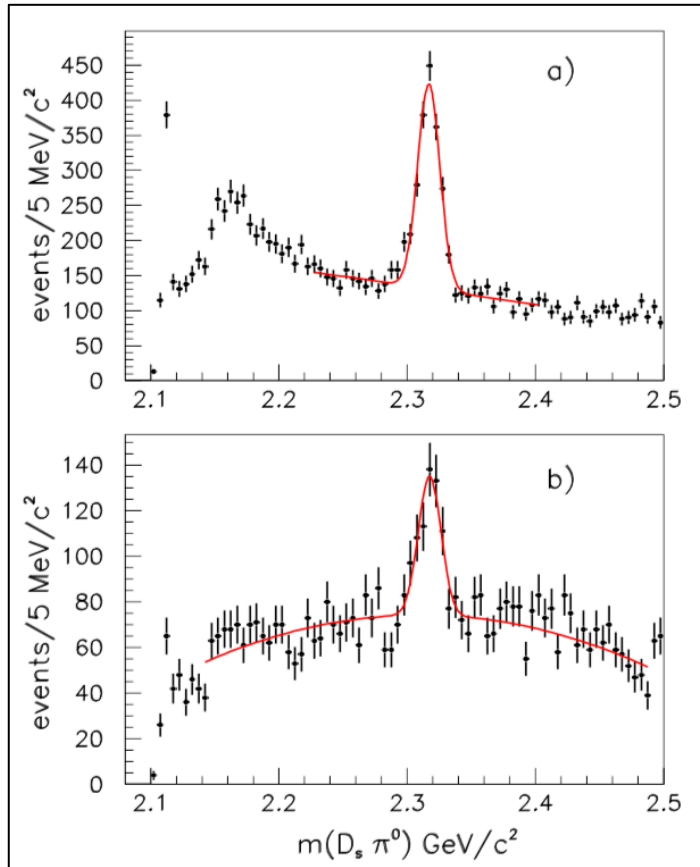


$D_{s0}^*(2317)$

$D_{s1}(2460)$

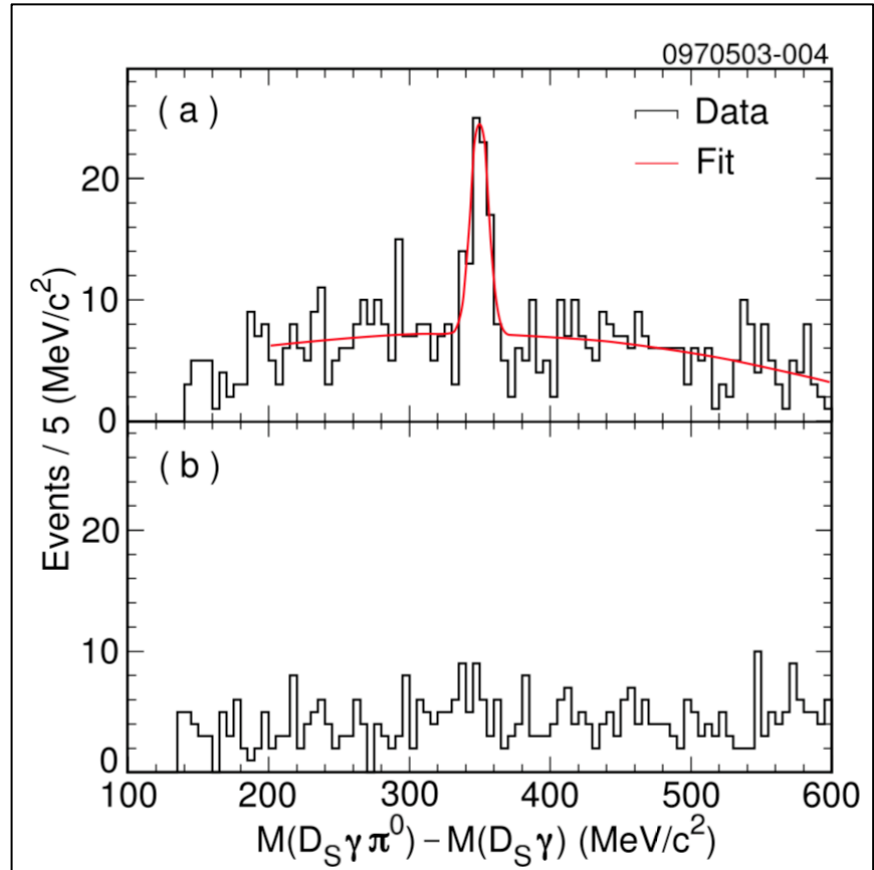
Two peculiar states

Ds0*(2317)



BaBar PRL90,242001(2003)

Ds1(2460)



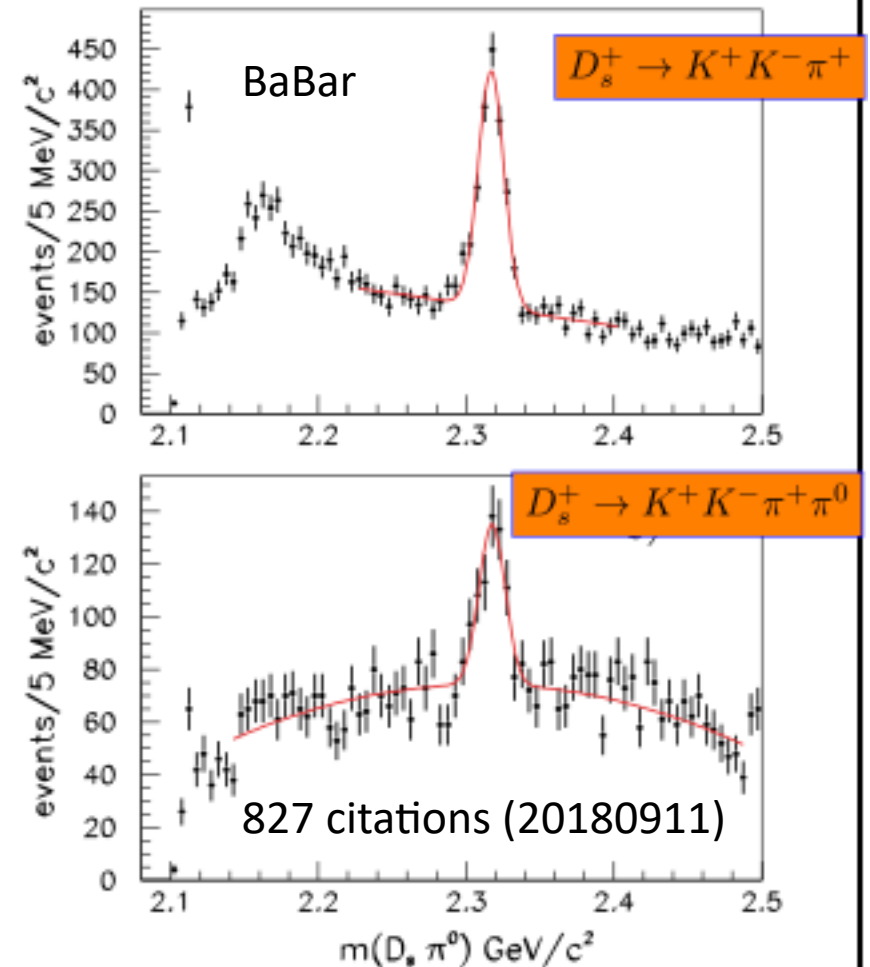
CLEO PRD68,032002(2003)

Two peculiar states

- $D_{s0}^*(2317)$, $D_{s1}(2460)$
- 160/70 MeV lower than quark model predictions--difficult to be understood as conventional $c\bar{s}$ bar states.
- “Dynamically generated” from strong DK interaction
 - ✓ E. E. Kolomeitsev 2004,
 - ✓ F. K. Guo 2006,
 - ✓ D. Gamermann 2007

$$m_{D_{s1}(2460)} - m_{D_{s0}^*(2317)} \approx m_{D^*} - m_D$$

$M = 2317.8 \pm 0.6$ and $\Gamma < 3.8$ MeV

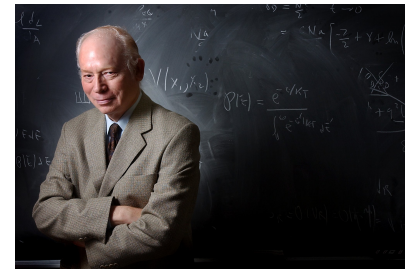


PRL90,242001(2003)

