Cluster virial expansion for quark-nuclear matter [1]

David Blaschke

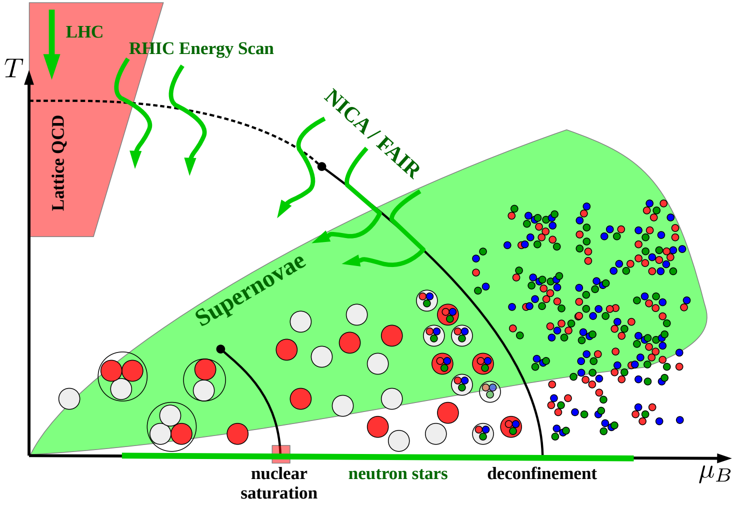
Institute of Theoretical Physics, University Wrocław, Poland
Bogoliubov Laboratory for Theoretical Physics, JINR Dubna, Russia

Workshop on INfinite and FInite NUclear Matter
Dubna, 20 March 2019

[1] Collab: N.-U. Bastian, A. Dubinin, A. Friesen, H. Grigorian, G. Röpke, L. Turko

QCD Phase Diagram with Clustering Aspects

From: N.-U. Bastian et al., Universe 4 (2018) 67; arxiv:1804.10178



Φ—Derivable Approach to the Cluster Virial Expansion

$$\Omega = \sum_{l=1}^{A} \Omega_l = \sum_{l=1}^{A} \left\{ c_l \left[\mathrm{Tr} \ln \left(-G_l^{-1} \right) + \mathrm{Tr} \left(\Sigma_l \; G_l \right) \right] + \sum_{\substack{i,j \\ i+j=l}} \Phi[G_i, G_j, G_{i+j}] \right\} ,$$

$$G_A^{-1} \quad = \quad G_A^{(0)^{-1}} - \Sigma_A \; , \; \Sigma_A(1 \ldots A, 1' \ldots A', z_A) = \frac{\delta \Phi}{\delta G_A(1 \ldots A, 1' \ldots A', z_A)}$$

Stationarity of the thermodynamical potential is implied

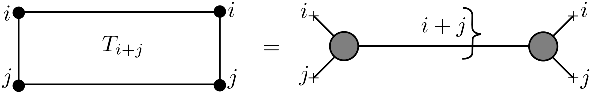
$$\frac{\delta \Omega}{\delta G_A(1 \ldots A, 1' \ldots A', z_A)} = 0 \; .$$

Cluster virial expansion follows for this $\Phi-$ functional

$\Phi =$ 

Figure: The Φ functional for $A-$particle correlations with bipartitions $A = i + j$.

Green's function and T-matrix: separable approximation



The T_A matrix fulfills the Bethe-Salpeter equation in ladder approximation

$$T_{i+j}(1,2,\ldots,A;1',2',\ldots A';z) = V_{i+j} + V_{i+j}G^{(0)}_{i+j}T_{i+j} \ ,$$

which in the separable approximation for the interaction potential,

$$V_{i+j} = \Gamma_{i+j}(1,2,\ldots,i;i+1,i+2,\ldots,i+j)\Gamma_{i+j}(1',2',\ldots,i';(i+1)',(i+2)',\ldots,(i+j)') \ ,$$

leads to the closed expression for the T_A matrix

$$T_{i+j}(1,2,\ldots,i+j;1',2',\ldots(i+j)';z) = V_{i+j}\left\{1-\Pi_{i+j}\right\}^{-1} \ ,$$

with the generalized polarization function

$$\Pi_{i+j} = \mathrm{Tr}\left\{\Gamma_{i+j}G^{(0)}_i\Gamma_{i+j}G^{(0)}_j\right\}$$

The one-frequency free $i-$particle Green's function is defined by the $(i-1)$-fold Matsubara sum

$$\begin{aligned} G^{(0)}_i(1,2,\ldots,i;\Omega_i) &= \sum_{\omega_1\ldots\omega_{i-1}}\frac{1}{\omega_1-E(1)}\frac{1}{\omega_2-E(2)}\cdots\frac{1}{\Omega_i-(\omega_1+\ldots\omega_{i-1})-E(i)} \\ &= \frac{(1-f_1)(1-f_2)\ldots(1-f_i)-(-)^i f_1f_2\ldots f_i}{\Omega_i-E(1)-E(2)-\ldots E(i)} \ . \end{aligned}$$

Useful relationships for many-particle functions

$$G_{i+j}^{(0)} = G_{i+j}^{(0)}(1,2,\ldots,i+j;\Omega_{i+j}) \quad = \quad \sum_{\Omega_i} G_i^{(0)}(1,2,\ldots,i;\Omega_i) G_j^{(0)}(i+1,i+2,\ldots,i+j;\Omega_j) \ .$$

Another set of useful relationships follows from the fact that in the ladder approximation both, the full two-cluster ($i+j$ particle) T matrix and the corresponding Greens' function

$$G_{i+j} = G_{i+j}^{(0)} \left\{1 - \Pi_{i+j}\right\}^{-1} \tag{1}$$

have similar analytic properties determined by the $i+j$ cluster polarization loop integral and are related by the identity

$$T_{i+j} G_{i+j}^{(0)} = V_{i+j} G_{i+j} \ . \tag{2}$$

which is straightforwardly proven by multiplying Equation for the $T_{i+j}-$ matrix with $G_{i+j}^{(0)}$ and using Equation (1). Since these two equivalent expressions in Equation (2) are at the same time equivalent to the two-cluster irreducible Φ functional these functional relations follow

$$T_{i+j} = \delta\Phi/\delta G_{i+j}^{(0)} \ ,$$
$$V_{i+j} = \delta\Phi/\delta G_{i+j} \ .$$

Next we prove the relationship to the Generalized Beth-Uhlenbeck approach!

