

Study of transport properties of QGP

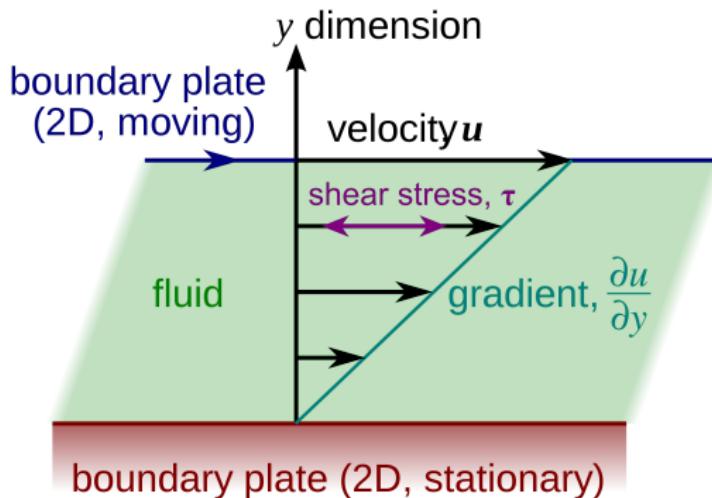
V.V. Braguta

ITEP

23 August, 2017

Outline:

- Temperature dependence of shear viscosity
- Temperature dependence of bulk viscosity



Shear viscosity

- $F_x = -\eta \cdot \frac{du}{dy} \cdot S$, η -viscosity
- Shear viscosity is connected with non-diagonal component of energy-momentum tensor: T_{xy}

Relativistic Hydrodynamics

- $T^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu} + (\eta \nabla^{\langle\mu} u^{\nu\rangle} + \zeta \Delta^{\mu\nu} \nabla_\alpha u^\alpha) + \dots$
 $\nabla^\alpha = \Delta^{\alpha\nu} \partial_\nu, \quad \Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$
 $\nabla^{\langle\mu} u^{\nu\rangle} = \nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} \nabla_\alpha u^\alpha$
- EOM $\partial_\mu T^{\mu\nu} = 0$
- Non-relativistic limit ($u^\mu = (1, \vec{v})$)
 - Continuity equation: $\partial_t \rho + \rho(\vec{\partial} \vec{v}) + \vec{v} \vec{\partial} \rho = 0$
 - Navier–Stokes equation: $\frac{\partial v^i}{\partial t} + v^k \frac{\partial v^i}{\partial x^k} = -\frac{1}{\rho} \frac{\partial p}{\partial x^i} - \frac{1}{\rho} \frac{\partial \Pi^{ki}}{\partial x^k}$
 - Viscous stress tensor: $\Pi^{ik} = -\eta \left(\frac{\partial v^i}{\partial x^k} + \frac{\partial v^k}{\partial x^i} - \frac{2}{3} \delta^{ik} \frac{\partial v^l}{\partial x^l} \right) - \zeta \delta^{ik} \frac{\partial v^l}{\partial x^l}$
- η –shear viscosity, ζ –bulk viscosity

Calculation of transport coefficients:

- Shear viscosity

$$\eta = \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^\infty dt \int d^3x e^{i\omega t} \langle [T_{12}(x), T_{12}(0)] \rangle$$

$$\rho_{12,12}(\omega) = -\frac{1}{\pi} \text{Im} G_R^{12,12}(\omega, \vec{k} = 0)$$

$$\eta = \pi \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{12,12}(\omega)$$

- Bulk viscosity

$$\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^\infty dt \int d^3x e^{i\omega t} \langle [T_\mu^\mu(x), T_\nu^\nu(0)] \rangle$$

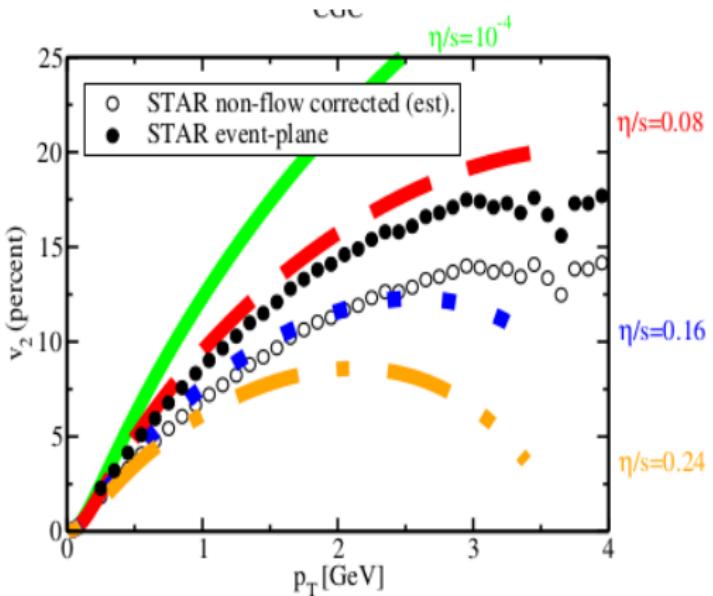
$$\rho_{\mu\mu,\nu\nu}(\omega) = -\frac{1}{\pi} \text{Im} G_R^{\mu\mu,\nu\nu}(\omega, \vec{k} = 0)$$

$$\zeta = \frac{\pi}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{\mu\mu,\nu\nu}(\omega)$$

- Analytic continuation

$$G_E(\omega, \vec{p}) = -G_R(i\omega, \vec{p}), \quad \omega > 0$$

$$C_E(t) = \int_0^\infty d\omega \rho(\omega) \frac{\text{ch}\left(\frac{\omega}{2T} - \omega t\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

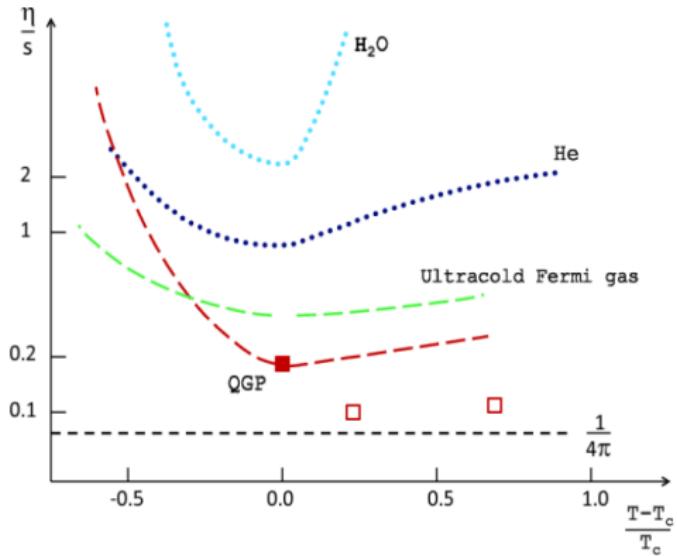


Elliptic flow from STAR experiment (Nucl. Phys. A 757, 102 (2005))

$$\frac{dN}{d\phi} \sim (1 + 2v_1 \cos(\phi) + 2v_2 \cos^2(\phi)), \text{ } \phi\text{-scattering angle}$$

Quark-gluon plasma is close to ideal liquid ($\frac{\eta}{s} = (1 - 3)\frac{1}{4\pi}$)

