

# Dark Matter Models (I)

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# Outline

- 1 Dark Matter in astrophysics
- 2 Dark Matter properties
- 3 Thermal Dark Matter
- 4 Summary

# Interplay: Standard Model and Cosmology

Gauge fields (interactions):  $\gamma, W^\pm, Z, g$

Three generations of matter:  $L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}, e_R; Q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}, d_R, u_R$

- SM Describes

- ▶ all experiments dealing with electroweak and strong interactions

- SM fails to describe (PHENO)

- ▶ Neutrino oscillations
- ▶ Dark matter ( $\Omega_{DM}$ )
- ▶ Baryon asymmetry ( $\Omega_B$ )
- ▶ Inflationary stage

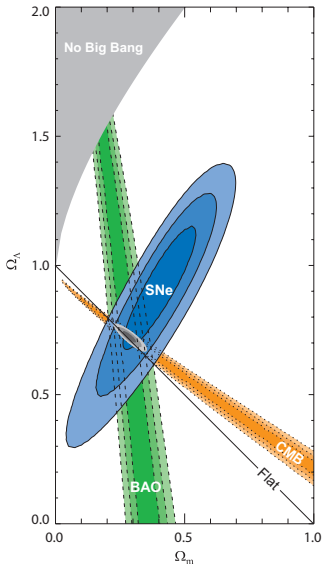
- (THEORY)

- ▶ Dark energy ( $\Omega_\Lambda$ )
- ▶ Strong CP-problem
- ▶ Gauge hierarchy
- ▶ Quantum gravity

Cosmology asks for new physics  
severely constrains many BSM

and limits neutrino mass  
relaxation..?

# Astrophysical and cosmological data are in agreement



$$\left(\frac{\dot{a}}{a}\right)^2 = H^2(t) = \frac{8\pi}{3} G \rho_{\text{density}}^{\text{energy}}$$

$$\rho_{\text{density}}^{\text{energy}} = \rho_{\text{radiation}} + \rho_{\text{matter}}^{\text{ordinary}} + \rho_{\text{matter}}^{\text{dark}} + \rho_{\Lambda}$$

$$\rho_{\text{radiation}} \propto 1/a^4(t), \quad \rho_{\text{matter}} \propto 1/a^3(t), \quad \rho_{\Lambda} = \text{const}$$

$$\frac{3H_0^2}{8\pi G} = \rho_{\text{density}}^{\text{energy}}(t_0) \equiv \rho_c \approx 0.53 \times 10^{-5} \frac{\text{GeV}}{\text{cm}^3}$$

Radiation:

$$\Omega_{\gamma} \equiv \frac{\rho_{\gamma}}{\rho_c} = 0.5 \times 10^{-4}$$

Baryons (H, He):

$$\Omega_B \equiv \frac{\rho_B}{\rho_c} = 0.05$$

Neutrino:

$$\Omega_{\nu} \equiv \frac{\sum \rho_{\nu i}}{\rho_c} < 0.01$$

$$N_{\nu} \simeq 3, \quad \sum m_{\nu} \lesssim 0.2 \text{ eV}$$

Dark matter:

$$\Omega_{\text{DM}} \equiv \frac{\rho_{\text{DM}}}{\rho_c} = 0.27$$

Dark energy:

$$\Omega_{\Lambda} \equiv \frac{\rho_{\Lambda}}{\rho_c} = 0.68$$

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# Observations we use in the analysis

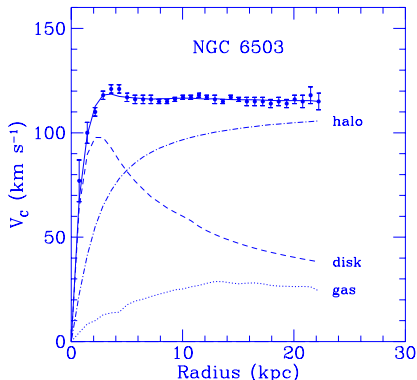
- **Astrophysical data** (favor Dark Matter)
  - ▶ Observations in galaxies: rotation curves, number of dwarfs
  - ▶ Observations in galaxy clusters: X-rays, strong lensing
- **Cosmological data** (favor Dark Matter and Dark Energy)
  - ▶ Observation of objects at cosmological distances (cosmic ladder): cefeids, SN Ia, LRG, RGB?
  - ▶ Baryonic Acoustic (Sakharov) Oscillations (BAO): two-point galaxy correlation function
  - ▶ Galaxy formation process: time-dependent galaxy spectrum
  - ▶ Evolution of galaxy clusters: X-rays, Sunyaev–Zeldovich effect
  - ▶ Anisotropy and Polarization of Cosmic Microwave Background (CMB): gaussianity, angular size of the sound horizon at recombination, ISW-effect, reionization, weak lensing, GW?, ...

