

Single Spin Asymmetries in a light-front quark-diquark model

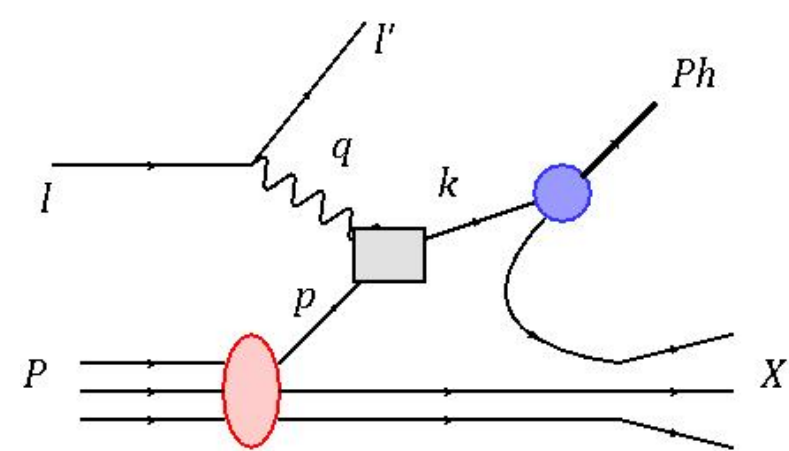
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Introduction

- In past few decades several experiment collaborations, e.g., HERMES, COMPASS, JLAB etc., have come up with one of the most exciting features of the nucleons: Single spin asymmetry(SSA).
- These asymmetries indicate existence of non-vanishing transverse momentum of interior partons and collinear picture is no longer sufficient to describe the transverse structure of nucleons.
- The information of asymmetries are encoded into the transverse momentum dependent distributions(TMDs) and fragmentation functions(FFs).
- We calculate the leading twist TMDs in a light-front quark-diquark model(LFQDM) with the wave functions adopted form soft-wall AdS/QCD and give model predictions to the single spin asymmetries. Particularly, here we present model result for Collins asymmetry $A_{UT}^{\sin(\phi_h+\phi_S)}(x, z, \mathbf{P}_{h\perp}, y)$ of proton in SIDIS($\ell P \rightarrow \ell' h X$) for the π^+ and π^- channels.

$$\ell(\ell) + N(P) \rightarrow \ell(\ell') + h(P_h) + X$$



At $P_{h\perp}^2 \simeq \Lambda_{QCD}^2 \ll Q^2$, the SIDIS cross-section is factorised as

$$d\sigma^{\ell N \rightarrow \ell' h X} = \sum_{\nu} \hat{f}_{\nu/P}(x, \mathbf{p}_{\perp}^2; Q^2) \otimes d\hat{\sigma}^{\ell q \rightarrow \ell' q} \otimes \hat{D}_{h/\nu}(z, \mathbf{k}_{\perp}^2; Q^2),$$

where the first term represents TMDs, the second term represents the hard scattering and the third term is for fragmentation functions(FFs).

- In the experimental measurement the azimuthal asymmetries are found to be non-zero and are defined as

$$A_{S_t S_N} = \frac{d\sigma^{\ell P \uparrow \rightarrow \ell' h X} - d\sigma^{\ell P \downarrow \rightarrow \ell' h X}}{d\sigma^{\ell P \uparrow \rightarrow \ell' h X} + d\sigma^{\ell P \downarrow \rightarrow \ell' h X}} \neq 0.$$

- The weighted asymmetry is projected out for the weight \mathcal{W} as

$$A_{S_t S_N}^{\mathcal{W}}(x, z, \mathbf{P}_{h\perp}, y) = 2 \frac{\int d\phi_h d\phi_S [d\sigma^{\ell P \uparrow \rightarrow \ell' h X} - d\sigma^{\ell P \downarrow \rightarrow \ell' h X}] \mathcal{W}}{\int d\phi_h d\phi_S [d\sigma^{\ell P \uparrow \rightarrow \ell' h X} + d\sigma^{\ell P \downarrow \rightarrow \ell' h X}]}$$

