## Rotating black holes in 5D Einstein-Maxwell-Chern-Simons theory with negative cosmological constant

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## Rotating black holes in 5D Einstein-Maxwell-Chern-Simons theory with negative cosmological constant

1. Introduction
2. Near-horizon formalism
3. Exploring the global solutions

EM-AdS vs EMCS-AdS $\lambda=1$ (SUGRA) vs EMCS-AdS $\lambda=1.5$
Global solutions and branch structure for $\lambda>2$

## 1. Introduction

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Black holes in $\mathrm{D}=5$ dimensions in
Einstein-Maxwell-Chern-Simons theory with negative cosmological constant

Asymptotically anti-de-Sitter space-times:
Interesting in the context of the AdS/CFT correspondence

Gravitating fields propagating in an AdS space-time


Fields propagating in a conformal field theory
Known analytical solutions:

- Myers-Perry black hole (uncharged)
- 5D Reissner-Nordström black hole (static)
- Cvetič-Lu-Pope black hole (rotating and charged, SUGRA) (PLB598 273)

What are the properties of black holes connecting these solutions?

## 1. Introduction

We are interested in the higher dimensional generalization of the Kerr-Newman black holes in 5D EMCS-AdS theory:

$$
I=\frac{1}{16 \pi G_{5}} \int d^{5} x\left[\sqrt{-g}\left(R-F^{2}-2 \Lambda\right)-\frac{2 \lambda}{3 \sqrt{3}} \varepsilon^{\mu \nu \alpha \beta \gamma} A_{\mu} F_{\nu \alpha} F_{\beta \gamma}\right]
$$

$$
R=\text { curvature scalar }
$$

$\mathrm{U}(1)$ electro-magnetic potential $A_{\mu}$
F = field strength tensor
$\Lambda=$ cosmological constant
$\lambda=$ Chern-Simons coupling parameter

## || 1. Introduction

$$
I=\frac{1}{16 \pi G_{5}} \int d^{5} x\left[\sqrt{-g}\left(R-F^{2}-2 \Lambda\right)-\frac{2 \lambda}{3 \sqrt{3}} \varepsilon^{\mu \nu \alpha \beta \gamma} A_{\mu} F_{\nu \alpha} F_{\beta \gamma}\right]
$$

Einstein-Maxwell-Chern-Simons theory in 5 dimensions

$$
G_{\mu \nu}+\Lambda g_{\mu \nu}=2\left(F_{\mu \rho} F_{\nu}^{\rho}-\frac{1}{4} F^{2}\right)
$$

Einstein equations

$$
G_{5}=1
$$

$$
\nabla_{\nu} F^{\mu \nu}+\frac{\lambda}{2 \sqrt{3}} \varepsilon^{\mu \nu \alpha \beta \gamma} F_{\nu \alpha} F_{\beta \gamma}=0
$$

Maxwell equations

## 1. Introduction

Ansatz constraints:

1. Axially symmetric and stationary: $U(1)^{N}$ symmetry

In D dimensions $\mathrm{N}=[(\mathrm{D}-1) / 2]$ (planes of rotation)
2. All angular momenta of equal magnitude: enhanced $U(N)$ symmetry

$$
\left|J_{(1)}\right|=\left|J_{(2)}\right|=\ldots=\left|J_{(N)}\right|=J
$$

3. Event horizon with spherical topology
4. Asymptotically AdS

## 1. Introduction

Ansatz for the metric (5D):

$$
\begin{aligned}
& d s^{2}=-b(r) d t^{2}+\frac{1}{u(r)} d r^{2}+g(r) d \theta^{2}+p(r) \sin ^{2} \theta\left(d \varphi_{1}-\frac{\omega(r)}{r} d t\right)^{2} \\
& +p(r) \cos ^{2} \theta\left(d \varphi_{2}-\frac{\omega(r)}{r} d t\right)^{2}+(g(r)-p(r)) \sin ^{2} \theta \cos ^{2} \theta\left(d \varphi_{1}-d \varphi_{2}\right)^{2}
\end{aligned}
$$

$$
\theta \in[0, \pi / 2], \varphi_{1} \in[0,2 \pi] \text { and } \varphi_{2} \in[0,2 \pi]
$$

Lewis-Papapetrou coordinates. The radial coordinate $\mathbf{r}$ is quasi-isotropic.
Ansatz for the gauge field:

$$
A_{\mu} d x^{\mu}=a_{0}(r) d t+a_{\varphi}(r)\left(\sin ^{2} \theta d \varphi_{1}+\cos ^{2} \theta d \varphi_{2}\right)
$$

System of second order ordinary differential equations + constraints

## | 1. Introduction

Global Charges:

Mass

$$
M=-\frac{\pi}{8} \frac{\beta-3 \alpha}{L^{2}}
$$

(Ashtekar-Magnon-Das conformal mass)

Angular
Momenum

$$
J_{(k)}=\int_{S_{\infty}^{3}} \beta_{(k)}
$$

$$
\beta_{(k) \mu_{1} \mu_{2} \mu_{3}} \equiv \epsilon_{\mu_{1} \mu_{2} \mu_{3} \rho \sigma} \nabla^{\rho} \eta_{(k)}^{\sigma}
$$

$$
\left|J_{(k)}\right|=J
$$

Electric charge

$$
Q=-\frac{1}{2} \int_{S_{\infty}^{3}} \tilde{F}
$$

$$
\tilde{F}_{\mu_{1} \mu_{2} \mu_{3}} \equiv \epsilon_{\mu_{1} \mu_{2} \mu_{3} \rho \sigma} F^{\rho \sigma}
$$