M. Luzum and P. Romatschke, Phys. Rev. C 78, 034915 (2008)

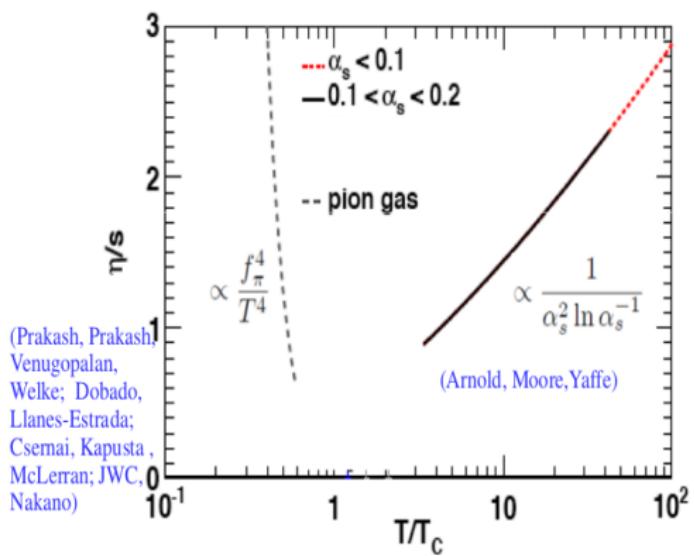


S.Cremonini, U.Gursoy, P.Szepietowski, JHEP 1208 (2012) 167

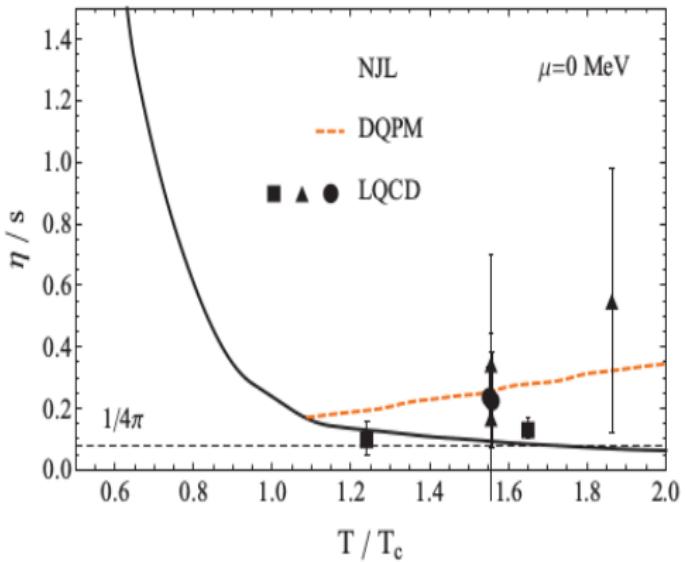
Comparison of different liquids

Quark-gluon plasma is the most ideal liquid

Study of shear viscosity (perturbative calculation + pion gas)



Study of shear viscosity (effective models)



R. Marty, E. Bratkovskaya, W. Cassing, J. Aichelin and H. Berrehrah, Phys. Rev. C 88, 045204 (2013)

Other works (SU(3) gluodynamics):

- Karsch, F. et al. Phys.Rev. D35 (1987)
- A. Nakamura, S. Sakai Phys. Rev. Lett. 94, 072305 (2005)
- H. B. Meyer, Phys.Rev. D76 (2007) 101701
- H. B. Meyer, Nucl.Phys. A830 (2009) 641C-648C

Results:

- $\frac{\eta}{s} = 0.134 \pm 0.033$ ($T/T_c = 1.65, 8 \times 28^3$)
- $\frac{\eta}{s} = 0.102 \pm 0.056$ ($T/T_c = 1.24, 8 \times 28^3$)
- $\frac{\eta}{s} = 0.20 \pm 0.03$ ($T/T_c = 1.58, 16 \times 48^3$)
- $\frac{\eta}{s} = 0.26 \pm 0.03$ ($T/T_c = 2.32, 16 \times 48^3$)

SU(2) gluodynamics:

- $\frac{\eta}{s} = 0.134 \pm 0.057$ ($T/T_c = 1.2, 16 \times 32^3$)

N.Yu. Astrakhantsev, V.V. Braguta, A.Yu. Kotov, JHEP 1509 (2015) 082

Lattice calculation of shear viscosity

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The first step:

Measurement of the correlation function:

$$C(t) = \langle T_{12}(t) T_{12}(0) \rangle$$

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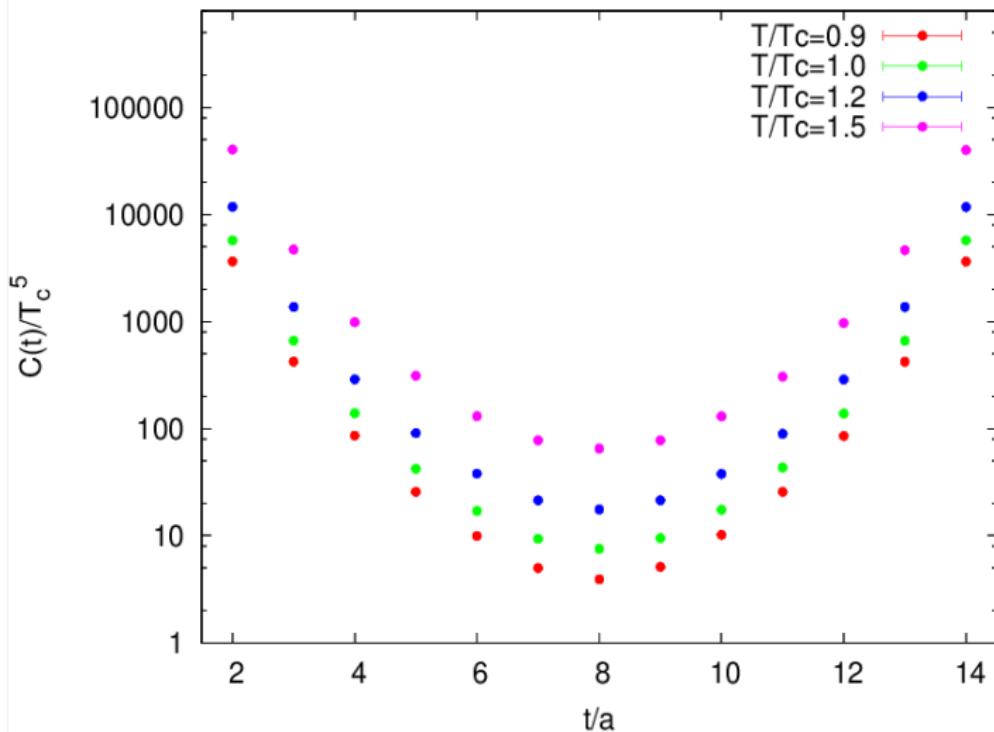
Calculation of the spectral function $\rho(\omega)$:

$$C(t) = \int_0^\infty d\omega \rho(\omega) \frac{\text{ch}\left(\frac{\omega}{2T} - \omega t\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$
$$\eta = \pi \lim_{\omega \rightarrow 0} \frac{\rho(\omega)}{\omega}$$

Details of the calculation

- SU(3) gluodynamics
- Two-level algorithm
- Lattice size $32^3 \times 16$
- Temperatures $T/T_c = 0.9, 0.925, 0.95, 1.0, 1.1, 1.2, 1.35, 1.5$
- Accuracy $\sim 2 - 3\%$ at $t = \frac{1}{2T}$
- $\langle T_{12}(x) T_{12}(y) \rangle \sim (\langle T_{11}(x) T_{11}(y) \rangle - \langle T_{11}(x) T_{22}(y) \rangle)$
- Clover discretization for the $\hat{F}_{\mu\nu}$
- Renormalization of EMT: F. Karsch, Nucl.Phys. B205 (1982) 285-300
- ...

Correlation functions



Spectral function

$$C(t) = \int_0^\infty d\omega \rho(\omega) \frac{ch\left(\frac{\omega}{2T} - \omega t\right)}{sh\left(\frac{\omega}{2T}\right)}$$

Properties of the spectral function:

- $\rho(\omega) \geq 0, \rho(-\omega) = -\rho(\omega)$
- Asymptotic freedom: $\rho(\omega)|_{\omega \rightarrow \infty}^{NLO} = \frac{1}{10} \frac{d_A}{(4\pi)^2} \omega^4 \left(1 - \frac{5N_c \alpha_s}{9\pi}\right)$
 $\sim 90\%$ of the total contribution $t = 1/2/T$
- Hydrodynamics: $\rho(\omega)|_{\omega \rightarrow 0} = \frac{\eta}{\pi} \omega$

Spectral function

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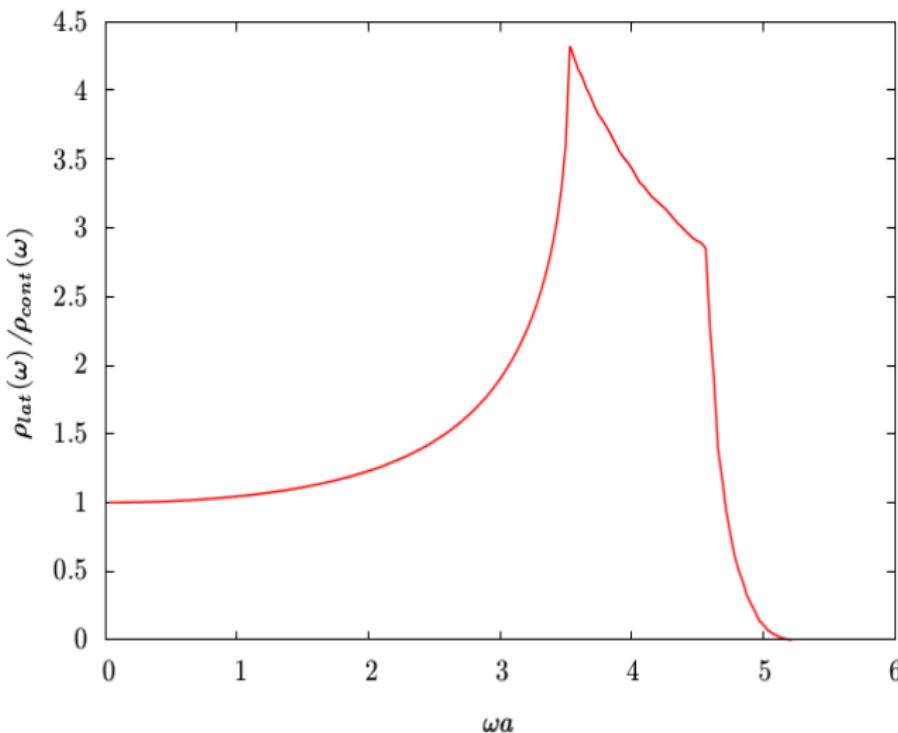
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Ansatz for the spectral function (QCD sum rules motivation)

$$\rho(\omega) = \frac{\eta}{\pi} \omega \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