GBU EoS from the $\Phi-$derivable approach

Consider the partial density of the $A-$particle state defined as

$$n_A(T,\mu) = -\frac{\partial \Omega_A}{\partial \mu} = -\frac{\partial}{\partial \mu} d_A \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega}{2\pi} \left[\ln\left(-G_A^{-1}\right) + \mathrm{Tr}\left(\Sigma_A\ G_A\right)\right] + \sum_{\substack{i,j\\ i+j=A}} \Phi[G_i, G_j, G_{i+j}]\ . \qquad (3)$$

Using spectral representation for $F(\omega)$ and Matsubara summation

$$F(iz_n) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{\mathrm{Im}F(\omega)}{\omega - iz_n}\ , \quad \sum_{z_n} \frac{c_A}{\omega - iz_n} = f_A(\omega) = \frac{1}{\exp[(\omega-\mu)/T] - (-1)^A}$$

with the relation $\partial f_A(\omega)/\partial\mu = -\partial f_A(\omega)/\partial\omega$ we get for Equation (3) now

$$n_A(T,\mu) = -d_A \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega}{2\pi} f_A(\omega) \frac{\partial}{\partial\omega} \left[\mathrm{Im}\ln\left(-G_A^{-1}\right) + \mathrm{Im}\left(\Sigma_A\ G_A\right)\right] + \sum_{\substack{i,j\\ i+j=A}} \frac{\partial \Phi[G_i,G_j,G_A]}{\partial\mu}\ ,$$

where a partial integration over ω has been performed For two-loop diagrams of the sunset type holds a cancellation[2] which we generalize here for cluster states

$$d_A \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega}{2\pi} f_A(\omega) \frac{\partial}{\partial\omega} \left(\mathrm{Re}\Sigma_A\ \mathrm{Im}G_A\right) - \sum_{\substack{i,j\\ i+j=A}} \frac{\partial \Phi[G_i, G_j, G_A]}{\partial\mu} = 0\ .$$

Using generalized optical theorems we can show that ($G_A = |G_A|\exp(i\delta_A)$)

$$\frac{\partial}{\partial\omega}\left[\mathrm{Im}\ln\left(-G_A^{-1}\right) + \mathrm{Im}\Sigma_A\ \mathrm{Re}G_A\right] = 2\mathrm{Im}\left[G_A\ \mathrm{Im}\Sigma_A \frac{\partial}{\partial\omega} G_A^*\ \mathrm{Im}\Sigma_A\right] = -2\sin^2\delta_A \frac{\partial\delta_A}{\partial\omega}\ .$$

The density in the form of a generalized Beth-Uhlenbeck EoS follows

$$n(T,\mu) = \sum_{i=1}^{A} n_i(T,\mu) = \sum_{i=1}^{A} d_i \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega}{2\pi} f_i(\omega) 2\sin^2\delta_i \frac{\partial\delta_i}{\partial\omega}\ .$$

[2] B. Vanderheyden & G. Baym, J. Stat. Phys. (1998), J.-P. Blaizot et al., PRD (2001)

Example: Deuterons in Nuclear Matter

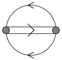
The $\Phi-$derivable thermodynamical potential for the nucleon-deuteron system reads

$$\Omega = -\mathrm{Tr}\,\{\ln(-G_1)\} - \mathrm{Tr}\{\Sigma_1 G_1\} + \mathrm{Tr}\,\{\ln(-G_2)\} + \mathrm{Tr}\{\Sigma_2 G_2\} + \Phi[G_1, G_2]\ ,$$

where the full propagators obey the Dyson-Schwinger equations

$$G_1^{-1}(1,z) = z - E_1(p_1) - \Sigma_1(1,z);\ \ G_2^{-1}(12,1'2',z) = z - E_1(p_1) - E_2(p_2) - \Sigma_2(12,1'2',z),$$

with selfenergies and Φ functional

$$\Sigma_1(1,1') = \frac{\delta\Phi}{\delta G_1(1,1')}\ ;\ \ \Sigma_2(12,1'2',z) = \frac{\delta\Phi}{\delta G_2(12,1'2',z)}\ ,\Phi =$$  ,

fulfilling stationarity of the thermodynamic potential $\partial\Omega/\partial G_1 = \partial\Omega/\partial G_2 = 0$.
For the density we obtain the cluster virial expansion

$$n = -\frac{1}{V}\frac{\partial\Omega}{\partial\mu} = n_{\mathrm{qu}}(\mu, T) + 2n_{\mathrm{corr}}(\mu, T)\ ,$$

with the correlation density in the generalized Beth-Uhlenbeck form

$$n_{\mathrm{corr}} = \int \frac{dE}{2\pi} g(E) 2\sin^2\delta(E)\frac{d\delta(E)}{dE}\ .$$

Example: Deuterons in Nuclear Matter

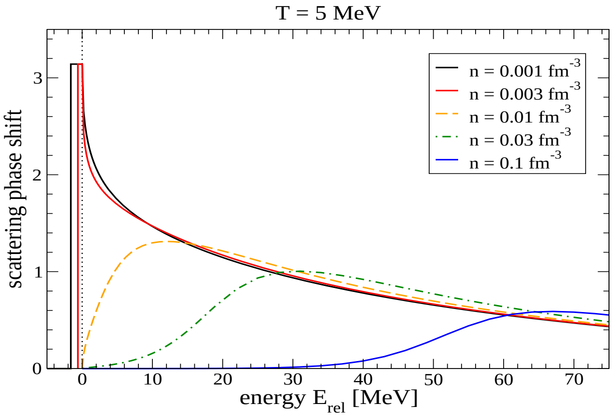


Figure: Integrand of the intrinsic partition function as function of the c.m.s. energy in the deuteron channel. Mott dissociation and Levinson's theorem! From G. Röpke, J. Phys. Conf. Ser. 569 (2014) 012014.