## Galactic dark halos:

## flat rotation curves

$$v(R) = \sqrt{G \frac{M(R)}{R}}$$

$$M(R) = 4\pi \int_0^R \rho(r) r^2 dr$$



observations:

$v(R) \simeq \text{const}$

visible matter:

internal regions  $v(R) \propto \sqrt{R}$   
 external ("empty") regions  $v(R) \propto 1/\sqrt{R}$

# Dark Matter in clusters

## X-rays from hot gas in clusters

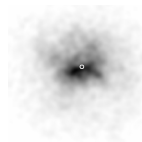
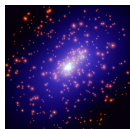
$$\frac{dP}{dR} = -\mu n_e(R) m_p \frac{GM(R)}{R^2}, \quad M(R) = 4\pi \int_0^R \rho(r) r^2 dr, \quad P(R) = n_e(R) T_e(R)$$

## galaxies in clusters

## virial theorem

$$U + 2E_k = 0$$

$$3M \langle v_r^2 \rangle = G \frac{M^2}{R}$$

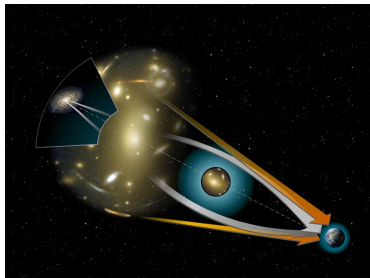


## Milky Way: Virgo infall

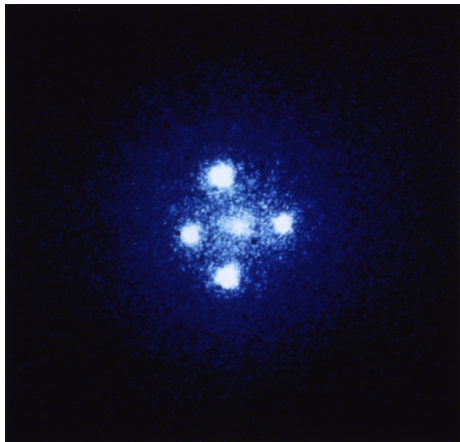


## Gravitational lensing in GR:

$$\alpha = 4GM/(c^2 b)$$

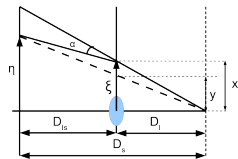


## Einstein Cross



source: quasar  $D_s = 2.4$  Gpc

lens: galaxy  $D_l = 120$  Mpc



$$\vec{\eta} = \frac{D_s}{D_l} \vec{\xi} - D_{ls} \vec{\alpha}(\vec{\xi})$$

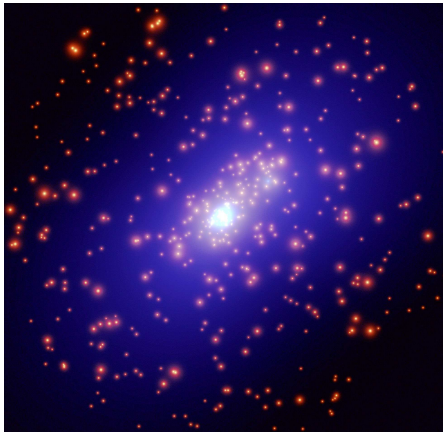
common lens  
with specific  
refraction  
coefficient

$$\vec{\alpha}(\vec{\xi}) = \frac{4G}{c} \int \frac{\vec{\xi} - \vec{\xi}'}{|\vec{\xi} - \vec{\xi}'|^2} d^2 \xi' \int \rho(\vec{\xi}', z) dz$$