- Where each weighted asymmetry has contribution from particular TMD and particular FF. For example:

$$\text{Sivers Asymmetry : } A_{UT}^{\sin(\phi_h - \phi_S)} \sim f_{1T}^{\perp} \otimes D_1 ; f_{1T}^{\perp} : \text{Sivers func.}$$

$$\text{Collins Asymmetry : } A_{UT}^{\sin(\phi_h + \phi_S)} \sim h_1^{\perp} \otimes H_1^{\perp} ; H_1^{\perp} : \text{Collins func.}$$

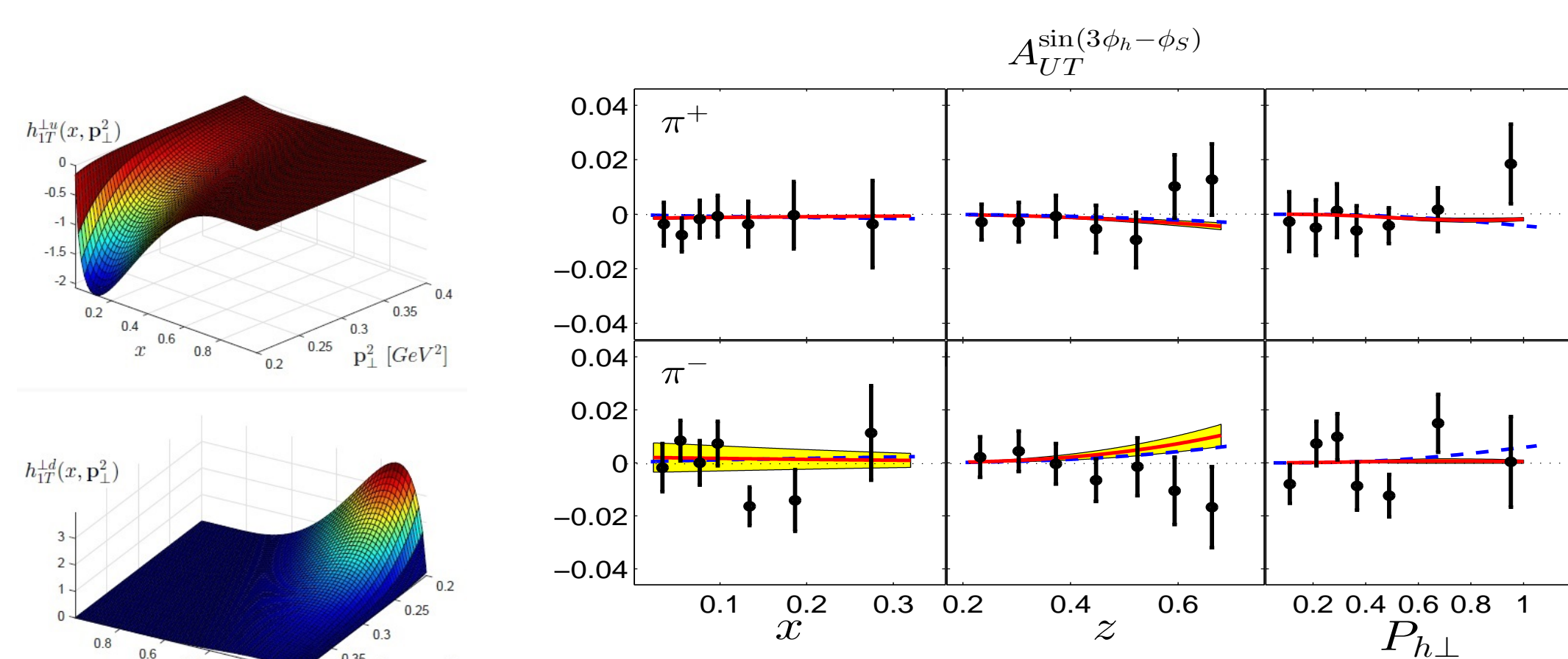
$$A_{UT}^{\sin(3\phi_h - \phi_S)} \sim h_{1T}^{\perp} \otimes H_1^{\perp}$$

$$A_{UL}^{\sin \phi_h} \sim h_{1L}^{\perp} \otimes H_1^{\perp} \dots$$

D_1 is the unpolarised FFs and $h_1, h_{1T}^{\perp}, h_{1L}^{\perp}$ are the three T-even TMDs.

