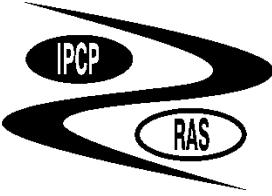


# The dipolar relaxation of multiple quantum coherences as a model for an investigation of decoherence processes in many-qubit clusters in multiple-quantum NMR

**G.A. Bochkin, E.B. Fel'dman, S.G. Vasil'ev**

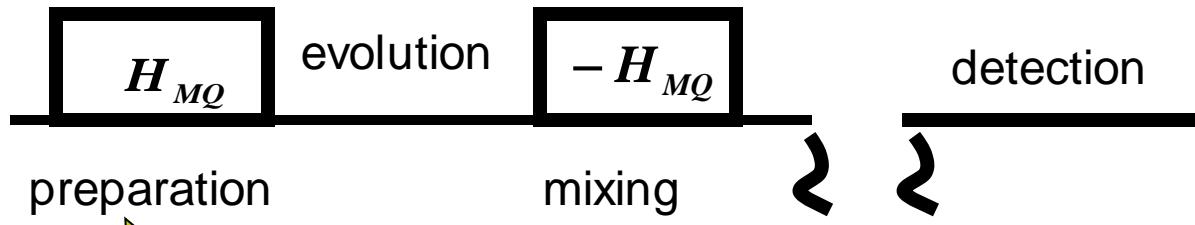
*Institute of Problems of Chemical Physics of Russian Academy of Sciences*