Cluster Virial Expansion for Quark-Hadron Matter within the Φ Derivable Approach

$$\Omega = \sum_{i=Q,M,D,B} c_i \left[\mathrm{Tr} \ln \left(-G_i^{-1}\right) + \mathrm{Tr}\left(\Sigma_i \; G_i\right)\right] + \Phi\left[G_Q, G_M, G_D, G_B\right] \;,$$

$$= \sum_{i=Q,M,D,B} d_i \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega}{2\pi} \left\{\omega + 2T \ln\left[1 - \mathrm{e}^{-\omega/T}\right]\right\} 2\sin^2\delta_i \frac{\partial \delta_i}{\partial \omega} \;.$$

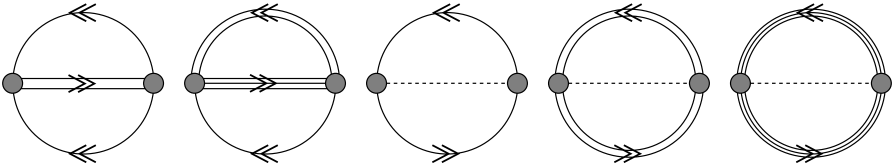


Figure: Φ functional for the quark-meson-diquark-baryon system in 2-loop approx.

$$\Sigma_i = \frac{\delta \; \Phi\left[G_Q, G_M, G_D, G_B\right]}{\delta \; G_i} \;.$$

The selfenergies ...

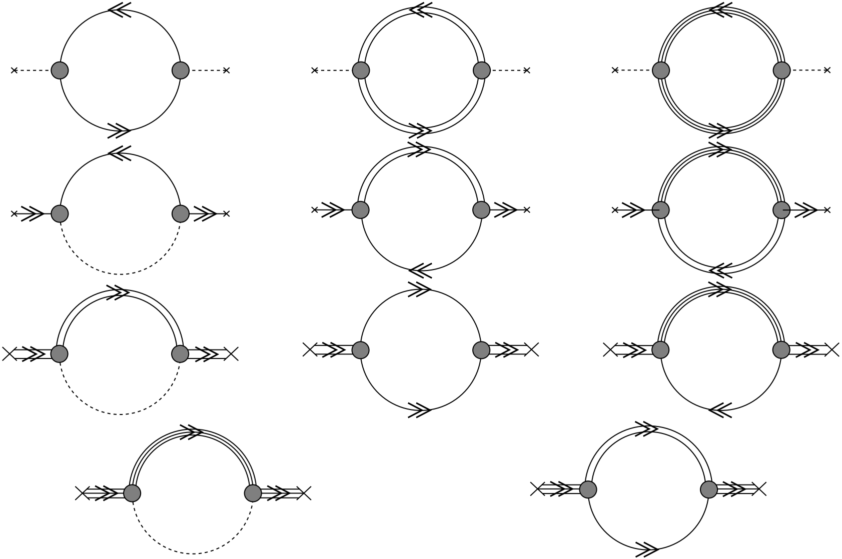


Figure: Selfenergies for Greens functions of Q-M-D-B system in 2-loop approx.

Mott Dissociation of Pions in Quark Matter

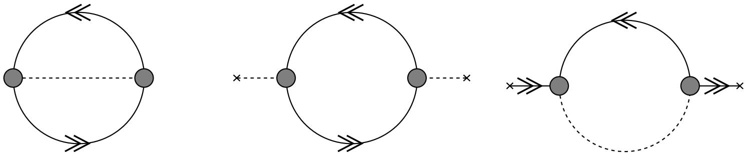


Figure: The Φ functional (left panel) for the case of mesons in quark matter, where the bosonic meson propagator is defined by the dashed line and the fermionic quark propagators are shown by the solid lines with arrows. The corresponding meson and quark selfenergies are shown in the middle and right panels, respectively.

Mott Dissociation of Pions in Quark Matter

The meson polarization loop $\Pi_M(q,z)$ enters the definition of the meson T matrix

$$T_M^{-1}(q,\omega+i\eta) = G_S^{-1} - \Pi_M(q,\omega+i\eta) = |T_M(q,\omega)|^{-1}\mathrm{e}^{-i\delta_M(q,\omega)} ,$$

which in the polar representation introduces a phase shift $\delta_M(q,\omega) = \arctan(\Im T_M/\Re T_M)$, that results in a generalized Beth-Uhlenbeck equation of state for the thermodynamics of the consistently coupled quark-meson system

$$\Omega = \Omega_{\rm MF} + \Omega_M ,$$

where the selfconsistent quark meanfield contribution is

$$\Omega_{\rm MF} = \frac{\sigma_{\rm MF}^2}{4G_S} - 2N_cN_f \int \frac{d^3p}{(2\pi)^3}\left[E_p + T\ln\left(1+\mathrm{e}^{-(E_p-\Sigma_+-\mu)/T}\right) + T\ln\left(1+\mathrm{e}^{-(E_p+\Sigma_-+\mu)/T}\right)\right] ,$$

The mesonic contribution to the thermodynamics is

$$\Omega_M = d_M \int \frac{d^3k}{(2\pi)^3}\int \frac{d\omega}{2\pi}\left\{\omega + 2T\ln\left[1-\mathrm{e}^{-\omega/T}\right]\right\} 2\sin^2\delta_M(k,\omega)\ \frac{\delta_M(k,\omega)}{d\omega} ,$$

Mott Dissociation of Pions in Quark Matter

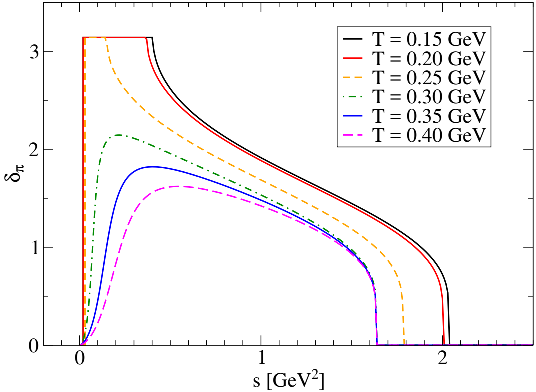


Figure: Phase shift of the pion as a quark-antiquark state for different temperatures, below and above the Mott dissociation temperature.
From D.B. et al., Ann. Phys. 348 (2014) 228.
See also D.B., A. Dubinin, D. Ebert, A.V. Friesen, PPN Lett. 15 (2018) 230.