Model result for $A_{UT}^{\sin(3\phi_h - \phi_S)}$

$A_{UT}^{\sin(3\phi_h - \phi_S)}$ is related to the pretzelosity TMD, h_{1T}^{\perp} . Our model results of $A_{UT}^{\sin(3\phi_h - \phi_S)}$ for the π^+ and π^- channels are compared with the HERMES data



Light-front quark-diquark Model(LFQDM)

- In this model proton is considered as a bound state of a quark and a diquark(of spin-0, singlet or spin-1,triplet) with effective mass and the proton state is written in spin-flavor SU(4) structure as

$$|P; \pm\rangle = C_S |u S^0\rangle^{\pm} + C_V |u A^0\rangle^{\pm} + C_{VV} |d A^1\rangle^{\pm}$$

- The two particle Fock-state expansion for $J^z = \pm 1/2$

$$|u S\rangle^{\pm} = \int \frac{dx d^2\mathbf{p}_{\perp}}{2(2\pi)^3 \sqrt{x(1-x)}} \sum_{\lambda} \psi_{\lambda}^{\pm(u)}(x, \mathbf{p}_{\perp}) |\lambda \Lambda_S; x P^+, \mathbf{p}_{\perp}\rangle \Big|_{\Lambda_S=0}$$

$$|\nu A\rangle^{\pm} = \int \frac{dx d^2\mathbf{p}_{\perp}}{2(2\pi)^3 \sqrt{x(1-x)}} \sum_{\lambda} \sum_{\Lambda_A} \psi_{\lambda\Lambda_A}^{\pm(\nu)}(x, \mathbf{p}_{\perp}) |\lambda \Lambda_A; x P^+, \mathbf{p}_{\perp}\rangle \Big|_{\Lambda_A=1,0,-1}$$

- The light-front wave functions:

$$\psi_{\lambda\Lambda}^{\pm(\nu)}(x, \mathbf{p}_{\perp}) = N^{\nu} f(x, \mathbf{p}_{\perp}, \lambda, \Lambda) \varphi_i^{\pm(\nu)}(x, \mathbf{p}_{\perp}) \Big|_{i=1,2}$$

- Modified soft-wall AdS/QCD wave function for two particle bound state:

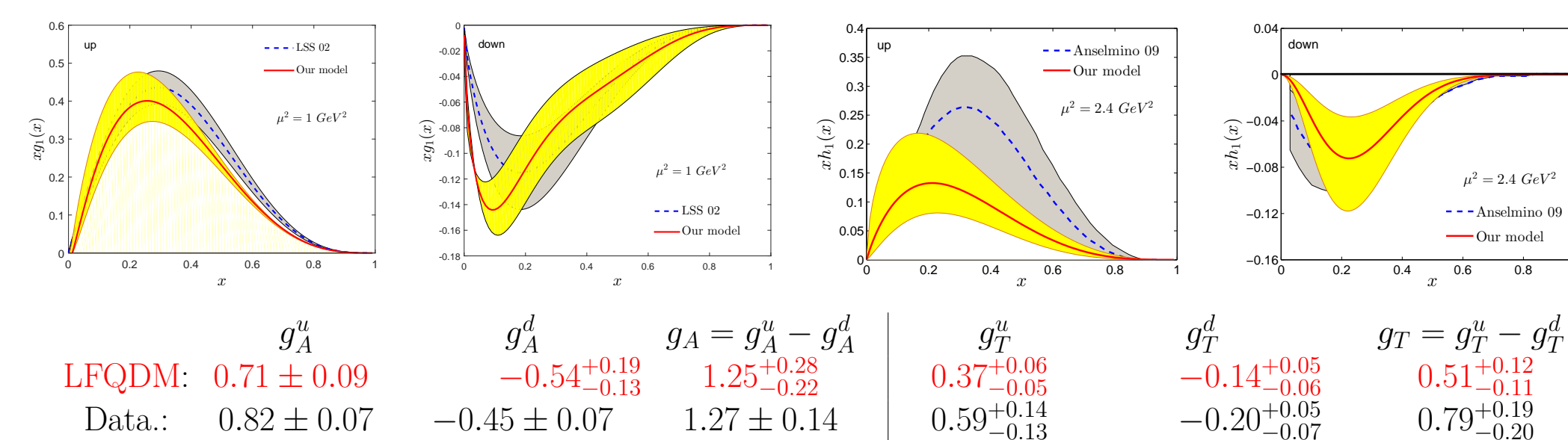
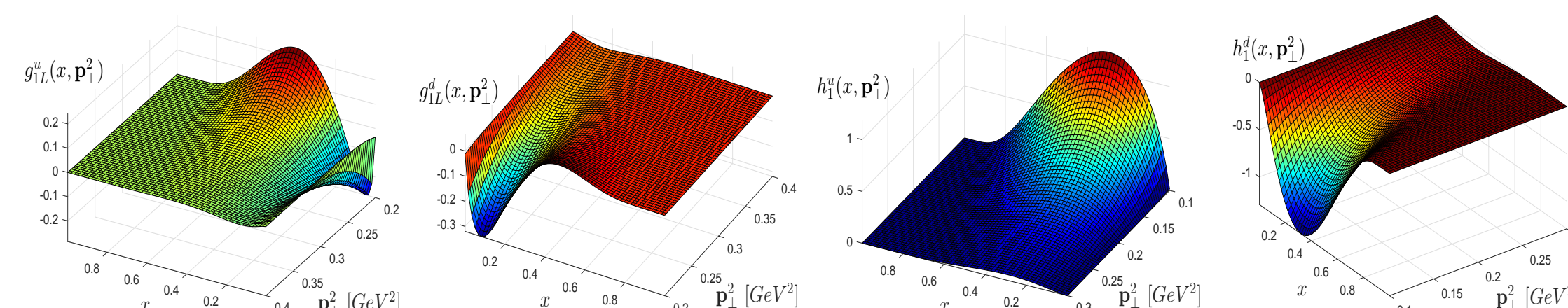
$$\varphi_i^{\pm(\nu)}(x, \mathbf{p}_{\perp}) = \frac{4\pi}{\kappa} \sqrt{\frac{\log(1/x)}{1-x}} x^{a_i^{\nu}} (1-x)^{b_i^{\nu}} \exp\left[-\delta^{\nu} \frac{\mathbf{p}_{\perp}^2 \log(1/x)}{2\kappa^2(1-x)^2}\right].$$

- We determine the parameters $a_i^{\nu}, b_i^{\nu}, \delta^{\nu}$ by fitting the experimental data of the Dirac $F_1(Q^2)$ and Pauli $F_2(Q^2)$ form factors.