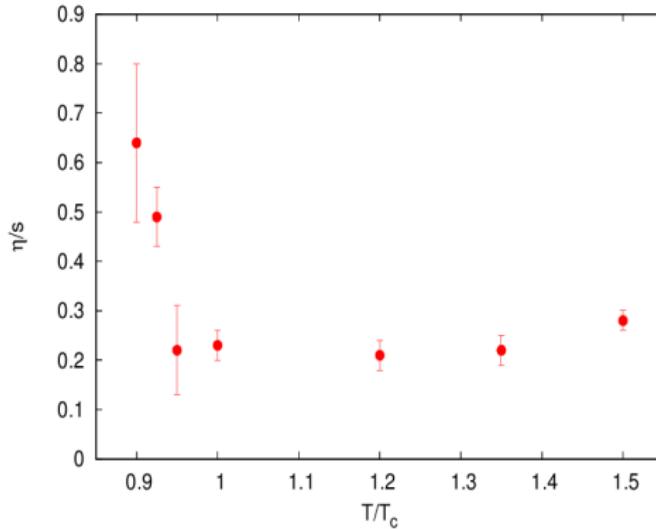
Lattice spectral function ρ_{lat}

Takes into account discretization errors in temporal direction



Spectral function

$$\rho_1(\omega) = \frac{\eta}{\pi} \omega \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$
$$\chi^2/dof \sim 1, A \sim 1, \omega_0/T \sim 7 - 8$$



Two additional ansatzs:

- $\rho_2(\omega) = \frac{1}{2} B \omega (1 + \tanh[\gamma(\omega_0 - \omega)]) + \frac{1}{2} A \rho_{lat}(\omega) (1 + \tanh[\gamma(\omega - \omega_0)])$
- $\rho_3(\omega) = B \omega (1 + C \omega^2) \theta(\omega_0 - \omega) + A \rho_{lat}(\omega) \theta(\omega - \omega_0)$

Properties of the spectral function

- Hydrodynamical approximation works well up to $\omega < \pi T \sim 1 \text{ GeV}$ (H.B. Meyer, arXiv:0809.5202)

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Properties of the spectral function

- Hydrodynamical approximation works well up to $\omega < \pi T \sim 1 \text{ GeV}$ (H.B. Meyer, arXiv:0809.5202)
- Asymptotic freedom works well from $\omega > 3 \text{ GeV}$
- Poor knowledge of the spectral function in the region $\omega \in (1, 3) \text{ GeV}$
⇒ Main source of uncertainty in the fitting procedure

Backus-Gilbert method for the spectral function

- Problem: find $f(\omega)$ from the integral equation

$$C(x_i) = \int_0^\infty d\omega f(\omega) K(x_i, \omega), \quad K(x_i, \omega) = \frac{\text{ch}\left(\frac{\omega}{2T} - \omega x_i\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

- Define an estimator $\tilde{f}(\bar{\omega})$ ($\delta(\bar{\omega}, \omega)$ - resolution function):

$$\tilde{f}(\bar{\omega}) = \int_0^\infty d\omega \hat{\delta}(\bar{\omega}, \omega) f(\omega)$$

- Let us expand $\delta(\bar{\omega}, \omega)$ as

$$\delta(\bar{\omega}, \omega) = \sum_i b_i(\bar{\omega}) K(x_i, \omega) \quad \tilde{f}(\bar{\omega}) = \sum_i b_i(\bar{\omega}) C(x_i)$$

- Goal: minimize the width of the resolution function

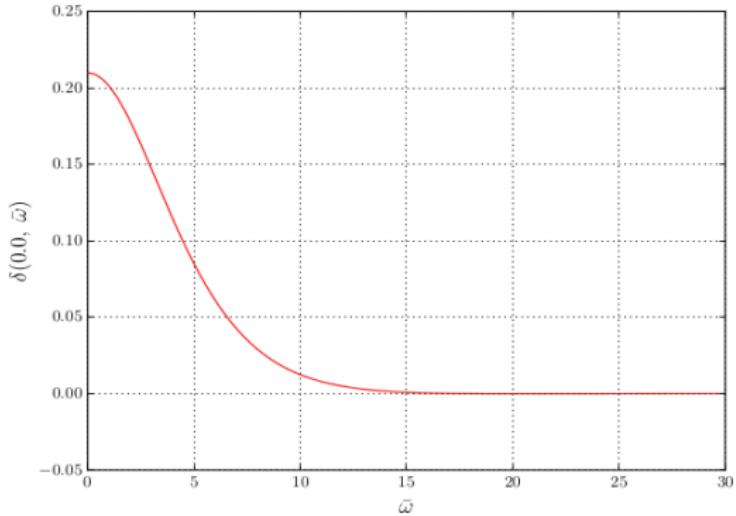
$$b_i(\bar{\omega}) = \frac{\sum_j W_{ij}^{-1} R_j}{\sum_{ij} R_i W_{ij}^{-1} R_j},$$

$$W_{ij} = \int d\omega K(x_i, \omega)(\omega - \bar{\omega})^2 K(x_j, \omega), R_i = \int d\omega K(x_i, \omega)$$

- Regularization by the covariance matrix S_{ij} :

$$W_{ij} \rightarrow \lambda W_{ij} + (1 - \lambda) S_{ij}, \quad 0 < \lambda < 1$$

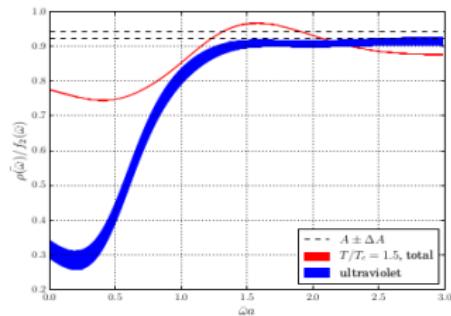
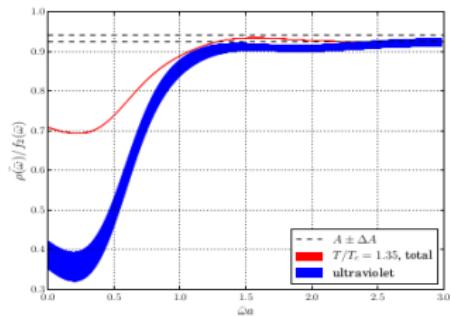
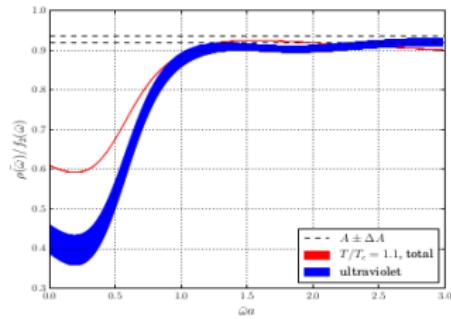
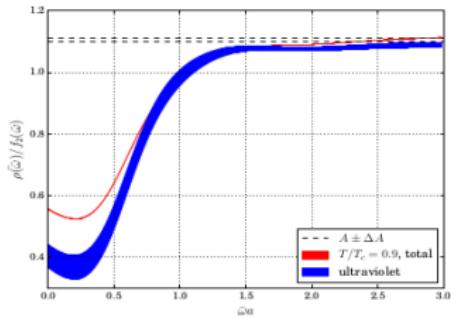
Resolution function $\delta(0, \omega)$ ($T/T_c = 1$, $\lambda = 0.001$)



- Width of the resolution function $\omega/T \sim 4$
- Hydrodynamical approximation works up to $\omega/T < \pi$
- Problem: large contribution from ultraviolet tail ($\sim 50\%$)