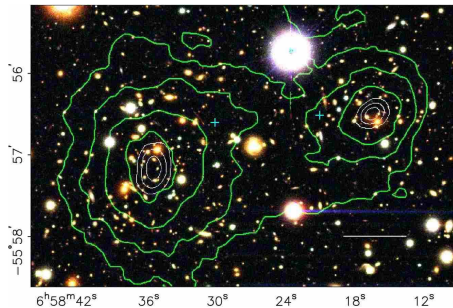
# Dark Matter in clusters

gravitational lensing

$$\rho_B \approx 0.25 \rho_{DM}$$



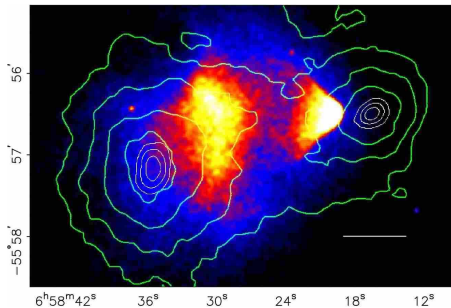
# Colliding clusters (Bullet clusters 1E0657-558)



gravitational lensing

scale is 200 kpc

clusters are at 1.5 Gpc



Observations in X-rays

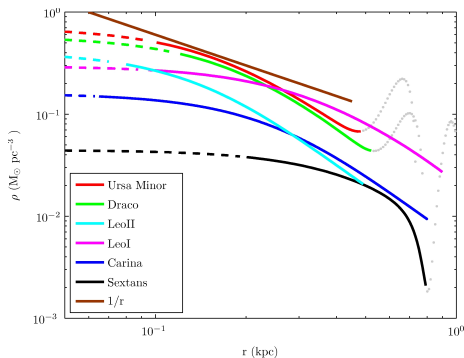
$M \simeq 10 \times m$

implies collisionless DM

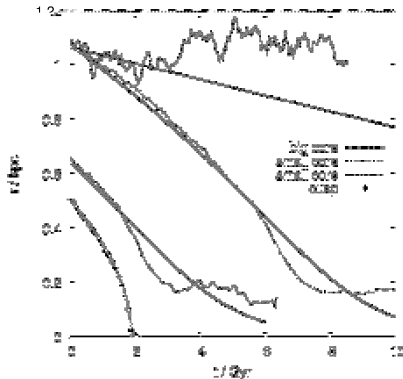
# CDM Problems at small-scales ... ?

- NFW profile fits nicely DM in galaxy clusters  $\rho \propto r^{-1}(r+r_c)^{-2}$
- Dwarf galaxy density profiles:  $\rho_M(r) \propto r^{-(0.5-1.5)}$  cusp  
most DM-dominated objects

## Cores observed (?)

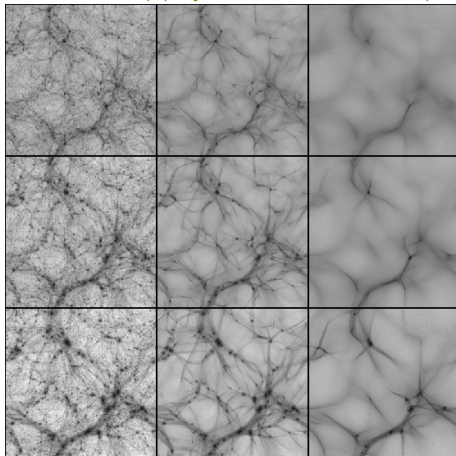


## 5 Clusters in the Fornax dSph



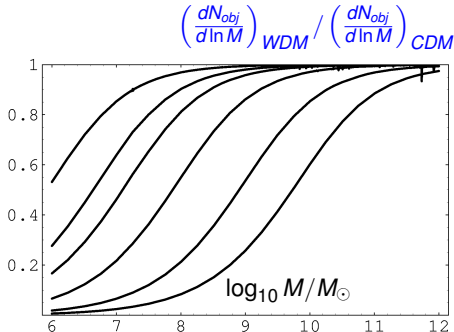
# CDM Problems ... ?

- Missing satellites:  $\frac{dN_{obj}}{d \ln M} \propto \frac{1}{M}$  no-scale 100 instead of 1000
- “Too big to fail” problem
- Solved (?) by Warm Dark Matter (sterile neutrino, gravitino) free-streaming



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Dark Matter Models (I)



DIAS Summer 2016

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# Matter perturbations

- CMB is isotropic, but “up to corrections, of course...”

