

Calculation of

Feynman diagrams

A.V. Kotikov (JINR, Dubna)

(11)

I Definitions:

$$\begin{array}{c} \rightarrow p \\ \text{---} \bullet \text{---} \\ d \\ \equiv \\ \frac{1}{[p^2]^d} \end{array}$$

$$\begin{array}{c} \rightarrow p \\ \text{---} \text{wavy} \text{---} \\ m^2 \\ \equiv \\ \frac{1}{[p^2 + m^2]^d} \end{array}$$

Euclidean space:
 $D = 4 - 2\epsilon$
 $\lambda = D/2 - 1 = 1 - \epsilon$

dimensional regularization
 $d^4k \rightarrow d^D k (\mu^2)^\epsilon$

(1.3)

(II) Tadpoles

$$A. \int_{\mathbb{R}^D} \frac{d^D k}{(k^2 + m^2)^\beta} \equiv \int_{\mathbb{R}^D} \frac{d^D k}{(k^2 + m^2)^\beta} \quad [k^2 = m^2 z]$$

$$d^D k \equiv \frac{1}{2} (k^2)^{\frac{D}{2}-1} dk^2 d\Omega$$

Polar coordinates:

$$\int_{\mathbb{R}^D} d\Omega = S_{D-1} = \frac{2\pi^{D/2}}{\Gamma(D/2)}$$

is the surface of the unit hypersphere in \mathbb{R}^D

↑ Euler Γ -function

$$\int_{\mathbb{R}^D} \frac{d^D k}{(k^2 + m^2)^\beta} = \int_0^\infty \frac{dz \cdot z^{\frac{D}{2}-1}}{(z+1)^\beta} \cdot \frac{1}{(m^2)^{\frac{D}{2}-\beta}} = \int_0^\infty \frac{dz \cdot z^{\frac{D}{2}-1}}{(z+1)^\beta} \cdot \frac{1}{(m^2)^{\frac{D}{2}-\beta}}$$

↑ Euler β -function

1.4

$$\int_0^\infty \frac{t^{\alpha-1} e^{-t}}{(1+t)^\alpha} dt = \int_0^\infty \frac{t^{\alpha-1} (1-t)^{\alpha-1}}{t^\alpha} dt = \int_0^1 \frac{p^\alpha (1-p)^{\alpha-1}}{p^\alpha} dp = \int_0^1 (1-p)^{\alpha-1} dp = \frac{1 - (1-1)^\alpha}{\alpha} = \frac{1}{\alpha}$$

$z+\alpha=t, t=\frac{1}{p}$

$$\left(\frac{x}{y}\right) = \frac{\pi^{\frac{d}{2}} \Gamma(\frac{d}{2}-\beta) \Gamma(d+\beta-\frac{d}{2})}{\Gamma(\frac{d}{2}) \Gamma(d)} = B(\beta, d) \quad (1)$$

$d-\beta-1 > 0, \beta > -1$

$B(0, 2) \equiv B(d)$

2. If $d+\beta < \frac{d}{2}$

$$\int_0^\infty \frac{t^{\alpha-1}}{(1+t)^\alpha} dt \xrightarrow{m^2 \rightarrow 0} 0$$

Dimensional regularization (convention): $\int_0^\infty \frac{t^{\alpha-1}}{(1+t)^\alpha} dt \xrightarrow{m^2 \rightarrow 0} 0$ for any d and β

3. Really $\int_0^\infty \frac{t^{\alpha-1}}{(1+t)^\alpha} dt \xrightarrow{m^2 \rightarrow 0} 0$ but the case $d+\beta = \frac{d}{2}$ is not relevant in dimensional regularization (it is possible to add small Δ when...)

(1.5)

4. Consider the integral

$$\int \frac{d^D k}{[k^2 - 2pk + m^2]^d} = \textcircled{x}$$

We introduce $k_1 = k - p$, then $[k^2 - 2pk + m^2] = [k_1^2 + m^2 - p^2]$

Moreover, $d^D k = d^D k_1$, then

$$\textcircled{x} = \int \frac{d^D k_1}{[k_1^2 + m^2 - p^2]^d} = \pi^{D/2} B(0, d) = \frac{\pi^{D/2}}{(m^2)^{d-D/2}} \frac{\Gamma(\frac{d-D/2}{2})}{\Gamma(d)}$$

m^2 $m^2 - p^2$

(2)

(1.6)

III Feynman parameterization

$$[k^2 + m^2]^{-1} = a_i$$

$$\frac{1}{\prod_{i=1}^N [k^2 + m_i^2]^{d_i}} = \frac{\Gamma(\sum_{i=1}^N d_i)}{\prod_{i=1}^N \Gamma(d_i)} \int_0^1 dx_1 \int_0^{1-x_1} dx_2 \dots \int_0^{1-\sum_{i=1}^{N-2} x_i} dx_{N-1} \frac{\prod_{i=1}^{N-1} x_i^{d_i-1}}{(\sum_{i=1}^N x_i a_i)^{\sum_{i=1}^N d_i}} \quad (3)$$

$$x_N = 1 - \sum_{i=1}^{N-1} x_i$$

If $N=2$

$$\frac{1}{a_1^{d_1} a_2^{d_2}} = \frac{\Gamma(d_1+d_2)}{\Gamma(d_1)\Gamma(d_2)} \int_0^1 dx_1 \frac{x_1^{d_1-1} (1-x_1)^{d_2-1}}{(x_1 a_1 + (1-x_1) a_2)^{d_1+d_2}} \quad (3a)$$

Feynman parameterization transforms products of propagators to a new propagator with effective mass, which is dependent on x_i .
 With this new propagator, we can apply formulas for tadpoles (which were done already)

I. Massless diagrams

1.6c

A.V.K., S. Teber "Multi-loop technique
for massless Feynman diagram
calculations", Phys. Part. Nucl. 50
(2019) 1. [[arXiv: 1805.05109 \[hep-th\]](https://arxiv.org/abs/1805.05109)]

(1.8)

Fourier transform

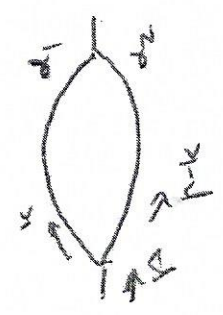
$$\boxed{\delta = \frac{1}{2} - \alpha}$$

2) Recover $\xi(\omega)$ using

$$\int \frac{d^2 p e^{i p x}}{p^2 \alpha} = \frac{2 \int \frac{d^2 \alpha}{\alpha^2 - \alpha} a(\alpha)}{\alpha^2 \alpha}$$

$$\int \frac{d^2 x e^{-i p x}}{\alpha^2 \alpha} = \frac{2 \int \frac{d^2 \alpha}{\alpha^2 - \alpha} a(\alpha)}{p^2 \alpha} \Rightarrow \frac{1}{p^2 \alpha} = \frac{1}{\pi^2} \frac{a(\alpha)}{2 \alpha}$$

Then



$$\int \frac{d^2 k}{k^2 \alpha} (p-k)^{2 \alpha} = \int \frac{d^2 k}{\pi^2} \frac{a(\alpha_1) a(\alpha_2)}{2^{2(\alpha_1 + \alpha_2)}} \int \frac{d^2 x_1 e^{-i p x_1}}{\alpha_1^2 \alpha_1} \int \frac{d^2 x_2 e^{-i(p-k)x_2}}{\alpha_2^2 \alpha_2} = \text{---}$$

$$\int d^2 x e^{i k(x_2 - x_1)} = (2\pi)^2 \delta^2(x_2 - x_1)$$

Note that

$$\text{---} = \frac{2 \int \frac{d^2 \alpha_1 e^{-i p x_1}}{\alpha_1^2 (2 + \alpha_2)} \int \frac{d^2 \alpha_2 e^{-i(p-k)x_2}}{\alpha_2^2 (\delta_1 + \delta_2)} a(\alpha_1) a(\alpha_2)}{\pi^2 \alpha} \frac{a(\alpha_1) a(\alpha_2)}{2^{2(\alpha_1 + \alpha_2)}}}{\frac{1}{(p^2)^{\alpha_1 + \alpha_2 + \alpha_2}}}$$

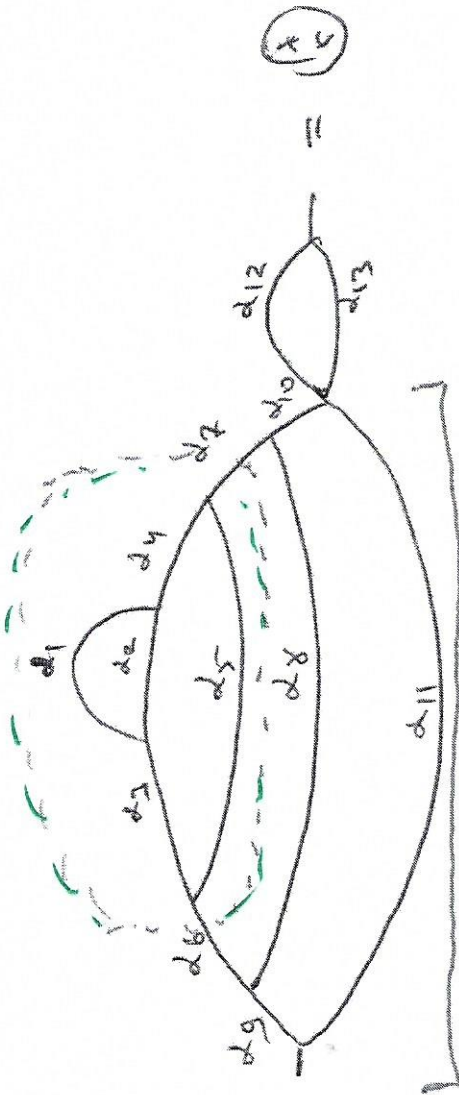
$$\delta = \alpha_1 + \alpha_2 - \alpha$$