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UChPT in Bethe-Salpeter equation

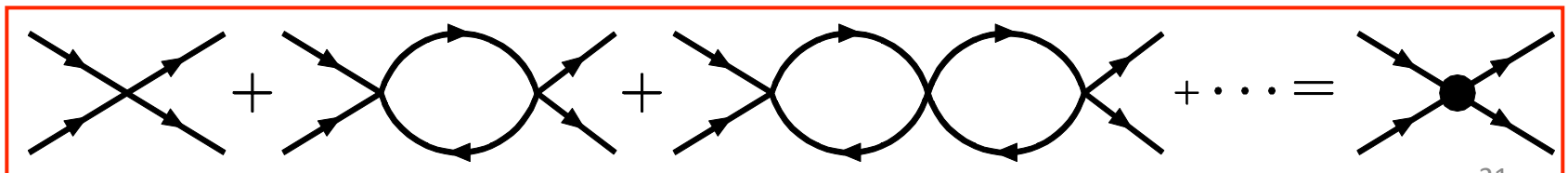


□ Model independent DK interaction from ChPT

$$\mathcal{V}_{\text{WT}}(P(p_1)\phi(p_2) \rightarrow P(p_3)\phi(p_4)) = \frac{1}{4f_0^2} \mathcal{C}_{\text{LO}} (s - u) \quad \text{Weinberg-Tomazawa}$$

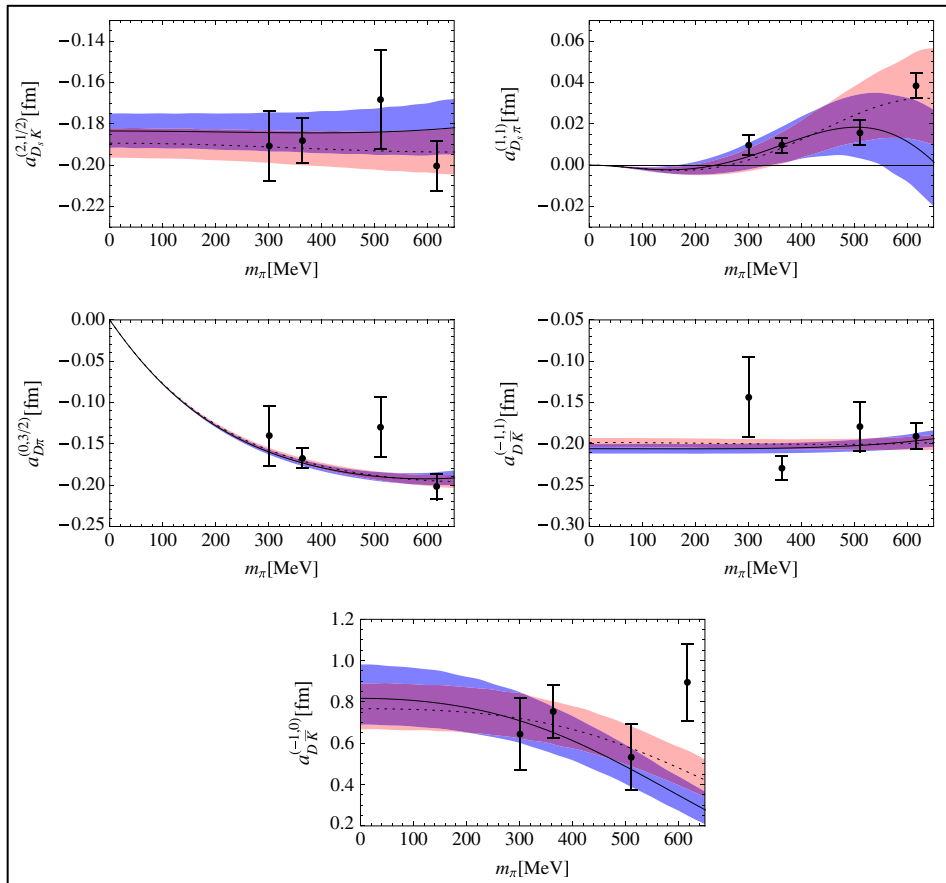
$$\begin{aligned} \mathcal{V}_{\text{NLO}}(P(p_1)\phi(p_2) \rightarrow P(p_3)\phi(p_4)) = & -\frac{8}{f_0^2} C_{24} \left(c_2 p_2 \cdot p_4 - \frac{c_4}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 + p_1 \cdot p_2 p_3 \cdot p_4) \right) \\ & -\frac{4}{f_0^2} C_{35} \left(c_3 p_2 \cdot p_4 - \frac{c_5}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 + p_1 \cdot p_2 p_3 \cdot p_4) \right) \\ & -\frac{4}{f_0^2} C_6 \frac{c_6}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 - p_1 \cdot p_2 p_3 \cdot p_4) \\ & -\frac{8}{f_0^2} C_0 c_0 + \frac{4}{f_0^2} C_1 c_1, \end{aligned} \quad (11)$$

□ Resummed in the Bethe-Salpeter equation (two-body elastic unitarity)

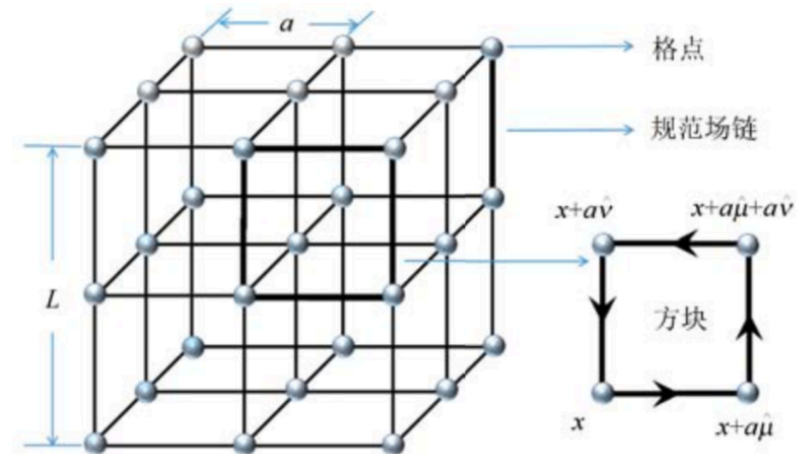


Fixing the LECs using latest LQCD* data

Liuming Liu et al., PRD87 (2013) 014508



- NLO ChPT kernel: 5 LECs
- A quite good description of the 20 Lattice **scattering lengths of pseudoscalar mesons and D mesons** (I=0 DK excluded) can be achieved.



Ds0 and Ds1 dynamically generated

● Charm sector

“Post-diction”

