

# Numerical algorithm for optimization of positive electrode in lead-acid batteries

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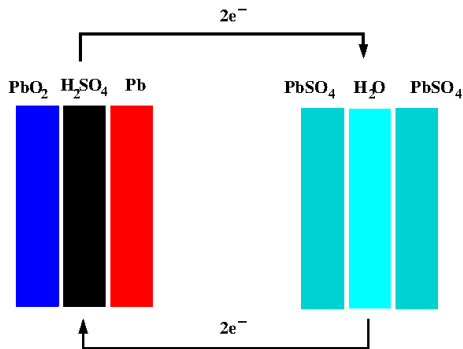
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# Lead-acid battery



- Fully charged: 2.14 V
- Reactions are reversible up to  $\approx 1.75$  V

# Electrode's structure

Electrode = metallic (Pb) support & porous active mass

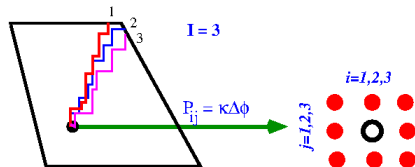


# The model

- hopping model: electrons are localized to specific sites; at each time step we have a probability of jump between current position and neighbours
- the jump probability  $\mathcal{P}_n$  is directly proportional to the potential gradient in the given point

$$\mathcal{P}_n = k \nabla \phi(\vec{r}) \quad (1)$$

- $n$  is the temporal index,  $k$  is a random number with uniform distribution between 0 and 1 and  $\nabla \phi(\vec{r})$  the gradient of the electric potential.



### What to do:

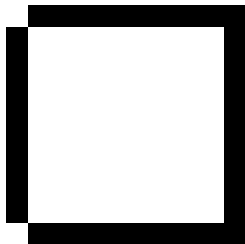
- Image analysis
- Poisson equation solver
- Time propagation
- Data analysis

### How to do:

- Use of the XPM format
- Relaxation method on grid
- Monte-Carlo subroutine
- Average values for pellets

## Software: handling the XPM format

```
/* XPM */
static char *noname[] = {
/* width height ncolors chars_per_pixel */
"11 11 2 1",
/* colors */
" c black",
". c white",
/* pixels */
".           ",
" .....  ",
" .....  ",
" .....  ",
" .....  ",
" .....  ",
" .....  ",
etc
};
```





## Mathematical model: potential during discharge

The potential distribution is extracted using the continuity equation for the electric density during the discharge

$$\vec{\nabla} j = -\frac{\partial \rho(\vec{r})}{\partial t} \quad (2)$$

$j$  is the current density while  $\rho$  is the charge density. Ohm's law in differential form:

$$j = \sigma \epsilon \epsilon = -\frac{d\phi(\vec{r})}{dn} \quad (3)$$

$\epsilon$  is the intensity of the electric field and  $\sigma$  is lead electric conductivity.

This leads to

$$\sigma \Delta \phi(\vec{r}) = -I \quad (4)$$

where  $I = \frac{\partial \rho}{\partial t}$  is the current generated by charge fluctuation (i.e. electrochemical reaction)

In a plane parallel to the electrode (no charge is created/destroyed) we get

$$\Delta \phi(\vec{r}) = 0 \quad (5)$$

# Mathematical model: potential during discharge

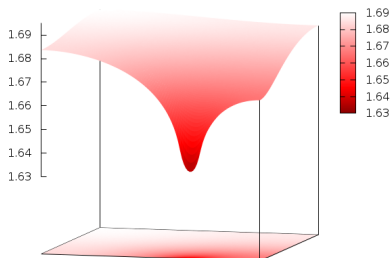
The boundary conditions for the equation are:

$$\frac{d\phi(\vec{r})}{dn} = 0 \quad (6)$$

at any edge points of the electrode, different from the collector, and

$$j_0 = -\sigma \frac{d\phi(\vec{r})}{dn} \quad (7)$$

for the collector region, and where  $\sigma$  is the lead conductivity (i.e.  $\sigma = 4550 \text{ S / mm}$ )



## Mathematical model: Poisson-Laplace equation

The relaxation method:  $\phi(x, y)$  is the potential at the coordinates  $(x, y)$  a Taylor expansion allows us approximate the values

$$\phi(x + h_x, y) = \phi(x, y) + h_x \frac{\partial \phi}{\partial x}(x, y) + \frac{1}{2!} \frac{\partial^2 \phi}{\partial x^2}(x, y) h_x^2 \quad (8)$$

By summing up the similar equation for the  $\phi(x - h_x, y)$  we get:

$$\phi(x + h_x, y) + \phi(x - h_x, y) = 2\phi(x, y) + h_x^2 \frac{\partial^2 \phi}{\partial x^2}(x, y) \quad (9)$$

A similar expression holds for  $\phi(x, y \pm h_y)$ .

## Mathematical model: Poisson-Laplace equation

If we use a grid representation for  $\phi(x, y)$  the  $h_{x/y}$  are defined by grid steps in discrete (grid) form we have

$$\frac{\phi(x + h_x, y) + \phi(x - h_x, y) - 2\phi(x, y)}{h_x^2} + E(y) = 0 \quad (10)$$

Next,  $\phi(x, y)$  is expressed as a function of the values of its neighbors points

$$\phi(x, y) = \beta(\phi(x+h_x, y)+\phi(x-h_x, y)+\alpha(\phi(x, h+h_y)+\phi(x, y-h_y))) \quad (11)$$

where

$$\alpha = \frac{h_x^2}{h_y^2} \beta = \frac{1}{2(1 + \alpha)} \quad (12)$$

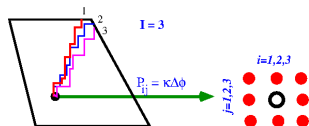
# Mathematical model: Poisson-Laplace equation

The equation:

$$\phi(x, y)^{(n+1)} = (1 - p)\phi(x, y)^{(n)} + p\phi(x, y) \quad (13)$$

is iterated until the convergence is reached. Here  $\phi(x, y)^{(n+1)}$  is the potential at iteration  $n + 1$  while  $\phi(x, y)^{(n)}$  is the potential at iteration  $n$ . Parameter  $p$  controls the convergence, typically taking values between 0.5 and 1.