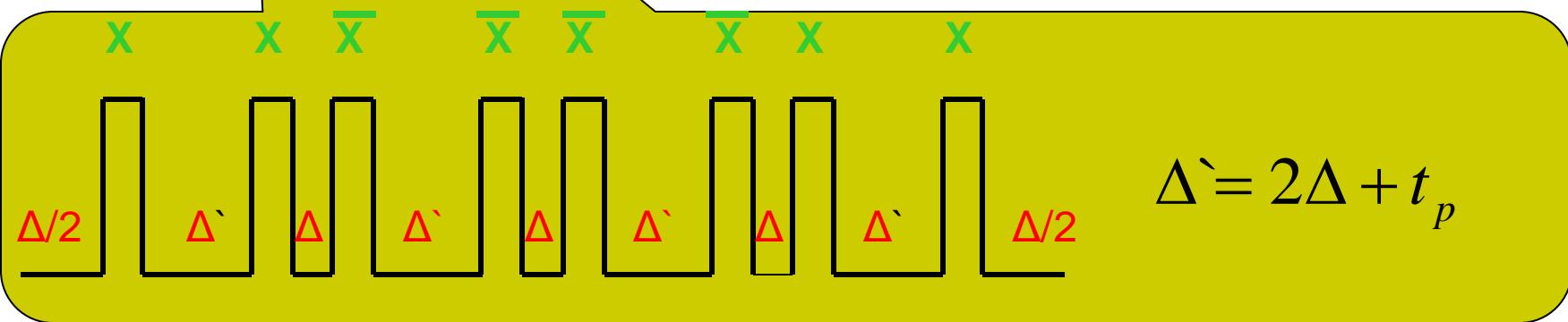


# MQ NMR EXPERIMENT IN SOLIDS

$$\rho_2 \sim I_{1z} I_2^+ I_3^+ I_{4z} I_5^+ I_6^-$$

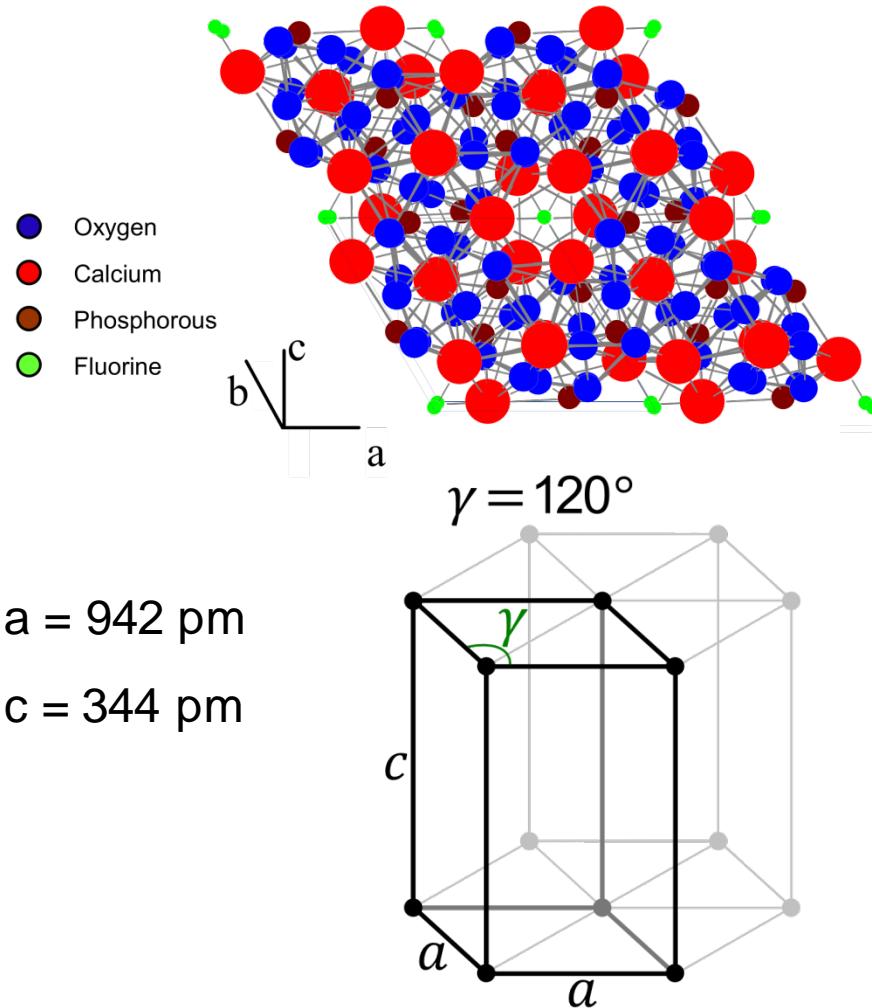


J.Baum, M.Munovitz, A.N.Garroway, A.Pines,  
J.Chem.Phys.83, 2015 (1985)

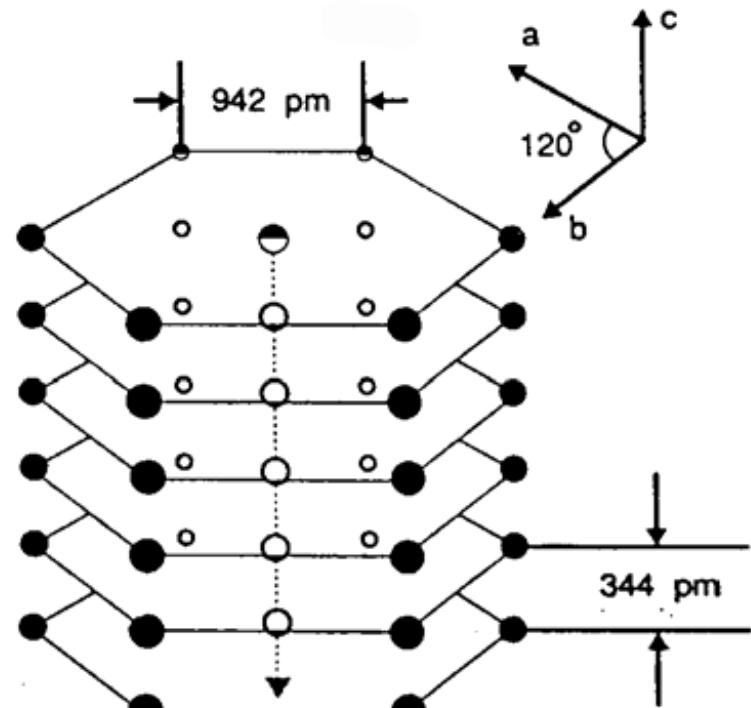


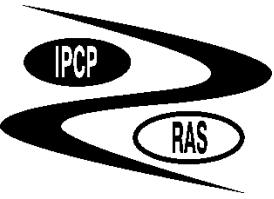
# One-dimensional spin chains

Fluorapatite –  $\text{Ca}_5(\text{PO}_4)_3\text{F}$



$\text{Ca}_5(\text{PO}_4)_3\text{F}$





# EXACT SOLUTION FOR MQ NMR DYNAMICS ON THE PREPARATION PERIOD

$$\sigma(\tau) = \sigma_0(\tau) + \sigma_2(\tau) + \sigma_{-2}(\tau)$$

$$\sigma_0(\tau) = \frac{1}{2} \sum_k \cos[2D\tau \sin k] (1 - a_k^+ a_k^-)$$

$$N \gg 1$$

$$k = \frac{2\pi n}{N} \left( n = -\frac{N}{2}, -\frac{N}{2} + 1, \dots, \frac{N}{2} - 1 \right)$$