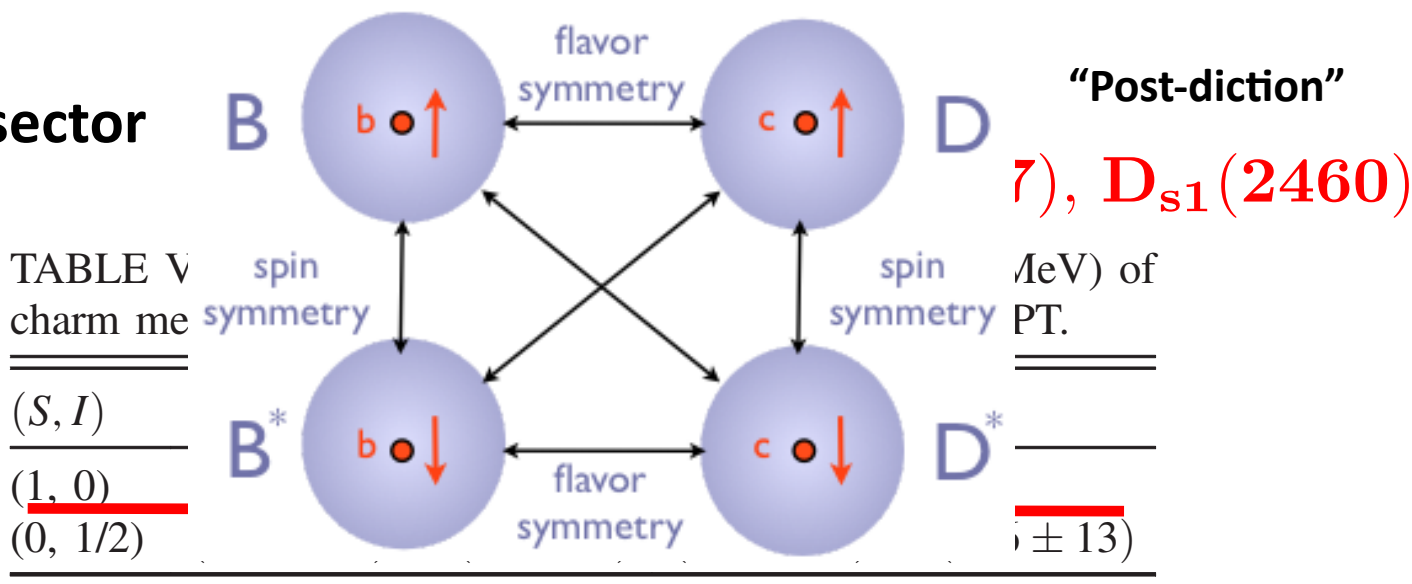
$D_{s0}^*(2317)$, $D_{s1}(2460)$

TABLE V. Pole positions $\sqrt{s} = M - i\frac{\Gamma}{2}$ (in units of MeV) of charm mesons dynamically generated in the HQS UChPT.

(S, I)	$J^P = 0^+$	$J^P = 1^+$
$(1, 0)$	2317 ± 10	2457 ± 17
$(0, 1/2)$	$(2105 \pm 4) - i(103 \pm 7)$	$(2248 \pm 6) - i(106 \pm 13)$

Ds0 and Ds1 dynamically generated

● Charm sector



● Bottom Sector

TABLE VI. Pole positions $\sqrt{s} = M - i\frac{\Gamma}{2}$ (in units of MeV) of bottom mesons dynamically generated in the HQS UChPT.

(S, I)	$J^P = 0^+$	$J^P = 1^+$
(1, 0)	5726 ± 28	5778 ± 26
(0, 1/2)	$(5537 \pm 14) - i(118 \pm 22)$	$(5586 \pm 16) - i(124 \pm 25)$

Predicted Bs0 and Bs1 states

Physics Letters B 750 (2015) 17–21



Contents lists available at ScienceDirect

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Predicting positive parity B_s mesons from lattice QCD



C.B. Lang^a, Daniel Mohler^{b,*}, Sasa Prelovsek^{c,d}, R.M. Woloshyn^e

^a Institute of Physics, University of Graz, A-8010 Graz, Austria

^b Fermi National Accelerator Laboratory, Batavia, IL 60510-5011, USA

^c Department of Physics, University of Ljubljana, 1000 Ljubljana, Slovenia

^d Jozef Stefan Institute, 1000 Ljubljana, Slovenia

^e TRIUMF, 4004 Wesbrook Mall, Vancouver, BC V6T 2A3, Canada

Table 5

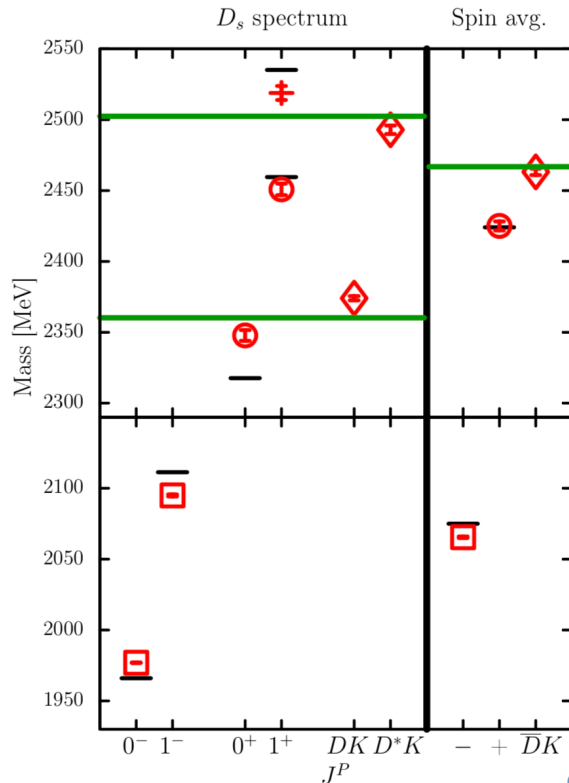
Comparison of masses from this work to results from various model based calculations; all masses in MeV.

J^P	0^+	1^+
Covariant (U)ChPT [24]	5726(28)	5778(26)
NLO UHMChPT [19]	5696(20)(30)	5742(20)(30)
LO UChPT [17,18]	5725(39)	5778(7)
LO χ -SU(3) [16]	5643	5690
HQET + ChPT [20]	5706.6(1.2)	5765.6(1.2)
Bardeen, Eichten, Hill [15]	5718(35)	5765(35)
rel. quark model [5]	5804	5842
rel. quark model [22]	5833	5865
rel. quark model [23]	5830	5858
HPQCD [30]	5752(16)(5)(25)	5806(15)(5)(25)
this work	5713(11)(19)	5750(17)(19)

In agreement with IQCD

Support from lattice QCD studies

- [G. S. Bali et al., arXiv:1706.01247 \[hep-lat\]](#).
- [C. B. Lang et al., arXiv:1403.8103 \[hep-lat\]](#).
- [D. Mohler et al., arXiv:1308.3175 \[hep-lat\]](#).



“DK components substantial”

FIG. 12. On the left, our final results for the lower lying D_s spectrum as detailed in Table VII. The short horizontal black lines indicate the corrected experimental values (see Section II) while the green horizontal lines give the positions of the DK and D^*K non-interacting thresholds. Our lattice results for the finite volume thresholds are labelled DK and D^*K , respectively. The errors indicated are statistical only. On the right, the negative parity spin-averaged $1S$ mass $m_- = \frac{1}{4}(m_{0^-} + 3m_{1^-})$ is shown and denoted $-$, while the same spin-average of the positive parity 0^+ and 1^+ states is labelled with $+$ and the weighted average of the threshold is labelled as \overline{DK} .

Further **tests** of the DK interaction

- Experiments, theory, and lattice QCD all show that DK or D^*K interaction is strong enough to form the $Ds0^*(2317)$ or $Ds1(2460)$
- A natural question is: if we add one more $D(\bar{D})$ or $D^*(\bar{D}^*)$, can they form molecules of three hadrons?
- This seems to **be a rather straightforward and naive question**, but **remains unexplored** until quite recently

Contents

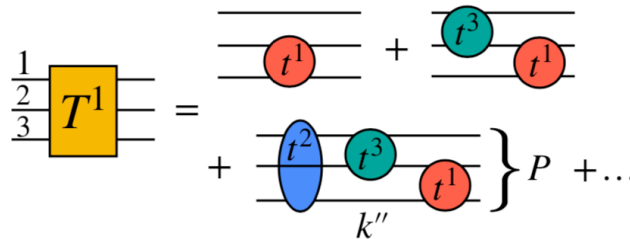
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An explicit three-body study of DDK

A. Martínez Torres, K. P. Khemchandani, and **LSG** 1809.01059

- Coupled-three-channel problem: $D(\text{DK} - D_S\pi - D_S\eta)$
- Three-body scattering matrix (Faddeev)

$$T = \sum_{i=1}^3 T^i$$



$$T^i = t^i \delta^3(\vec{k}'_i - \vec{k}_i) + \sum_{j \neq i=1}^3 T_R^{ij}, \quad i = 1, 2, 3,$$

$$T_R^{ij} = t^i g^{ij} t^j + t^i \left[G^{iji} T_R^{ji} + G^{ijk} T_R^{jk} \right],$$

A. Martínez Torres, K. P. Khemchandani, and E. Oset *PRC* **77**, 042203(R)

A. Martínez Torres, K.P. Khemchandani, **LSG**, M. Napsuciale, E. Oset, *PRD* **78** (2008) 074031

Two-body inputs

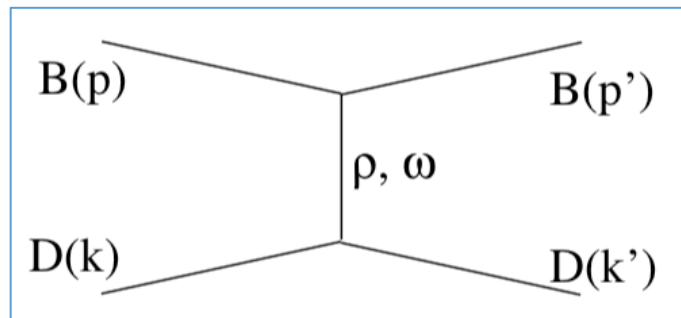
- **DK: leading order UChPT** DK , $D_s\eta$ and $D_s\pi$

$$V_{ij} = -\frac{C_{ij}}{4f^2}(s - u)$$

$$a(\mu) = -1.846, \mu = 1000 \text{ MeV} \Rightarrow \text{Pole} = 2318 \text{ MeV}$$

F.-K. Guo, P.-N. Shen, H.-C. Chiang, R.-G. Ping, and B.-S. Zou, PL B641, 278 (2006).