## || 1. Introduction

Horizon Charges:

Area

$$
A_{\mathrm{H}}=\int_{\mathcal{H}} \sqrt{\left|g^{(3)}\right|}=2 \pi^{2} r_{\mathrm{H}}^{3} \lim _{r \rightarrow r_{\mathrm{H}}} \sqrt{\frac{m^{2} n}{f^{3}}}
$$

Entropy
$S=4 \pi A_{H}$

Horizon
Mass

$$
M_{\mathrm{H}}=-\frac{3}{2} \int_{\mathcal{H}} \alpha=\lim _{r \rightarrow r_{\mathrm{H}}} 2 \pi^{2} r^{3} \sqrt{\frac{m n}{f^{3}}}\left[\frac{n \omega}{f}\left(\frac{\omega}{r}-\omega^{\prime}\right)+f^{\prime}\left(1+\frac{r^{2}}{L^{2}}\right)+\frac{2 r f}{L^{2}}\right]
$$

Horizon
Angular
Momenta

$$
J_{\mathrm{H}(k)}=\int_{\mathcal{H}} \beta_{(k)}=\lim _{r \rightarrow r_{\mathrm{H}}} \pi^{2} r^{3} \sqrt{\frac{m n^{3}}{f^{5}}}\left[\omega-r \omega^{\prime}\right]
$$

## 2. Near-horizon formalism

## || 2. Near Horizon Formalism

Properties of the near-horizon geometry of extremal black holes.
H. K. Kunduri and J. Lucietti, Living Reviews in Relativity 16 (2013)

- The near-horizon geometry of extremal black holes with spherical topology is the product of two independent spaces.


Isometries:

$$
\begin{aligned}
& S O(2,1) \times S O(D-1) \quad \text { static case (sphere) } \\
& S O(2,1) \times U(1)^{N} \quad \text { rotation (squashed sphere) }
\end{aligned}
$$

This factorization is obtained for all the known examples of topologically spherical black holes

## || 2. Near Horizon Formalism

Hence we can assume such factorization in our black holes (extremal case)
Metric:

$$
\begin{aligned}
& d s^{2}=v_{1}\left(d r^{2} / r^{2}-r^{2} d t^{2}\right)+v_{2}\left[4 d \theta^{2}+\sin ^{2} 2 \theta\left(d \phi_{2}-d \phi_{1}\right)^{2}\right] \\
& +v_{2} \eta\left[d \phi_{1}+d \phi_{2}+\cos ^{2} 2 \theta\left(d \phi_{2}-d \phi_{1}\right)-\alpha r d t\right]^{2}
\end{aligned}
$$

Gauge potential:

$$
A=-(\rho+p \alpha) r d t+2 p\left(\sin ^{2} \theta d \phi_{1}+\cos ^{2} \theta d \phi_{2}\right)
$$

- Field equations + Ansatz: algebraic relations for the Ansatz parameters
- Global charges can be calculated: (J, Q)
- Horizon charges: area, horizon angular momentum
- Parameters related to the asymptotical structure of the global solution cannot be calculated: Mass, angular velocity


## || 2. Near Horizon Formalism

Near-horizon geometry branch structure: EM flat


## || 2. Near Horizon Formalism

Near-horizon geometry branch structure: EM-AdS


## || 2. Near Horizon Formalism |

Near-horizon geometry branch structure: EMCS-AdS, Q=2.720699, $\lambda=5$


## 3. Exploring the global solutions

EM-AdS<br>VS<br>EMCS-AdS $\lambda=1$ (SUGRA)<br>VS<br>EMCS-AdS $\lambda=1.5$

## || 3. Exploring the global solutions ||

EM-AdS black holes with $\mathrm{Q}=0.044, \mathrm{~L}=1$


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## || 3. Exploring the global solutions ||

EMCS-AdS black holes with $\mathrm{Q}=0.044, \mathrm{~L}=1$
$\lambda=1$ (SUGRA)
$\lambda=1.5$



|| 3. Exploring the global solutions ||
EMCS-AdS black holes with $\mathrm{Q}=0.044, \mathrm{~L}=1$
$\lambda=1$ (SUGRA)
$\lambda=1.5$



EMCS-AdS black holes with $\mathrm{Q}=0.044, \mathrm{~L}=1, \lambda=1.5$


## 3. Exploring the global solutions

Global solutions and branch structure for $\lambda>2$

## || 3. Exploring the global solutions ||

Global and NH solutions, $\lambda>2$ scheme:


## || 3. Exploring the global solutions ||

Global extremal black holes, $\lambda>2$ scheme:


## || 3. Exploring the global solutions ||

Branch structure, $\lambda>2$ scheme:


J

## || 3. Exploring the global solutions |

$\mathrm{J}=0, \mathrm{Q}=2.720699, \Lambda=-0.06, \lambda=5$, Extremal

$\mathrm{J}=0, \mathrm{Q}=2.720699, \Lambda=-0.06, \lambda=5$, Extremal


## Thank you for your attention!

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