Φ-derivable cluster expansion for quark-nuclear matter

Cluster decomposition of the thermodynamic potential

$$\begin{aligned}\Omega &= \sum_i \Omega_i \\ &= \sum_i \Big\{ c_i \left[\mathrm{Tr}\ln\left(-G_i^{-1}\right) + \mathrm{Tr}\left(\Sigma_i\ G_i\right)\right] + \sum_{\substack{j,k \\ j+k=i}} \Phi[G_j, G_k, G_{j+k}] \Big\}\ ,\end{aligned}$$

generates the partial densities

$$\begin{aligned}n_j(T,\mu) = -\frac{\partial\Omega}{\partial\mu_j} &= \sum_{i=1} A_{ij} n_i(T,\mu) \\ &= \sum_{i=1} A_{ij} g_i \int \frac{\mathrm{d}^3 p}{(2\pi)^3} \int \frac{\mathrm{d}E}{2\pi} f_i(E) 2\sin^2\delta_i(E) \frac{\mathrm{d}\delta_i(E)}{\mathrm{d}E}\ ,\end{aligned}$$

Based on: N.-U. F. Bastian and D.B., arXiv:1812.11889

Φ-derivable cluster expansion for quark-nuclear matter

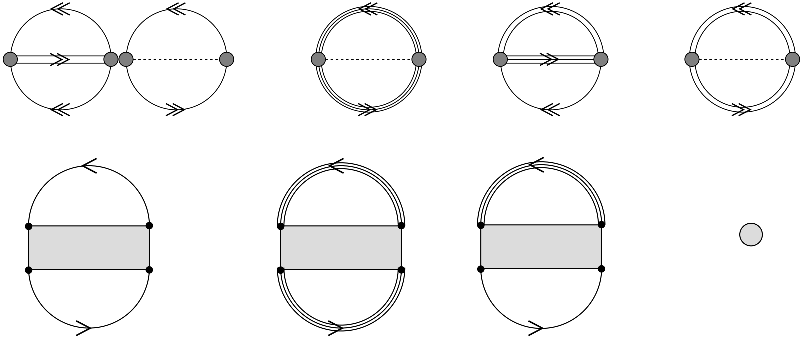


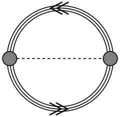
Figure:

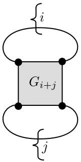
Upper row: full set of sunset type diagrams.

Lower row: after collapsing the diquark and meson-propagators ($\rightarrow$ mean fields).

Relativistic Density Functional Approach to Nuclear Matter

In case of color confinement all closed loop diagrams with Q- and D-lines vanish. The system reduces to the M-B- system. The $\Phi-$functional becomes a density functional.

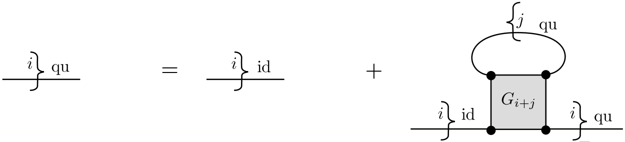
$\Phi =$ 

$\Longrightarrow \qquad \mathcal{U} =$ 

$$\Omega = T \sum_{i=n,p,\Lambda,\dots} c_i \left[\mathrm{Tr} \ln S_{i,qu}^{-1} + \sum_{j=S,V} n_{i,j} \Sigma_{i,j} \right] + U\left[\{n_{i,S}, n_{i,V}\}\right] \ ,$$

$$\frac{\partial \Omega}{\partial n_{i,S}} = \frac{\partial \Omega}{\partial n_{i,V}} = 0 \ , \quad i = n, p, \Lambda, \dots \ , \qquad \frac{\partial U}{\partial n_{i,S}} = \Sigma_{i,S} \ , \quad \frac{\partial U}{\partial n_{i,V}} = \Sigma_{i,V} \ .$$

The baryon quasiparticle propagators fulfill the Dyson equations $S_{i,qu}^{-1} = S_{i,0}^{-1} - \Sigma_{i,S} - \Sigma_{i,V}$,

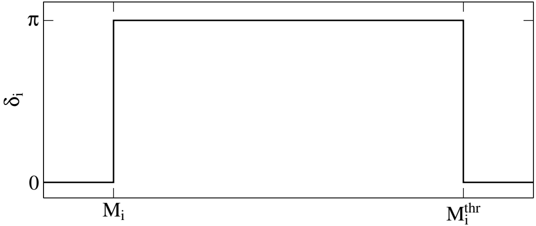


Φ-derivable cluster expansion for quark-nuclear matter

$$n_i^{\mathrm{free}} = g_i \int \frac{\mathrm{d}^3 p}{(2\pi)^3} \int \frac{\mathrm{d}M}{2\pi} f_i\left(\sqrt{p^2 + M^2} + V_i\right) 2\sin^2\delta_i(M) \frac{\mathrm{d}\delta_i(M)}{\mathrm{d}M}$$

Phase shift model:

$\delta_{i=\mathrm{p,n}}(M) = \pi\,\Theta(M - M_i)\,\Theta(M_i^{\mathrm{thr}} - M)$ with $M_{\mathrm{p}}^{\mathrm{thr}} = 2M_{\mathrm{u}} + M_{\mathrm{d}}$ etc.



$$n_{i=\mathrm{p,n}} = g_i \int \frac{\mathrm{d}^3 p}{(2\pi)^3} \left[f_i(\sqrt{p^2 + M_i^2} + V_i) - f_i(\sqrt{p^2 + (M_i^{\mathrm{thr}})^2} + V_i) \right] \Theta(M_i^{\mathrm{th}}$$

$$= \left(n_{\mathrm{N}}^{\mathrm{qu}} - n_{\mathrm{N}}^{\mathrm{thr}}\right) \Theta(M_i^{\mathrm{thr}} - M_i)$$