$$k = \alpha_1 + \alpha_2 + \alpha - 2(\delta_1 + \delta_2) = 0$$

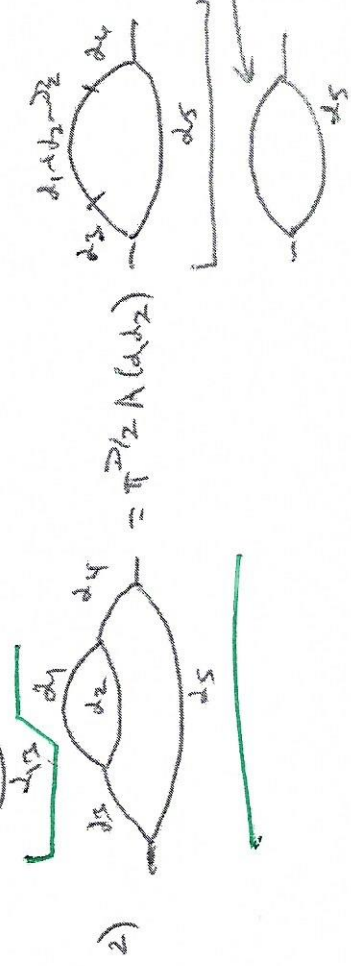
(1.9)

3) Chain $\xrightarrow{P} d_1 \cdot d_2 \xrightarrow{P} \frac{1}{(P^2)^{d_1}} \frac{1}{(P^2)^{d_2}} = \frac{1}{(P^2)^{d_1+d_2}} = \frac{1}{P^{2(d_1+d_2)}} \xrightarrow{P} d_1+d_2$ (S)

(V) So, Equations (4) and (5) give a possibility to calculate any |||| diagram, which is combination of loops and chains.

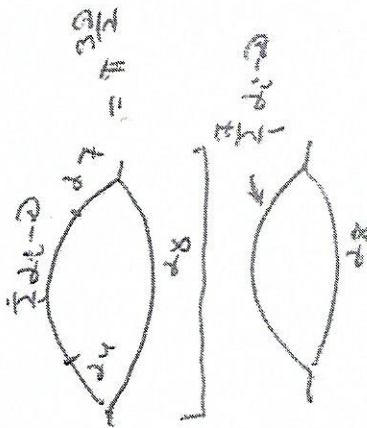


1) $\frac{d_{12}}{d_{13}} = \pi^{d_2} A(d_{12}, d_{13}) \frac{d_{12}+d_{13}-d_2}{d_{12}+d_{13}-d_2}$

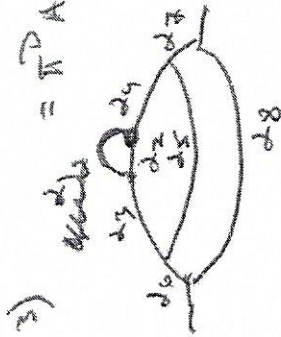


2) $\frac{d_{12}+d_{13}-d_2}{d_{12}+d_{13}-d_2} = \pi^{d_2} A(d_{12}, d_{13}) A(d_{12}, d_{13}) = \frac{\sum_i d_i - d_2}{\sum_i d_i - d_2}$

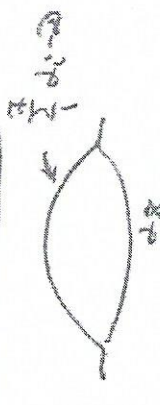
$$\frac{(2.20)}{8} \frac{\sum d_i - 3D}{2}$$



$$= \pi^D A(d_1, d_2) A\left(\sum_1^4 d_i, d_5\right)$$



$$= \pi^{\frac{3D}{2}} A(d_1, d_2) A\left(\sum_1^4 d_i, d_5, d_6, d_7, d_8\right)$$



number of loops
 $\sum_1^4 d_i - \frac{5D}{2}$

$$\textcircled{2} = \pi^{\frac{5D}{2}} A(d_1, d_2) A\left(\sum_1^4 d_i, d_5, d_6, d_7, d_8\right) \cdot A(d_{10}, d_{11})$$

4) To find "number" it is useful to use the relations $\Gamma(n+1) = n\Gamma(n)$ and the

$$\text{expansion for } \Gamma(1-\alpha\epsilon): \Gamma(1+\alpha\epsilon) = \exp\left[-\gamma_E \alpha\epsilon + \sum_{k=2}^{\infty} \frac{\zeta(k)}{k} (-\alpha\epsilon)^k\right]$$

number of integrations

Note: The terms $\pi^{\frac{D}{2}}$ should be not considered really in results,

because every integration gives:

$$\int \frac{\pi^{\frac{D}{2}} \Gamma(1+\epsilon)}{(2\pi)^D} \rightarrow \text{from measure of the integral}$$

$$\cdot g^2 (\mu^2)^2 = \frac{g^2}{16\pi^2} \Gamma(1+\epsilon) (\mu^2)^2 \equiv \frac{g^2}{16\pi^2}$$

dimensional regularization
 \uparrow charge squared

$$\frac{(4\pi)^{\epsilon} \Gamma(1+\epsilon) (\mu^2)^{\epsilon}}{(\mu^2)^{\epsilon}} \leftarrow \text{MS-scheme}$$

3) let $d_i = 1$

So, our result has the following form of arguments of Γ -function

$$\sum_1^4 d_i - \frac{D}{2} = 2 + 1, \sum_1^7 d_i - D = 3 + 2 + 1, \sum_1^{10} d_i - \frac{3D}{2} = 4 + 3 + 2$$

$$A^2(1,1) A(2+1,1) A(3+2,1) A(4+3,1) = \textcircled{A}$$

$$A(1,1) = \frac{\Gamma^2(1) \Gamma(2)}{\Gamma^2(1) \Gamma(2)} = \frac{\Gamma(1) \Gamma(2)}{\Gamma(1) \Gamma(2)}, \quad A(2+1,1) = \frac{\Gamma(1) \Gamma(2) \Gamma(3)}{\Gamma(1) \Gamma(2) \Gamma(3)}$$

$$\Gamma(a) = \Gamma(a-1) \Rightarrow \Gamma(a-1) = \frac{\Gamma(a)}{\Gamma(a)}$$

$$A(3+2,1) = \frac{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4)}{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4)} = -\frac{2}{3} \frac{\Gamma(1) \Gamma(2) \Gamma(3)}{\Gamma(1) \Gamma(2) \Gamma(3)}$$

$$A(4+3,1) = \frac{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)}{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)} = -\frac{5}{24} \frac{\Gamma(1) \Gamma(2) \Gamma(3)}{\Gamma(1) \Gamma(2) \Gamma(3)}$$

$$\textcircled{A} = -\frac{5}{36 \cdot 2} \frac{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)}{\Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)} = \frac{\Gamma^5(1) \Gamma(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)}{\Gamma^5(1) \Gamma(2) \Gamma(3) \Gamma(4) \Gamma(5)}$$

~~function~~
constant

Euler

$$\Gamma(1+\alpha\epsilon) = \exp\left[\alpha\gamma\epsilon + \sum_{\nu=2}^{\infty} \frac{(-\alpha\epsilon)^{\nu}}{\nu} \zeta(\nu)\right]$$

Euler constant

1.12

$$K_2 = \frac{\Gamma^7(1-\epsilon) \Gamma(1+\epsilon)}{\Gamma(1-2\epsilon) \Gamma(1-5\epsilon) \Gamma^4(1+\epsilon)} = \frac{\exp\left[-3\gamma\epsilon\right] + \sum_{\nu=2}^{\infty} \zeta(\nu) \frac{7 + (-1)^{\nu} 4^{\nu}}{\nu} \epsilon^{\nu}}{\exp\left[-3\gamma\epsilon\right] + \sum_{\nu=2}^{\infty} \zeta(\nu) \frac{2^{\nu} + 5^{\nu} + (-1)^{\nu} 4^{\nu}}{\nu} \epsilon^{\nu}}$$

$$= \exp\left[\sum_{\nu=2}^{\infty} \zeta(\nu) \frac{Q_1(\nu)\epsilon^{\nu}}{\nu}\right] = \mathcal{E}$$

$$Q_1(\nu) = 7 - 2^{\nu} - 5^{\nu} + (-1)^{\nu} (4^{\nu} - 4)$$

$$, \quad Q_1(\nu=2) = -10, \quad Q_1(\nu=3) = -186$$

$$\mathcal{E} = \exp\left[-5\zeta(2)\epsilon^2 - 62\zeta(3)\epsilon^3 + \mathcal{O}(\epsilon^4)\right]$$

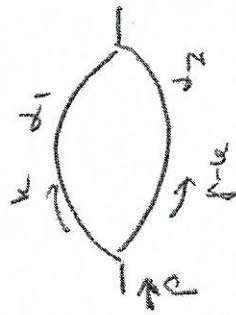
But already at two-loop



(VI)

Dual representation: all $p \rightarrow x \Rightarrow$

4.13



$$= \int \frac{d^D k}{k^{2+d_1} [(p-k)^2]^{d_2}}$$

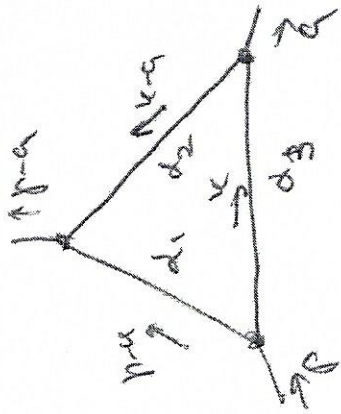
$$\int \frac{d^D x}{x^{2+d_1} [(y-x)^2]^{d_2}} =$$

$$= \frac{2^{d_1+d_2}}{\Gamma(d_1)\Gamma(d_2)} \int \frac{d^D y}{y}$$

\Rightarrow same integral but another graphic representation
 in partial, loops \leftrightarrow chains