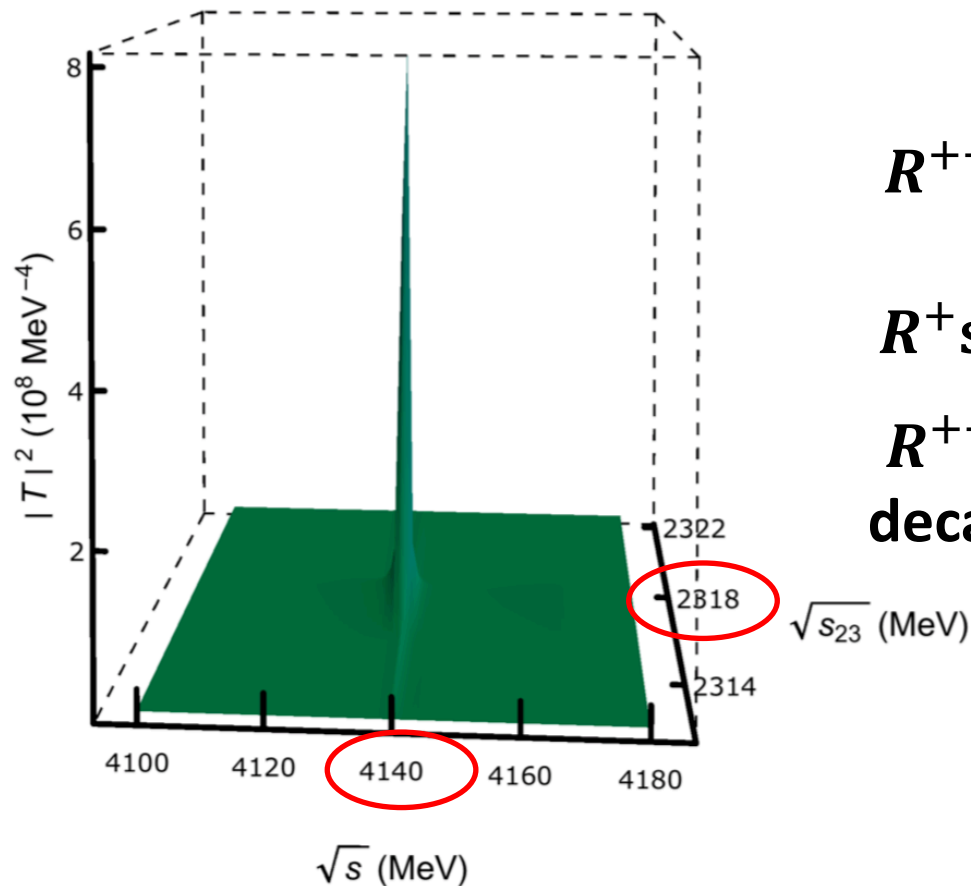
- **DD(Ds): local hidden gauge theory**



$$a(\mu) = -1.3 \sim -1.5, \mu = 1500 \text{ MeV} \Leftarrow \text{fixed from } D\bar{D}/D\bar{D}^* \rightarrow X(3700) / X(3872)$$

S. Sakai, L. Roca, and E. Oset, PRD96, 054023 (2017).

Three-body amplitudes



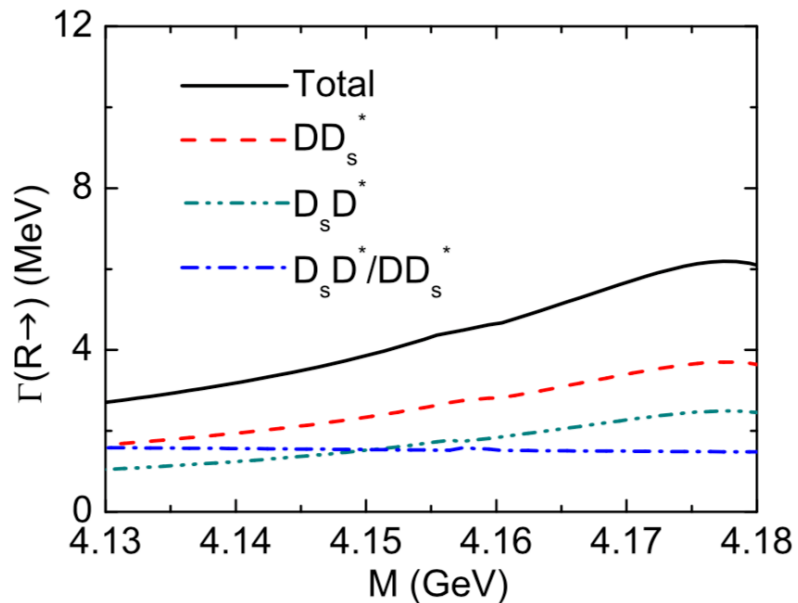
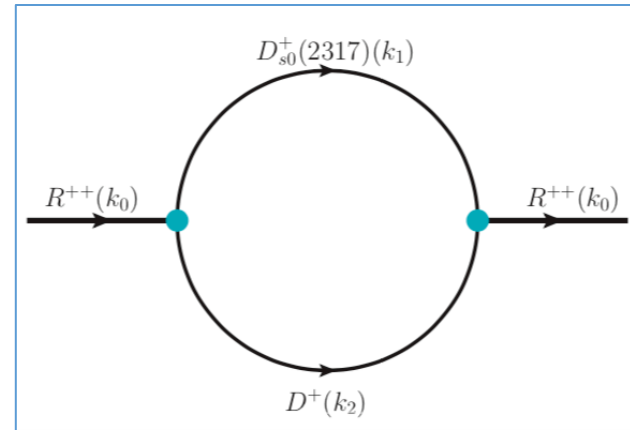
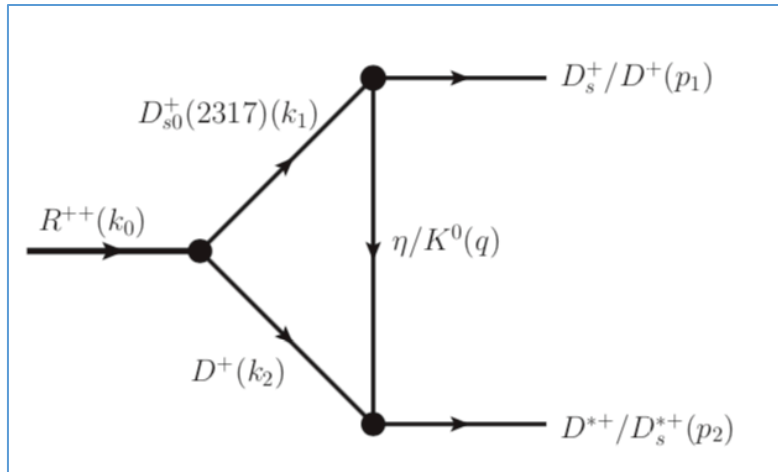
$$R^{++} = (I, I_{23}) = (\frac{1}{2}, 0)$$