Model for the ultraviolet contribution

$$\rho_{ultr} = A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

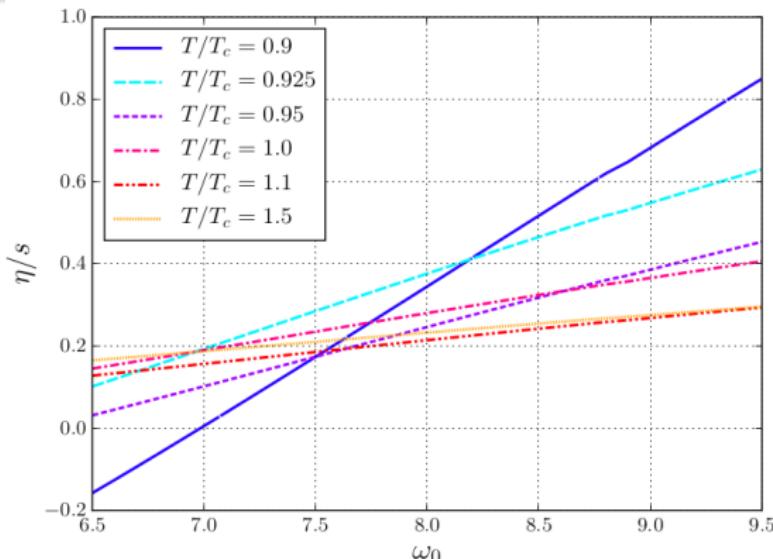


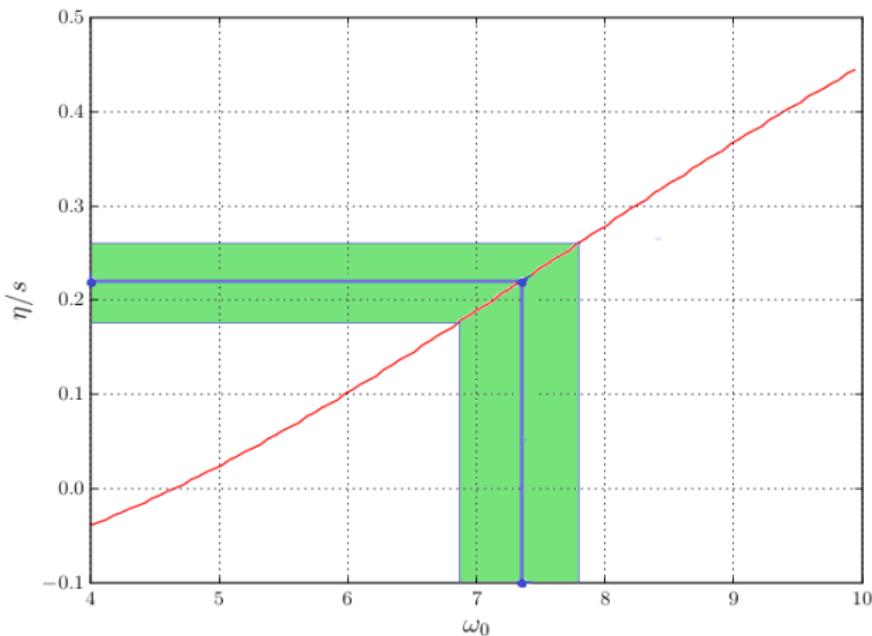
Solution:

- Take ultraviolet contribution in the form:

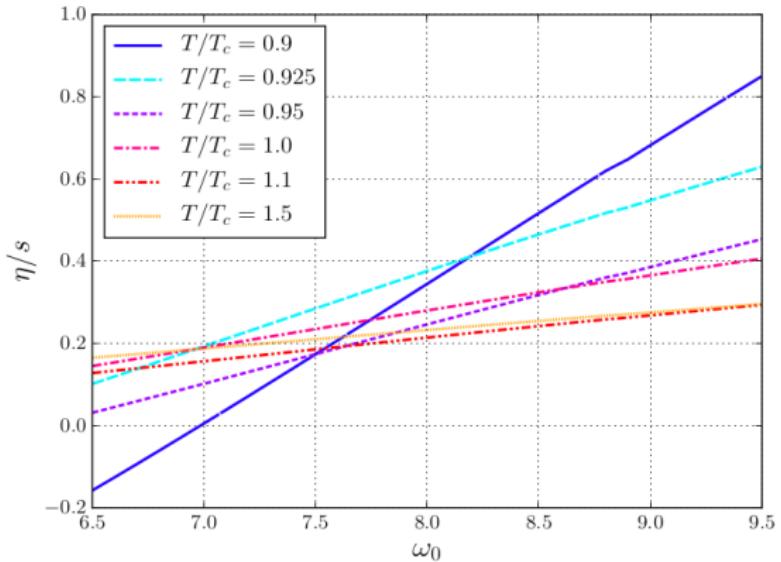
$$\rho_{ultr} = A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

- Determine the value of the A BG method
- Subtract ultraviolet contribution and obtain η/s as a function of ω_0