(VII)

Uniqueness relation

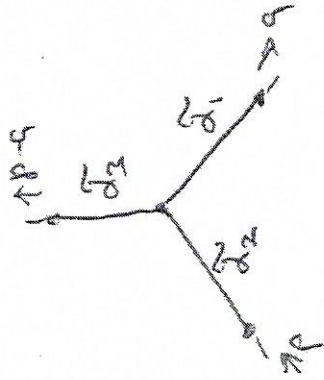


$$\sum_i d_i = D$$

$$\frac{3\pi}{2} A(d_1, d_2)$$

$$d_i = z_i - z_j$$

(1.14)



|||

$$\int \frac{d^D k}{k^{2\alpha} [(p-k)^2]^{d_1} [(k-q)^2]^{d_2}}$$

$$\frac{1}{[q-p]^2]^{d_2} p^{2d_2} q^{2d_1}}$$

from conformal invariance

Conformal transformations save angles between vectors (really ~~we~~ we will use inversion)

1) To prove uniqueness relation it is convenient to perform the inversion:

$$k_1 = \frac{1}{k_1}, p = \frac{1}{p_1}, q = \frac{1}{q_1}$$

$$k^2 = \frac{1}{k_1^2}, (p-k)^2 = p^2 - 2(pk) + k^2$$

Conformal
transformation \rightarrow

$$\frac{1}{p_1^2} - 2 \frac{(pk)}{p_1^2 k_1^2} + \frac{1}{k_1^2} = \left(\frac{p_1 - k_1}{p_1^2 k_1^2} \right)^2$$

1.15

$$d^D k = \frac{1}{k_1^2} d^D k_1$$

$$\int_0^\infty \int_0^\infty k_1^{2\alpha} dk_1^2 = \int_0^\infty \int_0^\infty k_1^{2\alpha} dk_1^2 \frac{1}{(k_1^2)^\alpha}, \text{ so } d^D k = \frac{d^D k_1}{(k_1^2)^\alpha}$$

$$\int \frac{d^D k}{k^{2\alpha_1} [(p-k)^2]^{\alpha_2}} = \int \frac{d^D k_1 \cdot k_1^{2\alpha_1} (k_1^2/p_1^2)^{\alpha_2} (k_1^2/q_1^2)^{\alpha_2}}{(k_1^2)^{\alpha_1} [(p_1 - k_1)^2]^{\alpha_2} [(q_1 - k_1)^2]^{\alpha_2}}$$

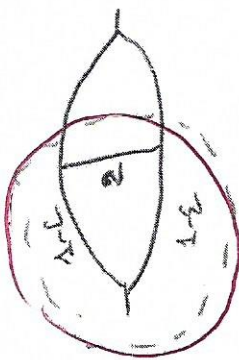
easy

$$\frac{(p_1^2)^{\alpha_1} (q_1^2)^{\alpha_2}}{[(p_1 - q_1)^2]^{\alpha_1 + \alpha_2}} = \frac{1}{p_1^{2\alpha_1} q_1^{2\alpha_2} [(p_1 - q_1)^2]^{\alpha_1 + \alpha_2}}$$

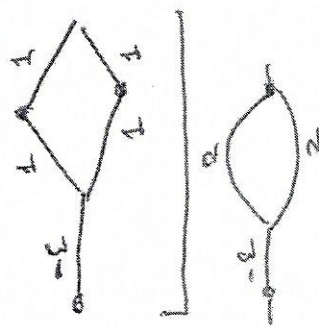
$$= \frac{1}{p_1^{2\alpha_1} q_1^{2\alpha_2}} \frac{(p_1^2 q_1^2)^{\alpha_1 + \alpha_2 - \alpha_2}}{[(p_1 - q_1)^2]^{\alpha_1 + \alpha_2 - \alpha_2}} = \frac{1}{p_1^{2\alpha_1} q_1^{2\alpha_2}} \frac{1}{[(p_1 - q_1)^2]^{\alpha_1 + \alpha_2 - \alpha_2}}$$

1.16

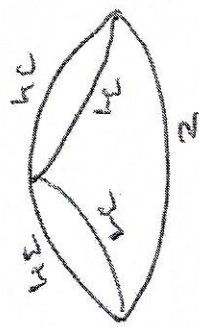
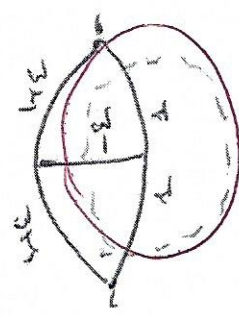
2) Examples



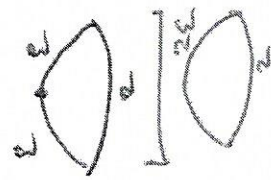
$$= \pi^2 A(\delta-\epsilon, \delta-\epsilon)$$



$$= \pi^2 A(\epsilon, \epsilon) A(2, 2)$$



$$= \frac{A^2(\delta-\epsilon, \delta-\epsilon)}{A(\epsilon, \epsilon)}$$



$$= \frac{A^2(\epsilon, \epsilon)}{A(\epsilon, \epsilon)} A(2, 2)$$

We introduce the additional integration



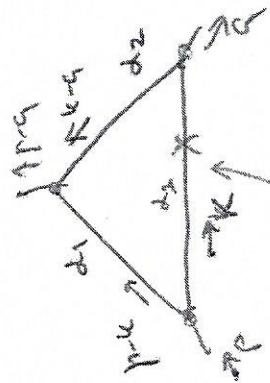
VIII

Chetyrin, Trachen, 1981 (1.17)
 Vassiliev et al., 1981

IBP (Integration by parts) procedure

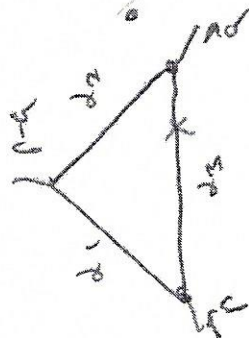
$$\int \frac{d^D k}{[k^{2d_2} (k-q)^{2d_2} (k-p)^{2d_1}]} = \int d^D k \left[\frac{d}{dk_p} \left(\frac{k^p}{[\]} \right) - k^p \frac{d}{dk_p} \frac{1}{[\]} \right]$$

= 0 Stora's theorem

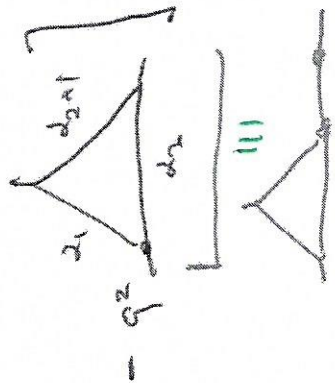
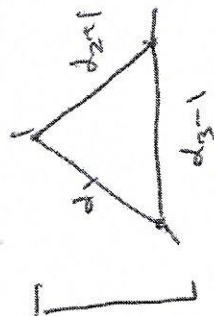


distinguished line

$$\left. \begin{aligned} -k^p \frac{d}{dk_p} \frac{1}{k^{2+2d_2}} &= \frac{2d_2}{k^{2+2d_2}} \\ -k^p \frac{d}{dk_p} \frac{1}{(k-q)^{2+2d_2}} &= 2d_2 \frac{(k, q-p)}{(k-q)^{2+2d_2+1}} = d_2 \frac{k^2 + (k-q) \cdot (-q)}{[(k-q)^2]^{d_2+1}} \end{aligned} \right\}$$



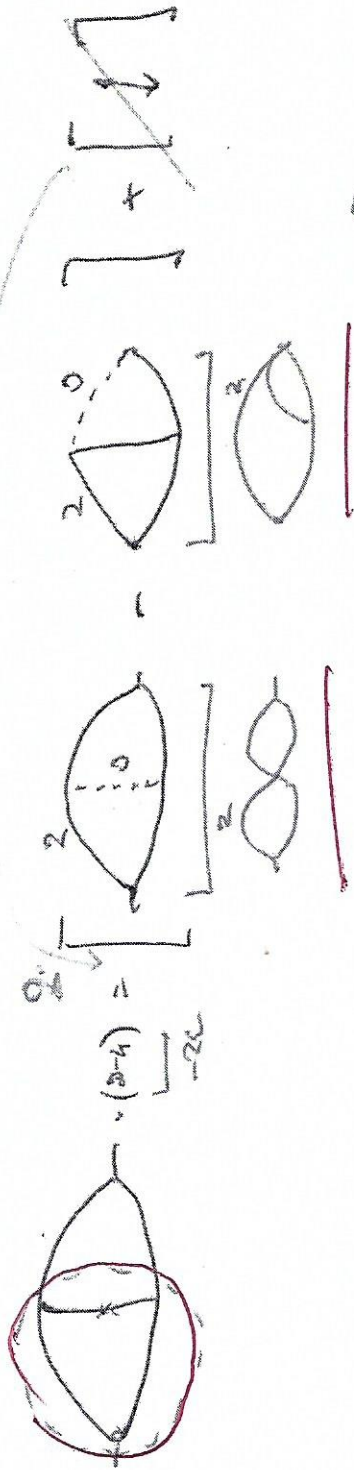
$$(D - 2d_2 - d_1 - d_2) = d_2$$



\leftrightarrow

Example:

1.18



$$S_0 = -\frac{1}{2} \left[\text{Diagram 1} - \text{Diagram 2} \right] = -\frac{A(1,1)}{2} \left[\text{Diagram 3} - \text{Diagram 4} \right] = -\frac{A(2,1) \epsilon}{2} \xrightarrow{z \rightarrow \infty}$$

$$c = -\frac{1}{2} \frac{\Gamma^2(\mu \epsilon) \Gamma(\epsilon)}{\Gamma^2(1) \Gamma(2\epsilon)} \left[\frac{\Gamma(\mu \epsilon) \Gamma(\mu \epsilon) \Gamma(\mu \epsilon)}{\Gamma(\epsilon) \Gamma(\mu \epsilon) \Gamma(\mu \epsilon)} - \frac{\Gamma(\epsilon) \Gamma(\mu \epsilon) \Gamma(\mu \epsilon)}{\Gamma(2) \Gamma(1) \Gamma(\mu \epsilon)} \right] = \frac{1}{2} \frac{\Gamma^2(\mu \epsilon) \Gamma^4(1-\epsilon)}{(1-2\epsilon) \Gamma^2(\mu \epsilon)} \left[1 - \frac{\Gamma^2(1-2\epsilon) \Gamma(1+2\epsilon)}{\Gamma^2(\mu \epsilon) \Gamma(1+3\epsilon) \Gamma(\mu \epsilon)} \right]$$

$$k_1 = \frac{\Gamma^4(\mu \epsilon)}{\Gamma^2(\mu \epsilon)} = \epsilon \kappa \Gamma \left[4\delta \epsilon + 4 \sum_{k=0}^{\infty} \frac{q_1(k)}{k} \epsilon^k \right] = \epsilon \kappa \Gamma \left[4 \sum_{k=2}^{\infty} \frac{q_1(k)}{k} (1-2^{k-1}) \epsilon^k \right]$$

$q_1(0) = 0, q_1(1) = -3, q_1(2) = -7, q_1(3) = -15, q_1(4) = -31$ (homework)

1.19

$$k_2 = \frac{r^2(1+c)(1+2c)}{r^2(1+c)(1+c)} = \frac{2x \left[2x + \sum_{k=2}^{\infty} \frac{r(1+c)}{k} [2^{k-1} + (-2)^k] \right] \epsilon^k}{2x \left[2x + \sum_{k=2}^{\infty} \frac{r(1+c)}{k} [2 \cdot (-1)^k + 3^k - 1] \right] \epsilon^k} \quad \oplus$$

$q_2(z) = 0, q_2(3) = -18, q_2(4) = -36, q_2(5) = -240, q_2(6) = -630$ homework

$$\textcircled{5} = 2x \left[-68(2) \epsilon^3 - 98(4) \epsilon^4 - 188(5) \epsilon^5 - 1058(6) \epsilon^6 + 0(\epsilon^7) \right] = (1-6)8(5) \epsilon^3 + 188(5) \epsilon^4$$

$$(1-9)8(4) \epsilon^4 - 3(17)8505 - 3(15)884 - 488(5) \epsilon^5 + 0(\epsilon^6) = 1(43)0 + 0(\epsilon^3) = 1-68(5) \epsilon^3 - 98(4) \epsilon^4 - 488(5) \epsilon^5$$

$$= (1058(6) - 189^2(3)) \epsilon^6$$

$$C = \frac{r^2(1+c)}{(1+c)} \left[68(3) + 98(4) \epsilon + 3(15)884 + (1058(6) - 189^2(3)) \epsilon^3 \right] \frac{r^2(1+c)(1+c)}{21(5^2(1) - 2^2(3)) \epsilon^2}$$

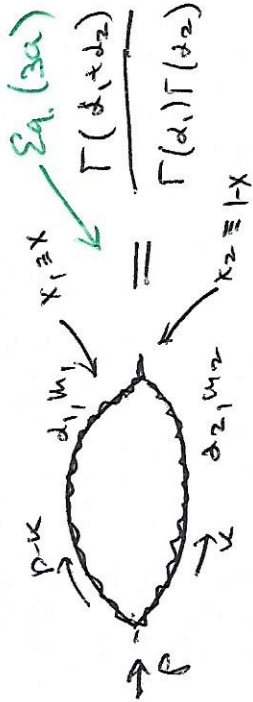
2.1

II Evaluating diagrams

with masses

One-loop

Using Feynman parameters



$$[\Gamma]_1 = X \left\{ \underbrace{p^2 - k^2 + m_1^2}_{\mu^2} + (1-X) \left\{ k^2 + m_2^2 \right\} \right\} = k^2 + p^2 X - 2(p \cdot k) X + m_1^2 X + m_2^2 (1-X)$$

$$m_{1,2}^2 = m_{eff}^2 - p^2 X^2 = p^2 X(1-X) + m_1^2 X + m_2^2 (1-X) = X(1-X) \left[p^2 + \frac{m_1^2 X + m_2^2}{1-X} \right]$$

$$\textcircled{3} = \pi^{\frac{d_1+d_2-2}{2}} \int_0^1 dx \cdot X^{\frac{d_1+d_2-2}{2}} (1-X)^{\frac{d_2-1-d_1}{2}} \frac{1}{(p^2 + \mu^2)^{d_1+d_2-2}}$$

$$\mu^2 = \frac{m_1^2}{1-X} + \frac{m_2^2}{X}$$

Loop as the integral from the tadpole propagator with the effective mass μ :
 \Rightarrow L-loop as integral from (L-1) loop with the effective mass μ .

$$\text{Eq. (2)} \quad \pi^{\frac{d_1+d_2-2}{2}} \frac{\Gamma(d_1+d_2-2)}{\Gamma(d_1+d_2)} \frac{1}{(m_{eff}^2)^{d_1+d_2-2}} = \int \frac{d^D K}{[\Gamma]^{d_1+d_2}} = \textcircled{3}$$

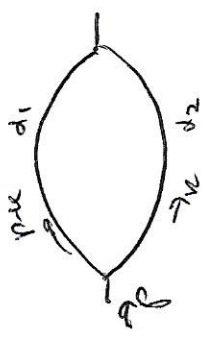
$$\int_0^1 dx \cdot X^{d_1-1} (1-X)^{d_2-1}$$

$$\frac{X^{d_1-1} (1-X)^{d_2-1}}{[X(1-X)]^{d_1+d_2-2}}$$

$$= \pi^{\frac{d_1+d_2-2}{2}} \frac{\Gamma(d_1+d_2-2)}{\Gamma(d_1)\Gamma(d_2)} \int_0^1 dx \cdot X^{\frac{d_1-1-d_2}{2}} (1-X)^{\frac{d_2-1-d_1}{2}}$$

1) massless case: $p^2=0$ (we considered already)

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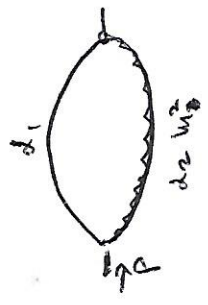


$$= \int \frac{d^2x}{(p^2)^{d_1+d_2-2}} \frac{\Gamma(d_1+d_2-2)}{\Gamma(d_1)\Gamma(d_2)} \int_0^1 dx \cdot x^{\frac{d_2-1}{2}-1} (1-x)^{\frac{d_1-1}{2}-1} \frac{\Gamma(\frac{d_2-d_1}{2})\Gamma(\frac{d_2-d_2}{2})}{\Gamma(d_1-d_1-d_2)}$$

$$= \pi^2 A(d_1, d_2) \frac{1}{(p^2)^{d_1+d_2-2}} \quad (4)$$

$$A(d_1, d_2) = \frac{a(d_1)a(d_2)}{a(d_1+d_2-2)}, \quad a(d) = \frac{\Gamma(d)}{\Gamma(d)}, \quad d = \frac{d_2-d_1}{2}$$

2) the case $m_1^2=0, m_2^2=m^2$

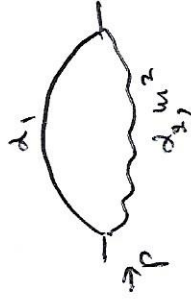


$$\int_0^1 dx \cdot x^{\frac{d_2-1}{2}-1} (1-x)^{\frac{d_2-1}{2}-1} \frac{\Gamma(\frac{d_2-d_1}{2})\Gamma(\frac{d_2-d_2}{2})}{\Gamma(d_1-d_1-d_2)}$$

$$= \int \frac{d^2x}{(p^2)^{d_1+d_2-2}} \frac{\Gamma(d_1+d_2-2)}{\Gamma(d_1)\Gamma(d_2)} \int_0^1 dx \cdot x^{\frac{d_2-1}{2}-1} (1-x)^{\frac{d_2-1}{2}-1} \frac{\Gamma(\frac{d_2-d_1}{2})\Gamma(\frac{d_2-d_2}{2})}{\Gamma(d_1-d_1-d_2)}$$

~~Engelmann~~
~~Vektor~~
 ~~$\vec{p} = m^2$~~
 ~~$\vec{k} = m^2$~~
 ~~$\vec{p} = m^2$~~

2a Important on-shell case: $p^2 = -m^2 \Rightarrow (p^2 + m^2) = m^2 (1-x)$



$$= \int_0^1 dx \frac{\Gamma(d_1 + d_2 - 2\epsilon)}{\Gamma(d_1) \Gamma(d_2)} \frac{1}{(m^2)^{d_1 + d_2 - 2\epsilon}} \int_0^1 dx \cdot x^{d_1-1} (1-x)^{d_2-1-2d_1-d_2} = \frac{\Gamma(d_1) \Gamma(d_2 - 2d_1 - d_2)}{\Gamma(2d_1 - d_2)}$$

$$= \int_0^1 dx \frac{\Gamma(d_1 + d_2 - 2\epsilon) \Gamma(2d_1 - d_2)}{\Gamma(d_2) \Gamma(2d_1 - d_2)} \frac{1}{(m^2)^{d_1 + d_2 - 2\epsilon}}$$