R^+ should also exist

R^{++} is a bound state, but can decay strongly

Two-body decay width

Yin Huang, *LSG*, et al., in preparation



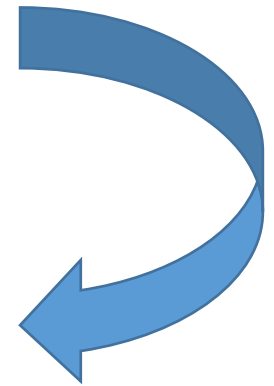
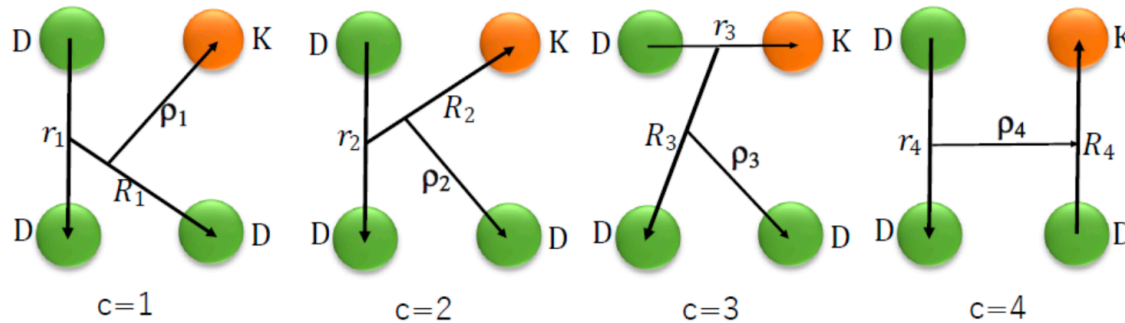
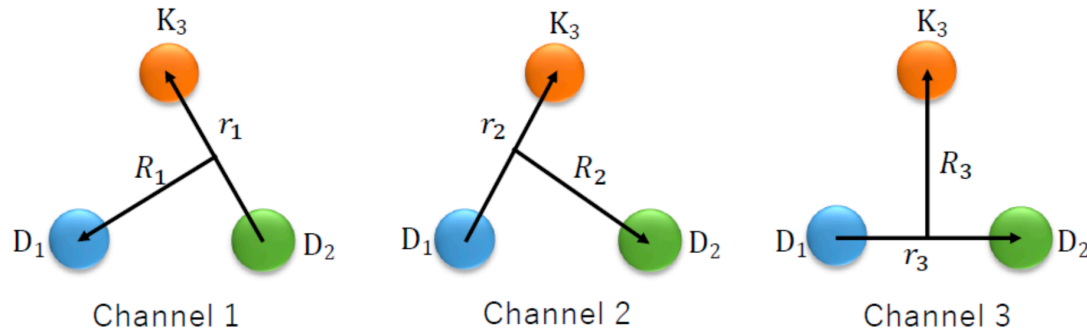
Preliminary result:
 $\Gamma \sim 10$ MeV

A DDDK state

$1(0^+)$

Gaussian Expansion Method

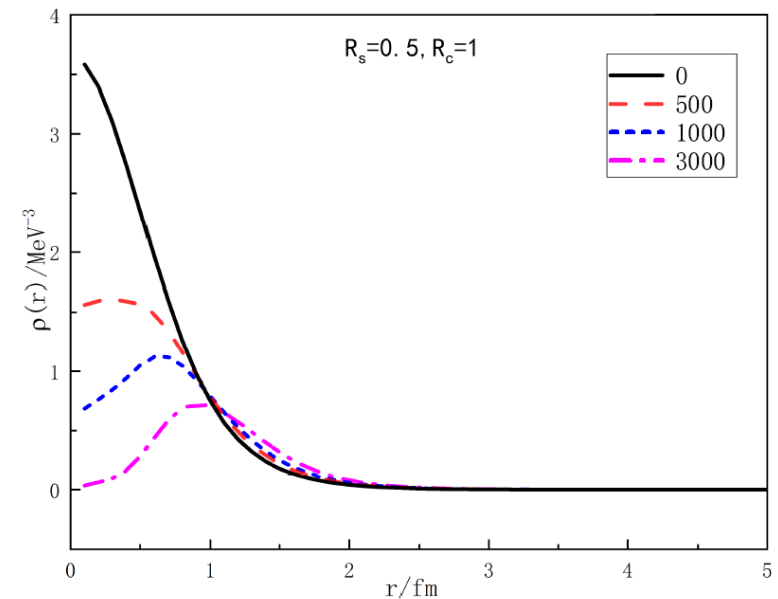
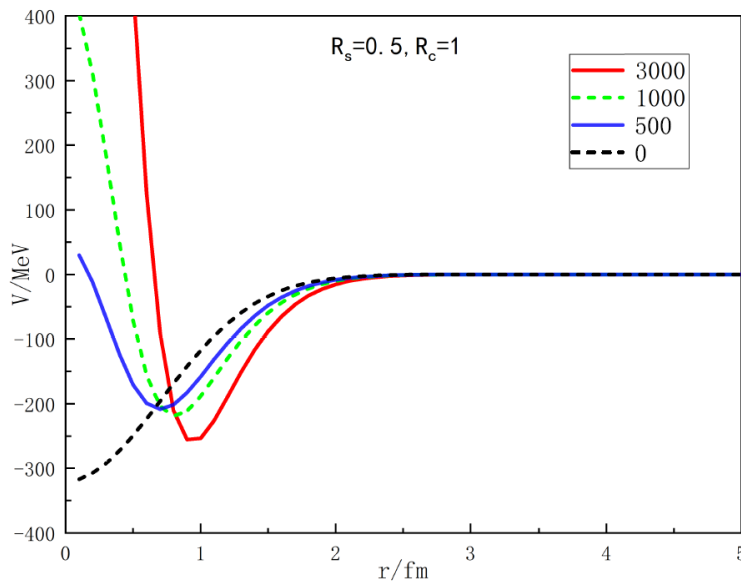
What if we add one more D?



A DDDK state $1(0^+)$

Gaussian Expansion Method

What if we add one more D? **Preliminary results** show that such a state exists as well



Uncertainties are at **the order of 10-20 MeV**

$$V_{DK}(\vec{r}; R_c) = C_S \frac{e^{-(r/R_S)^2}}{\pi^{3/2} R_S^3} + C(R_C) \frac{e^{-(r/R_c)^2}}{\pi^{3/2} R_c^3},$$

	DK*	DDK	DDDK
Binding	45 MeV	(67-71) MeV	91-107 MeV

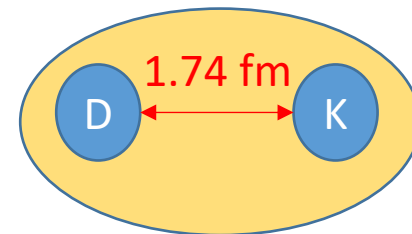
DD interactions play a minor role

$\frac{C_S}{\pi R_S^3}$	$\frac{C(R_c)}{\pi R_c^3}$	E_2	$E_3(\text{only } V_{DK})$	$E_3(V_{DK} + V_{DD})$	$E_4(\text{only } V_{DK})$	$E_4(V_{DK} + V_{DD})$
		$R_S = 0.5\text{fm}$		$R_c = 1\text{fm}$		
0	-320.1	-45.0	-65.8	-71.2	-89.4	-106.8
500	-455.4	-45.0	-65.8	-70.4	-89.2	-103.5
1000	-562.6	-45.0	-65.7	-69.7	-88.8	-101.4
3000	-838.7	-45.0	-65.0	-68.4	-87.0	-97.3
		$R_S = 0.5\text{fm}$		$R_c = 2\text{fm}$		
0	-149.1	-45.0	-66.0	-68.8, -45.1	-88.7, -66.3	-97.6, -70.7
500	-178.4	-45.0	-65.9	-68.2, -45.5	-88.5, -66.7	-95.5, -70.9
1000	-195.0	-45.0	-65.8, -45.2	-67.9, -45.8	-88.2, -66.9	-94.5, -71.2
3000	-225.9	-45.0	-65.3, -45.6	-67.2, -46.6	-87.0, -67.0	-92.6, -71.7
		$R_S = 0.5\text{fm}$		$R_c = 3\text{fm}$		
0	-107.0	-45.0	-66.2, -47.3	-68.0, -48.3	-88.8, -70.2	-94.4, -74.3
500	-119.4	-45.0	-66.2, -48.2	-67.7, -49.3	-88.7, -71.0	-93.2, -74.8
1000	-125.6	-45.0	-66.1, -48.7	-67.5, -49.8	-88.4, -71.3	-92.5, -75.2
3000	-136.2	-45.0	-65.8, -49.4	-67.1, -50.7	-87.6, -71.7	-91.4, -75.7