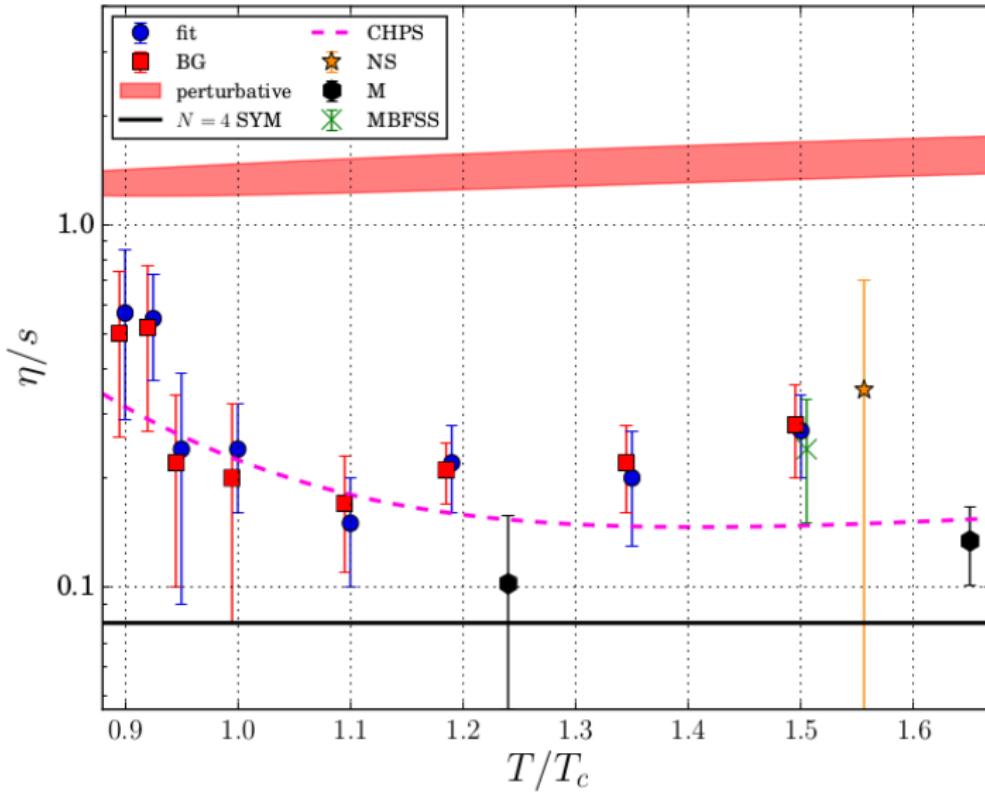




- For $T/T_c = 1$. $\omega_0/T = 7.33 \pm 0.47$
- $\eta/s = 0.22 \pm 0.04$



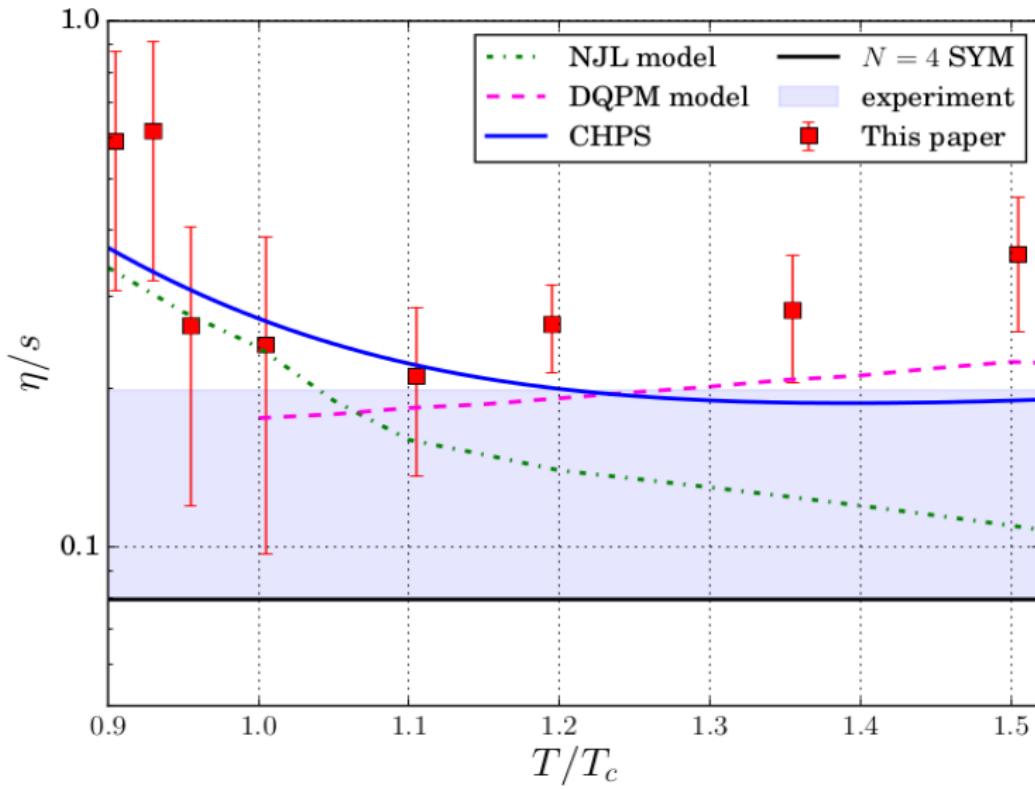
Our results

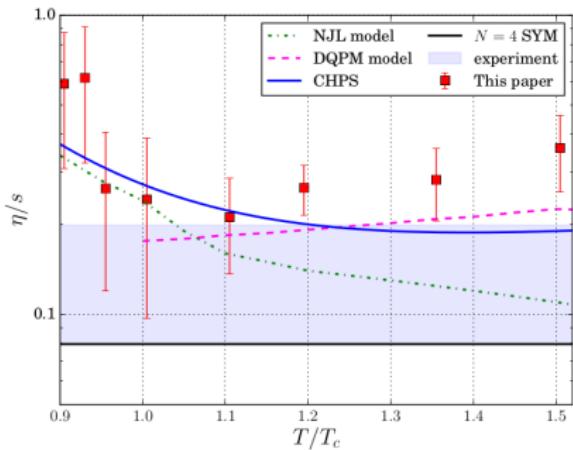
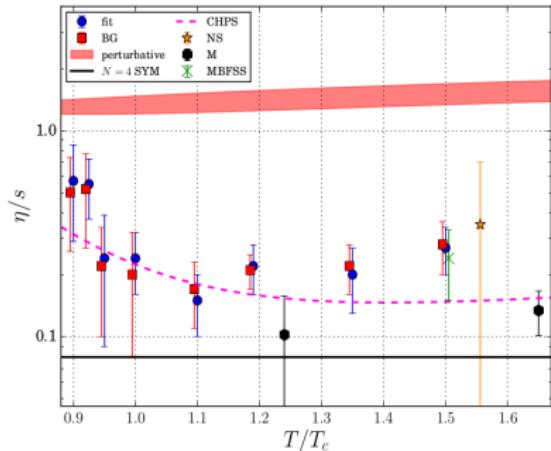


Estimation of the η/s for QCD with dynamical quarks ($N_f = 3$):

$$\left(\frac{\eta}{s}\right)_{QCD} = \left(\frac{(\eta/s)_{QCD}}{(\eta/s)_{Gluodynamics}} \right)_{pert} \times \left(\frac{\eta}{s}\right)_{Gluodynamics}$$

Our results





Conclusion:

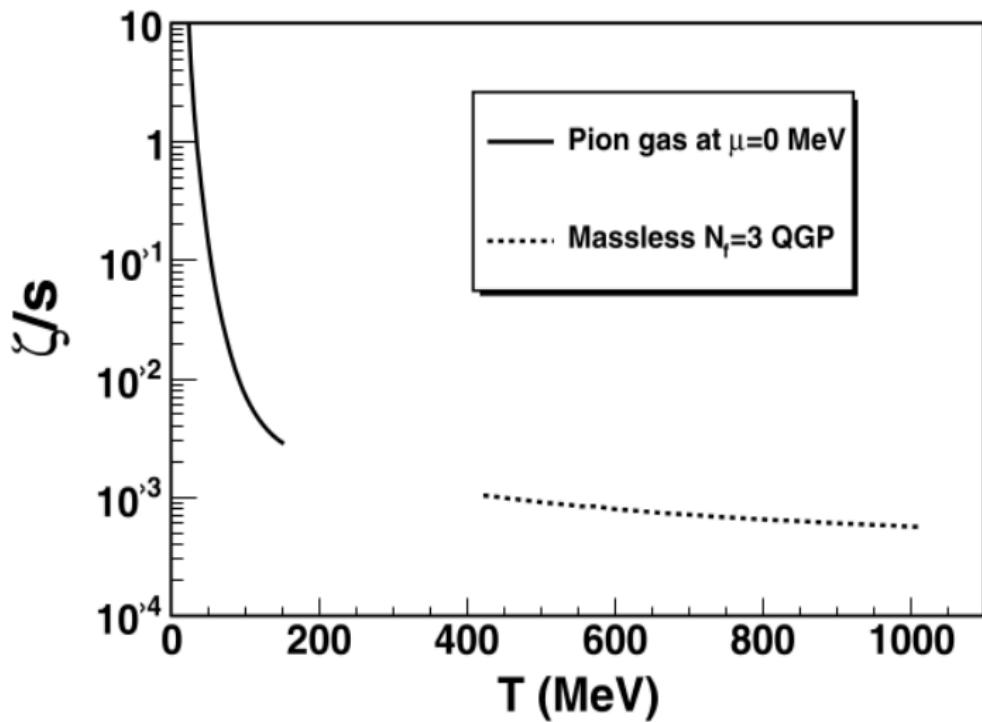
- We calculated η/s for set of temperatures $T/T_c \in (0.9, 1.5)$
- Applied fitting procedure and Backus-Gilbert method for the SF
- η/s is close to $N=4$ SYM and in agreement with experiment
- Large deviation from perturbative results

Bulk viscosity

- $\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int_0^\infty dt \int d^3x e^{i\omega t} \langle [T_\mu^\mu(x), T_\nu^\nu(0)] \rangle$
- $C_E(t) = \int_0^\infty d\omega \rho_{\mu\mu,\nu\nu}(\omega) \frac{ch\left(\frac{\omega}{2T} - \omega t\right)}{sh\left(\frac{\omega}{2T}\right)}$
- $\zeta = \frac{\pi}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \rho_{\mu\mu,\nu\nu}(\omega)$

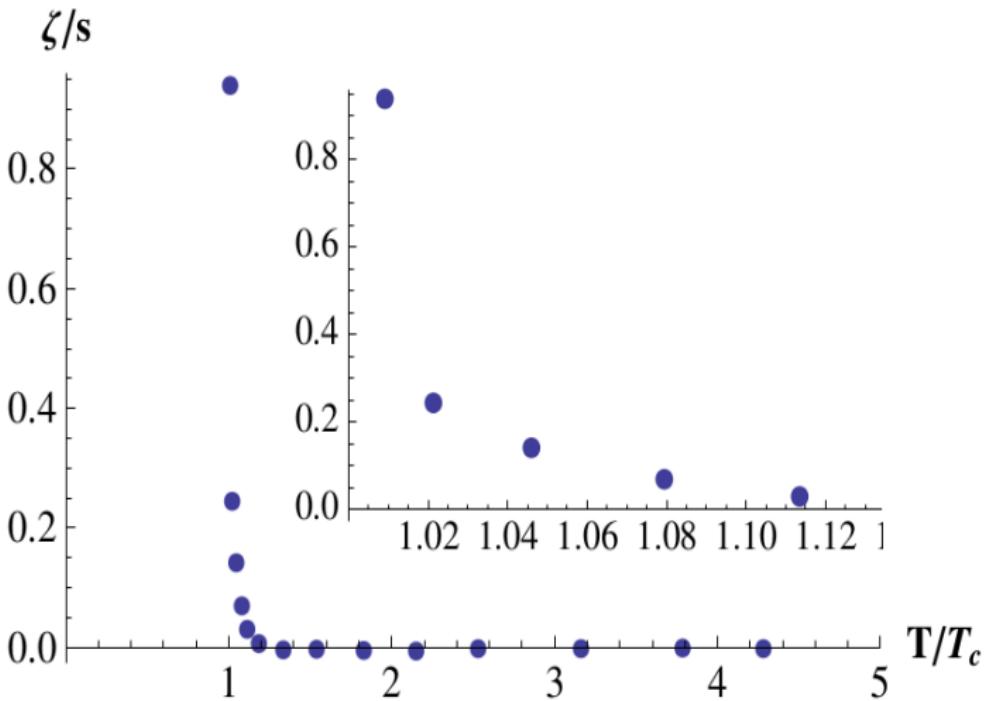
Violation of conformal symmetry in QCD

- In conformal theory $T_\mu^\mu = 0$
- Conformal symmetry is violated by quantum corrections
- Trace anomaly: $T_\mu^\mu = \frac{\beta(g)}{2g} Tr[G^{\mu\nu} G_{\mu\nu}] + \sum_i m_i \bar{q} q$