2b Inverse mass expansion:

$$\frac{1}{[p^2 + m^2]^{d_1 + d_2 - 2\epsilon}} = \frac{1}{(m^2)^{d_1 + d_2 - 2\epsilon}} \sum_{k=0}^{\infty} \frac{\Gamma(h + d_1 + d_2 - 2\epsilon)}{h! \Gamma(d_1 + d_2 - 2\epsilon)} (-1)^k \left(\frac{p^2}{m^2}\right)^k$$

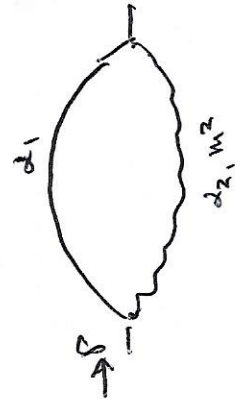
$$\int_0^1 dx \cdot x^{d_1-1} (1-x)^{d_2-d-d_1} \frac{1}{[p^2 + m^2]^{d_1 + d_2 - 2\epsilon}} = \frac{1}{(m^2)^{d_1 + d_2 - 2\epsilon}} \sum_{k=0}^{\infty} \frac{\Gamma(h + d_1 + d_2 - 2\epsilon)}{h! \Gamma(d_1 + d_2 - 2\epsilon)} \frac{\Gamma(d_1) \Gamma(2\epsilon - d_1)}{\Gamma(2\epsilon + h)} (-1)^k \left(\frac{p^2}{m^2}\right)^k =$$

$$= \frac{1}{(m^2)^{d_1 + d_2 - 2\epsilon}} \mathbb{F}_2 \left[\begin{matrix} d_1, d_1 + d_2 - 2\epsilon \\ 2\epsilon \end{matrix} \middle| -\frac{p^2}{m^2} \right] \cdot \frac{\Gamma(d_1) \Gamma(2\epsilon - d_1)}{\Gamma(2\epsilon)}$$

\mathbb{F}_2 - hypergeometric function

(2.5)

(1.5)

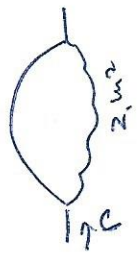


$$F_{z_1} \left[\frac{d_1, d_1 + d_2 - 2\alpha_2}{z_1} \mid -\frac{p^2}{w^2} \right]$$

$$\frac{\Gamma(\frac{\alpha_2}{2} - d_1) \Gamma(d_1 + d_2 - 2\alpha_2)}{\Gamma(d_2) \Gamma(2\alpha_2)} \frac{1}{(w^2)^{d_1 + d_2 - 2\alpha_2}}$$

$B(d_1, d_2) \neq$ (see tadpole diagram)

Let $d_1 = 1, d_2 = 1$



$$F_{z_1} \left[\frac{d_1, m_2}{z_1} \mid -\frac{p^2}{w^2} \right] = \left(\frac{x}{x-1} \right)$$

$$F_{z_1} \left[\frac{d_1, m_2}{z_1} \mid -x \right] = \sum_{k=0}^{\infty} \frac{\Gamma(k+1+k)}{k! \Gamma(k)} \frac{\Gamma(2-k)}{\Gamma(k+k)} (-x)^k = (1-x) \sum_{k=2}^{\infty} \frac{(-x)^k}{k+1-k} \frac{\Gamma(k+1+k)}{\Gamma(k+k)}$$

$$\frac{(1-x)}{(-x)} \sum_{k=0}^{\infty} \frac{(-x)^k}{k-k} \frac{\Gamma(k+k)}{\Gamma(k+k)} = \frac{(1-x)}{(-x)} \sum_{m=2}^{\infty} \epsilon^m \phi_m(-x)$$

$$\left(\frac{x}{x-1} \right) = \frac{\Gamma(\frac{\alpha_2}{2})}{(w^2)^{\frac{\alpha_2}{2}}} \frac{\Gamma(m)}{(-p^2)} \sum_{m=2}^{\infty} \epsilon^m \phi_m(-x)$$

2.6

2.6

harmonic sums

$$\frac{\Gamma(2k)\Gamma(2l)}{\Gamma(k)\Gamma(l)} = \sum_{m=1}^{\infty} \left[- \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \frac{\Gamma(1-n)}{\Gamma(1-n)} \right] \frac{\Gamma(1-n)}{\Gamma(1-n)} = S_m$$

$$S_m(k) = \sum_{r=1}^k \frac{(-1)^{r+m}}{r^m} \sum_{\pm m_1, \pm m_2, \dots} \frac{(-1)^{m_1}}{r^{m_1}} \sum_{\pm m_2, \dots} \frac{(-1)^{m_2}}{r^{m_2}} \dots$$

nested sums

$$\frac{\Gamma(2k)\Gamma(2l)}{\Gamma(k)\Gamma(l)} = \exp \left[- \sum_{m=1}^{\infty} \frac{(-1)^m}{m} S_m \right] \frac{\Gamma(1-2k)\Gamma(1-2l)}{\exp \left[- \sum_{m=1}^{\infty} \frac{(-1)^m}{m} S_m \right]}$$

$$= \exp \left[\sum_{m=1}^{\infty} \frac{(1-(-1)^m)}{m} S_m \right] = \exp \left[2 \sum_{m=2k+1}^{\infty} \frac{(-1)^m}{m} S_m \right]$$

(34)

$$= \exp \left[2S_1 + \frac{2}{3}S_3 + O(\epsilon^4) \right] = (1+2S_1 + \frac{2}{3}S_3 + \dots) \left(1 + \frac{2}{3}S_3 + \dots \right) = (1+2S_1 + 2S_1^2 + \frac{2}{3}S_3 + \frac{2}{3}S_3^2 + \dots) + O(\epsilon^4)$$

(33) 0 +

$$\sum_{k=1}^{\infty} \frac{(-1)^k}{k} \frac{\Gamma(1-k)\Gamma(2k)}{\Gamma(k)\Gamma(2k)} = \sum_{k=1}^{\infty} \frac{(-1)^k}{k} \left(1 + \frac{2}{3}S_3 + \frac{1}{k} + \frac{2S_1^2}{k} + \frac{2S_1}{k^2} + \frac{2S_1^2}{k} + \frac{2S_1}{k^2} + \frac{1}{k^2} \right) + O(\epsilon^4)$$

$$\frac{1}{k-2} = \frac{1}{k} + \frac{1}{k^2} + \frac{1}{k^3} + \dots$$

$$= \sum_{m=2}^{\infty} (-1)^m \psi_m(z)$$

Flensburg, A.V.K., Venetian,
1998

2.17

2.17

$$\sum_{n=1}^{\infty} \frac{z^n}{n^a} = \zeta_a(z), \quad \sum_{n=1}^{\infty} \frac{z^n}{n^a} = S_{a-1,2}(z)$$

$$\text{Li}_2(z) = -\text{Li}_1(z) \quad S_0(z) = \sum_{n=1}^{\infty} \ln^2(1-z) = \sum_{n=1}^{\infty} \ln^2(z) = L_1^2(z)$$

Nielsen
Polylogarithmic
functions

$$S_{a+1,b}(z) = \frac{(-1)^{a+b}}{a!b!} \int_0^1 \ln^a(x) \ln^b(1-zx) dx, \quad \text{Li}_a(z) = S_{a-1,1}(z)$$

$$\sum_{n=1}^{\infty} \frac{z^n}{n^2} = L_1 L_2 + \frac{1}{2} L_1^2 - 2S_{1,2}, \quad \sum_{n=1}^{\infty} \frac{z^{2n}}{n^2} = \frac{1}{2} L_2^2 + 2S_{1,2} - 2S_{2,2}, \quad \sum_{n=1}^{\infty} \frac{z^{4n}}{n^2} = L_1 L_2 - \frac{1}{2} L_2^2$$

$$\sum_{n=1}^{\infty} \frac{z^{2n}}{n^3} = \frac{1}{4} L_1^4 + L_1 L_2 - \frac{1}{2} L_2^2 + \frac{3}{2} L_1^2 L_2 - 3L_1 S_{1,2}$$

$$\sum_{n=1}^{\infty} \frac{z^{2n}}{n^4} = L_1 L_2 = -\text{Li}_1(L_2) = \psi_0(z)$$

$$\sum_{n=1}^{\infty} \frac{z^{2n}}{n^4} (2S_{1,2} - \frac{1}{2} L_2^2) = L_1^2(z) + \psi_0(z) = \psi_1(z)$$

$$\sum_{n=1}^{\infty} \frac{z^{2n}}{n^4} (2S_{1,2}^2 + \frac{2S_{1,2}}{n} + \frac{1}{n^2}) = 2[L_1 L_2 + \frac{1}{2} L_1^2 - 2S_{1,2}] + 2S_{1,2} + L_4 = 2[\frac{1}{2} L_1^2 + L_1 L_2 - S_{1,2}] + L_4 = \psi_2(z)$$

$$\sum_{n=1}^{\infty} \frac{z^{2n}}{n^4} \left[\frac{2}{3} (2S_{1,2}^2 + S_{2,2}) + \frac{2S_{1,2}}{n} + \frac{1}{n^2} \right] = \frac{2}{3} \left[2 \left(\frac{1}{2} L_1^4 + L_1 L_2 - \frac{1}{2} L_2^2 + 3 \frac{1}{2} L_1^2 L_2 - 3 L_1 S_{1,2} \right) + L_1 L_2 - \frac{1}{2} L_2^2 \right] +$$

$$+ 2 \left\{ \frac{1}{2} L_2^2 + 2S_{1,2} - 2S_{2,2} \right\} + 2S_{2,2} + L_4 = \frac{2}{3} \left[\frac{1}{2} L_1^4 + 3L_1^2 L_2 + 3L_1 L_2 - 6L_1 S_{1,2} - \frac{3}{2} L_2^2 \right] + 2 \left[\frac{1}{2} L_2^2 + 2S_{1,2} - S_{2,2} \right] + L_4 = \psi_3(z)$$