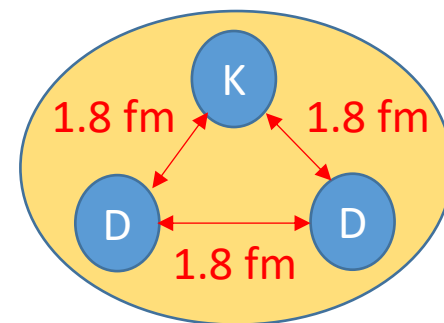
Spatial distributions

$\frac{C_S}{\pi R_S^3}$	$\frac{C(R_c)}{\pi R_c^3}$	$r_2(DK)$	$r_3(DK)$	$r_3(DD)$	$\langle T \rangle$	$\langle V_{DK} \rangle$	$\langle V_{DD} \rangle$
$R_S = 0.5\text{fm } R_c = 1\text{fm}$							
0	-320.1	1.28	1.32	1.36	124.37	-189.61	-5.98
500	-455.4	1.39	1.44	1.47	99.51	-164.83	-5.03
1000	-562.6	1.46	1.53	1.54	91.43	-156.67	-4.51
3000	-838.7	1.61	1.69	1.68	93.24	-157.80	-3.82
$R_S = 0.5\text{fm } R_c = 2\text{fm}$							
0	-149.1	1.74	1.80	1.80	60.20	-125.74	-3.23
500	-178.4	1.91	1.98	1.96	51.00	-116.59	-2.64
1000	-195.0	1.99	2.07	2.04	50.63	-116.12	-2.43
3000	-225.9	2.13	2.22	2.15	53.61	-118.59	-2.24
$R_S = 0.5\text{fm } R_c = 3\text{fm}$							
0	-107.0	2.13	2.19	2.17	39.49	-105.35	-2.13
500	-119.4	2.31	2.38	2.34	34.80	-100.73	-1.77
1000	-125.6	2.37	2.47	2.42	34.90	-100.77	-1.65
3000	-136.2	2.53	2.61	2.53	36.66	-102.24	-1.54

Ds0*(2317)



R(4140)

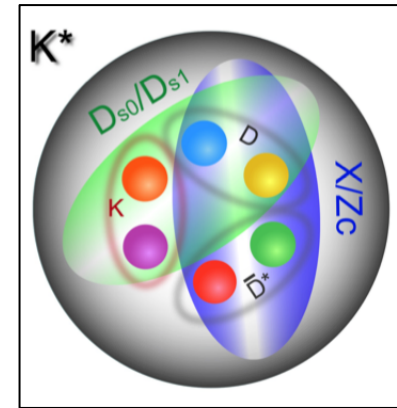


Contents

- Motivation: new types of clusters of color singlets in addition to nuclei
- $Ds0^*(2317)$ and $Ds1(2460)$ as DK/ D^*K molecules
- Explicit studies of the **DDK & DDDK** systems
- **A $K^*(4307)$ with hidden charm as $KX(3872)/Zc(3900)$ molecule**
- Summary and outlook

$K^*(4307)$

Instead of a D , what happens if we add a \bar{D}^* to the DK pair



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K^* mesons with hidden charm arising from $KX(3872)$ and $KZ_c(3900)$ dynamics



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K*(4307)

- Fixed center approximation (FCA):

$$K(D\bar{D}^* + \bar{D}D^*) \sim KX(3872)/Zc(3900)$$

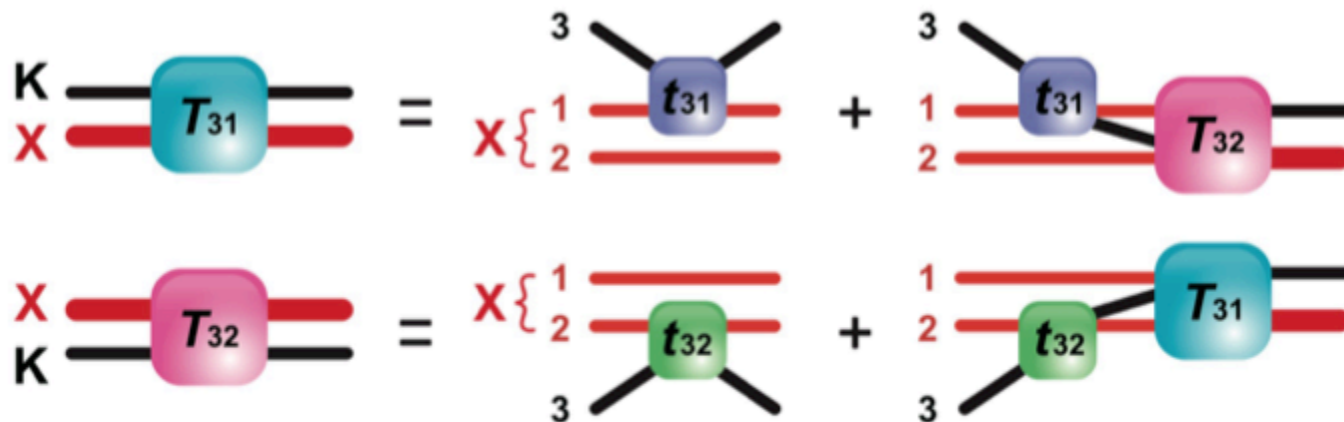
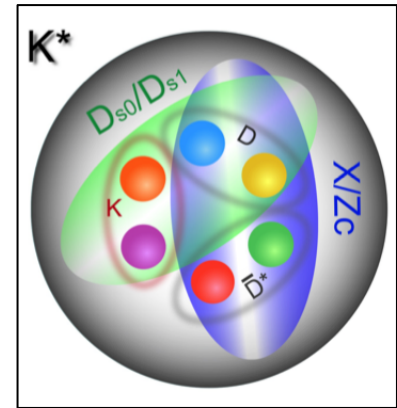
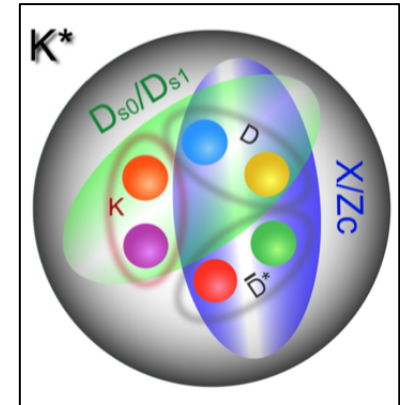
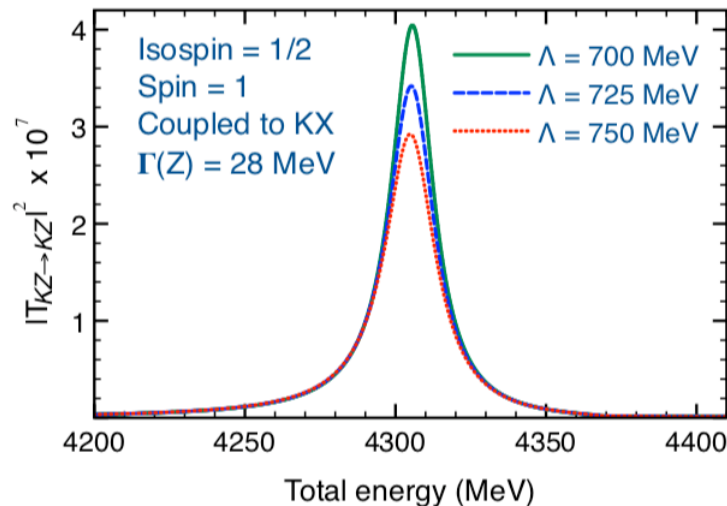
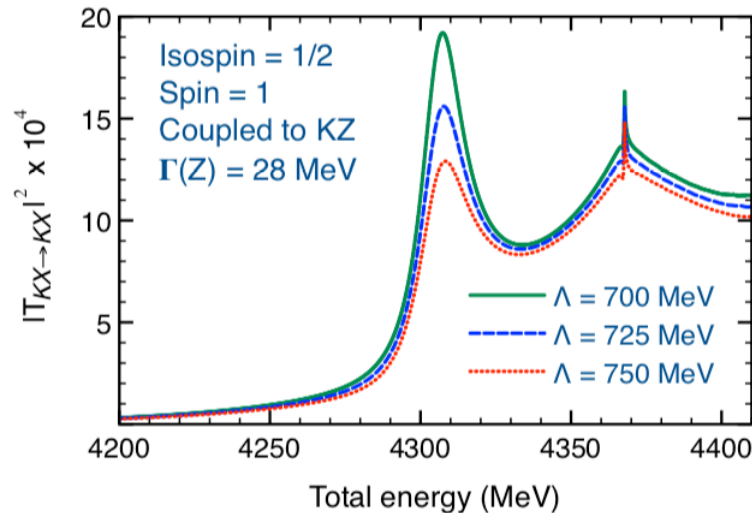


Figure 2: Diagrams showing the scattering of the particle labeled “3” (K) on a cluster (X) made of particles 1 (D) and 2 (\bar{D}^*).