Selfenergies: Relativistic Density Functionals

Quark Matter: $U_{\rm SFM} = D(n_{\rm v}) n_{\rm q,s}^{2/3} + a n_{\rm q,v}^2$.

$$S_{\rm SFM} = \frac{2}{3} D(n_{\rm q,v})\, n_{\rm q,s}^{-1/3}\ , \quad D(n_{\rm v}) = D_0 \exp(-\alpha n_{\rm q,v}^2) \tag{5}$$

Nuclear Matter:

$$U_{\rm DD2} = -\frac{1}{2}\frac{\Gamma_\sigma^2}{m_\sigma^2} n_{\rm s}^2 + \frac{1}{2}\frac{\Gamma_\omega^2}{m_\omega^2} n_{\rm v}^2 + \frac{1}{2}\frac{\Gamma_\rho^2}{m_\rho^2} n_{\rm vi}^2\ , \tag{6}$$

Scalar self-energy: $S_{\rm DD2} = -(\Gamma_\sigma^2/m_\sigma^2) n_{\rm s}$, Vector self-energy:

$$V_{{\rm DD2},i} = \frac{\Gamma_\omega^2}{m_\omega^2} n_{\rm v} + \tau_i \frac{\Gamma_\rho^2}{m_\rho^2} n_{\rm vi} - \frac{\Gamma_\sigma \Gamma_\sigma'}{m_\sigma^2} n_{\rm s}^2 + \frac{\Gamma_\omega \Gamma_\omega'}{m_\omega^2} n_{\rm v}^2 + \frac{\Gamma_\rho \Gamma_\rho'}{m_\rho^2} n_{\rm vi}^2\ . \tag{7}$$

Thermodynamics of quark-nuclear matter

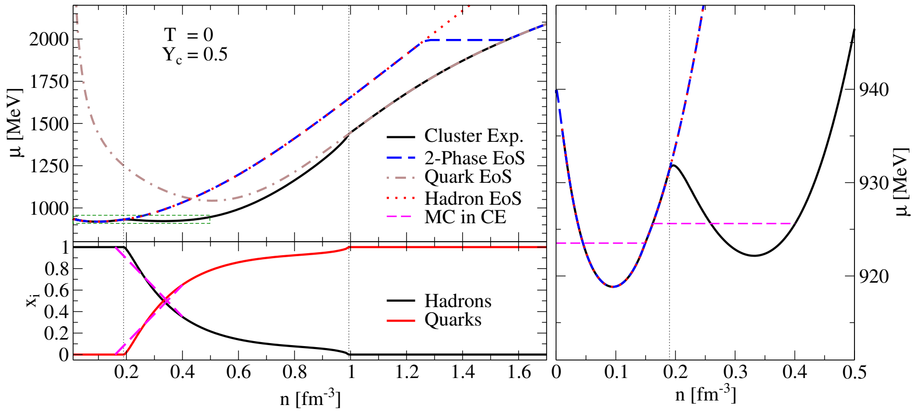
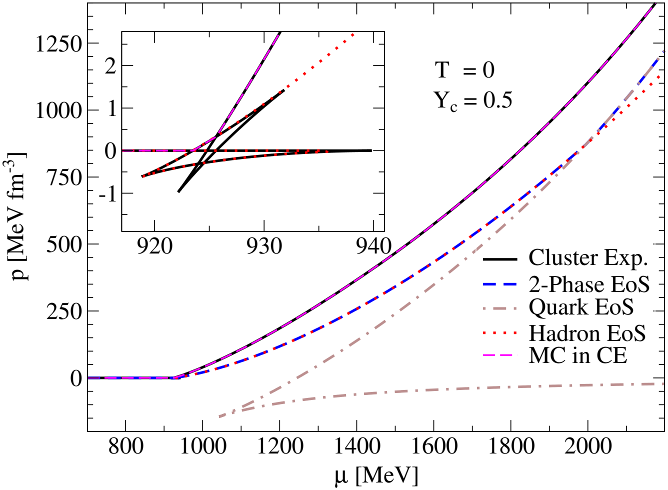


Figure: Van der Waals type thermodynamic behaviour with two first order phase transitions: liquid-gas and deconfinement. Note the wide density range for quark-hadron mixture after the deconfinement transition.

Maxwell construction for quark-nuclear matter EoS



Phase diagram of quark-nuclear matter

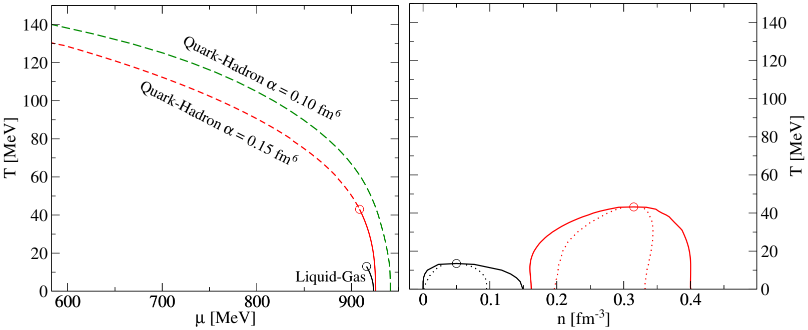
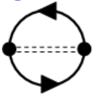


Figure: Phase diagram of the cluster expansion, with the liquid-gas phase transition (black) and two parametrizations of quark-hadron transition. One features a first order phase transition (red) with mixed phase and a critical endpoint at $T_C = 43$ MeV and $\mu_C = 909$ MeV. The other represents a crossover all over scenario (green) without critical endpoint. Right panel: the borders of the phase coexistence regions (binodals) for liquid-gas (black) and quark-hadron (red) transitions; dotted lines show the spinodal region of thermodynamic instability ($\partial\mu/\partial n < 0$).