- 1 Earth movement with respect to CMB

$$\frac{\Delta T_{\text{dipole}}}{T} \sim 10^{-3}$$

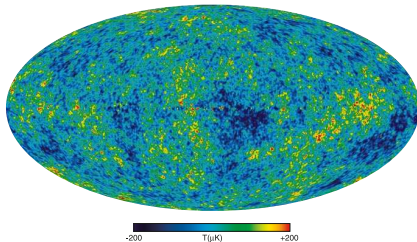
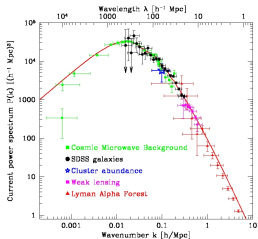
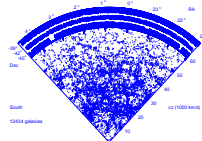
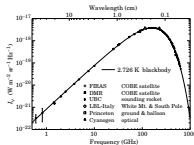
- 2 More complex anisotropy:  $\frac{\Delta T}{T} \sim 10^{-4}$

- There were matter inhomogeneities  $\Delta\rho/\rho \sim \Delta T/T$  at the stage of recombination ( $e + p \rightarrow \gamma + H^*$ )  $\Rightarrow$

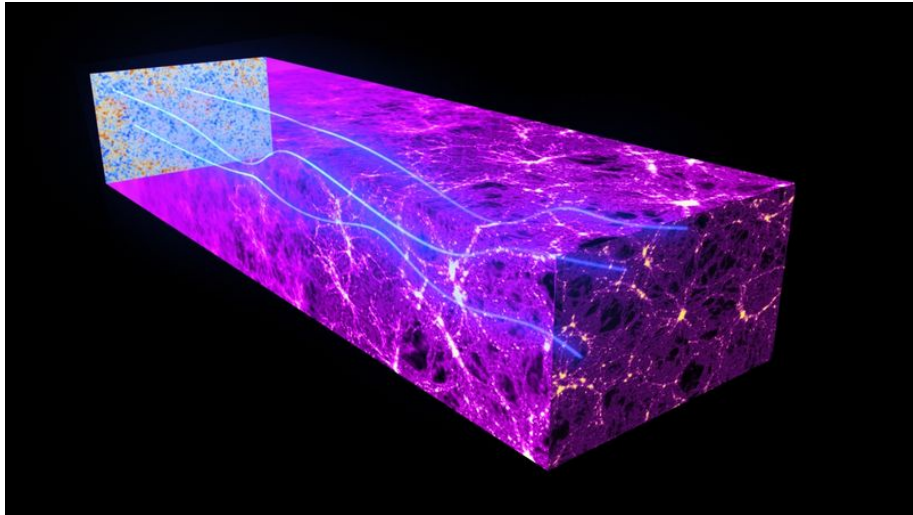
Jeans instability in the system of gravitating particles at rest  $\Rightarrow \Delta\rho/\rho \nearrow$  galaxies (CDM halos)

- $\Delta\rho_{DM}/\rho_{DM} \propto a \propto 1/T$  from  $T = 0.8 \text{ eV}$ ,  
while  $\Delta\rho_B/\rho_B \propto a \propto 1/T$  only after recombination  
 $T = 0.25 \text{ eV}$

without DM total growth factor would be 1100  
not enough to explain structures!



# On top of that: propagation in expanding Universe



# So far only gravitational evidence for DM

$$\left(\frac{\dot{a}}{a}\right)^2 = H^2(t) = \frac{8\pi}{3} G \rho_{\text{density}}^{\text{energy}}$$

$$\rho_{\text{density}}^{\text{energy}} = \rho_{\text{radiation}} + \rho_{\text{matter}}^{\text{ordinary}} + \rho_{\text{matter}}^{\text{dark}} + \rho_{\Lambda}$$

$$\rho_{\text{radiation}} \propto 1/a^4(t) \propto T^4(t), \quad \rho_{\text{matter}} \propto 1/a^3(t)$$

$$\rho_{\Lambda} = \text{const}$$

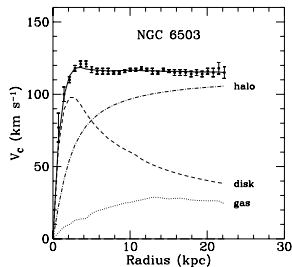
Why do we think it is most probably new particle physics  
(new gravity if any is not enough) ?