- CHPT: A. Dobado, F.J. Llanes-Estrada, J.M. Torres-Rincon, Physics Letters B 702 (2011) 43
- Perturbative QCD: P. Arnold, C. Dogan, G. Moore , Physical Review D 74, 085021 (2006)



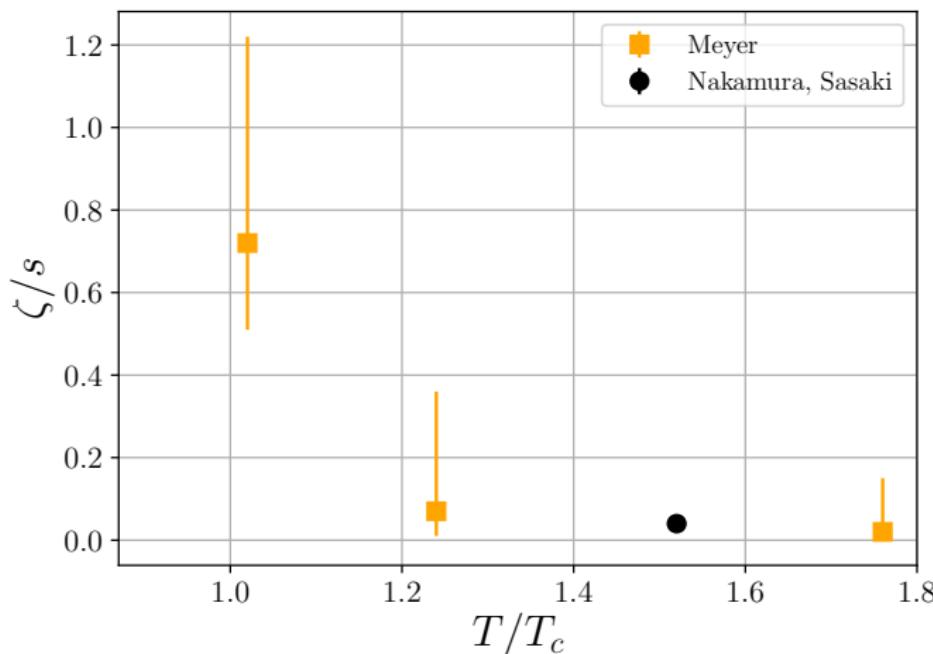


- Low energy theorems of QCD: $\zeta = \frac{1}{9\omega_0} \left(T^5 \frac{\partial}{\partial T} \frac{\epsilon - 3P}{T^4} + 16\epsilon_V \right)$

D. Kharzeev, K. Tuchin, JHEP 0809 (2008) 093,
D. Kharzeev, F. Karsch, K. Tuchin, Phys.Lett. B663 (2008) 217

Previous works (SU(3) gluodynamics):

- A. Nakamura, S. Sakai Phys. Rev. Lett. 94, 072305 (2005)
- H. B. Meyer, Phys.Rev.Lett. 100 (2008) 162001

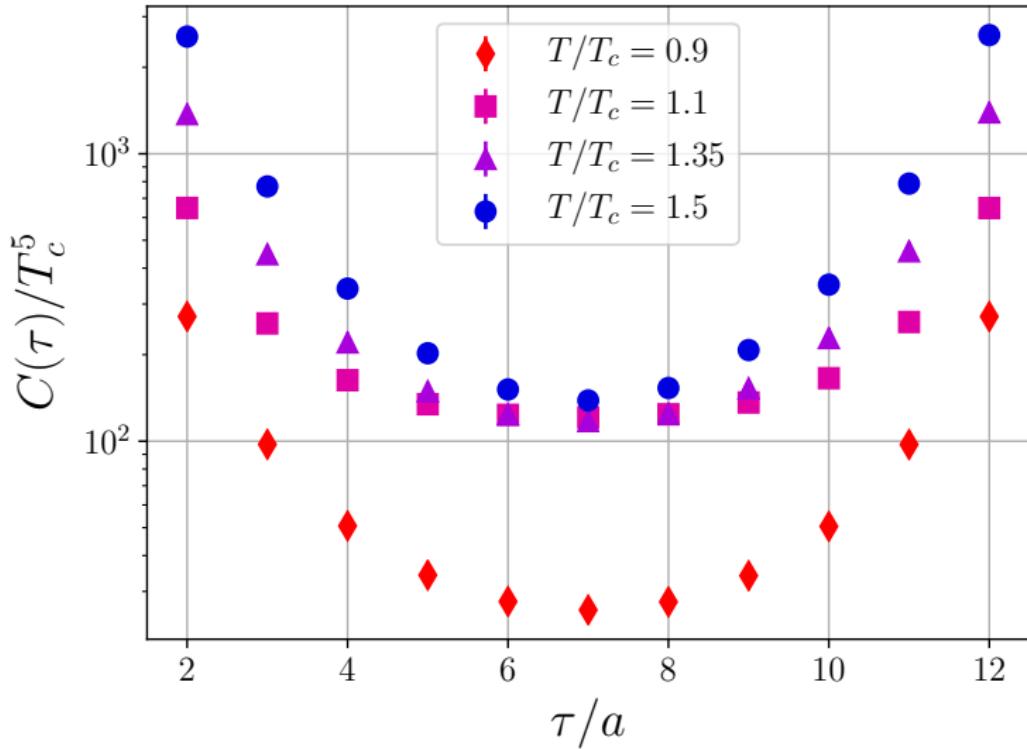


Details of the calculation

- SU(3) gluodynamics
- Two-level algorithm
- Lattice size $32^3 \times 16$
- Temperatures $T/T_c = 0.9, 0.925, 0.95, 1.0, 1.1, 1.2, 1.35, 1.5$
- Accuracy $\sim 3\%$ at $t = \frac{1}{2T}$
- Clover discretization for the $\hat{F}_{\mu\nu}$
- Renormalization of EMT: F. Karsch, Nucl.Phys. B205 (1982) 285-300
- ...

Our results are preliminary !

Correlation functions

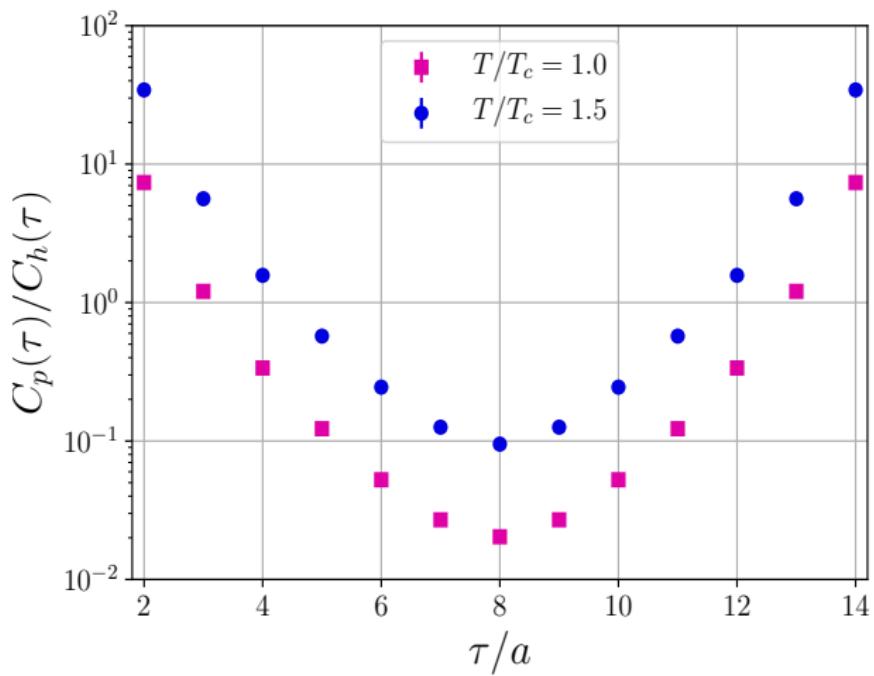


Spectral function

$$C(t) = \int_0^\infty d\omega \rho(\omega) \frac{\text{ch}\left(\frac{\omega}{2T} - \omega t\right)}{\text{sh}\left(\frac{\omega}{2T}\right)}$$

Properties of the spectral function:

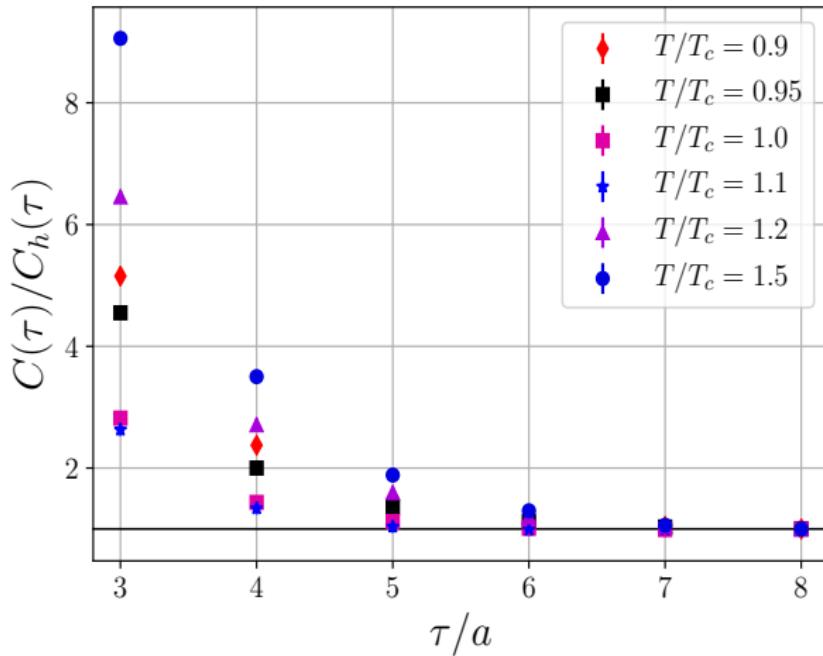
- $\rho(\omega) \geq 0, \rho(-\omega) = -\rho(\omega)$
- Asymptotic freedom: $\rho(\omega)|_{\omega \rightarrow \infty}^{NLO} = d_A \left(\frac{11\alpha_s}{(4\pi)^2} \right)^2 \omega^4$
compare with shear channel $\sim d_A \frac{1}{10(4\pi)^2} \omega^4$
- Hydrodynamics: $\rho(\omega)|_{\omega \rightarrow 0} = \frac{9}{\pi} \zeta \omega$

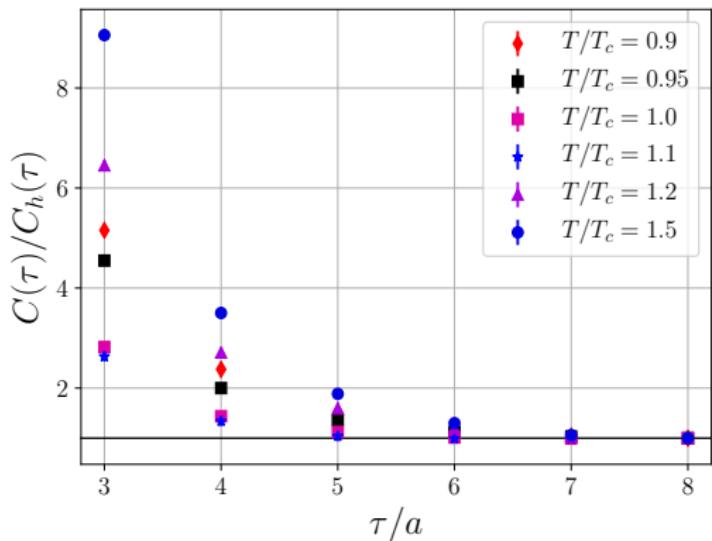


- In the region $\tau/a \sim \frac{\beta}{2}$ hydrodynamics is dominant
- In the region $\tau/a \sim \text{few}$ perturbative contribution is dominant

Hydrodynamical approximation

$$C_h(\tau) = \int_0^\infty d\omega \rho_h(\omega) \frac{ch\left(\frac{\omega}{2T} - \omega\tau\right)}{sh\left(\frac{\omega}{2T}\right)}, \quad \rho_h(\omega) = \frac{9}{\pi} \zeta \omega \theta(\omega_0 - \omega)$$

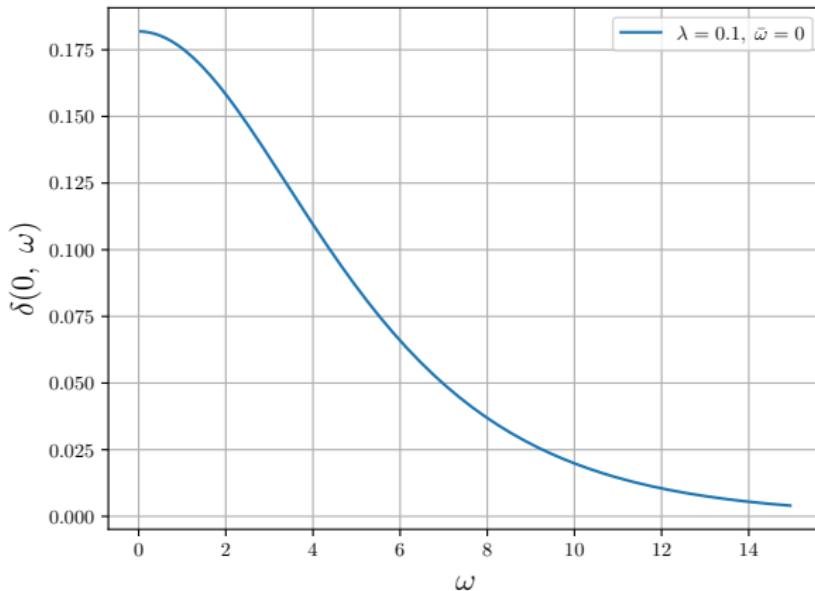




Middle point estimation of bulk viscosity

- In the vicinity of the phase transition hydrodynamics very works well!
- $C_h\left(\frac{\beta}{2}\right) = \frac{9}{\pi} \zeta \int_0^{\omega_0} d\omega \frac{\omega}{sh\left(\frac{\omega}{2T}\right)}$
- ω_0 is varied within the interval $1.5 - 3$ GeV

Resolution function $\delta(0, \omega)$ ($T/T_c = 1.5$, $\lambda = 0.1$)



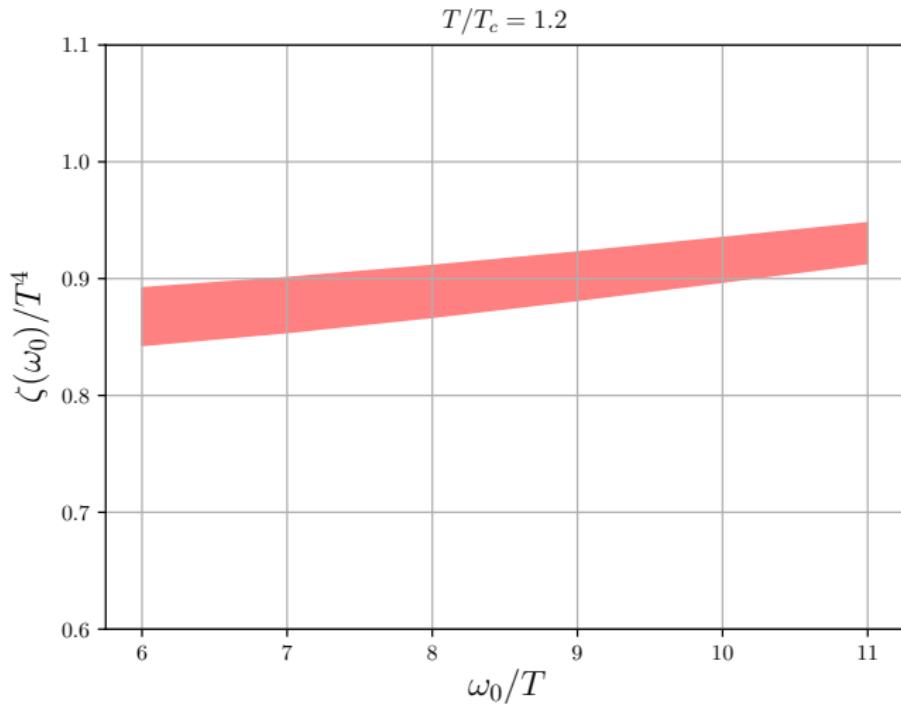
- Width of the resolution function $\omega/T \sim 5$

