Design of nanomechanical sensors based on carbon nanoribbons and nanotubes in a distributed computing system

C.M. Visan, T.L. Mitran, G.A. Nemnes, I.T. Vasile, M.A. Dulea

Computational Physics and Information Technologies Department - dfcti.ifin.ro Horia Hulubei National Institute for R&D in Physics and Nuclear Engineering (IFIN-HH) Reactorului Str. no.30, P.O.BOX MG-6, Bucharest - Magurele, ROMANIA



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Presentation outline:

- Research strategy and computing infrastructure
- *Ab initio* simulation of nanostructures SIESTA software package
- Density Functional Theory SIESTA implementation
- Numerical tricks for simulation speedup
- Parallel SIESTA
- Nanomechanical resonators based on carbon nanoribbons and nanotubes
- Conclusions

Research strategy and computing infrastructure

The main strategic directions for R&D of the CPIT Department are:

- Numerical investigation of mesoscopic quantum systems
- High-performance computing for applications in nuclear and condensed matter physics
- Study of physical properties of nanostructures through ab initio computing methods
- Computational modeling and simulation of condensed systems with many interacting particles
- Modeling and simulations in biophysics

Computing infrastructure:

- IBM BladeCentre H Infiniband Voltaire 4X (40Gbps)
- 10 nodes LS22 + 16 nodes QS22
 80 cores AMD Opteron, 80GB RAM
 - 32 PowerXCell 8i, 512GB
- 56 nodes HS22: 2x Intel Xeon X5650 @ 2.67GHz (6 Core) - 672 cores,
 - 2412 GB RAM @ 1333 MHz
 - 500GB SATA3 HDD per node



- 112 cores
- 420 GB RAM @ 1600MHz
- 500GB SAS HDD per node
- 4x NVIDIA Tesla M2090 GPU Graphics clock: 650 MHz Processor clock: 1300 MHz Memory size: 6GB GDDR5 CUDA cores: 512
- Storage 35TB Raid6 w GbE Storage 14TB Raid6 w IB Storage 66TB Raid6 w IB







Ab initio simulation of nanostructures – SIESTA software package

• SIESTA - Spanish Initiative for Electronic Simulations with Thousands of Atoms

"The SIESTA method for ab initio order-N materials simulation" José M Soler, Emilio Artacho, Julian D Gale, Alberto García, Javier Junquera, Pablo Ordejón and Daniel Sánchez-Portal **J. Phys.: Condens. Matter 14 2745 (2002)**

- SIESTA 4.0 (GPL license) departments.icmab.es/leem/siesta/
- Most important feature of SIESTA: linear scaling, while usual DFT methods scale as N³ !

What can you compute with it ?

What else can it do ?

- Electronic band structure
- Electron density
- Total and partial density of states
- Electric dipole moment
- Mulliken population
- Spin polarized calculations
- Atomic forces
- Structural relaxation
- Stress tensor
- Moleculat dynamics
- Phonon spectra



Ballistic electronic transport with **TRANSIESTA**

"Density-functional method for nonequilibrium electron transport" Mads Brandbyge, Jose-Luis Mozos, Pablo Ordejon, JeremyTaylor, and Kurt Stokbro **Phys. Rev. B 65 165401 (2002)**



• **Traditional approach:** find approximate solutions for the Schrödinger equation

$$\left\{-\frac{h^2}{2m}\sum_{j}\nabla_j^2 - \sum_{j,i}\frac{e_0^2 Z_i}{|r_j - R_i|} + \frac{1}{2}\sum_{j\neq i}\frac{e_0^2}{|r_j - r_i|} - E\right\}\psi(r_1,...,r_N) = 0$$

• Density Functional Theory (DFT): Hohenberg & Kohn & Sham

Main ideas:

- replace the N electron wavefunction with a density functional for non interacting electrons
- include all multi-particle effects in a separate energy term known as the exchange correlation energy (E_{xc})
- Kohn-Sham equations:

$$\left(-\frac{\hbar^2}{2m}\nabla^2 + V_{KS}(\vec{r})\right)\psi_i(\vec{r}) = E_i\psi_i(\vec{r}) \qquad E[\psi_i(\vec{r})] = T_e + U_{nn} + U_{ne} + U_{ee} + E_{xc}$$

Schrödinger-like equation for system of non-interacting particles that generates the **same electron** density as the real system of interacting particles.

• Solve the KS equations by minimizing the energy of the system – variational approach.

Not trivial even for N=2 !