DM phenomena happen at various spatial and time scales



# Dark Matter in astrophysics

## Rotational curves



## Gravitational lensing

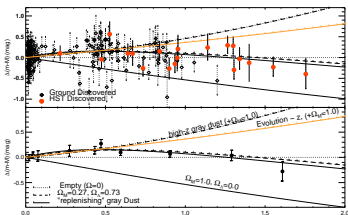


## X-rays from centers of galaxy clusters

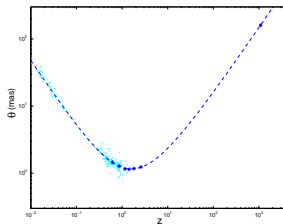
## “Bullet” cluster

# Dark matter in cosmology

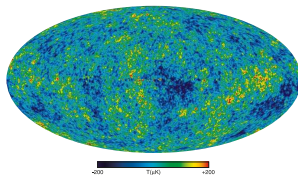
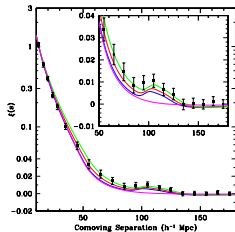
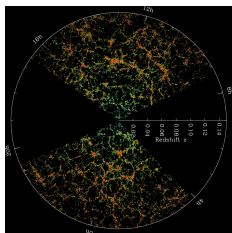
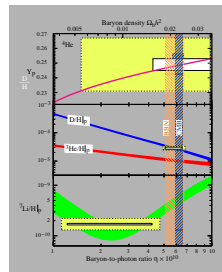
## Standard candles



## Angular distance



## Nucleosynthesis



Large Scale Structures

Baryon acoustic oscillations

CMB anisotropy

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# Dark Matter properties from cosmology: $p = 0$

(If) particles:

- 1 **stable** on cosmological time-scale  
requires new (almost) **conserved quantum number**
  - 2 **produced in the early Universe**  
some time before RD/MD-transition ( $T = 0.8 \text{ eV}$ )
  - 3 **nonrelativistic** particles long before RD/MD-transition ( $T = 0.8 \text{ eV}$ )  
(either **Cold** or **Warm**,  $v_{RD/MD} \lesssim 10^{-3}$ )  
Otherwise no small-size structures, like dwarf galaxies:  
smoothed out by free streaming
- If were in **thermal equilibrium**:  $M_X \gtrsim 1 \text{ keV}$
- 4 (almost) **collisionless**  $p = 0, v_{\text{sound}} = 0$
  - 5 (almost) electrically **neutral** CMB distortion
  - 6 **all matter inhomogeneities (perturbations) are adiabatic**:

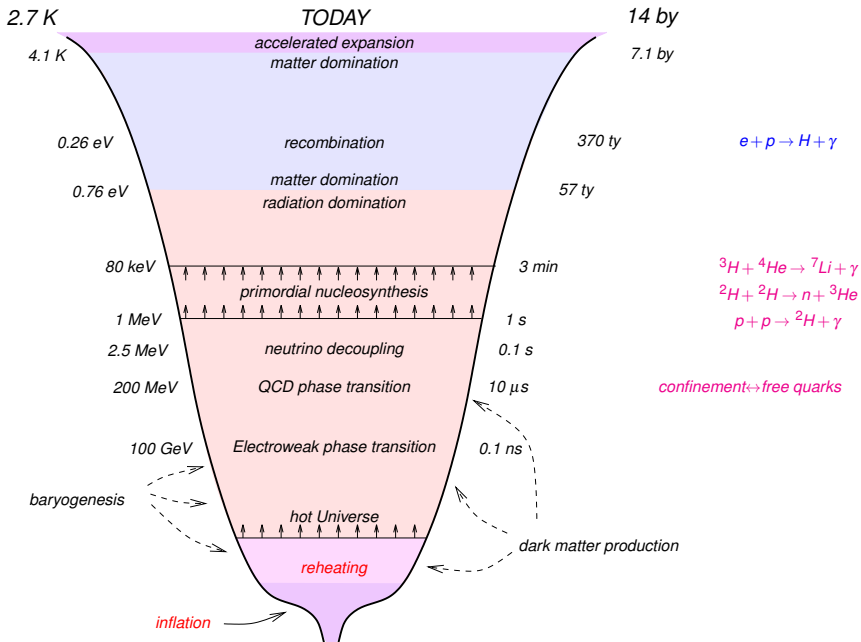
$$\delta \left( \frac{n_B}{n_{DM}} \right) = \delta \left( \frac{n_B}{n_\gamma} \right) = \delta \left( \frac{n_\nu}{n_\gamma} \right) = 0$$