Removal of the ultraviolet contribution

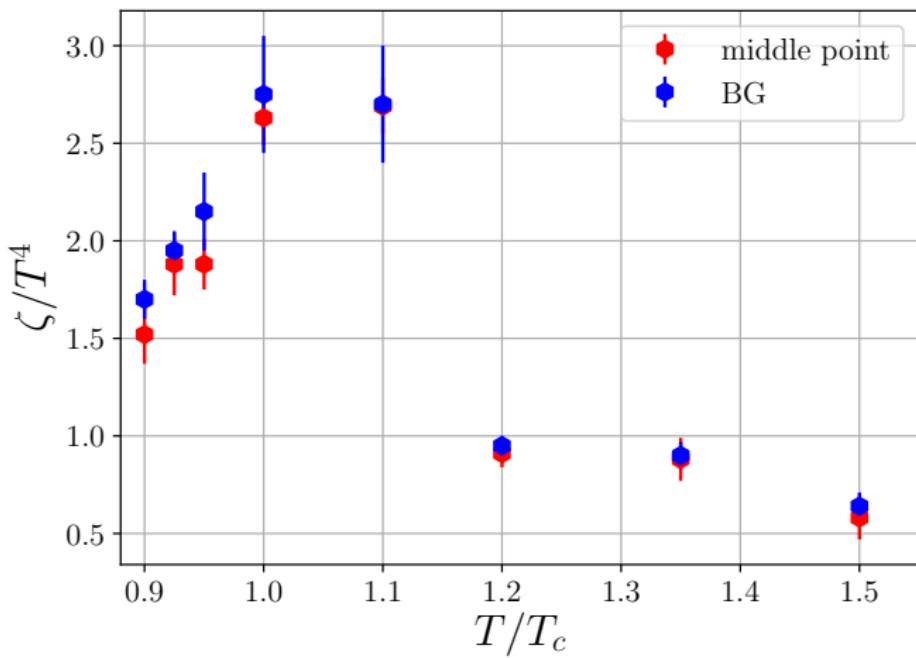
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$$\rho_{ultr} = A \rho_{lat}(\omega) \theta(\omega - \omega_0)$$

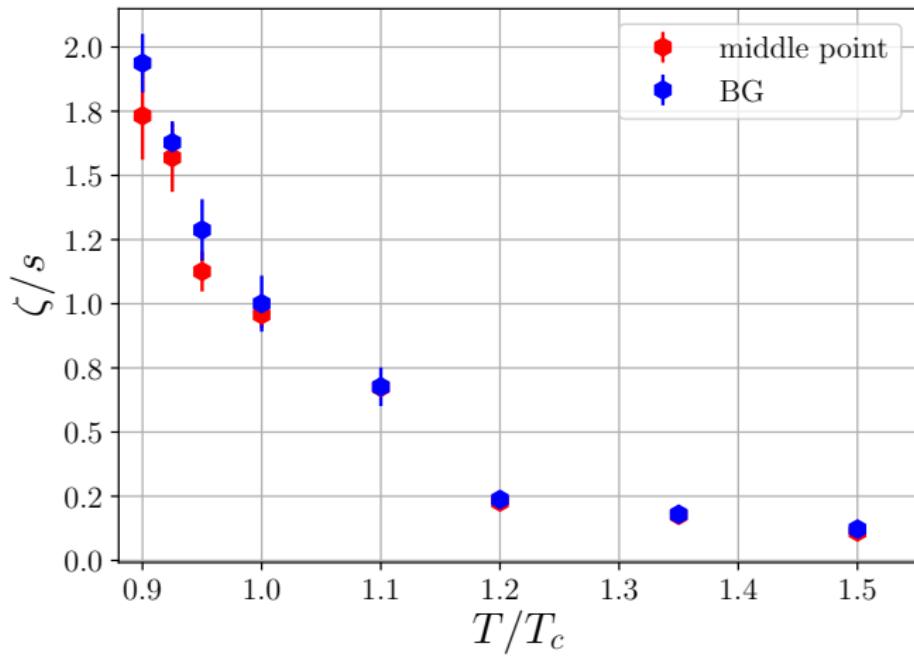
- Determine the value of the constant A from the $C(\tau/a = 2)$
- Subtract ultraviolet contribution and obtain ζ/T^4 as a function of ω_0



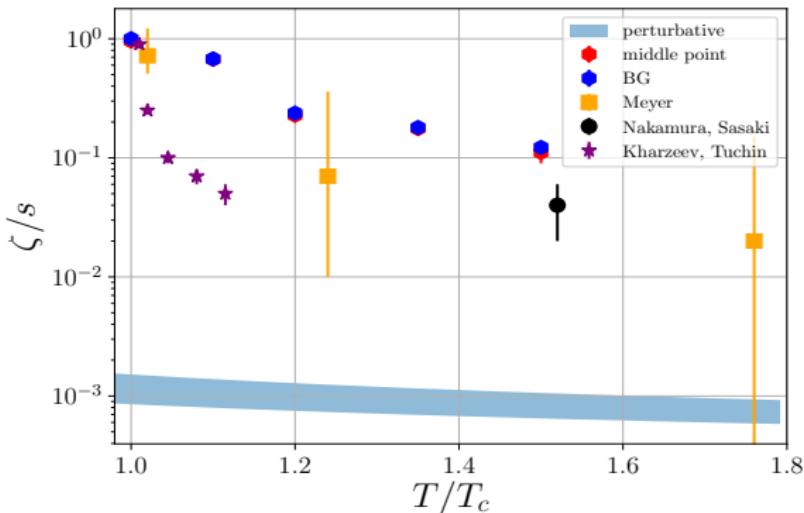
Our results (preliminary!)



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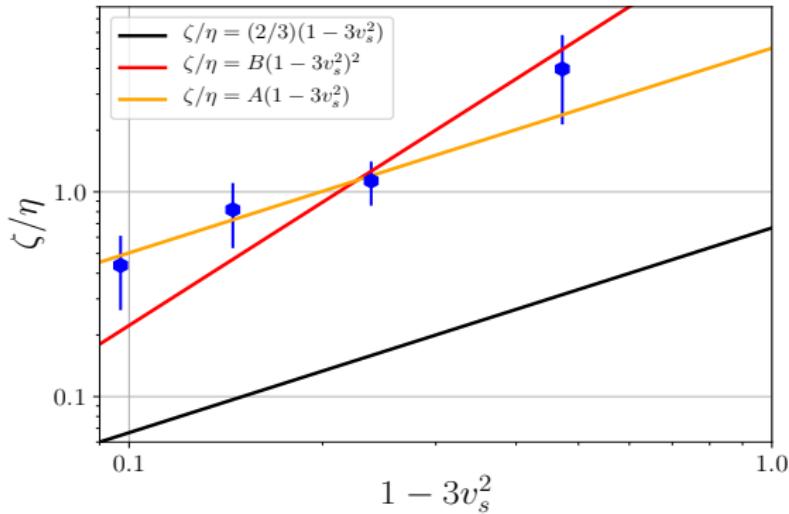


Comparison with other approaches

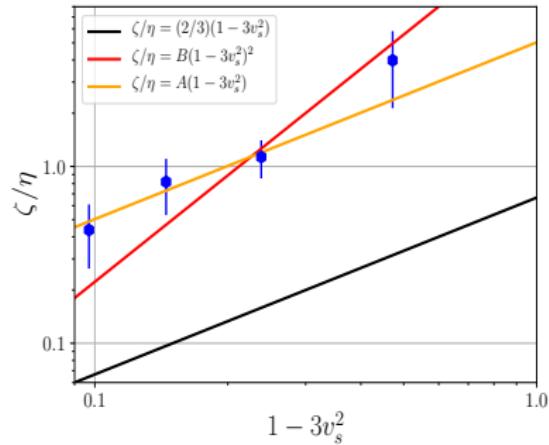
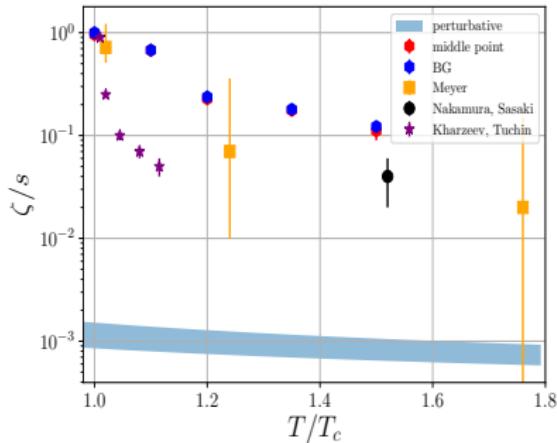


- Agreement with other lattice studies
- Large deviation from perturbative results

Weakly or strongly coupled QGP?



- Weakly coupled system $\zeta/\eta \sim (1 - 3v_s^2)^2$ ($\chi^2/dof \sim 4$)
- Strongly coupled system $\zeta/\eta \sim (1 - 3v_s^2)$ ($\chi^2/dof \sim 1$)
- $\zeta/\eta \geq \frac{2}{3}(1 - 3v_s^2)$ (A. Buchel, Physics Letters B663, 286 (2008))



Conclusion:

- We calculated ζ/s for set of temperatures $T/T_c \in (0.9, 1.5)$
- Agreement with previous lattice results
- Large deviation from perturbative calculation
- QGP reveals the properties of strongly coupled system