- P. Hohenberg and W. Kohn
 Phys. Rev., 136 B864 (1964)
- W. Kohn and L. J. Sham
 Phys. Rev., 140 A1133 (1965)

Numerical tricks for simulation speedup

- Use **pseudopotentials**: replace core electrons and nuclear potential with an effective pseudo-potential.
- Benefits of using pseudopotnetials:
 - Get rid of core electrons (not chemically active)
 - Iron out strongly varying wavefunction (near core)
- Use soft pseudopotentials for transferability (compute once and use in any system)





• Numerical orbitals: SIESTA uses a **localized basis set** which leads to **sparse Hamiltonian and overlap matrices + empty space has almost no additional computational cost**

Reduce the number of electrons and the basis

set that will be required.

- **Periodic Born von Karman boundary conditions:** easily solve infinite systems but can also treat non-periodic systems if the supercell is large enough
- Monkhorst-Pack scheme for Brillouin zone integration

SIESTA software stack:

- MPI communication layer
- BLACS basic linear algebra library with integrated MPI communication
- SCALAPACK linear algebra package: uses BLAS + LAPACK for serial operations and MPI+ BLACS for parallel operations



Parallel SIESTA

• Main parallelization options:

- Parallel over K: solve problem for each separate reciprocal space point
- Solve eigenvalue problem by parallel matrix diagonalization
- Spatial decomposition based on atomic proximity
- Difficult to test scaling because of a large variety of physical systems and numerical requirements: 1D, 2D, 3D, periodic, non-periodic, number of atoms, different numerical accuracy required for different investigated parameters, etc.

Personal user-case results:

- Optimal configuration: 8 cores (on same processor) per instance while running multiple independent instances
- Possible speedup from library and architecture optimization: over 2.7 times
- HPC scaling for DFT simulations:

- "Performance Analysis of Electronic Structure Codes on HPC Systems: A Case Study of SIESTA" F. Corsetti, PLoS ONE 9 e95390 (2014)

- "Linear-scaling density-functional theory with tens of thousands of atoms: Expanding the scope and scale of calculations with ONETEP" N.D.M. Hine, P.D. Haynes, A.A. Mostofi, C.-K. Skylaris, M.C. Payne, Comp. Phys. Comm. 180 1041–1053 (2009)



Georg Huhs - "Parallelization issues" (handout)

Nanomechanical resonators based on carbon nanoribbons and nanotubes

• Motivation:

- multiple possible practical uses such as: high-frequency oscillators or ultra-sensible acceleration detectors

- **Possible practical implementation:** measure transmission current between oscillator tip and contact to estimate position and frequency
- Systems studied:

- monoatomic thickness nanoribbons of length 17.3 Å and 2.8 Å diameter made out of 32 C atoms and passivated with 18 H atoms

- nanotubes of length 35.9 Å and 4.2 Å diameter made out of 92 C atoms with closed ends

• Simulation parameters:

- basis type: double zeta polarized
- local density approximation functional (LDA) + Ceperley Alder exchange correlation parametrization

- 300 Ry mesh cut-off, 10^{-3} density matrix tolerance, 0.01 eV/Ang maximum force tolerance

Goals:

- structural relaxation + molecular dynamics with thermostat at 300 and 150 K
- transmission function over gap (between nanotube and contact) for different separations (TRANSIESTA NEGF)



1000 (fs) 2000 (fs) 3000 (fs) Time: 0 (fs) Initial angle = 20° T = 300 K Initial angle = 20° T = 150 K Initial angle = 10° T = 300 K Initial angle = 10° T = 150 K

Nanomechanical resonators based on carbon nanoribbons and nanotubes

0 (fs) 200 (fs) 400 (fs) 600 (fs) Time: Initial angle = 20° T = 300 K Initial angle = 20° T = 150 K Initial angle = 10° T = 300 K Initial angle = 10° T = 150 K

Nanomechanical resonators based on carbon nanoribbons and nanotubes

Nanomechanical resonators based on carbon nanoribbons and nanotubes



Position of tip atom in oscillation plane as a function of time for the first 3000 fs: *nanoribbon vs. nanotube*



Transmission from -1 to 1 eV as a function of *nanotube*-contact distance (initial and after relaxation)



Position of tip atom of *nanotubes* in oscillation plane as a function of time for the first 5000 fs



Resonant frequency of *nanotubes* as a function of the initial angle

Conclusions:

- Because of the reduced rigidity, the studied nanoribbons have a much higher period of oscillation and are more susceptible to thermal noise than the nanotubes. This makes nanotubes better suited as nanomechanical resonators.
- The **nanotubes show a non-linear behavior**, with a resonant frequency that depends on both the initial amplitude (angle) and also on temperature.
- The nanotube-contact gap has an optimal value of ~3 Å
- Small diameter nanotubes prove to be suitable candidates for nano-resonators in the THz region

Thank you for your attention !