# Dark Matter properties from astrophysics

- 1 **stable** on cosmological time-scale
  - 2 (almost) **collisionless** to form ellipsoidal halos
  - 3 (almost) electrically **neutral** to be Dark
  - 4 **stability of globular stellar clusters**  $M_X \lesssim 10^3 M_\odot \approx 10^{61} \text{ GeV}$  otherwise too strong tidal forces
  - 5 **confinement in a galaxy:** quantum physics!
- de Broglie wavelength:  $\lambda = 2\pi / (M_X v_X) < l_{\text{galaxy}}$ , for bosons
- in a galaxy  $v_X \sim 0.5 \cdot 10^{-3}$   $\longrightarrow$   $M_X \gtrsim 3 \cdot 10^{-22} \text{ eV}$  for fermions
- Pauli blocking:  $M_X \gtrsim 750 \text{ eV}$

$$f(\mathbf{p}, \mathbf{x}) = \frac{\rho_X(\mathbf{x})}{M_X} \cdot \frac{1}{\left(\sqrt{2\pi} M_X v_X\right)^3} \cdot e^{-\frac{\mathbf{p}^2}{2M_X^2 v_X^2}} \Bigg|_{\mathbf{p}=0} \leq \frac{g_X}{(2\pi)^3}$$

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# Decoupling of relativistic Dark Matter

## Assumptions

- DM particles are in equilibrium in plasma
- DM decouple from plasma at temperature  $T_d \gtrsim M_X$ ,  
so they are **relativistic** (e.g. neutrino)

Later on

$$n_X(T_d) = g_X \cdot \left(\frac{1}{4}\right) \cdot \frac{\zeta(3)}{\pi^2} T_d^3$$

$$n_X a^3 = \text{const}, \quad sa^3 = \text{const} \quad \Rightarrow \quad \frac{n_X}{s} = \text{const} = \# \frac{g_X}{g_*(T_d)}$$

useful

DM particle mass  $M_X$  fixes  $\Omega_X$ :

$$\Omega_X = \frac{M_X \cdot n_{X,0}}{\rho_c} = \frac{M_X \cdot s_0}{\rho_c} \frac{n}{s} \approx 0.2 \times \frac{M_X}{100 \text{ eV}} \left(\frac{g_X}{2}\right) \cdot \left(\frac{100}{g_*(T_d)}\right)$$

– NO heavy stable feebly coupled to SM particles !

– NO realistic DM models:

Pauli blocking prevents fermionic DM

$$\frac{p_X}{M_X} \propto \frac{a_d}{a} \sim \frac{3T}{M_X} \left(\frac{g_*(T)}{g_*(T_d)}\right)^{1/3}$$

too energetic for the proper structure formation



# Decoupling of relativistic Dark Matter

Can we save the relativistic Dark Matter ??

one can try, say, nonstandard cosmological evolution

with entropy production

- 1 hot stage (radiation domination)  $\rho \propto 1/a^4$
- 2 add new nonrelativistic particles decoupled from plasma  $\rho \propto 1/a^3$
- 3 later they start to dominate  
intermediate stage of matter domination terminates before BBN !!
- 4 both relativistic DM density and entropy density drop
- 5 new nonrelativistic particles decay reheating the Universe  $T > 3 \text{ MeV}$   
entropy production

# Decoupling of nonrelativistic Dark Matter

## Assumptions:

- 1 no  $X - \bar{X}$  asymmetry
  - 2 @  $T \lesssim M_X$  in thermal equilibrium with plasma
- either  $X = \bar{X}$  or  $n_X = n_{\bar{X}}$  (e.g. neutrons)

$$n_X = n_{\bar{X}} = g_X \left( \frac{M_X T}{2\pi} \right)^{3/2} e^{-M_X/T}$$

$X\bar{X} \rightarrow$  light particles

freeze-out temperature  $T_f$

$$H \equiv T^2 / M_{\text{Pl}}^*$$

$$n_X \langle \sigma_{\text{ann}} v \rangle = H(T_f) \rightarrow T_f = \frac{M_X}{\ln \left( \frac{g_X M_X M_{\text{Pl}}^* \sigma_0}{(2\pi)^{3/2}} \right)}.$$

Bethe formula:

$$\text{s-wave: } \sigma_{\text{ann}} = \frac{\sigma_0}{v}$$

# Weakly Interacting Massive Particles

density after freeze-out:

$$n_x(T_f) = \frac{T_f^2}{M_{\text{Pl}}^* \sigma_0}$$

present density:

$$n_x(T_0) = \left(\frac{a(T_f)}{a(T_0)}\right)^3 n_x(T_f) = \left(\frac{s_0}{s(T_f)}\right) n_x(T_f) \propto \frac{1}{T_f}$$