Now there are popular the Remiddi-Verwaseven (RV) Polynomials and the GonZarov V Polynomials

$$H_0(z) = \int \frac{dz}{z} \equiv \ln z; \quad H_{-1}(z) = \int \frac{z dz'}{1-z'} \equiv -\ln(1-z), \quad H_{-1}(z) = \int \frac{dz'}{1-z'} = \ln(1+z)$$

RV polynomials:

$$H_{0, b_1, \dots}(z) = \int \frac{z dz'}{z} H_{b_1, \dots}(z); \quad H_{-1, b_1, \dots}(z) = \int \frac{z dz'}{1-z'} H_{b_1, \dots}(z)$$

$$H_{-1, b_1, \dots}(z) = \int \frac{dz'}{1+z'} H_{b_1, \dots}(z) \text{ ?}$$

Properties $H_{-1, \dots, -1}(z) = \frac{1}{m!} \ln^m(1-z)$, $H_{-1, \dots, -1}(z) = \frac{1}{m!} \ln^m(1+z)$

G polylogarithms

$$G_{a, b_1, \dots}(z) = \int \frac{dz'}{z'^{1-a}} G_{b_1, \dots}(z')$$

RV polylogarithms

$$L_{11} z(z-1) = L_1(z) \quad , \quad L_2 z(z-1) = - \int_0^z \frac{d\tau}{\tau} L_1(1-\tau) = - \int_0^z \frac{d\tau}{\tau} L_1(\tau) = \int_0^z \frac{d\tau}{\tau} H_1(\tau) = K_{0,1}(z)$$

$$L_{13}(z) \quad || \int_0^z \frac{d\tau}{\tau} L_1(\tau) = K_{0,1}(z)$$

$$L_3 = \int_0^z \frac{d\tau}{\tau} L_2(z-1) = K_{0,1}(z) \quad , \quad L_4 = K_{0,0,0,1}(z) \Rightarrow L_4 = K_{0,0,0,1}(z)$$

$$S_{0,2}(z) = \frac{1}{2} L_2(z) = - \int_0^z \frac{d\tau}{\tau} L_1(1-\tau) = \int_0^z \frac{d\tau}{\tau} L_1(\tau) = K_{1,1}(z) \quad (OK, probably)$$

$$\int_0^z \frac{d\tau}{\tau} L_1(1-\tau) = \int_{1-\tau}^1 \frac{d\tau}{\tau} L_1(\tau) = - \int_{1-\tau}^1 L_1(\tau) = - \frac{1}{2} L_2(1-\tau)$$

$$S_{1,2}(z) = \int_0^z \frac{d\tau}{\tau} S_{0,2}(1-\tau) = K_{0,1}(z) \quad , \quad S_{1,2}(z) = \int_0^z \frac{d\tau}{\tau} S_{1,2}(z) = K_{0,0,1,1}(z)$$

$$S_{0,3}(z) = \int_0^z \frac{d\tau}{\tau} L_2(z-1) = \int_0^z \frac{d\tau}{\tau} L_2(1-\tau) = \int_0^z \frac{d\tau}{\tau} L_2(\tau) = K_{1,1,1}(z) \quad (OK)$$

$$\int_0^z \frac{d\tau}{\tau} L_2(1-\tau) = \int_{1-\tau}^1 \frac{d\tau}{\tau} L_2(\tau) = \int_{1-\tau}^1 L_2(\tau) = - \frac{1}{3} L_3(1-\tau)$$

$$S_{1,3}(z) = \int_0^z \frac{d\tau}{\tau} S_{0,3}(z) = K_{0,1,1,1}(z)$$

So, we expressed all elements of our results through RV polylogarithms !!!

$$p^2 x^{\gamma} = \frac{1}{x^{(d_1)}} [p^2 x^{(d_1)} + w^2]$$

$$w^2 = \frac{w^2}{x} + \frac{w^2}{x} = \frac{w^2}{x^{(d_1)}}$$

3) the case $w_1^2 = w_2^2 = w^2$

$$d_1, w_2 \quad d_2, w_2$$

$$\int \frac{dx \cdot x^{d_1-1} (wx)^{d_2-1}}{[p^2 x^{(d_1)} + w^2]^{d_1+d_2-2d_2}} = \frac{\Gamma(d_1+d_2-2d_2)}{\Gamma(d_1)\Gamma(d_2)} = \frac{1}{x}$$

$$\frac{1}{[p^2 x^{(d_1)} + w^2]^{d_1+d_2-2d_2}} = \frac{1}{(w^2)^{d_1+d_2-2d_2}} \sum_{n=0}^{\infty} \frac{\Gamma(n+d_1+d_2-2d_2)}{n! \Gamma(d_1+d_2-2d_2)} (-1)^n \left[\frac{p^2 x^{(d_1)}}{w^2} \right]^n$$

$$\int_0^1 dx \cdot x^{d_1-1} (wx)^{d_2-1} \left[\frac{p^2 x^{(d_1)} + w^2}{p^2 x^{(d_1)} + w^2} \right]^{d_1+d_2-2d_2} = \frac{1}{(w^2)^{d_1+d_2-2d_2}} \sum_{n=0}^{\infty} \frac{\Gamma(n+d_1+d_2-2d_2)}{n! \Gamma(d_1+d_2-2d_2)} (-1)^n \left(\frac{p^2}{w^2} \right)^n =$$

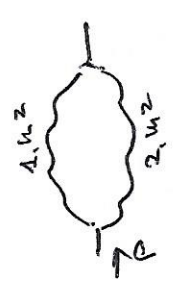
$d_1 + d_2 = d$

$$\frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2-2d_2)} \frac{\Gamma\left(\frac{d_1+d_2}{2} + n\right)}{\sqrt{\pi}} \Gamma\left(\frac{d+1}{2} + n\right)$$

$$\Gamma(2d) = \frac{2^{2d-1}}{\sqrt{\pi}} \Gamma(d) \Gamma\left(d + \frac{1}{2}\right) = \int_0^1 (w^2)^{d_1+d_2-2d_2} \frac{\Gamma(d_1)\Gamma(d_2) \sqrt{\pi}}{2^{d_1+d_2-2d_2} \Gamma\left(\frac{d_1}{2}\right) \Gamma\left(\frac{d_2}{2}\right)} F_2 \left[\begin{matrix} d_1, d_2 \\ d_1+d_2-2d_2 \end{matrix} \middle| -\frac{p^2}{4w^2} \right]$$

$$\frac{\Gamma\left(\frac{d_1}{2}\right) \Gamma\left(\frac{d_2}{2}\right) \sqrt{\pi}}{2^{d_1+d_2-2d_2} \Gamma\left(\frac{d_1}{2}\right) \Gamma\left(\frac{d_2}{2}\right)} F_2 \left[\begin{matrix} d_1, d_2 \\ d_1+d_2-2d_2 \end{matrix} \middle| -\frac{p^2}{4w^2} \right]$$

$d_1 = 1, d_2 = 2 \Rightarrow d = 3$



$$= \frac{\pi^{d/2}}{(u_2)^{d/2}} \frac{\Gamma(1+d/2) \sqrt{\pi}}{2^2 \Gamma(3/2) \Gamma(2)} \frac{\Gamma(4+d)}{\Gamma(4+d)} \frac{1}{\Gamma(d/2)} \left[-\frac{p^2}{4u^2} \right] = \frac{\Gamma(4+d)}{\Gamma(4+d)} \frac{1}{\Gamma(d/2)} \sum_{m=0}^{\infty} \varepsilon^m \bar{\varphi}_m \cdot \rho$$

$$F_2 \left[\begin{matrix} 1, 1+d/2 \\ 3/2 \end{matrix} \middle| -\frac{x}{4} \right] = \sum_{k=0}^{\infty} \frac{\Gamma(1+k) \Gamma(1+k)}{\Gamma(1+k)} \frac{\Gamma(3/2)}{\Gamma(1+k/2)} \left(-\frac{x}{4}\right)^k = \frac{1}{(-x)^k} \sum_{k=2,1}^{\infty} \frac{\Gamma(1+k)}{\Gamma(1+k)} \frac{\Gamma(1+k)}{\Gamma(2k)} \left(-\frac{x}{4}\right)^k$$

$$\Gamma(2k+2) = \frac{2^{2k+1}}{\sqrt{\pi}} \Gamma(k+1) \Gamma(k+3/2) \Rightarrow \frac{1}{\Gamma(k+3/2)} = \frac{2^{2k+1}}{\sqrt{\pi}} \frac{\Gamma(k+1)}{\Gamma(2k+2)} = \frac{2^k \Gamma(k+1)}{\Gamma(2k+2)}$$

$$\left(\frac{k}{2k}\right)^{-1} = \frac{\Gamma(k+1)^2}{\Gamma(2k+1)} = \frac{\Gamma(k)}{\Gamma(2k)} \cdot \frac{k}{2} \quad \left| \right. = \frac{2}{(-x)^{-1}} \sum_{k=1}^{\infty} \left(\frac{k}{2k}\right)^{-1} \frac{1}{k} \frac{\Gamma(k+1)}{\Gamma(k)} \Gamma(4k) \left(-\frac{x}{4}\right)^k = \frac{2}{(-x)^k} \sum_{m=0}^{\infty} \varepsilon^m \bar{\varphi}_m \cdot \rho$$

$$\frac{\Gamma(4+d)}{\Gamma(4+d) \Gamma(4+d)} = \exp \left[- \sum_{m=1}^{\infty} \frac{(-1)^m}{m} S_m(u_1) \right] = \exp \left[S_{1,2} - \frac{\varepsilon^2}{2} S_2 + O(\varepsilon^3) \right] = 1 + S_{1,2} + \frac{\varepsilon^2}{2} [S_{1,2} - S_2] + O(\varepsilon^3)$$

$$y = \frac{1 + \sqrt{\frac{x}{x+4}}}{1 + \sqrt{\frac{x}{x+4}}} \quad , \quad \rho = \frac{1-y}{1+y}$$