$$\sigma_2(\tau) = -\frac{1}{2} \sum_k \sin[2D\tau \sin k] a_k a_{-k}$$

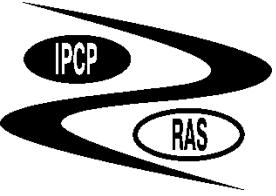
$$\sigma_{-2}(\tau) = \frac{1}{2} \sum_k \sin[2D\tau \sin k] a_k^+ a_{-k}^+$$

$$a_k = \frac{1}{\sqrt{N}} \sum_m \Psi_m e^{-ikm}$$

$$a_k^+ = \frac{1}{\sqrt{N}} \sum_m \Psi_m^+ e^{ikm}$$

Jordan-Wigner transformation

$$\Psi_m = 2^{m-1} I_1^z I_2^z \dots I_{m-1}^z I_m^+; \quad \Psi_m^+ = 2^{m-1} I_1^z I_2^z \dots I_{m-1}^z I_m^-$$



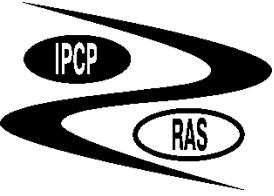
# ZZ MODEL FOR RELAXATION OF MQ NMR COHERENCES

$$H_{dz} = \sum_{i < j} D_{ij} (3I_i^z I_j^z - \vec{I}_i \vec{I}_j) = \sum_{i < j} D_{ij} (2I_i^z I_j^z - I_i^x I_j^x - I_i^y I_j^y)$$

$$H_{zz} = 2 \sum_{i < j} D_{ij} I_i^z I_j^z = \sum_{i \neq j} D_{ij} I_i^z I_j^z$$

$$F_0(\tau, t) = \frac{\text{Tr}\left[e^{-iH_{zz}t} \sigma_0(t) e^{iH_{zz}t} \sigma_0(t)\right]}{\text{Tr}(I_z^2)}$$

$$F_{\pm 2}(\tau, t) = \frac{\text{Tr}\left[e^{-iH_{zz}t} \sigma_2(t) e^{iH_{zz}t} \sigma_{-2}(t)\right]}{\text{Tr}(I_z^2)}$$