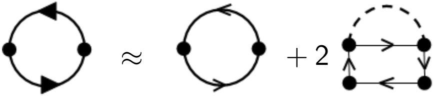
More details in: N.-U. F. Bastian and D.B., arXiv:1812.11889

Precursor to Mott Dissociation: Quark Pauli Blocking

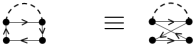
$\Omega = \mathrm{Tr}\left\{\ln S_q^{-1} - \Sigma_q S_q - \frac{1}{2}\ln D_\pi^{-1} + \frac{1}{2}D_\pi \Pi_\pi\right\} + \Phi[S_q, D_\pi]$, $\Phi[S_q, D_\pi] =$ 

Dyson-Schwinger equations $S_q^{-1} = S_{q,MF}^{-1} - \Sigma_q$ and $D_\pi^{-1} = G_\pi^{-1} - \Pi_\pi$, resp.

$$\Sigma_q = \frac{\delta\Phi}{\delta S_q} =$$  $+ \ldots$;

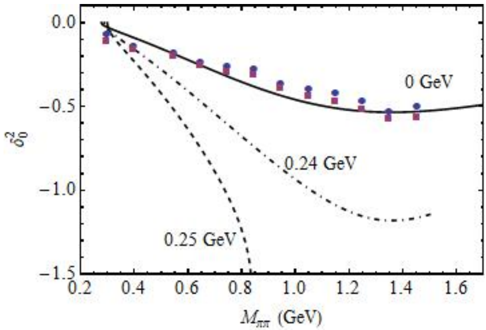
$$\Pi_\pi = \frac{\delta\Phi}{\delta D_\pi} =$$  $+ \ldots$.

Perturbation around selfconsistent meanfield reveals quark exchange contribution to $\pi\pi$ scattering in isospin=2 channel

 ,

Quark Pauli Blocking in the Pion Gas

Nonrelativistic potential model evaluation of quark exchange diagrams
$\Longrightarrow \delta_0^2$ phase shift for $\pi^+\pi^+$ scattering [Barnes & Swanson, PRD (1992)]

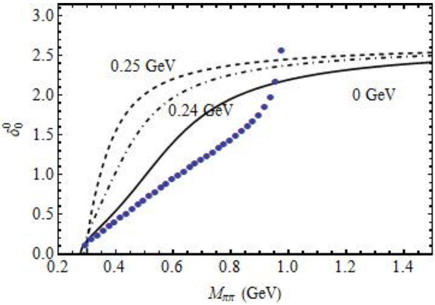


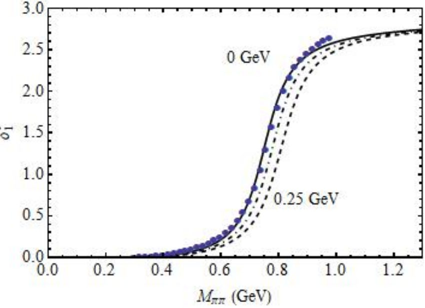
$$\sin\delta_0^2 = -\frac{\alpha_s}{9\lambda m_q^2}\sqrt{\frac{s}{\xi}}\left(1 - \mathrm{e}^{-\lambda\xi/2} + \frac{4\lambda\xi}{3\sqrt{3}}\mathrm{e}^{-\lambda\xi/3}\right) , \quad \xi = s - 4M_\pi^2$$

Let m_q depend on temperature $\longrightarrow$ chiral enhancement!

Quark Pauli Blocking: T-Dependence of Phase Shifts

Simple Breit-Wigner parametrization of $\pi - \pi$ phase shifts in the sigma-meson (δ_0^0) and rho-meson δ_1^1 channels

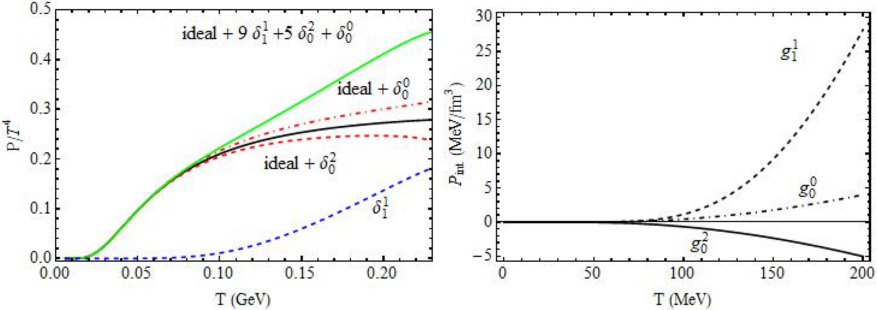




Let meson masses and widths depend on temperature $\longrightarrow$ chiral enhancement!

Virial expansion of the Pion Gas Pressure

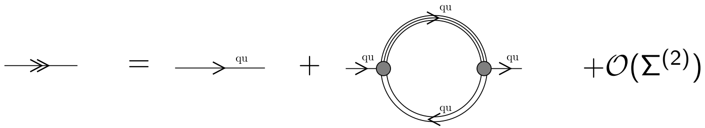
Second virial coefficient in Beth-Uhlenbeck formula with $\pi - \pi$ phase shifts

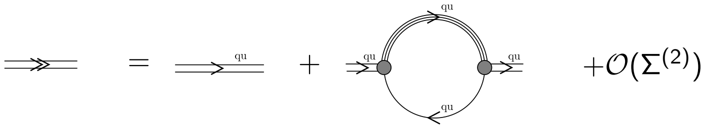


Compensation of sigma-meson and Pauli blocking contributions keeps intact, despite strong temperature dependent chiral enhancement!

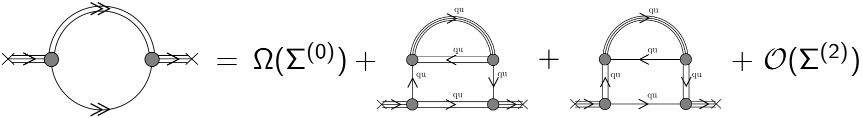
Quark Pauli Blocking in Hadronic Matter

Perturbative expansion around the quasiparticle Q- and D- propagators



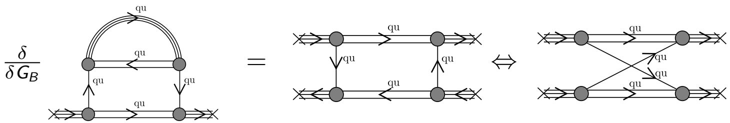


Insertion into the Q-D loop diagram defining the baryon



Quark Pauli Blocking in Hadronic Matter

The "new" baryon selfenergy diagrams contain one closed baryon loop, are proportional to baryon density. Functional derivative w.r.t. the baryon propagator yields effective interaction



For the diagrammatic expansion, see also K. Maeda, Ann Phys. **326** (2011) 1032. Quark Pauli blocking has been evaluated, e.g. in nonrelativistic quark models, with constant (constituent) quark mass [G. Röpke et al., PRD **34** (1986) 3499]. Here, effects of chiral symmetry restoration in a hadronic medium are taken into account. They lead to a strong enhancement of the Pauli blocking energy shift and drive the system into dissociation/deconfinement!