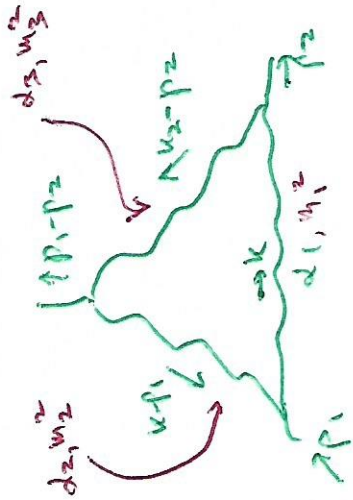
$$\bar{\varphi}_0 = \ln y \equiv \ell, \quad \bar{\varphi}_1 = \frac{1}{2} \ell^2 - \eta_2 + 2\ell \bar{\Gamma}_1 - 2\bar{\Gamma}_2, \quad \bar{\Gamma}_1 \equiv \ln(-y), \quad \bar{\Gamma}_2 \equiv \ln_2(-y) = \ln(-y) \ln(-y)$$

$$\bar{\varphi}_2 = \frac{1}{3} \ell^3 + 2L_1 \ell^2 + 2\ell [2\bar{\Gamma}_1^2 - \eta_2] - 4L_1 [\eta_2 - L_2] - 4\eta_3 - 4L_2 + 8S_{1,2}(-y), \quad L_3 \equiv \ln_3(-y)$$

Davydychev, Kalmykov, 2004

IBP relation:

Consider the triangle:



$$\int \frac{d^D k}{\{ \}} \left[D = \frac{D}{2} k_\mu \right]$$

$$\{ \} = (k+u_1)^2 (k-p_1)^2 + u_2^2 \int_{d_2}^{d_2} ((k-p_2)^2 + u_3^2)^{d_3}$$

IBP relation: $\frac{1}{\{ \}} \frac{\partial}{\partial k_\mu} k_\mu = \frac{\partial}{\partial k_\mu} \left\{ \frac{k_\mu}{\{ \}} \right\} - k_\mu \frac{\partial}{\partial k_\mu} \frac{1}{\{ \}}$
 boundary integral = 0

$$k_\mu \frac{\partial}{\partial k_\mu} \frac{1}{(k^2 - m^2)^d} = -2d \frac{k_\mu k_\mu}{(k^2 - m^2)^{d+1}} = -2d \frac{k^2}{(k^2 - m^2)^{d+1}} = -2d \left[\frac{1}{(k^2 - m^2)^{d+1}} - \frac{m^2}{(k^2 - m^2)^{d+1}} \right]$$

$$k_\mu \frac{\partial}{\partial k_\mu} \frac{1}{(k-P)^2 + m^2} = -2d \frac{(k-P)_\mu (k-P)_\mu}{(k-P)^2 + m^2} = -d \left[\frac{1}{(k-P)^2 + m^2} + \frac{k^2 + P^2}{(k-P)^2 + m^2} - \frac{(P^2 + m^2 + P^2)}{(k-P)^2 + m^2} \right]$$

$$2(k-P) = k^2 + (k-P)^2 - P^2$$

$$D \frac{\partial}{\partial k_\mu} \frac{1}{(k-P)^2 + m^2} = 2d_1 \left[\frac{1}{(k-P)^2 + m^2} - m^2 \frac{1}{(k-P)^2 + m^2} \right] + d_2 \left[\frac{1}{(k-P)^2 + m^2} + \frac{1}{(k-P)^2 + m^2} - (P^2 + m^2 + m^2) \frac{1}{(k-P)^2 + m^2} \right] + d_3 [\leftrightarrow]$$

resting line

$$(D - 2d_1 - d_2 - d_3) \frac{1}{(k-P)^2 + m^2} = d_2 \left[\frac{1}{(k-P)^2 + m^2} - \frac{1}{(k-P)^2 + m^2} - (m^2 + m^2) \frac{1}{(k-P)^2 + m^2} \right] + d_3 [\leftrightarrow] - 2d_1 m^2 \frac{1}{(k-P)^2 + m^2}$$

massive propagators

additional terms in massive case !!!

Q.14

$(d-2d_1-d_2-d_3)$
 $= d_2 \left[\text{Diagram} \right] + d_3 \left[\text{Diagram} \right] - (m_1^2 + m_2^2) \left[\text{Diagram} \right] - 2d_1 m_1^2 \left[\text{Diagram} \right]$

(A) One-loop with $m_1^2=0$, and $m_2^2=m^2$

$(d-3) = \frac{\Gamma(d)}{\Gamma(m)^2} \left[\text{Diagram} \right] - (p^2+m^2) \left[\text{Diagram} \right]$

 $(d-3) = \cancel{\text{Diagram}} - (p^2+m^2) \left[\text{Diagram} \right] - 2m^2 \left[\text{Diagram} \right]$

one-loop
 Only one \checkmark diagram is independent
 It is water diagram
 we calculate $\left[\text{Diagram} \right] \leftarrow$ finite

For water-integral IBP gives differential equation (DE) with the inhomogeneous term containing only tadpole

$(d-3-d) = d \left[\text{Diagram} \right] - (p^2+m^2) \left[\text{Diagram} \right] - \frac{\Gamma(d)}{\Gamma(m)^2} \left[\text{Diagram} \right]$

 , because $\frac{d}{(k^2+m^2)^{d+1}} = -\frac{2}{2m^2} \frac{1}{(k^2+m^2)^d}$

$\left[(d-3-d) \left[\text{Diagram} \right] \frac{\partial}{\partial m^2} \right] - \text{Diagram} = d \left[\text{Diagram} \right]$

 $d-4 = -2\epsilon$ for $d=2$

(D) One-loop with $k_1^2 = k_2^2 = k^2$

$$\begin{aligned}
 & \frac{\Gamma(k)}{(k^2)^k} \\
 & = \text{Diagram} - (k^2)^{k-1} \text{Diagram} \\
 & \quad \left[\text{only one master integral} \right] \text{DE for } \text{Diagram} \\
 & \quad - \frac{1}{2} \frac{\partial}{\partial k^2} \text{Diagram}
 \end{aligned}$$

If we line \rightarrow master-integral:
 L (finite)

$$\text{Diagram} (D-4) = 2 \text{Diagram} - 2 (k^2)^{2k-2} \text{Diagram} - 2k^2 \text{Diagram} \quad (1)$$

$$\text{Diagram} (D-6) = - \text{Diagram} - (k^2)^{2k-2} \text{Diagram} - 4k^2 \text{Diagram} \quad (2)$$

Sum of (1) and (2)

$$\text{Diagram} (D-5) = 2 \text{Diagram} - (k^2)^{2k-2} \left[2 \text{Diagram} + \text{Diagram} \right] - \frac{\partial}{\partial k^2} \text{Diagram}$$

J. HEWITT, 2013

Put all master integrals in a column and solve the DE system

The most popular method now.

Two-loops with masses (only sunsets, for simplicity)



(A) two massless lines

$$\int \frac{d^D k}{(2\pi)^D} A(d_1, d_2) \quad \text{one loop}$$

(B) one massless line and $m_1^2 = m_2^2 = m^2$



$$\bar{d} = d_1 + d_2 + d_3$$

Tadpole $\int \frac{d^D k}{(2\pi)^D} \frac{1}{k^2} \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2-2)} = \int \frac{d^D x \cdot x^{d_1-1} (1-x)^{d_2-1}}{[x(1-x)]^{d_1+d_2-2}} = \int \frac{dx \cdot x^{d_1-1} (1-x)^{d_2-1}}{[x(1-x)]^{d_1+d_2-2}}$

$$\int \frac{d^D k}{(2\pi)^D} \frac{1}{k^2} \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2-2)} = \frac{\Gamma(d_1+d_2-2) \Gamma(d_1+d_2-2)}{\Gamma(d_1+d_2+2d_3-2)}$$

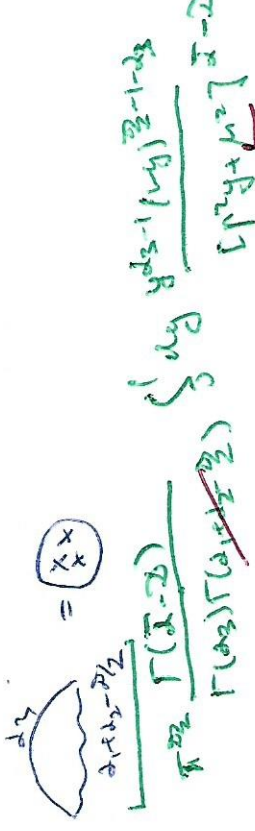
$$= \pi^D \frac{\Gamma(\frac{D}{2}-d_3) \Gamma(d_1+d_2+d_3-2) \Gamma(d_2+d_3-2) \Gamma(d_2+d_3-2)}{\Gamma(d_1) \Gamma(d_2) \Gamma(\frac{D}{2}) \Gamma(d_1+d_2+2d_3-2) (m^2)^{d_1+d_2}}$$

$$= \pi^D \frac{\Gamma(d_1+d_2-2)}{\Gamma(d_1) \Gamma(d_2)} \int \frac{dx \cdot x^{d_1-1} (1-x)^{d_2-1}}{[x(1-x)]^{d_1+d_2-2}} = \frac{1}{[x(1-x)]^{d_1+d_2-2}} \frac{\Gamma(d_1+d_2-2)}{[\Gamma(x) \Gamma(1-x)]^{d_1+d_2-2}} \int \frac{dx \cdot x^{d_1-1} (1-x)^{d_2-1}}{[x(1-x)]^{d_1+d_2-2}}$$

$$\mu^2 = \frac{m^2}{x(1-x)}$$

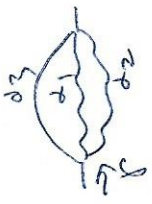
General case (i.e. $p \neq 0$)