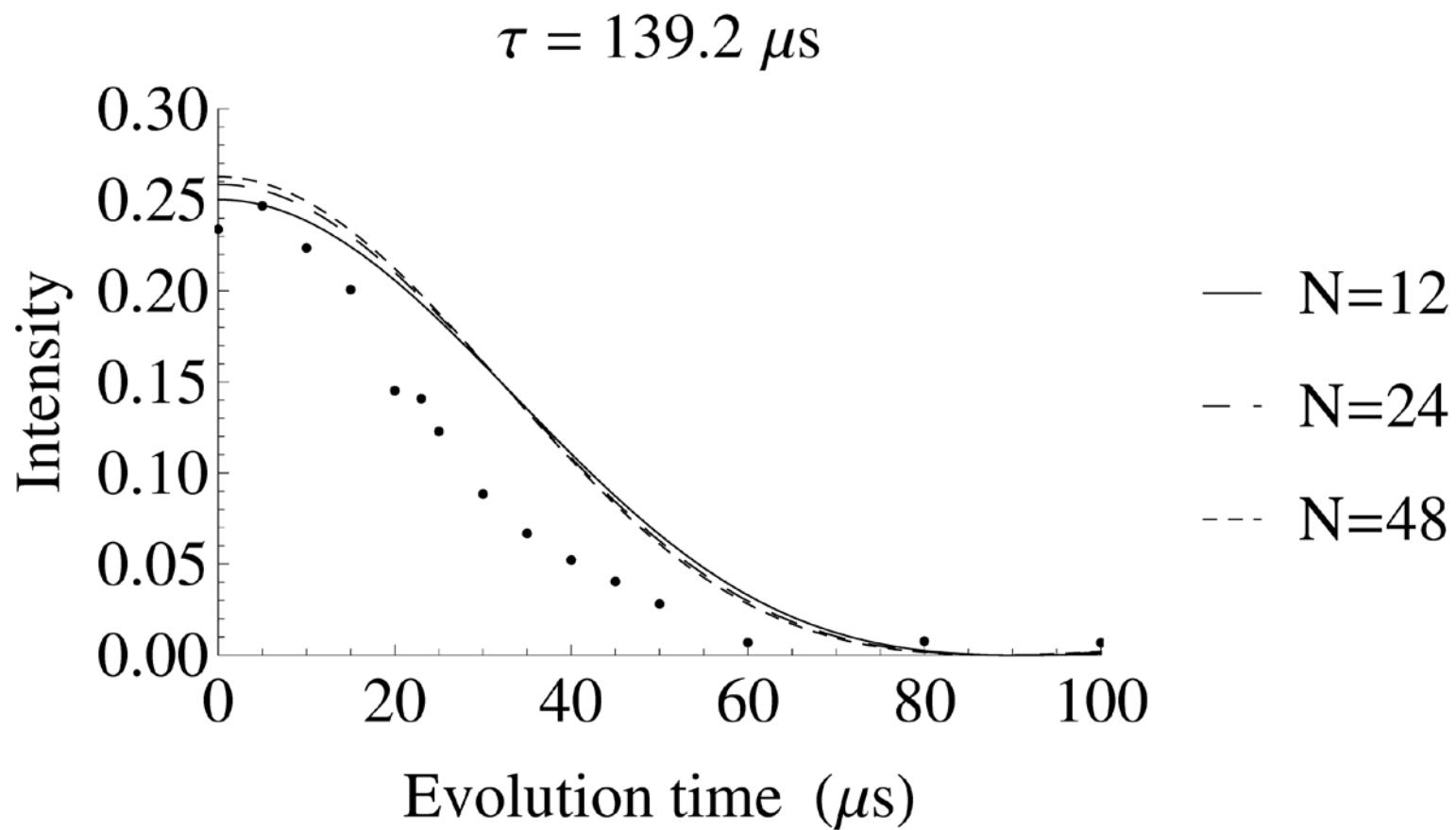


# RELAXATION OF THE MQ NMR COHERENCES OF THE SECOND ORDER

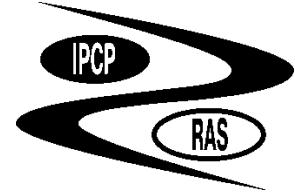
$$F_2(\tau, t) = -\frac{4}{N(N+1)^2} \sum_{l=1,3,\dots,N-1} \prod_{\substack{l'=2,4,\dots,N \\ m \neq l, m \neq l'}} \cos[(D_{ml} + D_{ml'})t] \cdot \left( \sum_k \sin kl \sin kl' e^{-2iD\tau \cos k} \right)^2$$

$$M_2(\tau) = -\frac{1}{G_2(\tau)} \frac{d^2 F_2(\tau, t)}{dt^2} \Bigg|_{t=0}$$

$$S(t) = e^{-\frac{1}{2} M_2(\tau) t^2} \quad t_e = \sqrt{\frac{2}{M_2(\tau)}}$$



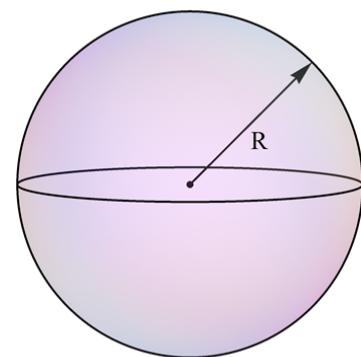
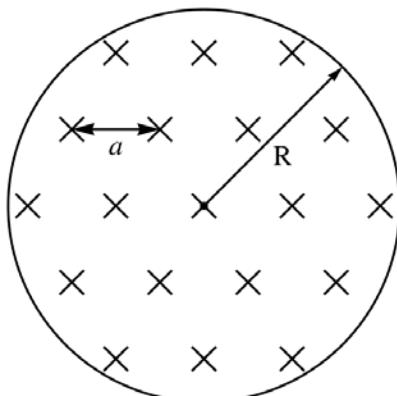
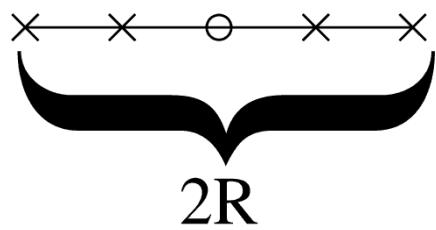
The experimental and theoretical intensities of the MQ NMR coherence of the second order on the evolution period  $t$  for the duration of the preparation period  $\tau=139.2 \mu\text{s}$ . The lines are the theoretical plots for various  $N$ .