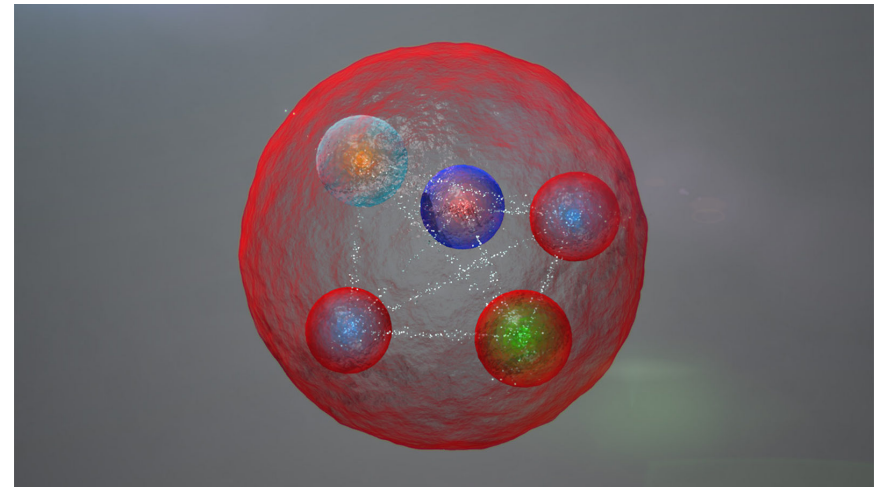
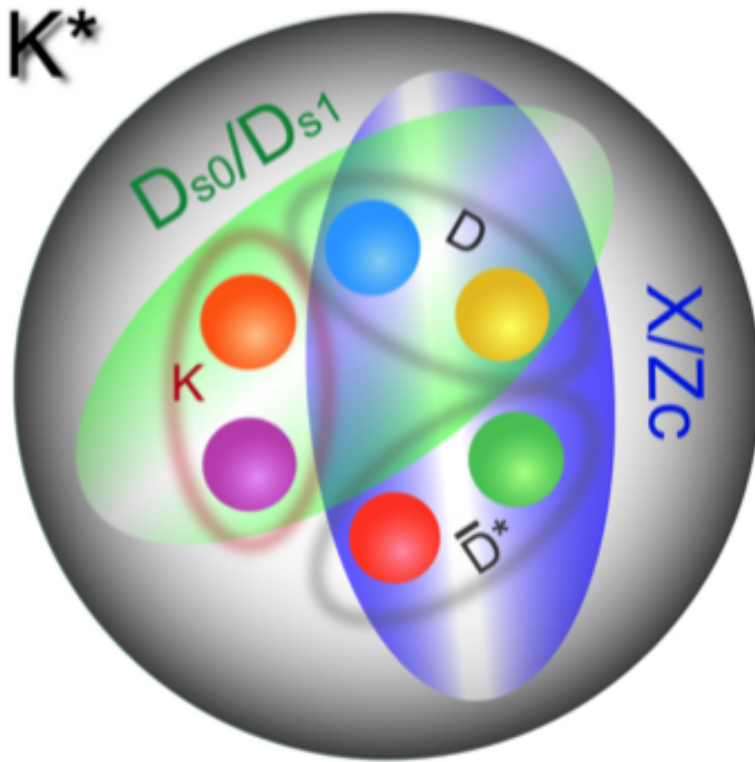
K*(4307)



- Treating KX and KZ as coupled channel systems
- A resonance with $M=(4307 \pm 2) - i(9 \pm 2)$ MeV with $I(J^P) = 1/2(1^-)$

In agreement with Li Ma, Qian Wang, [Ulf-G. Meißner](#), 1711.06143, but with completely different dynamics

$K^*(4307)$ —bosonic counterpart of P_c



Pentaquark (N^*) by LHCb

Phys.Rev.Lett. 115 (2015) 072001

Prediction of narrow N^* and Λ^* resonances with hidden charm above 4 GeV,
Jia-Jun Wu, R. Molina, E. Oset, B.S. Zou, 1007.0573

Analogy between KD and Kbar N

$D_{s0}^*(2317)$

- DK bound state
- **Dynamically generated**--
Unitary heavy hadron
chiral perturbation theory
- Coupled channels

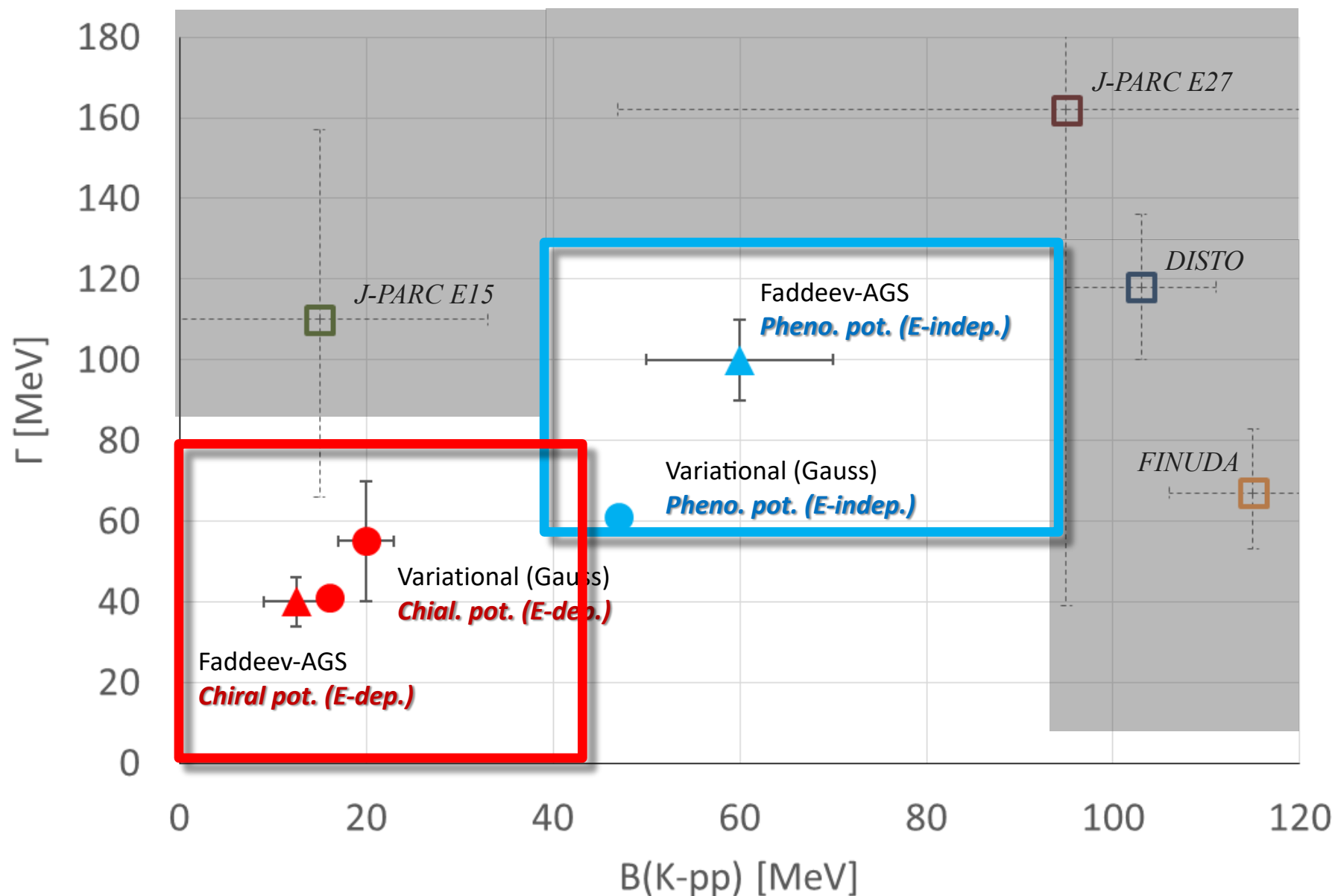
$\Lambda(1405)$

- N-Kbar bound state
- **Dynamically generated**--
Unitary baryon chiral
perturbation theory
- Coupled channels