$\bar{d} = d_1 + d_2 + d_3$



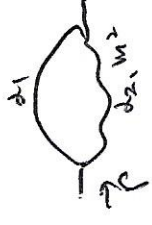
$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$= \pi^{d_2} \frac{\Gamma(d_1+d_2-d_2)}{\Gamma(d_1) \Gamma(d_2)}$$



$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1+d_2-d_2)}{\Gamma(d_1) \Gamma(d_2)}$$

$$= \pi^{d_2} \frac{\Gamma(d_1+d_2-d_2)}{\Gamma(d_1) \Gamma(d_2)}$$



$$\int_{d_2} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_3} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_2} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_3} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$= \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_2} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_3} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_2} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_3} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_2} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

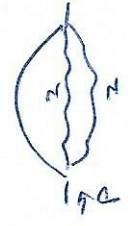
$$\int_{d_3} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

$$\int_{d_1} x^{d_1-1} (1-x)^{d_2-1-d_1} dx = \frac{\Gamma(d_1) \Gamma(d_2)}{\Gamma(d_1+d_2)}$$

2.18

IBP gives relations between diagrams with various indices di

let $d_3=1, d_1=d_2=2 \Rightarrow F = S \quad 2X = 6-2 = 2+2C \Rightarrow X=4$



$$F = \frac{\int \frac{d^3k}{(2\pi)^3} \frac{1}{(k^2)^2 (k^2)^2 (k^2)^2} = \frac{1}{(4\pi)^3} \int \frac{d^3k}{k^4} = \frac{1}{(4\pi)^3} \int_0^\infty \frac{4\pi k^2 dk}{k^4} = \frac{1}{(4\pi)^3} \int_0^\infty \frac{4\pi}{k^2} dk = \frac{1}{(4\pi)^3} \left[-\frac{4\pi}{k} \right]_0^\infty = \frac{1}{(4\pi)^3} \left(\frac{1}{\epsilon} + \dots \right)$$

$$\Gamma(1+2\epsilon) = \frac{2^{2\epsilon}}{\Gamma(\epsilon)} \Gamma(2\epsilon) \Gamma(\epsilon)$$

$$\frac{\Gamma(1+2\epsilon) \sqrt{k}}{2^{1+2\epsilon} \Gamma(3/2+\epsilon)} = \frac{1}{2} \frac{\Gamma(3/2+\epsilon) \Gamma(\epsilon)}{\Gamma(3/2+\epsilon) \Gamma(3/2+\epsilon)} \frac{\Gamma(2\epsilon)}{\Gamma(2\epsilon)}$$

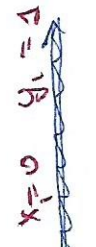
$$= \frac{\Gamma(2\epsilon)}{\Gamma(\epsilon)^2} \frac{\Gamma(1+\epsilon)}{\Gamma(1-\epsilon)} \frac{\Gamma(2\epsilon)}{\Gamma(2\epsilon)} = \frac{\Gamma(2\epsilon)}{\Gamma(\epsilon)^2} \frac{\Gamma(1+\epsilon)}{\Gamma(1-\epsilon)} = \frac{\Gamma(2\epsilon)}{\Gamma(\epsilon)^2} \frac{\Gamma(1+\epsilon)}{\Gamma(1-\epsilon)}$$

$$= \frac{1 - \sqrt{\frac{x}{x+4}}}{1 + \sqrt{\frac{x}{x+4}}}$$

$$F_3 \left[\begin{matrix} 1, 1, 1 \\ 2, 2, 2 \end{matrix} \middle| -\frac{x}{4} \right] = \sum_{k_1=0}^{\infty} \sum_{k_2=0}^{\infty} \sum_{k_3=0}^{\infty} \frac{\Gamma(1-k_1) \Gamma(1-k_2) \Gamma(1-k_3)}{\Gamma(2+k_1) \Gamma(2+k_2) \Gamma(2+k_3)} \left(-\frac{x}{4} \right)^{k_1+k_2+k_3} = \frac{1}{(4\pi)^3} \int \frac{d^3k}{k^4} \dots$$

$$= \frac{1 - \sqrt{\frac{x}{x+4}}}{1 + \sqrt{\frac{x}{x+4}}} \frac{4\pi}{(4\pi)^3} \frac{1}{1-\epsilon} = \frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon}$$

$$\frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon} = \frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon} \frac{\Gamma(1+\epsilon) \Gamma(1-\epsilon)}{\Gamma(1+\epsilon) \Gamma(1-\epsilon)} = \frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon} \frac{\Gamma(1+\epsilon) \Gamma(1-\epsilon)}{\Gamma(1+\epsilon) \Gamma(1-\epsilon)}$$

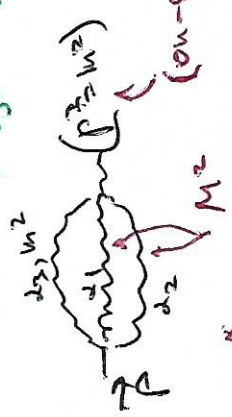


$$F_3 \left[\begin{matrix} 1, 1, 1 \\ 2, 2, 2 \end{matrix} \middle| -\frac{x}{4} \right] = \frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon} \frac{\Gamma(1+\epsilon) \Gamma(1-\epsilon)}{\Gamma(1+\epsilon) \Gamma(1-\epsilon)} = \frac{1}{(4\pi)^3} \frac{4\pi}{1-\epsilon}$$

$$R_{mp} \frac{\Gamma(x-1)}{\Gamma(x)} = \frac{1}{\Gamma(x-1)} \int_{-2\epsilon}^{2\epsilon} \int_{-2\epsilon}^{2\epsilon} \int_{-2\epsilon}^{2\epsilon} \dots = \frac{1}{\Gamma(x-1)} \int_{-2\epsilon}^{2\epsilon} \int_{-2\epsilon}^{2\epsilon} \int_{-2\epsilon}^{2\epsilon} \dots$$

Last example

$d_2 = 1, d_2 = d_1 = 2$



(on-shell $p^2 = -m^2$)

$S_0(u) = \sum_{l=1}^u \frac{1}{l}$

$S_1 = S_1(u-1), \bar{S}_1 = S_1(2u-1), \bar{\bar{S}}_1 = S_1(4u-1)$

$$M^2 N_{122} = \frac{1}{x} \sum_{h=1}^{\infty} (-x^2)^h \binom{2h}{h} \binom{4h}{2h} \left(-\frac{1}{2u} \ln x + \frac{S_1}{2u} - \frac{1}{2} \frac{1}{h} + \frac{1}{2} \frac{1}{u} \right) + \dots$$

$x = \frac{m^2}{M^2}$

$$+ \frac{1}{x^2} \sum_{h=1}^{\infty} \frac{(-16x^2)^h}{\binom{2h}{h} \binom{4h}{2h}} \left(-\frac{1}{2h-1} + \frac{1}{2h} - \frac{1}{4h} \right) = O(\epsilon)$$

$$= \frac{1}{dS} \left\{ \frac{x^{1+\delta}}{(3+2\delta)(2+\delta)} \sqrt{3} \left[\begin{matrix} 1, 3/2, 1/2+\delta, 1+\delta/2 \\ 2+\delta/2, S_1+\delta/2, 1/4+\delta/2 \end{matrix} \right] - \frac{x^2}{4} \right\} \Big|_{\delta=0}$$

$$+ \frac{\sqrt{3}}{4} \left[\begin{matrix} 1, 1, 1, 1/2 \\ 3/2, 3/2, 3/2 \end{matrix} \right] - \frac{x^2}{4}$$

NRQED, NRQCD

(Kniehl, A.V.K., Quirchelenko, Vorefin 2005)

Compact form of the integral representation:

$$M^2 \Omega_{1,2,2} = -\frac{1}{2\pi} \int_0^1 \frac{dt}{t\sqrt{1-t}} \left(\frac{1}{\sqrt{1+4A^2}} L(A) - 2\ln A \right) \rightarrow \text{integrals of elliptic integrals}$$

$A = \frac{\alpha t}{1-t}$
 \Rightarrow Elliptic integral (in the case if $L(A) = 1$)

$$L_1(A) = \ln \frac{\sqrt{1+4A^2} - 1}{\sqrt{1+A^2}} \quad ; \quad L_2 = \ln \frac{\sqrt{1-2A}}{\sqrt{1+2A}} \quad ; \quad L(A) = L_1(A) + L_2(A)$$

Elliptic polilogarithms

- [Paris group [Bloch, Vanhove, 2013]
- [Mainz group [S. Weinzierl et al, started 2017]
- [Broedel, Duhr, Dulat, Penante, Taroni [started 2017]

(A. Levin, 2007) \Leftarrow Mathematics

Conclusion

I have shown rather effective methods of calculation of massless

Feynman diagrams and also

Feynman integrals containing massive propagators