Note:Pauli blocking effect in a pion gas completely analogous!

Jean-Paul Blaizot Kenji Maeda D.B.

Helmholtz International Summer School
“Dense QCD Phases in Heavy-Ion Collisions”
Dubna, August 21 – September 4, 2010

Example: Quark Pauli Blocking in Nuclear Matter

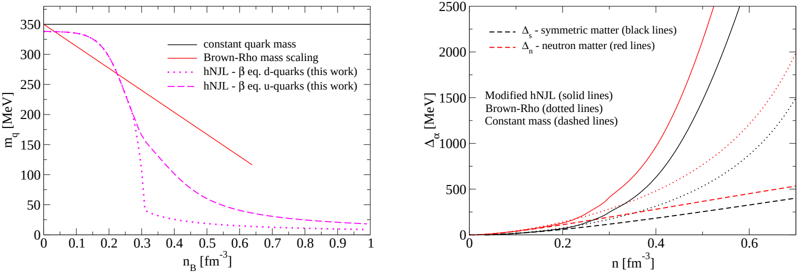


Figure: Left panel: Different quark mass dependences on the density; Right panel: Resulting Pauli blocking energy shift in symetric matter (black lines) and in pure neutron matter (red lines).

Example: Quark Pauli Blocking in Nuclear Matter

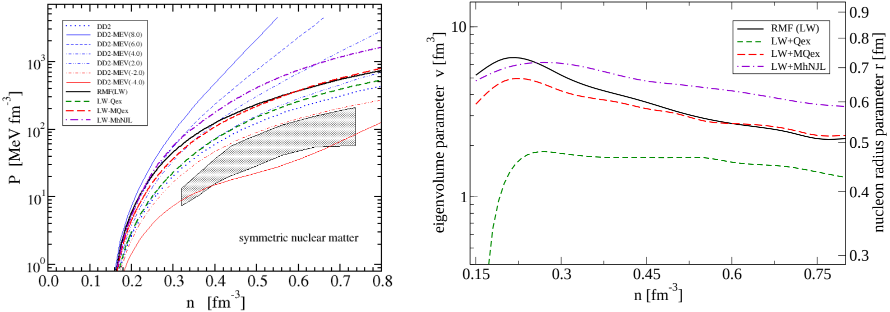


Figure: Left Panel: Pressure vs. density for chirally enhanced quark Pauli blocking within a linear Walecka model scheme. Diffent line colors stand for the quark mass scalings. For comparison the DD2 RMF model with modified excluded volume [S. Typel, EPJA 52 (2016)] is shown by blue lines (positive $v-$ parameter) and red lines (negative $v-$parameter).
Right Panel: Density-dependent excluded volume parameter and corresponding nucleonic hard-core radius.

Example: Quark Pauli Blocking in Nuclear Matter

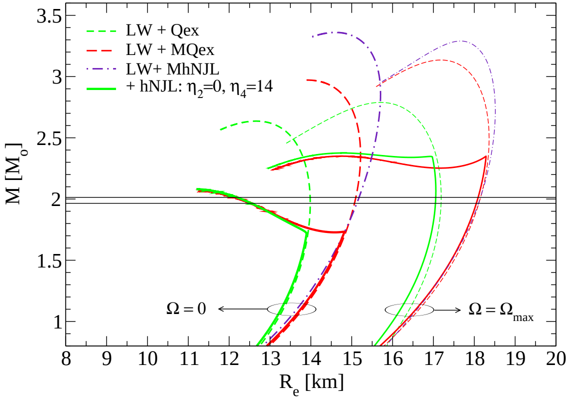


Figure: Mass vs. radius for hybrid stars resulting from a hadronic EoS with quark Pauli blocking and a higher order NJL model for quark matter.
D. Blaschke, H. Grigorian, G. Röpke, in preparation for MDPI Particles (2019).

Relativistic Density Functional Approach to Quark Matter

$$\mathcal{Z} = \int \mathcal{D}\bar{q}\mathcal{D}q \exp\left\{\int_0^\beta d\tau \int_V d^3x \left[\mathcal{L}_{\rm eff} + \bar{q}\gamma_0\hat{\mu}q\right]\right\} ~,$$

$$U(\bar{q}q, \bar{q}\gamma_0 q) = U(n_{\rm s}, n_{\rm v}) + (\bar{q}q - n_{\rm s})\Sigma_{\rm s} + (\bar{q}\gamma_0 q - n_{\rm v})\Sigma_{\rm v} + \dots ~,$$

$$\Omega = -T \ln \mathcal{Z} = \Omega^{\rm quasi} + U(n_{\rm s}, n_{\rm v}) - n_{\rm s}\Sigma_{\rm s} - n_{\rm v}\Sigma_{\rm v} ~.$$

The quasi-particle term (for the case of isospin symmetry and degenerate flavors)

$$\Omega^{\rm quasi} = -2N_c N_f T \int \frac{d^3p}{(2\pi)^3} \left\{\ln\left[1 + {\rm e}^{-\beta(E^* - \mu^*)}\right] + \ln\left[1 + {\rm e}^{-\beta(E^* + \mu^*)}\right]\right\}$$

can be calculated by using the ideal Fermi gas distribution for quarks with the quasiparticle energy $E^* = \sqrt{p^2 + M^2}$, the effective mass $M = m + \Sigma_s$ and effective chemical potential $\mu^* = \mu - \Sigma_v$. The self energies are determined by the density derivations

$$\Sigma_{\rm s} = \frac{\partial U(n_{\rm s}, n_{\rm v})}{\partial n_{\rm s}} ~, \quad \text{and}$$

$$\Sigma_{\rm v} = \frac{\partial U(n_{\rm s}, n_{\rm v})}{\partial n_{\rm v}} ~.$$

In this approach the stationarity of the thermodynamical potential

$$0 = \frac{\partial \Omega}{\partial n_{\rm s}} = \frac{\partial \Omega}{\partial n_{\rm v}}$$

is always fulfilled.