# Software: Monte Carlo time propagation



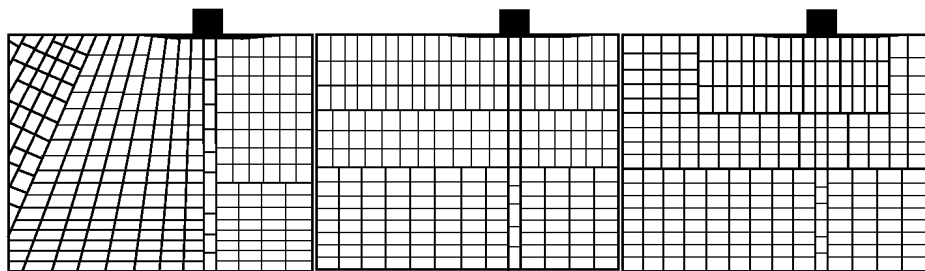
- At each iteration,  $n$  for each point/pelet the charge is propagated with probability  $\mathcal{P}_n$

- 

$$\mathcal{P}_n = k \nabla \phi(\vec{r}) \quad (14)$$

- $k$  is a random number with constant distribution, ranging from 0 to 1 and  $\nabla \phi(\vec{r})$  is the gradient of the potential.
- if the gradient is negative, no jump takes place
- if the metallic part is reached, we add  $+1$  at total charge generate by the pellet.

## Geometrical models



**Figure:** Geometric shapes of the grids studied: from left to right we note them as G1, G2 and G3.

- All models were drawn using Xfig (vectorial format fig, export to xpm)
- The same size for the electrode is used in all cases
- The same potential file is used for propagation
- Results are given in pixels - for the electrode drawings and iterative cycles, respectively

# Time to collect charge

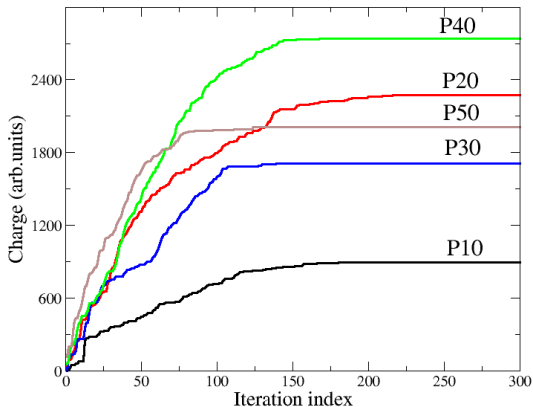
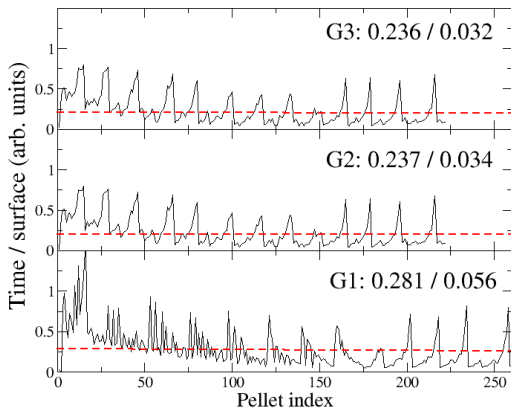


Figure: Example: time dependence of the current for selected pellets in G1.

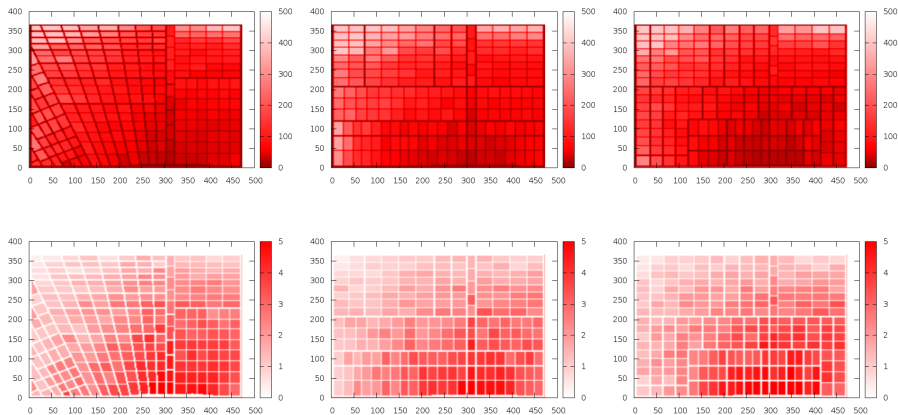


# Time to collect charge



**Figure:** Time needed to collect the charge for grids G1, G2 and G3: average and MSQ for all pellets

# Pellet usage at discharge



**Figure:** Top: time needed to collect the charge for grids G1, G2 and G3. Bottom: Percent of usage for each pellet after 1/10 of the discharge time for grids

# Summary

- We propose an algorithm to analyse the quality of the lead-acid collector in the positive electrode of lead-acid battery
- The idea is to use a hopping model and a Monte-Carlo procedure to analyse the system's state during battery discharge
- A dedicated software was developed to implement all the features
- This includes image analysis, Poisson solver and Monte Carlo time propagation of the charge generated by electrochemical reaction.
- Data analysis is complex
- A "qualitative index" may be assigned to each support of the positive electrode

Thank you!