$$\omega_d \tau = 1$$

$$\omega_d = \frac{\gamma^2 \hbar}{R^3}$$

$$V_d = \frac{\pi^{D/2} R^D}{\Gamma(D/2 + 1)}$$

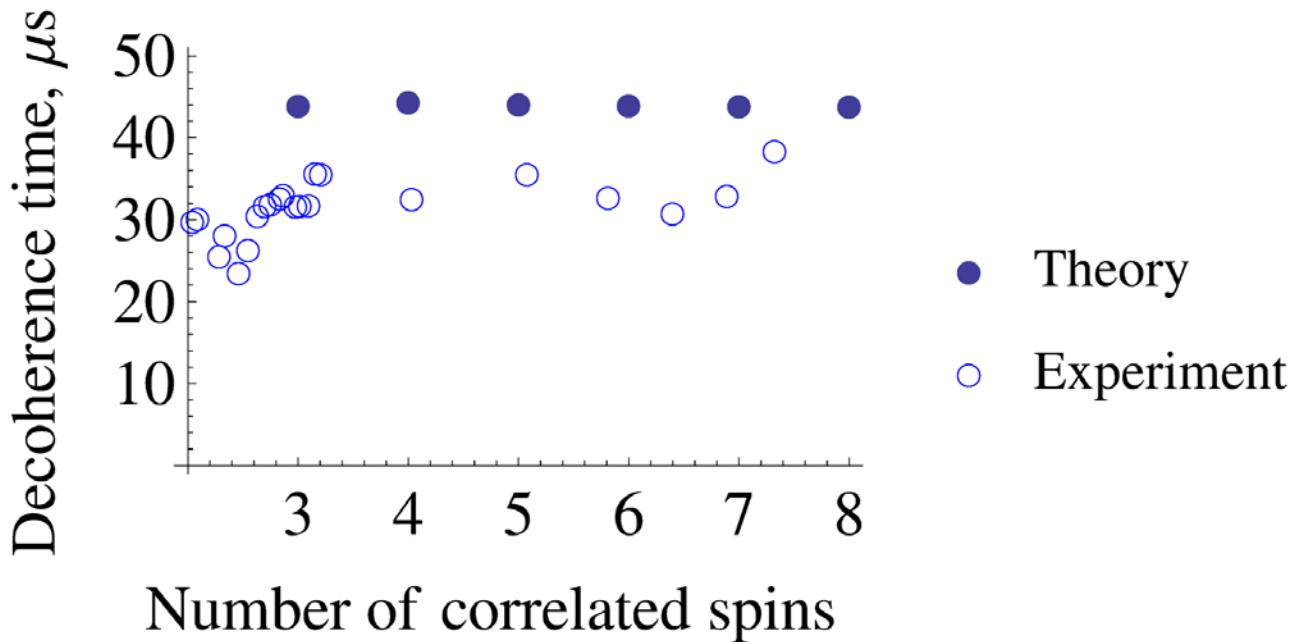


$$V_1 = 2R$$

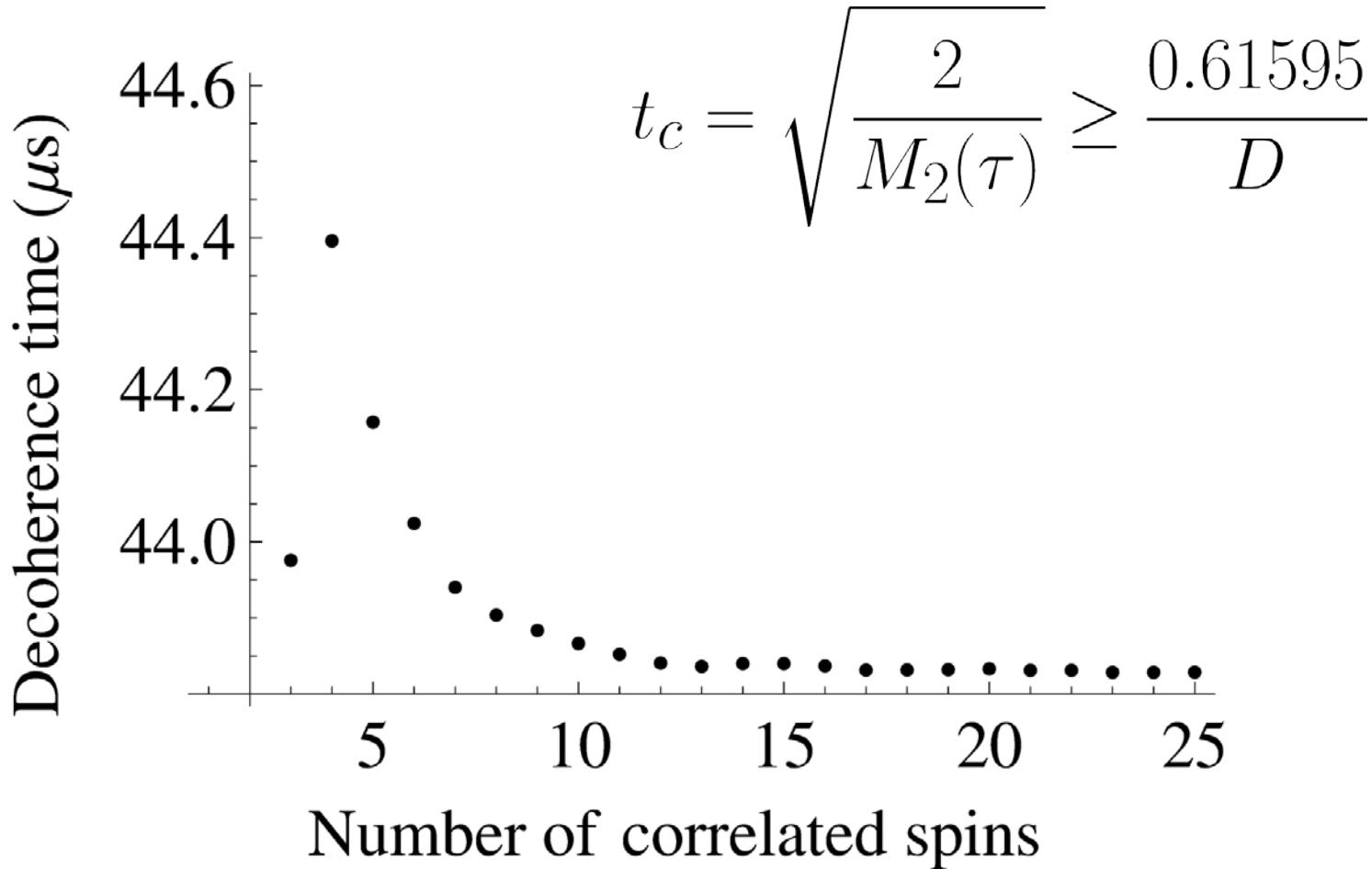
$$V_2 = \pi R^2$$

$$V_3 = \frac{4}{3}\pi R^3$$

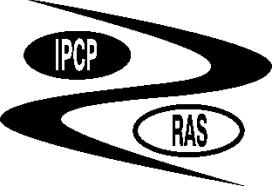
$$N = \frac{V_d}{a^D} = \frac{\pi^{D/2} R^D}{a^D \Gamma(D/2 + 1)} = \underbrace{\frac{\pi^{D/2} (\gamma^2 \hbar)^{D/3}}{a^D \Gamma(D/2 + 1)}}_k \tau^{D/3}$$



The experimental and theoretical times of dipolar relaxation of the MQ NMR coherence of the second order versus the cluster size.

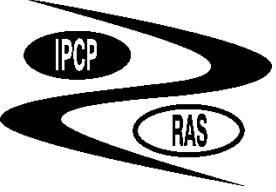


The theoretical times of dipolar relaxation of the MQ NMR coherence of the second order versus the cluster size, for larger cluster sizes.



# Conclusions

- A theory of the dipolar relaxation of multiple quantum NMR coherences is developed for one dimensional spin systems.
- It is shown that relaxation of the multiple quantum NMR coherence of the second order can be considered as a model for investigations of decoherence processes
- The decoherence (relaxation) time of the MQ NMR coherence of the second order decreases only slowly when the cluster size increased in our model.
- A lower bound of decoherence time does not depend on the number of qubits in correlated clusters.



Thank you for  
attention!