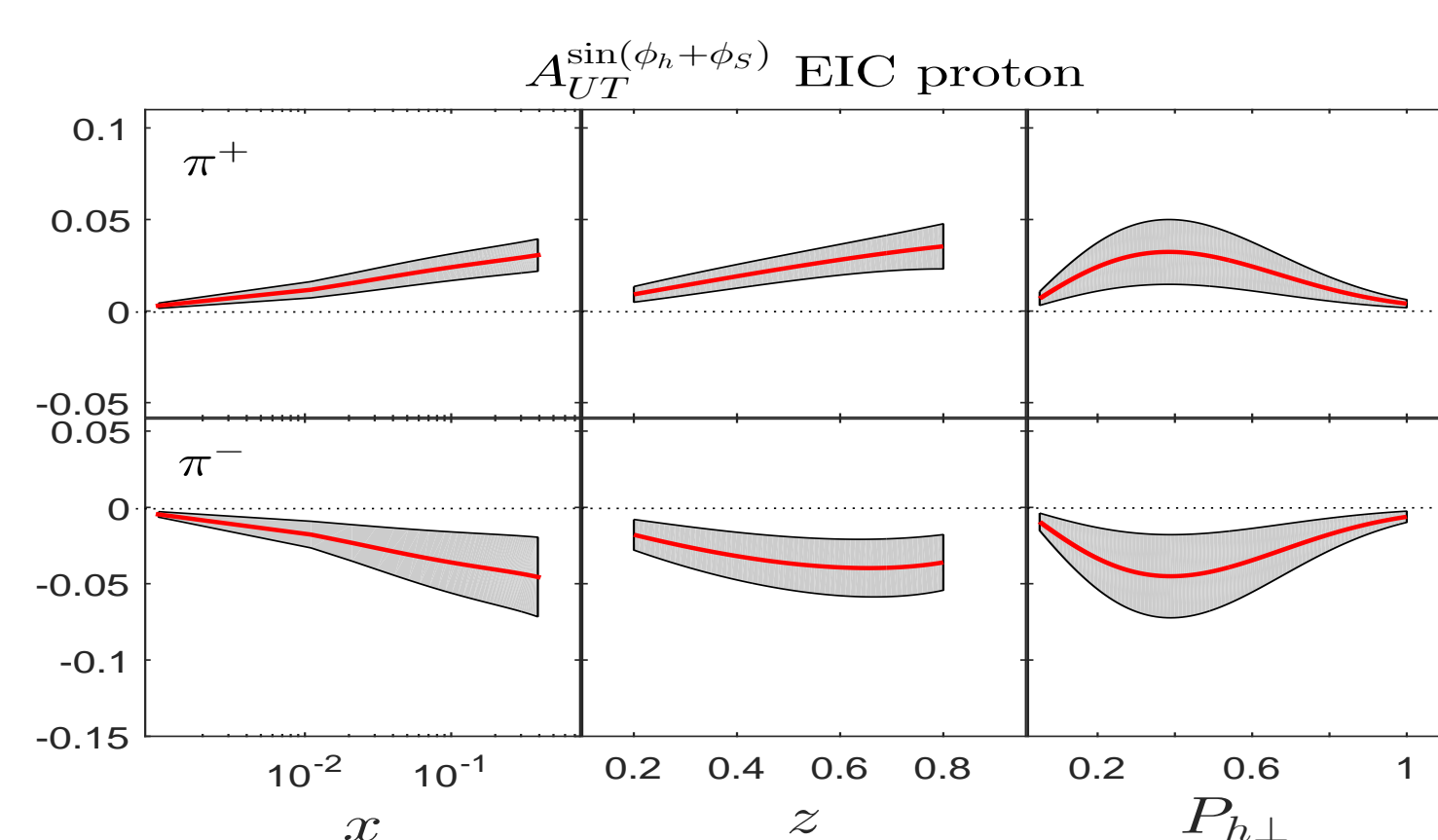
Helicity and Transversity

In the SIDIS process, the TMDs correlator is defined as (at leading twist $\Gamma = \gamma^+, \gamma^+ \gamma^5$ and $i\sigma^j + \gamma^5$)

$$\Phi^{\nu[\Gamma]}(x, \mathbf{p}_{\perp}; S) = \frac{1}{2} \int \frac{dz^- d^2z_T}{2(2\pi)^3} e^{ip \cdot z} \langle P; S | \bar{\psi}^{\nu}(0) \Gamma \mathcal{W}_{[0,z]} \psi^{\nu}(z) | P; S \rangle$$



