

# Chiral heat-vortical wave in cold Fermi liquid and the deformation of zero sound

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# The plan

- Anomalous transport effects
- Chiral waves
- Chiral kinetic theory formalism for Fermi liquid
- Results

- blah-blah

- Anomalous transport effects should reveal themselves in heavy ion collisions

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- Namely, chiral waves should cause generation of quadrupole moment in quark-gluon plasma

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# Anomalous transport effects

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- Chiral vortical effect

$$\mathbf{j}_A = \left[ \frac{(\mu_V^2 + \mu_A^2)}{2\pi^2} + \frac{T^2}{6} \right] \boldsymbol{\Omega}$$

$$\mathbf{j}_V = \frac{\mu_V \mu_A}{\pi^2} \boldsymbol{\Omega}$$



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$$\omega \chi \delta \mu_{V/A} + \frac{(\mathbf{k} \cdot \mathbf{B}) \delta \mu_{A/V}}{2\pi^2} = 0$$

## Dispersion relation

$$v_{CMW} = \frac{B}{2\pi^2 \chi}$$

# Chiral vortical wave

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- We need finite vector charge density in the background

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## Dispersion relation

$$v_{CVW} = \frac{\mu_V \Omega}{2\pi^2 \chi}$$

- Anomalous energy transport due to rotation:

$$\mathbf{j}_E = \frac{\mu_A}{3} \left[ \frac{3\mu_V^2 + \mu_A^2}{\pi^2} + T^2 \right] \boldsymbol{\Omega}$$



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- For small fluctuations of vector charge and energy density above equilibrium

$$\delta \rho_V = \chi \delta \mu_V + \alpha \delta T$$

$$\delta \epsilon = C \delta T + \gamma \delta \mu_V$$

## Dispersion relation

$$v_{CHW} = \sqrt{\frac{T^3}{C\chi - \alpha\gamma} \frac{\Omega}{3}}$$

# Mixed chiral heat-vortical wave

- Non-zero background charge density
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## Dispersion relation

$$v_{CHVW} = \sqrt{\frac{\Omega^2 |9\mu_V^2(C - \alpha\mu_V) - 3\mu_V T \pi^2(\alpha T + \gamma - \chi\mu_V) + \pi^2 \chi T^3|}{9\pi^4 \chi |C\chi - \alpha\gamma|}}$$

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# Berry Fermi liquid

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- Not far from equilibrium is described in terms of quasiparticles
- The anomalous effects are captured by Berry connection in momentum space



$$\frac{\partial n_{R/L}}{\partial t} + \dot{\mathbf{x}}_{R/L} \cdot \frac{\partial n_{R/L}}{\partial \mathbf{x}} + \dot{\mathbf{p}}_{R/L} \cdot \frac{\partial n_{R/L}}{\partial \mathbf{p}} = C_{R/L}[n_R, n_L]$$

- $n_{R/L}(t, \mathbf{x}, \mathbf{p})$  are right and left quasiparticles distribution functions
- $C_{R/L}$  are respective collision integrals

$$\begin{aligned}\sqrt{G_{R/L}}\dot{\mathbf{x}}_{R/L} &= \mathbf{v}_{R/L} + 2\epsilon_{R/L}\boldsymbol{\Omega}(\mathbf{v}_{R/L} \cdot \mathbf{b}_{R/L}) + \boldsymbol{\mathcal{E}}_{R/L} \times \mathbf{b}_{R/L} \\ \sqrt{G_{R/L}}\dot{\mathbf{p}}_{R/L} &= \boldsymbol{\mathcal{E}}_{R/L} + 2\epsilon_{R/L}\mathbf{v}_{R/L} \times \boldsymbol{\Omega} + (\boldsymbol{\mathcal{E}}_{R/L} \cdot \boldsymbol{\Omega})2\epsilon_{R/L}\mathbf{b}_{R/L}\end{aligned}$$

- Here  $\epsilon_{R/L}$  are quasiparticles energy functionals,
- $\mathbf{v}_{R/L} = \frac{\partial\epsilon_{R/L}}{\partial\mathbf{p}}$ ,  $\boldsymbol{\mathcal{E}}_{R/L} = -\frac{\partial\epsilon_{R/L}}{\partial\mathbf{x}}$
- $\mathbf{b}_{R/L} = \pm\frac{\hat{\mathbf{p}}}{2p^3}$  are Berry connections curvature in momentum space ( $p = |\mathbf{p}|$ )
- $\sqrt{G_{R/L}} = 1 + 2\epsilon_{R/L}(\mathbf{b}_{R/L} \cdot \boldsymbol{\Omega})$  modify phase space volume

- We are considering the small fluctuations above the equilibrium distribution functions:

$$n_{R/L} = n_{R/L}^0 + \delta n_{R/L} e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)}$$

where

$$n_{R/L}^0 = \frac{1}{e^{\beta(\epsilon^0 - \mu)} + 1}$$

- Here  $\beta = T^{-1}$  and we assume  $\mu \gg T$

# Collective excitations in kinetic theory formalism

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where

$$n_{R/L}^0 = \frac{1}{e^{\beta(\epsilon^0 - \mu)} + 1}$$

- Here  $\beta = T^{-1}$  and we assume  $\mu \gg T$
- Two regimes:  $\omega \ll \tau^{-1}$  (hydrodynamic regime)  $\omega \gg \tau^{-1}$  (collisionless regime). Here  $\tau$  is characteristic relaxation time of collision integral

## Dispersion relation

$$\omega = \pm \frac{\mu(\mathbf{k} \cdot \boldsymbol{\Omega})}{2\pi^2\chi} \sqrt{F_1 F_2} \left[ 1 + \frac{T^2 \pi^2}{6\mu^2 F_1} (1 - 4A + 4A^2) \right]$$

- $A = \frac{\mu}{v_{FPF}}$
- $F_1$  and  $F_2$  are taken from linearised energy functionals

## Non-modified implicit zero sound dispersion relation

$$\operatorname{arccotanh} s_0 = \frac{1}{s_0} \left( \frac{1}{2(F_S \pm F_A)} + 1 \right)$$

- Here  $s_0 = \frac{\omega}{v_f k}$ ,  $F_S$  and  $F_A$  are taken from linearised energy functionals

## Modification due to the chiral heat-vortical wave

$$\delta s = s - s_0 \approx \mp \frac{\omega^2 \mu^2}{p_F^4} \left( 1 - \frac{2\pi^2 T^2}{3v_F^2 p_F^2} \right).$$

$$\frac{F_S L_2(s_0) + [L_1(s_0)^2 - 2L_0(s_0)L_2(s_0)](F_S^2 - F_A^2)}{F_A [\text{arccotanh } s_0 - \frac{s_0}{2((s_0)^2 - 1)}]}$$

$$L_0(s_0) = s_0 \text{arccotanh } s_0 - 1$$

$$L_1(s_0) = 3s_0(s_0 \text{arccotanh } s_0 - 1)$$

$$L_2(s_0) = 2s_0[-3s_0 + (3s_0^2 - 1) \text{arccotanh } s_0]$$

- The result for velocity of chiral heat-vortical wave coincides with the one known from hydrodynamics



- The result for velocity of chiral heat-vortical wave coincides with the one known from hydrodynamics
- There turn out to be two branches of modified zero sound with the correction to velocity being quadratic in angular velocity

THANK YOU FOR YOUR ATTENTION