The interaction between a **kaon and a heavy particle** seems to play an important role

Current status on “ K^-pp ”

A. Dote, Menu2019



Summary and outlook

- From nucleons, we can build nuclei, based on which the whole visible universe is formed
- If the $Ds_0^*(2317)$ is indeed a molecule of DK , then new forms of matter may be built upon them
- We have performed explicit few-body studies—demonstrating that indeed both DDK and $DDDK$ bound states **exist**
- Now we need experimental or lattice QCD confirmations and further theoretical studies on their production and decay mechanisms



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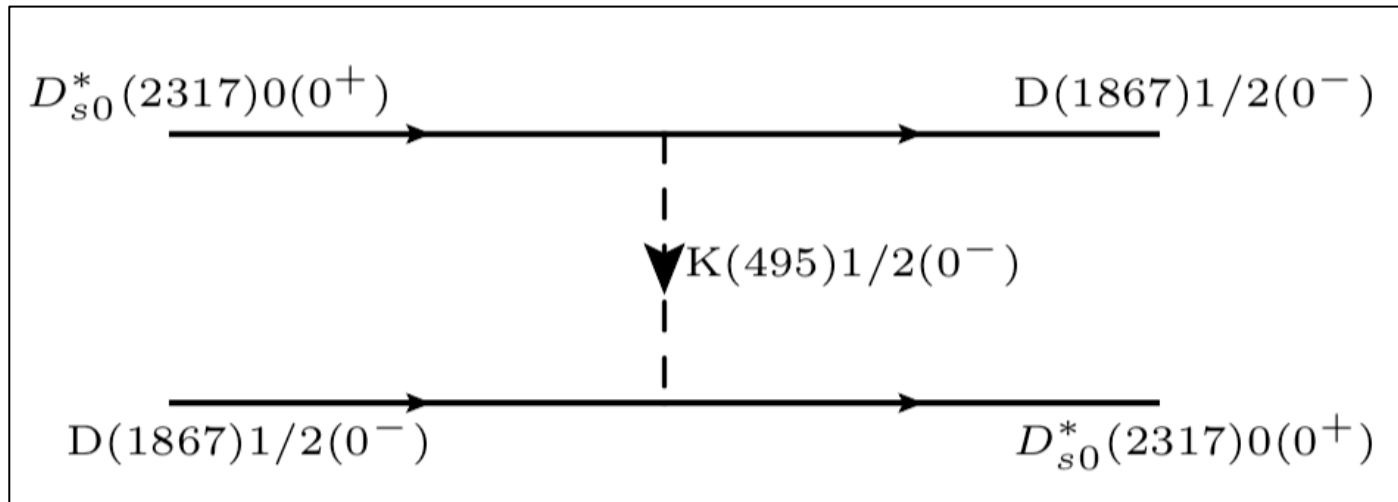


Thanks for your attention !

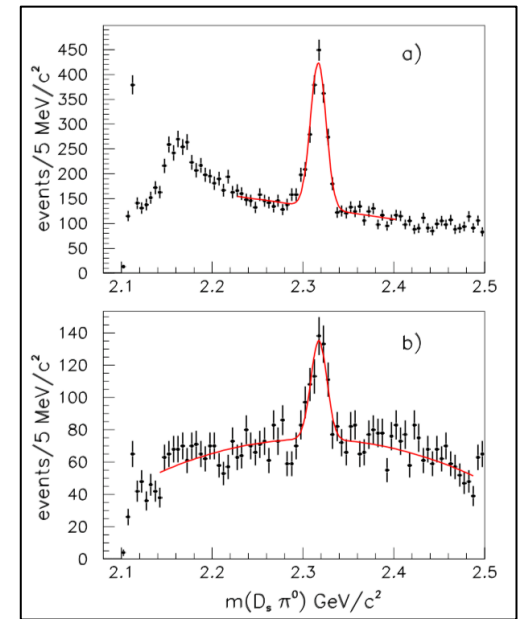
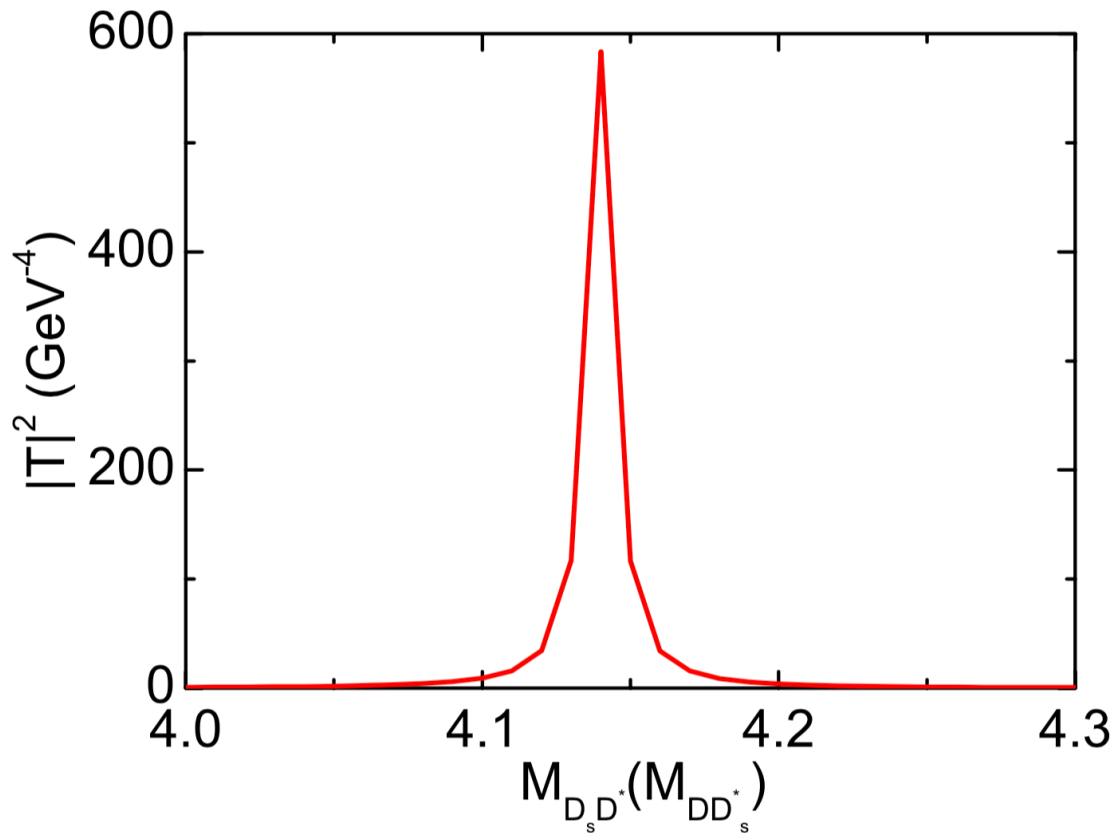
July 18, 2019

Even if the D_{s0} is a $c\bar{s}$ state, D_{s0} D binds

Interestingly, the explicit three-body result is consistent with the **quasi two body** study, where one treats the DK pair as the D_{s0}^* and describes the interaction between the D_{s0}^* with the D using one-kaon exchange



Where to look for the R^{++}/R^+



□ Belle/BelleII

$$e^+e^- \rightarrow X R \rightarrow X D^+ D_s^+ \pi^0$$

□ LHCb

$$p\bar{p} \rightarrow X R \rightarrow X D^+ D_s^+ \pi^0$$