$X + \bar{X}$  contribution to critical density:

$$\begin{aligned} \Omega_x &= 2 \frac{M_x n_x(T_0)}{\rho_c} = 7.6 \frac{s_0 \ln\left(\frac{g_x M_{\text{Pl}}^* M_x \sigma_0}{(2\pi)^{3/2}}\right)}{\rho_c \sigma_0 M_{\text{Pl}} \sqrt{g_*(T_f)}} \\ &= 0.1 \cdot \left(\frac{(10 \text{ TeV})^{-2}}{\sigma_0}\right) \frac{10}{\sqrt{g_*(T_f)}} \ln\left(\frac{g_x M_{\text{Pl}}^* M_x \sigma_0}{(2\pi)^{3/2}}\right) \cdot \frac{1}{2h^2} \end{aligned}$$

# WIMPs: discussion

$$\Omega_X = 0.1 \cdot \left( \frac{(10 \text{ TeV})^{-2}}{\sigma_0} \right) \frac{10}{\sqrt{g_*(T_f)}} \ln \left( \frac{g_X M_{\text{Pl}}^* M_X \sigma_0}{(2\pi)^{3/2}} \right) \cdot \frac{1}{2h^2}$$

- **natural DM: subweak-scale cross section**  $\sigma_0 \sim 0.01 \times \sigma_W$   
say,  $M_X \sim 1 \text{ TeV}$  or  $X$  is not a weak gauge eigenstate
- **naturally "light"**                      unitarity                       $\sigma_0 \lesssim \frac{4\pi}{M_X^2} \rightarrow M_X \lesssim 100 \text{ TeV}$
- **all stable particles with smaller  $\sigma_0$  are forbidden !!**
- WIMPs remain in kinetic equilibrium with plasma till  $T \sim 10 \text{ MeV}$   
this is **Cold Dark Matter**,  $v_{RD/MD} \ll 10^{-3}$   
WIMPs may form dark halos (clumps) much lighter than  
dwarf galaxies

# Weakly IMPs are mostly welcome (e.g. LSP in SUSY)

## We can fully explore the model !!

lectures by S.Demidov

- Direct searches for Galactic Dark Matter ( $v \sim 10^{-3}$ )

$$X + \text{nuclei} \rightarrow X + \text{nuclei} + \Delta E$$

- Can search for WIMPs in cosmic rays: products of WIMPs annihilation (in Galactic center, dwarf galaxies, Sun)

$$X + \bar{X} \rightarrow p\bar{p}, e^+e^-, \nu, \gamma, \dots$$

- Can search for WIMPs in collision experiments (LHC):

$$X + \bar{X} \leftrightarrow \text{SM} + \text{SM}' + \dots$$

# If thermal CDM but not **Weakly** IMPs?

We still can study the model if DM annihilates (partly) into SM particles

- But DM particle  $X$  can be light and feebly coupled ( $t$ -channel)

$$\sigma_0 \sim \frac{\xi^4}{M_X^2}$$

$\xi$  is not a gauge coupling within GUT !

- With small  $\sigma_0$  one needs entropy production
- $\sigma_0$  may be increased by **s-channel resonance**,  $M_Y \approx 2M_X$
- annihilation can be amplified by **co-annihilation channels**,  $X + A \rightarrow SM$
- With light messengers between Dark and Visible sectors many estimates change, say  $\sigma_0 = \sigma_0(\nu)$
- DM interaction at freeze-out and now are not the same  
say, **Sommerfeld enhancement** of the annihilation of slow particles  $v \sim 10^{-3}$

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# Summary (I)

- 1 We need DM both in past (cosmology)  
and at present (astrophysics)
- 2 For stability a symmetry is needed
- 3 There are claimed discrepancies between CDM simulations and observations of small scale structures, observations of central regions of dwarf galaxies
- 4 WDM? selfinteracting DM? no proof
- 5 Structures: DM cannot be hot (e.g. SM neutrinos can not help)
- 6 WIMPs (neutralino) are natural candidates for Cold Dark Matter
- 7 Much more options for WIMP-like candidates. . .
- 8 Generally, heavy and/or feebly coupled thermal relics  
are forbidden !!