Relativistic Density Functional Approach to Quark Matter

To capture the phenomenology of a confining meanfield (string-flip model), the following density functional of the interaction is adopted,

$$U(n_{\rm s}, n_{\rm v}) = D(n_{\rm v}) n_{\rm s}^{2/3} + a n_{\rm v}^2 + \frac{b n_{\rm v}^4}{1 + c n_{\rm v}^2} \ .$$

The first term captures aspects of (quark) confinement through the density dependent scalar self-energy,

$$\Sigma_{\rm s} = \frac{2}{3} D(n_{\rm v}) n_{\rm s}^{-1/3} \ ,$$

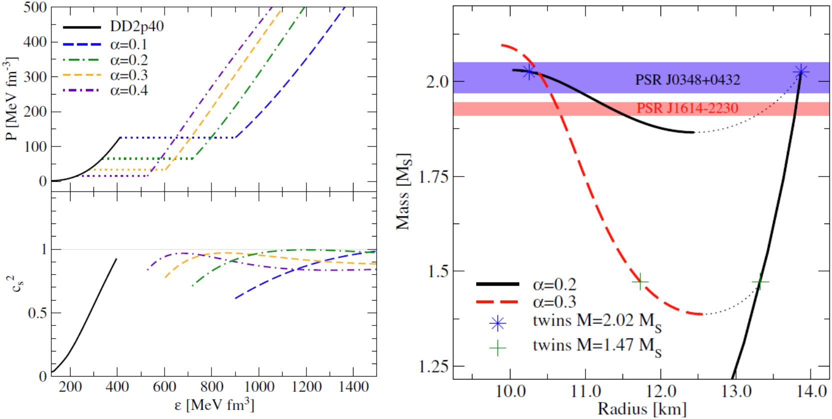
defining the effective quark mass $M = m + \Sigma_{\rm s}$. We also employ higher-order quark interactionsto obey the observational constraint of 2 $M_\odot$. The denominator in the last term of Equation (33) guarantees that the speed of sound $c_{\rm s} = \sqrt{\partial P / \partial \varepsilon}$ does not exceed the speed of light). All terms in Equation (33) that contain the vector density contribute to the shift defining the effective chemical potentials $\mu^* = \mu - \Sigma_{\rm V}$, where

$$\Sigma_{\rm v} = 2 a n_{\rm v} + \frac{4 b n_{\rm v}^3}{1 + c n_{\rm v}^2} - \frac{2 b c n_{\rm v}^5}{(1 + c n_{\rm v}^2)^2} + \frac{\partial D(n_{\rm v})}{\partial n_{\rm v}} n_s^{2/3} \ .$$

The reduction of the string tension $D(n_{\rm v}) = D_0 \phi(n_{\rm v}; \alpha)$ is modeled via a Gaussian function of the baryon density $n_{\rm v}$,

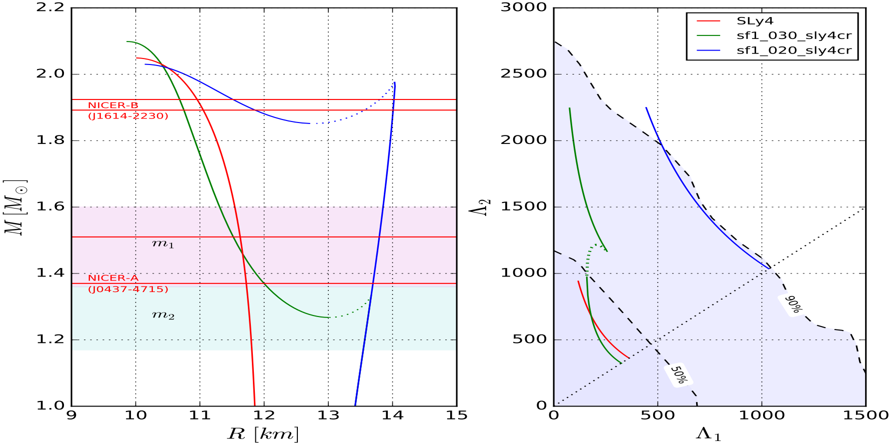
$$\phi(n_{\rm v}; \alpha) = \exp\left[-\alpha (n_{\rm v} \cdot {\rm fm}^3)^2\right] \ ,$$

Hybrid EoS: Third Family of Compact Stars & Mass Twins



KALTENBORN, BASTIAN, and BLASCHKE PHYSICAL REVIEW D **96,** 056024 (2017)

Was GW170817 NOT a neutron star (NS) - NS merger ?



GW170817 can be explained as a hybrid star (HS) - HS merger for a low-mass twin EoS as well as a NS - NS merger for a soft nuclear matter EoS. If NICER measures $R_{J0437} \geq 14$ km $\Longrightarrow$ evidence for a strong phase transition !

Summary

- cluster virial expansion developed for sunset-type Φ functionals made of cluster Green's functions and a cluster T-matrix
- cluster Φ functional approach to quark-meson-diquark-baryon system developed and example for meson dissociation outlined
- quark Pauli blocking in hadronic matter is contained in the approach
- selfconsistent density-functional approach to quark matter with confinement and chiral symmetry breaking obtained as limiting case
- applications to nuclear clustering and quark deconfinement in the astrophysics of supernovae and compact stars as well as in HIC

Outlook

- full hadron resonance gas description of the hadronic phase
- cluster virial corrections (hadron-hadron and hadron-quark scattering)
- inclusion of gluons $\rightarrow$ density functional for QGP a la Khvorostukin et al., Eur. Phys. J. C 48 (2006) 531