Prediction for EIC kinematics



$$0.001 < x < 0.4, \quad 0.2 < z < 0.8, \quad 0.05 < P_{h\perp} < 1, \quad 0.01 < y < 0.95$$

$$\sqrt{s} = 45 \text{ GeV} \quad \mu^2 = 100 \text{ GeV}^2$$

Collins Asymmetry

- Collins asymmetry provides correlation between the transversely polarised quark in a transversely polarised nucleon and transverse momentum of the produced hadron.
- Collins asymmetry (has contribution from transversity TMDs, $h_1^{\nu}(x, \mathbf{p}_{\perp})$) is defined as

$$A_{UT}^{\sin(\phi_h + \phi_S)}(x, z, \mathbf{P}_{h\perp}, y) = \frac{4\pi^2 \alpha^2 (1-y) C}{sxy^2} \left[\frac{P_{h\perp} - z(\hat{\mathbf{P}}_{h\perp} \cdot \mathbf{p}_{\perp})}{zM_h} h_1^{\nu}(x, \mathbf{p}_{\perp}^2) H_1^{\perp\nu}(z, \mathbf{k}_{\perp}) \right]$$

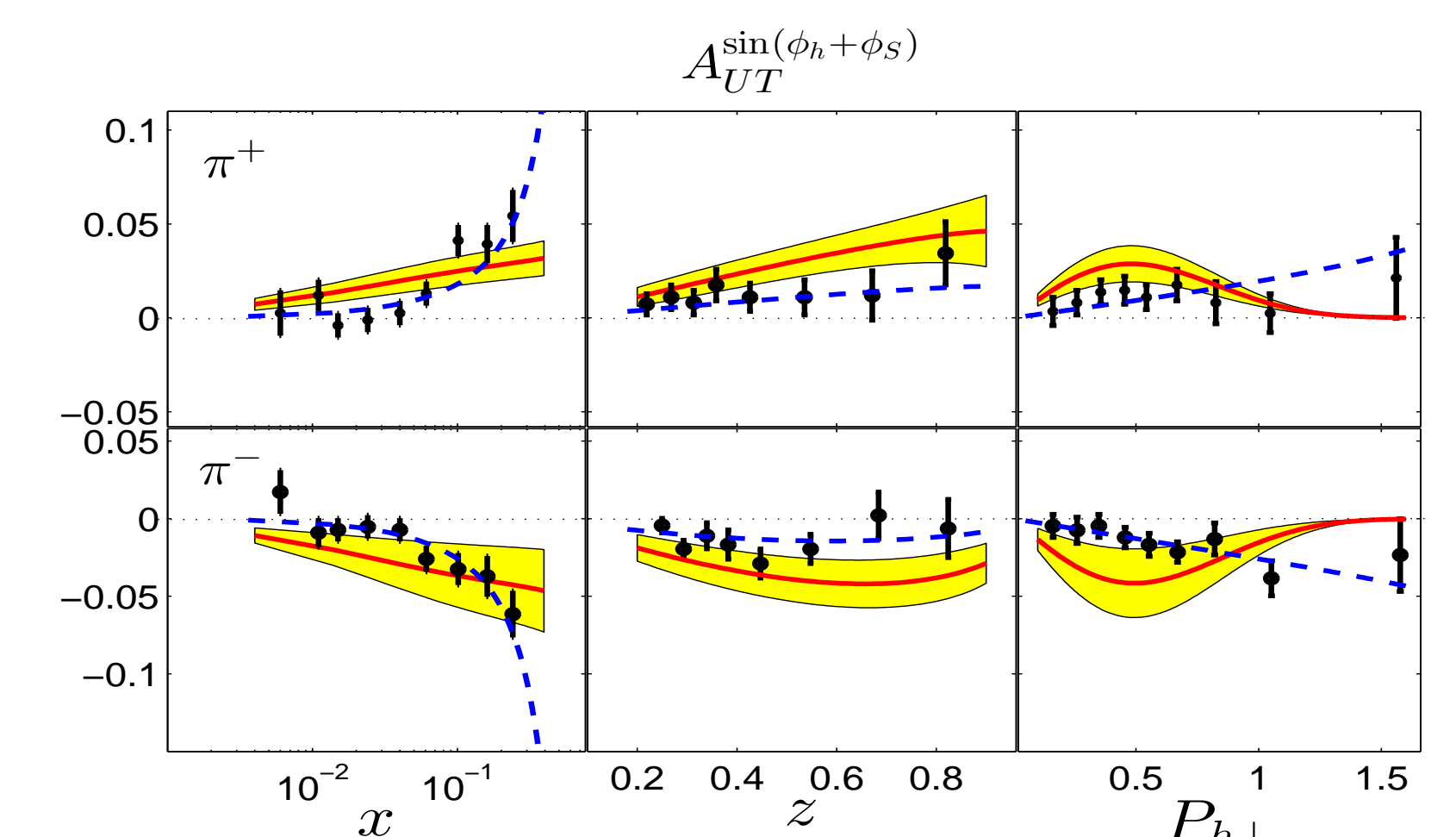
$$= \frac{2\pi^2 \alpha^2 (1-y)^2 C [f_1^{\nu}(x, \mathbf{p}_{\perp}^2) D_1^{h/\nu}(z, \mathbf{k}_{\perp})]}{sxy^2}$$

- In the LFQDM:

$$A_{UT}^{\sin(\phi_h + \phi_S)} = \frac{\frac{2(1-y) P_{h\perp} \sqrt{2e}}{sxy^2} \sum_{\nu} e_{\nu}^2 \hat{h}_1^{\nu}(x)}{\frac{1+(1-y)^2}{sxy^2} \sum_{\nu} e_{\nu}^2 N_{f_1} \frac{\ln(1/x)}{\pi\kappa^2} \left[T_1^{\nu}(x) - \frac{\langle m_{\perp}^2 \rangle}{M^2} T_2^{\nu}(x) \right]}$$

$$\times \frac{\mathcal{N}_{\nu}^C(z) D_1^{h/\nu}(z) \frac{\langle k_{\perp}^2 \rangle_C \langle p_{\perp}^2 \rangle_x e^{-\mathbf{P}_{h\perp}^2 / (P^2_{h\perp}) C}}{\langle k_{\perp}^2 \rangle_C \langle P^2_{h\perp} \rangle_C} \frac{\langle P^2_{h\perp} \rangle_C}{D_1^{h/\nu}(z) \langle p_{\perp}^2 \rangle_x \frac{e^{-\mathbf{P}_{h\perp}^2 / (P^2_{h\perp})}}{\langle P^2_{h\perp} \rangle_C}}}{}$$

The Collins asymmetry in this model is shown below and compared with the COMPASS data.



Where, "-" represent QCD evolution of $f_1(x, \mathbf{p}_{\perp}^2)$, and "- - -" is for parameter evolution of $f_1(x, \mathbf{p}_{\perp}^2)$.

- A complete QCD evolutions of all the leading twist TMDs are not known. h_1^{ν} is kept at initial scale.
- Parameter evolution approach assumes that the evolution information encoded into the parameters(PRD94,094020(2016)) of this model generates the TMD evolution.
- The fragmentation functions $H_1^{\perp\nu}(z, \mathbf{k}_{\perp})$ and $D_1^{h/\nu}(z, \mathbf{k}_{\perp})$ are taken as phenomenological inputs.

Conclusions and future directions

- The light-front quark-diquark model(LFQDM) inspired by soft-wall AdS/QCD predicts the TMDs whose PDF limits satisfy the phenomenological extractions quite well.
- Model prediction to the Collins asymmetry and $A_{UT}^{\sin(3\phi_h - \phi_S)}$ show well agreement with the COMPASS and HERMES data.
- Work in progress to predict Sivers asymmetry in this model.

References

- T. Maji, D.Chakrabarti, Phys. Rev. D94,094020(2016)
- R. Jakob, p.j. Mulders, j. Rodrigues Nucl. Phys. A 626(1997)937
- G. F. de Teramond and S. J. Brodsky, arXiv:1203.4025 [hep-ph].
- S. M. Aybat, T. C. Rogers, PRD83, 114042 (2011); S. M. Aybat, J. C. Collins et.al. PRD85, 034043 (2012); M. Anselmino et.al. PRD86, 014028 (2012)
- E. Leader et.al., Phys. Rev. D 67 (2003) 037503
- A. Martin [COMPASS Collaboration], Phys. Part. Nucl. 45, 141 (2014)
- G. Schnell [HERMES Collaboration], PoS DIS 2010, 247 (2010)

